A Positron Physics Program for the Jefferson Laboratory

 e^+ @ JLab 12 GeV e^+ @ JLEIC

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and the JLab Positron Working Group

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⁴Istituto Nazionale di Fisica Nucleare, Genova, Italia

⁵Università di Genova, Genova, Italia

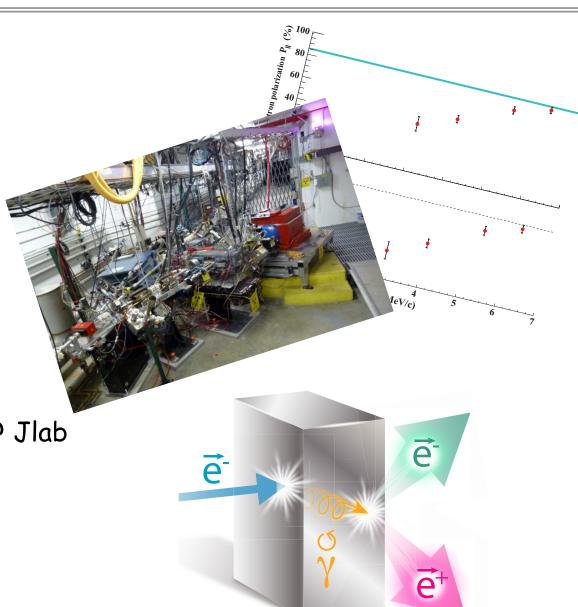
⁶Argonne National Laboratory, Argonne, IL, USA

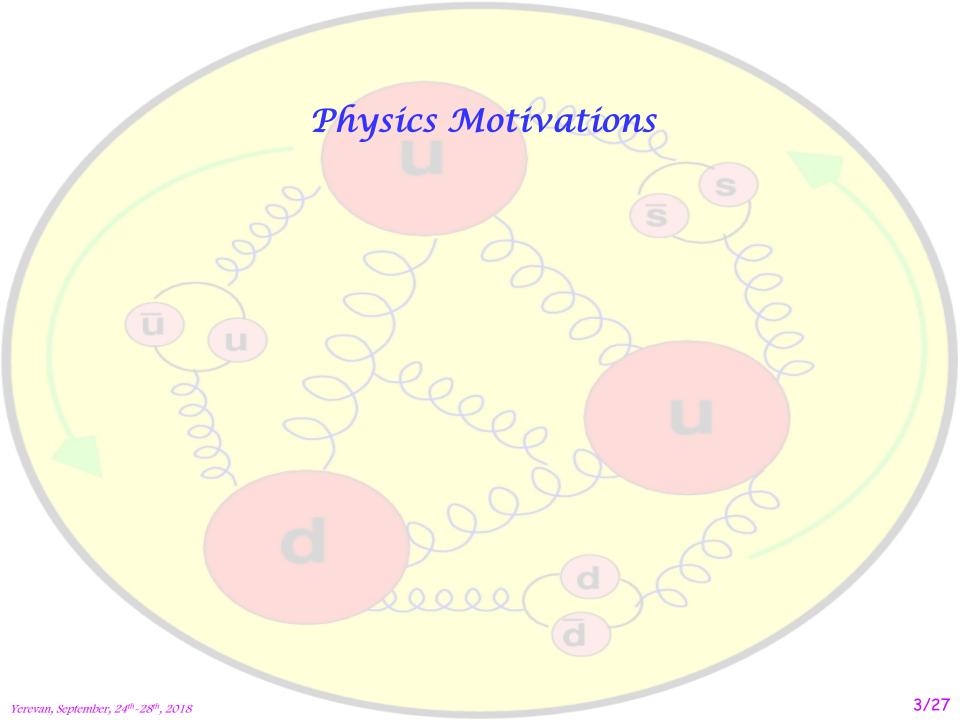
⁷University of Connecticut, CT, USA

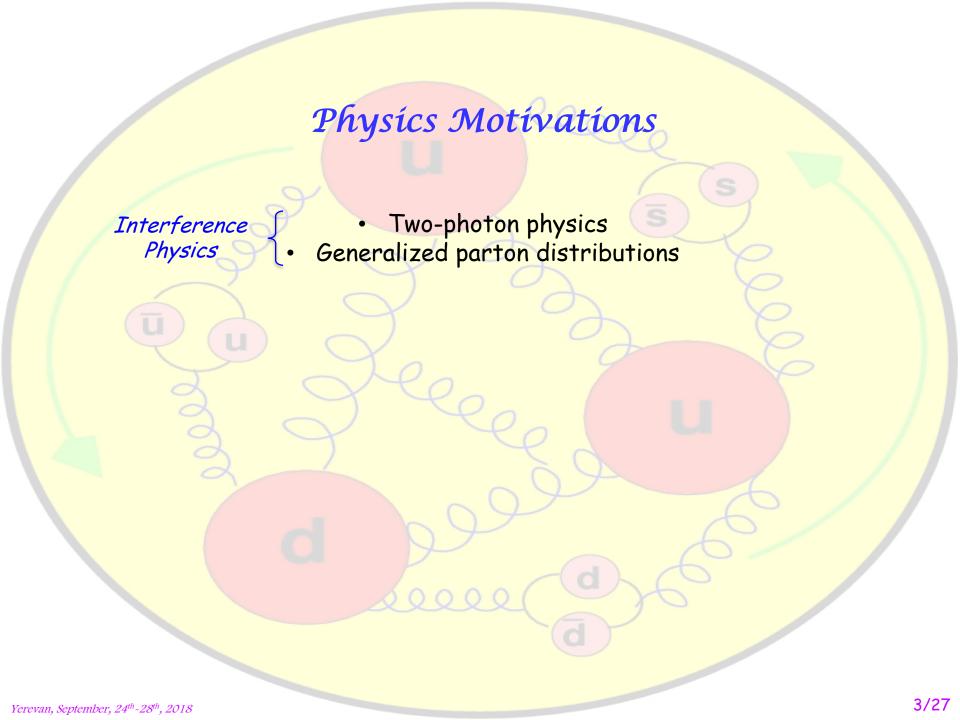
⁸University of North Georgia, Dahlonega, GA, USA



- (i) Physics motivations
- (ii) Interference physics
- (iii) Structure functions
- (iv) Sandard Model tests
- (v) Polarized positrons @ Jlab
- (vi) Summary







Physics Motivations

Interference Physics

- Two-photon physics
- Generalized parton distributions

Structure Functions

- Neutral and charged DIS
 - Charm production
 - Pion and kaon structure

Physics Motivations

Interference Physics

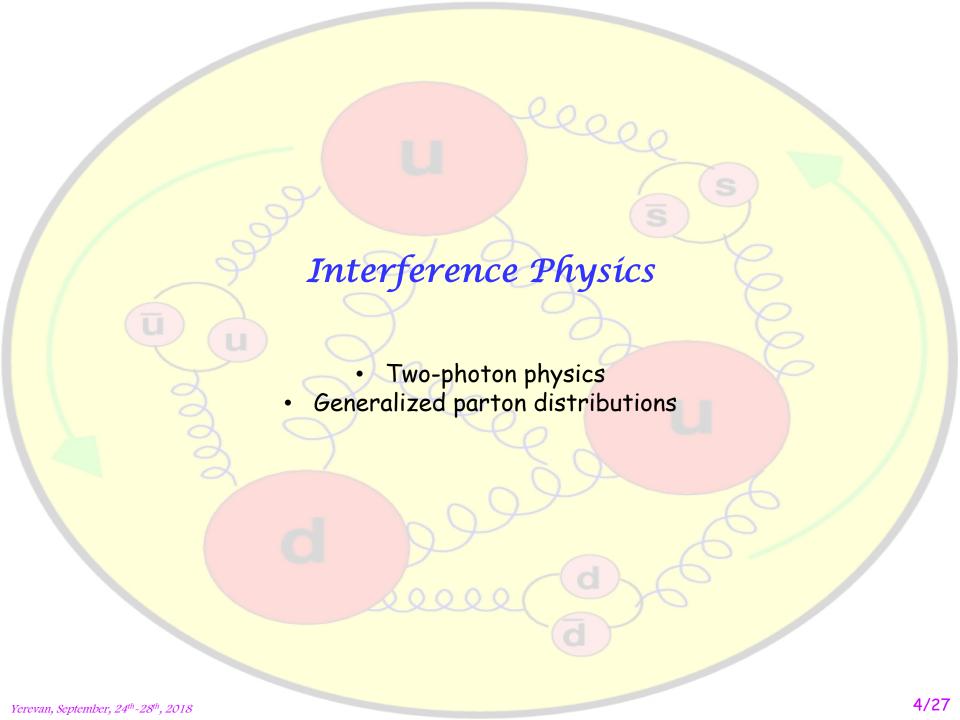
- Two-photon physics
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Structure Functions

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Standard Model Tests

- Charge conjugation violation
 - Right-handed W-bosons
 - Dark photon
 - Leptoquarks, leptogluons

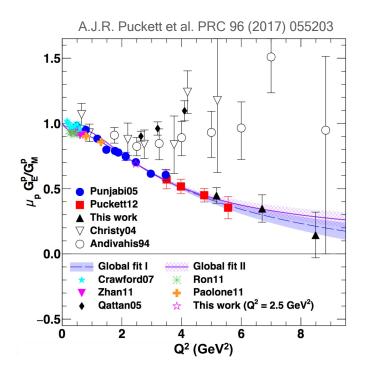




Electromagnetic Form Factors

P.A.M. Guichon, M. Vanderhaeghen, PRL 91 (2003) 142303 P.G. Blunden, W. Melnitchouk, J.A. Tjon, PRL 91 (2003) 142304

Measurements of polarization transfer observables in electron elastic scattering off protons question the validity of the 1γ exchange approximation of the electromagnetic interaction.



Within the 2γ exchange hypothesis, the electromagnetic structure of the nucleon may be parametrized by 3 generalized form factors, corresponding to 8 unknow quantities.

$$\widetilde{G}_{M} = G_{M}(Q^{2}) + e_{l} \delta \widetilde{G}_{M}(Q^{2}, \varepsilon)$$

$$\widetilde{G}_{E} = G_{E}(Q^{2}) + e_{l} \delta \widetilde{G}_{E}(Q^{2}, \varepsilon)$$

$$\widetilde{F}_{3} = e_{l} \delta \widetilde{F}_{3}(Q^{2}, \varepsilon)$$



Experimental Observables

M.P. Rekalo, E. Tomasi Gustafsson, NPA 742 (2004) 322 C. Carlson, M. Vanderhaeghen, ARNPS 57 (2007) 171

> Unpolarized e[±] elastic scattering and polarization transfert observables off the nucleon involve up to 5 unknown quantities.

$$\sigma_R = G_M^2 + \frac{\varepsilon}{\tau} G_E^2 \pm 2G_M \Re \Big[f_0 \Big(\delta \widetilde{G}_M, \delta \widetilde{F}_3 \Big) \Big] \pm 2 \frac{\varepsilon}{\tau} G_E \Re \Big[f_1 \Big(\delta \widetilde{G}_E, \delta \widetilde{F}_3 \Big) \Big]$$

$$\sigma_R P_t = -\sqrt{\frac{2\varepsilon(1-\varepsilon)}{\tau}} \left(G_E G_M \pm G_E \Re \Big[\delta \widetilde{G}_M \Big] \pm G_M \Re \Big[f_1 \Big(\delta \widetilde{G}_E, \delta \widetilde{F}_3 \Big) \Big] \right)$$

$$\sigma_R P_l = \sqrt{1-\varepsilon^2} \left(G_M^2 \pm 2G_M \Re \Big[f_2 \Big(\delta \widetilde{G}_M, \delta \widetilde{F}_3 \Big) \Big] \right)$$
Polarization Transfert

5 unknown contributions for 6 independent observables



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$$\sigma_R P_t = -\sqrt{\frac{2\mathcal{E}(1-\mathcal{E})}{\tau}} \left(G_E G_M \pm G_E \Re \Big[\delta \widetilde{G}_M \Big] \pm G_M \Re \Big[f_1 \Big(\delta \widetilde{G}_E, \delta \widetilde{F}_3 \Big) \Big] \right)$$

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Polarization Transfert

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Combining polarized electrons and positrons allows a model independent separation of the electromagnetic form factors of the nucleon.



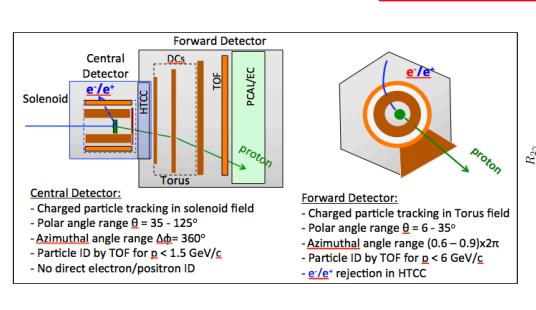
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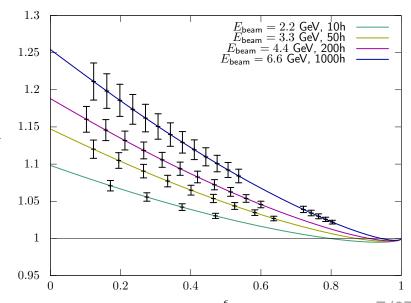
e+@JLab12GeV - J. Grames, E. Voutier et al.

J. Arrington, J. Bernauer, V. Burkert, A. Schmidt et al.

 \triangleright A modified CLAS12 hosting an electromagnetic calorimeter in place of the Central Neutron Detector would allow to map out the 2γ -effects in the (Q^2, ϵ) space, providing a conclusive answer about the relevance of 2γ -effects.

$$R_{2\gamma} = \frac{\sigma_{e^+}}{\sigma_{e^-}} \approx 1 + \delta_{2\gamma}$$





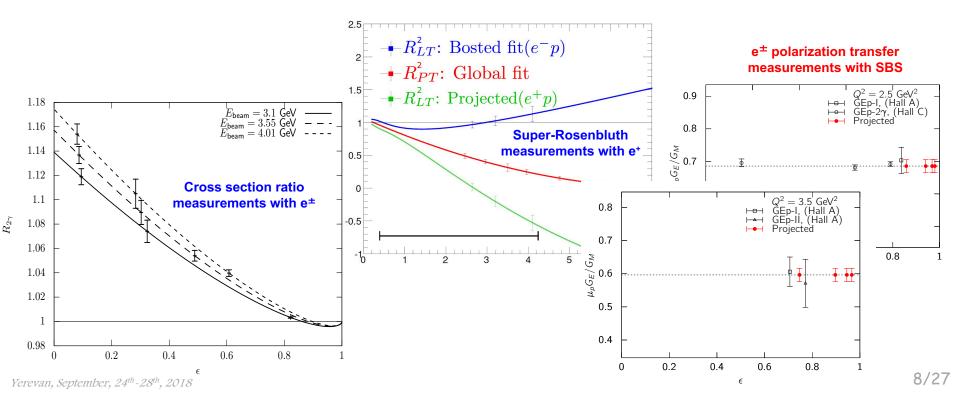


LOI12-18-004

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ightharpoonup High impact measurements are also feasible in Hall A & C to sign the existence of 2γ -effects, better control systematics with the Super-Rosenbluth technique, and to provide unique polarization transfer data.

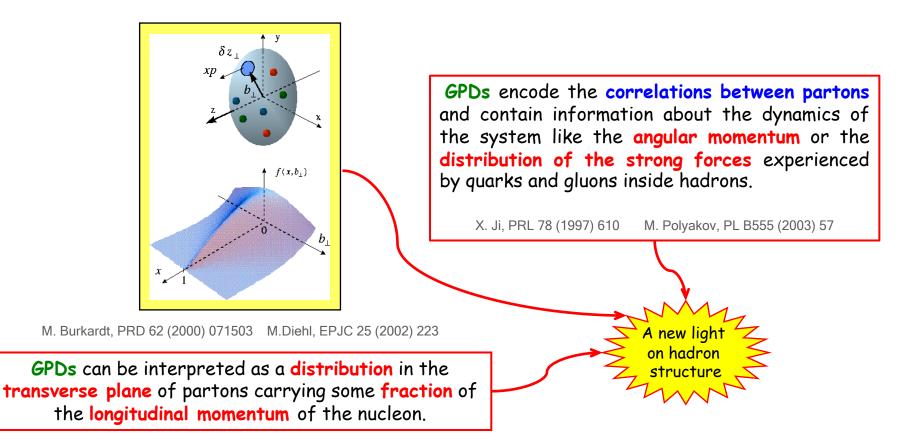




Parton Imaging

D. Müller, D. Robaschik, B. Geyer, F.M. Dittes, J. Horejsi, FP 42 (1994) 101 X. Ji, PRD 55 (1997) 7114 A. Radyushkin, PRD 56 (1997) 5524

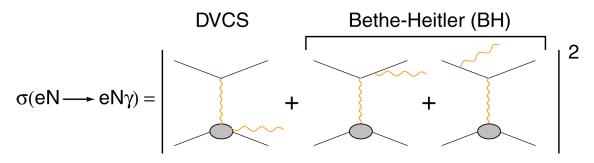
GPDs parameterize the partonic structure of hadrons and offer the unprecedented possibility to access the spatial distribution of partons.



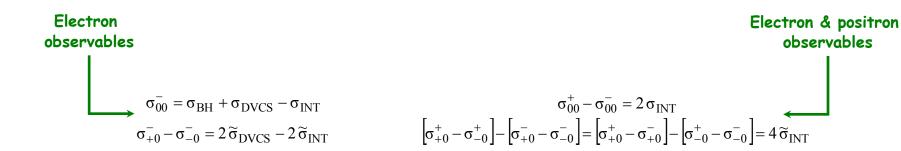


N(e,e'yN) Differential Cross Section

M. Diehl at the CLAS12 European Workshop, Genova, February 25-28, 2009



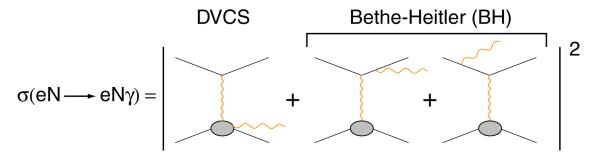
$$\sigma_{P0}^{e} = \sigma_{BH} + \sigma_{DVCS} + P_{l} \widetilde{\sigma}_{DVCS} + e_{l} (\sigma_{INT} + P_{l} \widetilde{\sigma}_{INT})$$



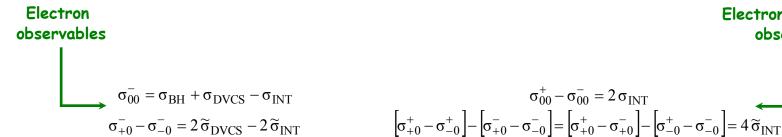


N(e,e'\gammaN) Differential Cross Section

M. Diehl at the CLAS12 European Workshop, Genova, February 25-28, 2009



$$\sigma_{PS}^{e} = \sigma_{P0}^{e} + S \left[P_{l} \Delta \sigma_{BH} + \left(\Delta \widetilde{\sigma}_{DVCS} + P_{l} \Delta \sigma_{DVCS} \right) + e_{l} \left(\Delta \widetilde{\sigma}_{INT} + P_{l} \Delta \sigma_{INT} \right) \right]$$

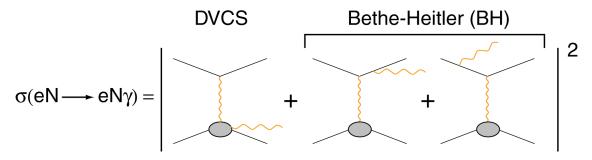


Electron & positron observables



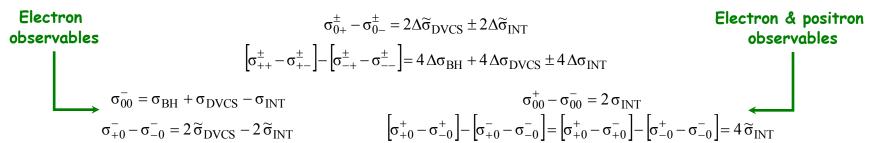
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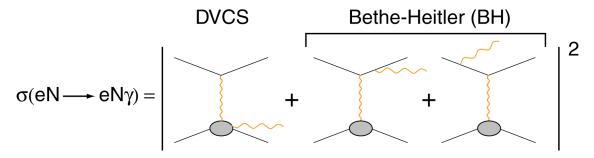
Additional observables





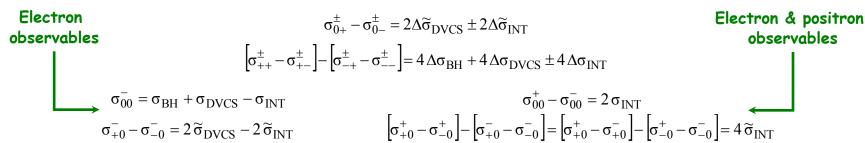
N(e,e'\gammaN) Differential Cross Section

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Additional observables

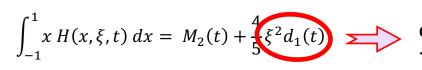


Polarized electrons and positrons allow to separate the unknown amplitudes of the cross section for electro-production of photons.



Nucleon Internal Pressure

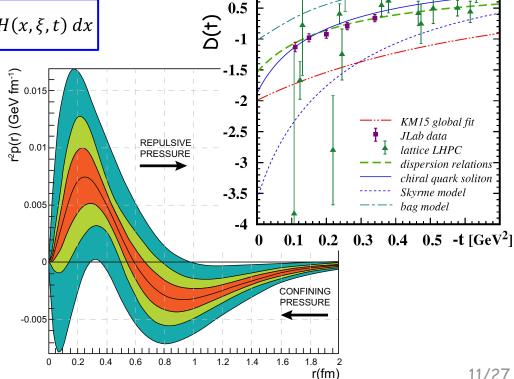
V. Burkert, L. Elouadrhiri, F.-X. Girod, Nature 557 (2018) 396 M.V. Polyakov, P. Schweitzer, Int. J. Mod. Phys. A33 (2018) 1830025



The 2nd order Mellin moment of GPDs allow to access the pressure distribution inside hadrons through the skewness dependency of GPDs ... DDVCS.

CFF
$$\mathcal{H}(\xi,t) = \int_{-1}^{1} \left[\frac{1}{\xi - x - i\epsilon} - \frac{1}{\xi + x - i\epsilon} \right] H(x,\xi,t) dx$$

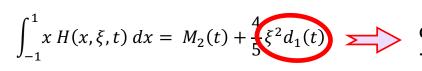
$$\Re[\mathcal{H}(\xi,t)] \stackrel{\text{LO}}{=} D(t) \\ + \mathcal{P}\left\{ \int_{-1}^{1} \left[\frac{1}{\xi - x} - \frac{1}{\xi + x} \right] \Im[\mathcal{H}(x,t)] dx \right\} \stackrel{\text{D}}{=} 0.015 \\ D(t) = \frac{1}{2} \int_{-1}^{1} \frac{D(z,t)}{1 - z} dz \\ D(z,t) = (1 - z^2) \left[d_1(t) C_1^{3/2}(z) + \dots \right]$$





Nucleon Internal Pressure

V. Burkert, L. Elouadrhiri, F.-X. Girod, Nature 557 (2018) 396 M.V. Polyakov, P. Schweitzer, Int. J. Mod. Phys. A33 (2018) 1830025



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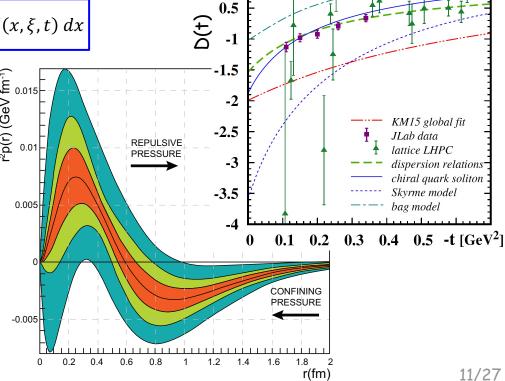
DDVCS.

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$$D(z,t) = (1-z^2) \left[d_1(t) C_1^{3/2}(z) + \dots \right]$$

Real part of Compton form factors.





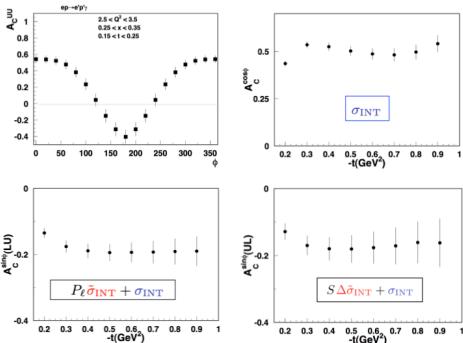
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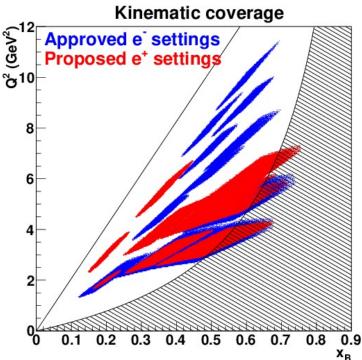
e+@ JLab 12 GeV - J. Grames, E. Voutier et al.

V. Burkert, L. Elouadrhiri, F.-X. Girod, C. Muñoz Camacho et al.

> It is proposed to measure polarized and unpolarized beam charge asymmetries off protons and neutrons (CLAS12), and unpolarized p-DVCS cross section with positrons (HMS+NPS).





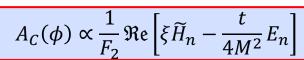




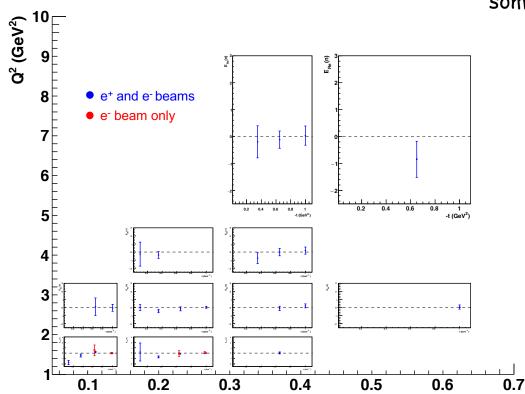
LOI12-18-004

e+@ JLab 12 GeV - J. Grames, E. Voutier et al.

S. Niccolai, E. Voutier et al.



The BCA on the neutron accesses the real part of the CFF E, and is sensitive to H at some kinematics.

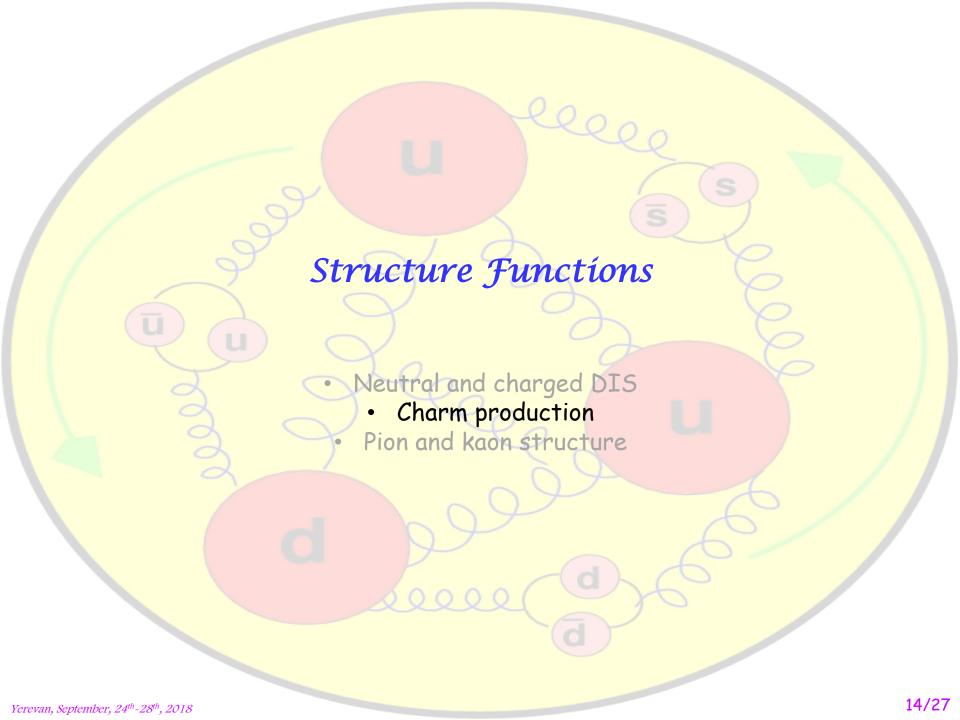


$$\Delta_{LU}^{C} = \frac{(d^{4}\sigma_{+0}^{+} - d^{4}\sigma_{-0}^{+}) - (d^{4}\sigma_{+0}^{-} - d^{4}\sigma_{-0}^{-})}{d^{4}\sigma_{+0}^{+} + d^{4}\sigma_{-0}^{+} + d^{4}\sigma_{-0}^{-} + d^{4}\sigma_{-0}^{-}}$$

$$\Sigma_{LU}^{C} = \frac{(d^{4}\sigma_{+0}^{+} - d^{4}\sigma_{-0}^{+}) + (d^{4}\sigma_{+0}^{-} - d^{4}\sigma_{-0}^{-})}{d^{4}\sigma_{+0}^{+} + d^{4}\sigma_{-0}^{+} + d^{4}\sigma_{-0}^{-} + d^{4}\sigma_{-0}^{-}}$$

$$A_{LU} = \frac{\Sigma_{LU}^C}{1 - A_C} - \frac{\Delta_{LU}^C}{1 - A_C}$$

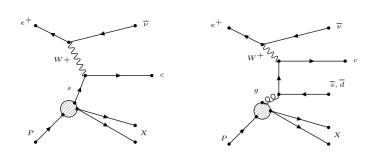
Polarized positron beams provide a different access to the imaginary part, and probe the importance of higher twists.



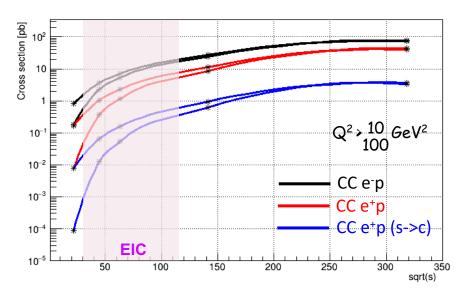


Strangeness Tagging

The modest center-of-mass energy at the EIC can be compensated by the high luminosity and polarization degree of the lepton beam, to help for precise measurements of the small charge current cross section.



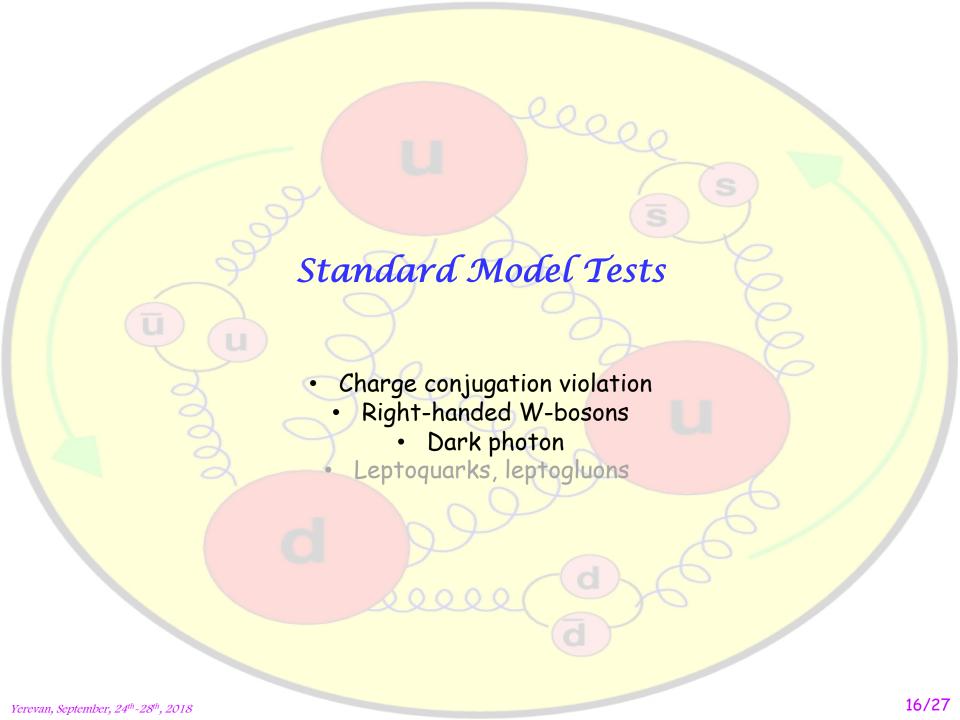
- Charm production via charge current exchange preferentially couples to the strange content of the nucleon.
- The boson-gluon-fusion mechanism is a source of background, manageable with good PID and vertex reconstruction.



$$e^+ + p \rightarrow \overline{v_e} + c + X$$

 $e^- + p \rightarrow v_e + \overline{c} + X$

An integrated luminosity of 10 fb⁻¹/year would provide ~1500 events/year.



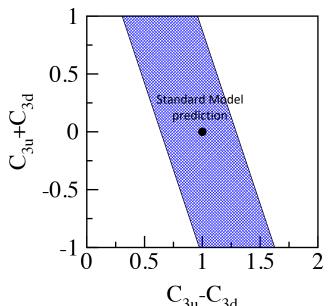


Weak Neutral Current Couplings

S.M. Berman, J.R. Primack, Phys. Rev. D 9 (1974) 2171 X. Zheng, (JPos09) AIP CP 1160 (2009) 160

> The comparison of polarized electron and polarized position scatterings provides access to the charge conjugation-violation coupling C_{3q} from the interference between the weak neutral and electromagnetic currents.

$$\mathcal{L} = \frac{G_F}{\sqrt{2}} \sum_{q} \left[C_{1q} \, \bar{\ell} \gamma^{\mu} \gamma_5 \ell \bar{q} \gamma_{\mu} q + C_{2q} \, \bar{\ell} \gamma^{\mu} \ell \bar{q} \gamma_{\mu} \gamma_5 q + C_{3q} \, \bar{\ell} \gamma^{\mu} \gamma_5 \ell \bar{q} \gamma_{\mu} \gamma_5 q \right] g_A^e g_A^q$$



The combination of C_{3q} couplings is **poorly known**; have only been measured at CERN using **muon** and **anti-muon** beams on a carbon target.

$$0.81 (2C_{2u}-C_{2d}) + 2C_{3u}-C_{3d} = 1.53 \pm 0.45$$

D. Wang et al. Phys. Rev. C 91 (2015) 045506

 $2C_{3u} - C_{3d} = 1.65 \pm 0.45$

 C_{3q} are known only within 30%

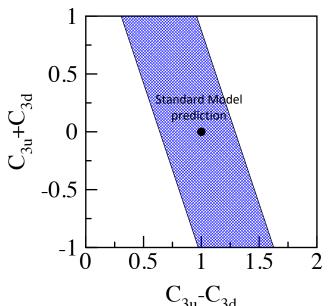


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$$A^{e_L^- - e_R^+} = \frac{d\sigma(e_L^- N) - d\sigma(e_R^+ N)}{d\sigma(e_L^- N) + d\sigma(e_R^+ N)}$$

$$A_d^{e_L^- - e_R^+} = \left(\frac{3G_F Q^2}{2\pi\alpha\sqrt{2}}\right) \frac{y(2-y)}{2} \frac{2C_{2u} - C_{2d} + 2C_{3u} - C_{3d}}{5} \frac{u_V + d_V}{u + \bar{u} + d + \bar{d}}$$

$$A_d^{e_L^- - e_R^+} \approx 108 \frac{y(2-y)}{2} (2C_{3u} - C_{3d}) Q^2 \text{ ppm}$$

There exists a unique opportunity for polarized electron and positron beams at EIC to improve the C_{3q} knowledge.



Polarized Charge Current Cross Section

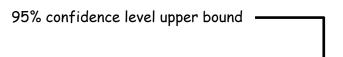
Y. Furletova, S. Mantry, (JPos17) AIP CP 1970 (2018) 030005

> The polarization dependence of the charge current cross section can be measured to potentially reveal deviations from the Standard Model prediction.

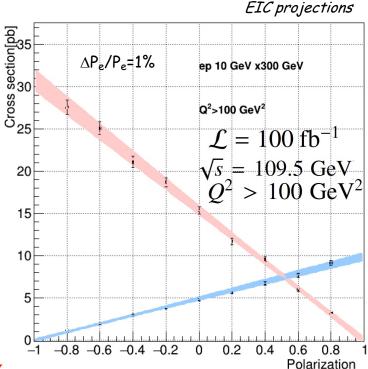
$$\frac{d^2 \sigma_{SM}^{e^+ p}}{dx \ dQ^2} = (1 + P_e) \frac{G_F^2}{2\pi} \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \left[\bar{u} + \bar{c} + (1 - y^2)(d + s) \right]$$

$$\frac{d^2 \sigma_{SM}^{e^- p}}{dx \ dQ^2} = (1 - P_e) \frac{G_F^2}{2\pi} \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \left[u + c + (1 - y^2)(\bar{d} + \bar{s}) \right]$$

$$\sigma^{e^{\pm}p} (P_e = \mp 1) \neq 0 \rightarrow \begin{array}{c} \text{Limit on the mass of} \\ \text{right-handed W-boson} \end{array}$$



$$\sigma^{e^{\pm}p} (P_e = \pm 1) < 0.0776 \text{ pb } \rightarrow M_R \ge 285 \text{ GeV}$$



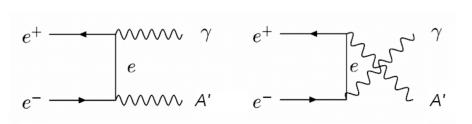
18/27 Yerevan, September, 24th ~28th, 2018



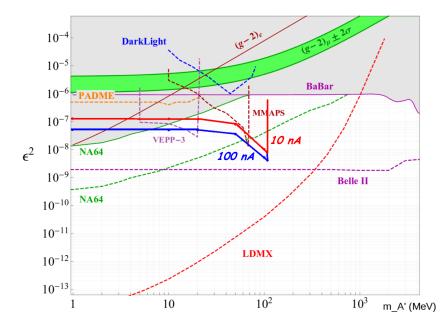
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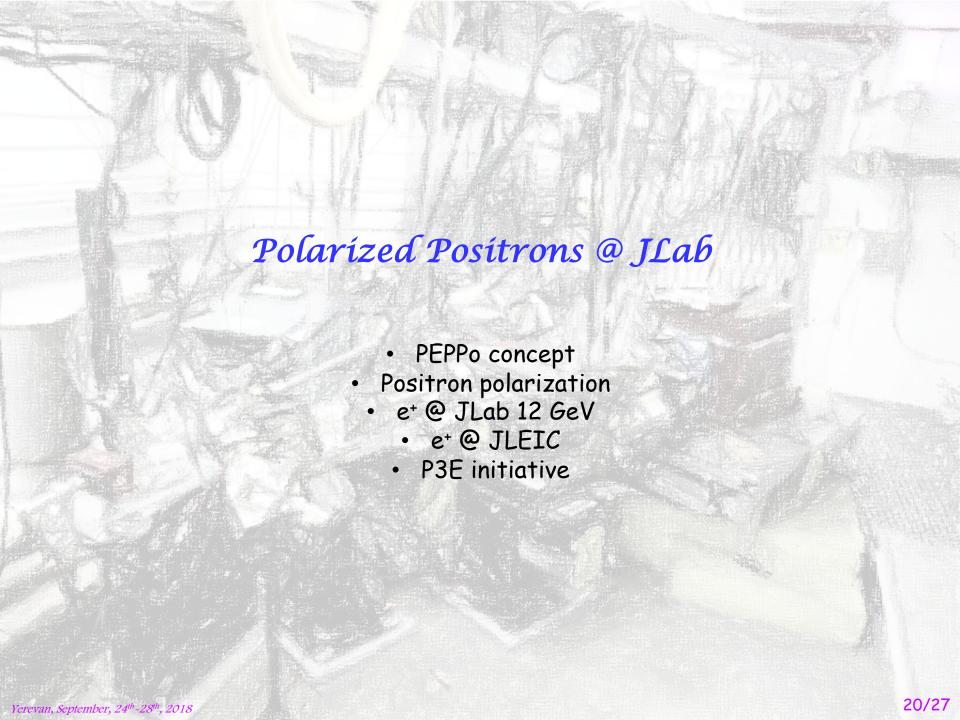


The e⁺e⁻-annihilation</sup> offers an alternative way to probe the existence of a dark photon by measuring the photon spectra in the A' invisible decay chanel.



P. Valente, J. Alexander, (JPos17) AIP CP 1970 (2018) 020007 L. Marsicano, (JPos17) AIP CP 1970 (2018) 020008

180 days data taking with
a high energy positron beam
would extend the
A' mass reach up to 100 MeV
and would improve by a
factor 10 the coupling strength sensitivity.



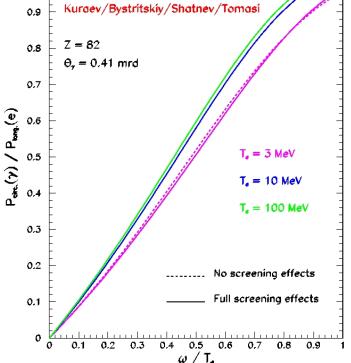


Bremsstrahlung and Pair Creation

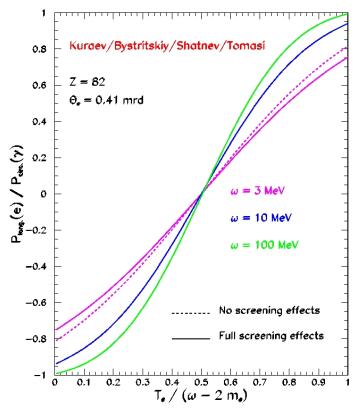
H. Olsen, L. Maximon, PR 114 (1959) 887 E.A. Kuraev, Y.M. Bystritskiy, M. Shatnev, E. Tomasi-Gustafsson, PRC 81 (2010) 055208

Finite lepton mass calculations of the bremsstrahlung and pair creation processes predict very efficient polarization transfers.

BREMSSTRAHLUNG Kuraev/Bystritskiy/Shatnev/Tomasi



PAIR CREATION



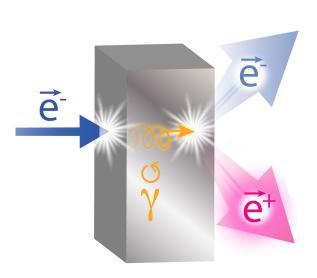
21/27

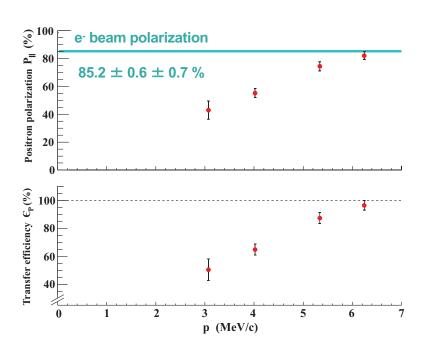


PEPPo Technique

(PEPPo Collaboration) D. Abbott et al., Phys. Rev. Lett. 116 (2016) 214801

> PEPPo demonstrated efficient polarization transfer from 8.2 MeV/c electrons to positrons, expanding polarized positron capabilities from GeV to MeV accelerators.





The PEPPo technique can achieve up to 100% transfer of the electron.

20

Yerevan, September, 24th - 28th, 2018

60

E_Pos (MeV)

80

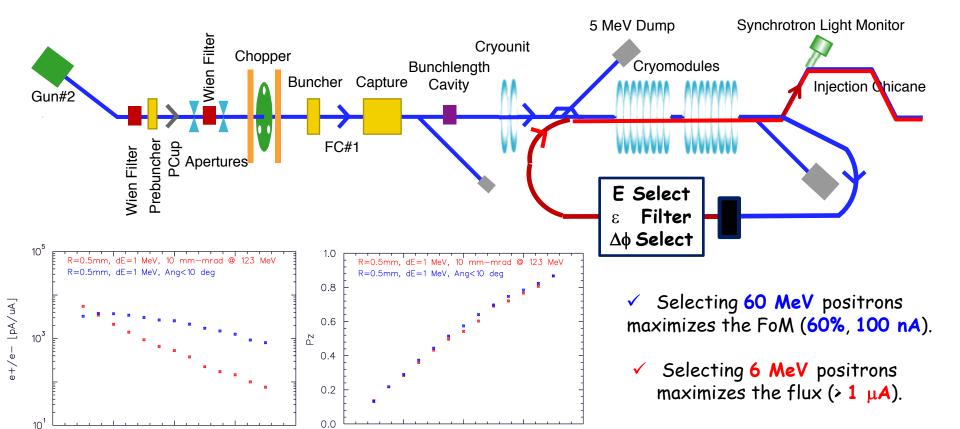


LOI12-18-004

e+@JLab12GeV - J. Grames, E. Voutier et al.

L. Cardman et al. (JPos17) AIP CP 1970 (2018) 050001

Positrons would be created at the CEBAF injector, using the 123 MeV electron beam.



60

E_Pos (MeV)

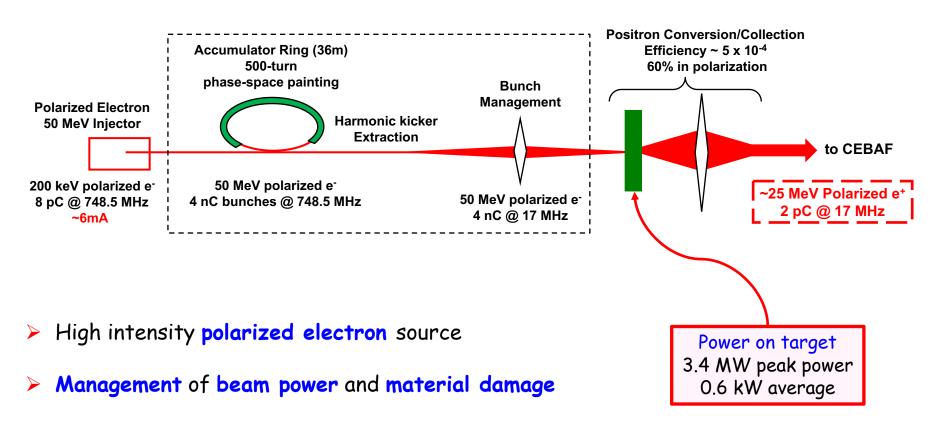
80

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Positron Production Scheme

F. Lin, J. Grames, J. Guo, V. Morozov, Y. Zhang, (JPos17) AIP CP 1970 (2018) 050005



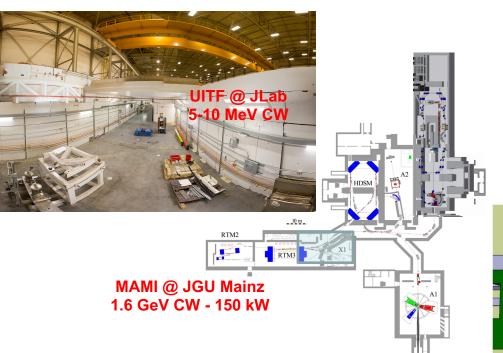
Optimized collection system



Technical Concept Validation

K. Aulenbacher, N. Berger, I. Chaivoska, J. Enders, J. Grames, G. Moortgat-Pick, S. Riemann, E. Voutier et al.

Technical solutions for each component of the positron source must be experimentally evaluated and optimized.







The P3E (Polarized Electrons, Positrons, and Polarimetry) initiative, part of the STRONG-2020 proposal (H2020-INFRAIA-2018-2019 call), intends to address these issues.



e+ @ JLab

Polarized and unpolarized positron beams provide a unique opportunity to enhance the physics reach of JLab and a future EIC.

Interferences, neutral and charged currents, test of the Standard Model...

JPos17 proceedings, AIP CP 1970 (2018)

https://aip.scitation.org/toc/apc/1970/1?size=all&expanded=1970

LOI12-18-004 to JLab PAC46

« ... These measurements all have significant physics interest. The proposers should carefully evaluate feasibility and present the best case possible in a future proposal. The justification must be very strong to enable the significant changes needed in the accelerator, both in equipment and in schedule... »

Work is progressing to assess the quantitative impact of a polarized positron beam at JLab12 and EIC.

White Paper (planned 2019), Conceptual Design Report (goal 2020)



Polarízed e+ @ ALTO, MAMI, MESA, JLab, EIC, ILC, CEPC...

- The P3E european/US initiative has been proposed to evaluate/address some of the technical challenges of polarized positron beams driven by high intensity polarized electron beams.
 - □ The P3E initiative will also investigate the constraints/abilities for the production of low energy polarized positron beams (Material Science...) particularly in the context of the upgrade of the ALTO facility in Orsay.



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Consider contributing to this adventure!!

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Subscribe to pwg@jlab.org