

# QCD EVOLUTION

## Workshop

May 20 - 24, 2018  
Santa Fe, NM,  
Drury Plaza Hotel

### TOPICS:

- Hadron structure: theory & experiment
- Transverse momentum dependent distributions
- Generalized parton distributions
- Effective theories, SCET
- Lattice QCD
- Multi-parton interactions
- Resummation techniques
- Nuclear effects, small-x

### Organizing Committee

- |                     |                    |
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| Martha Constantinou | Anatoly Radyushkin |
| Leonard Gamberg     | Matt Sievert       |
| Chris Lee           | Ivan Vitev (Chair) |
| Duff Neill          | Shinsuke Yoshida   |

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[www.jlab.org/conferences/qcd-evolution2018](http://www.jlab.org/conferences/qcd-evolution2018)

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# First extraction of Transversity from a global analysis of lepton-hadron scattering and hadronic collisions



**Marco Radici**  
INFN - Pavia

in collaboration with  
A. Bacchetta (Univ. Pavia)



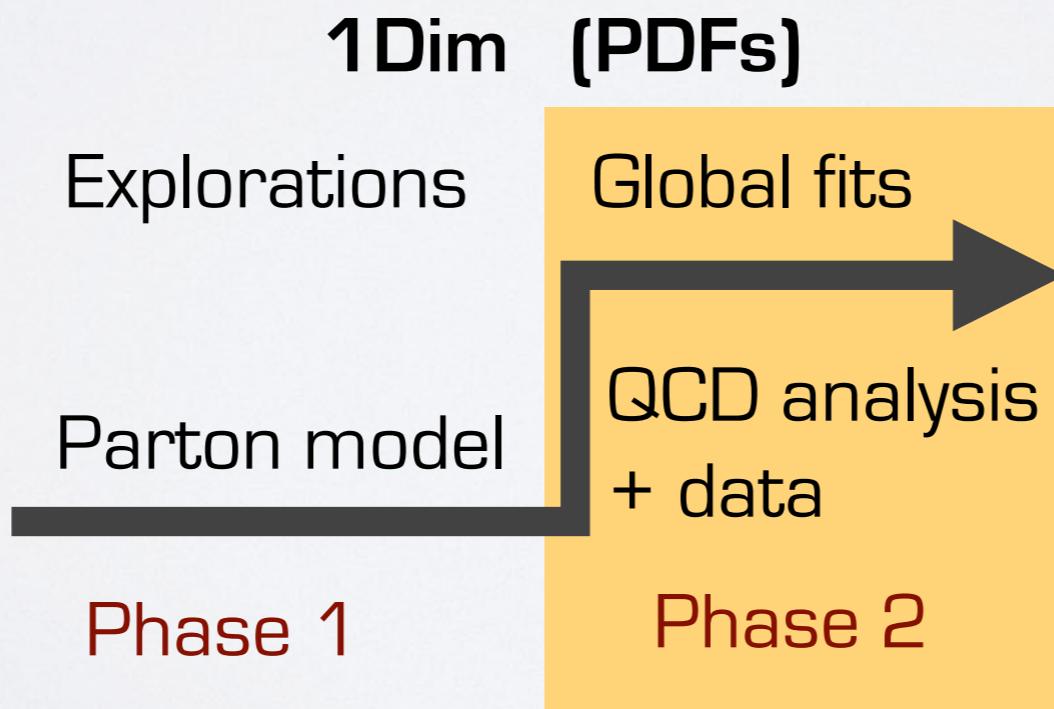
based on  
**P.R.L. 120** (2018) 192001  
arXiv:1802.05212  
**plus updates**

# a phase transition

		quark polarization		
		U	L	T
nucleon polarization	U	$f_1$		$h_{1^\perp}$
	L		$g_{1L}$	$h_{1L^\perp}$
	T	$f_{1T^\perp}$	$g_{1T}$	$h_1$ $h_{1T^\perp}$

chiral-odd  $\rightarrow$  SIDIS

first global fit  
(= lepton-hadron scatt.  
and hadron collisions)  
of PDF  $h_1$

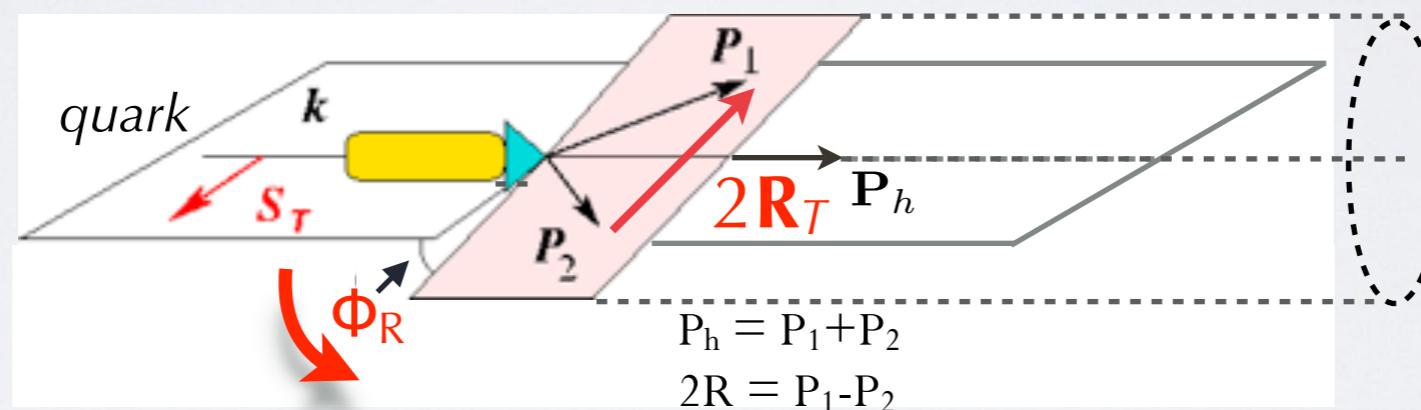


# 2-hadron-inclusive production

Collins, Heppelman, Ladinsky,  
N.P. **B420** (94)

$$R_T \ll Q \quad H_1^{\triangleleft}$$

↑  
 $M_h$



correlation  $S_T$  and  $R_T \rightarrow$  azimuthal asymmetry

invariant mass

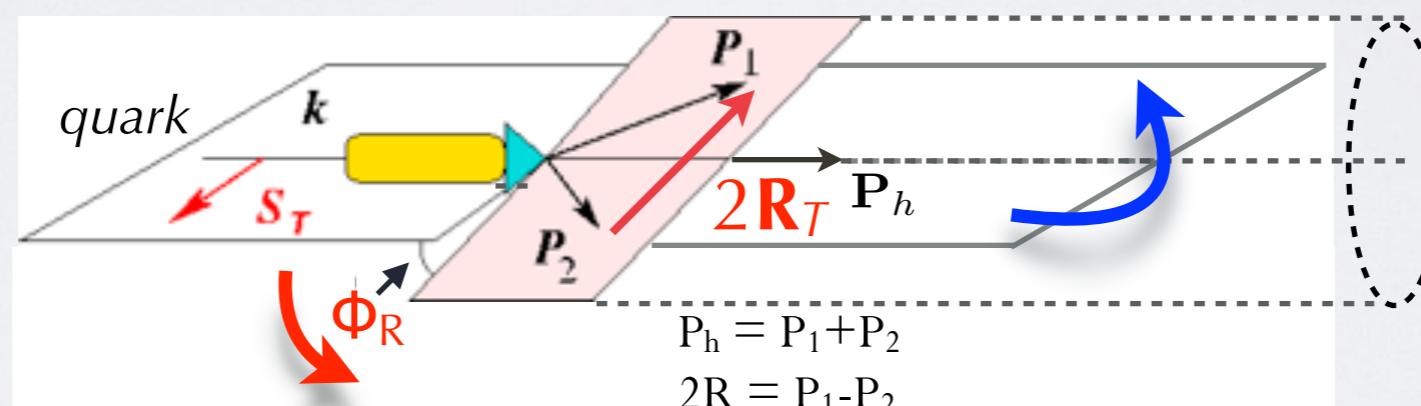
# 2-hadron-inclusive production

framework  
collinear  
factorization

Collins, Heppelman, Ladinsky,  
N.P. **B420** (94)

$$R_T \ll Q \quad H_1^{\triangleleft}$$

↑  
 $M_h$

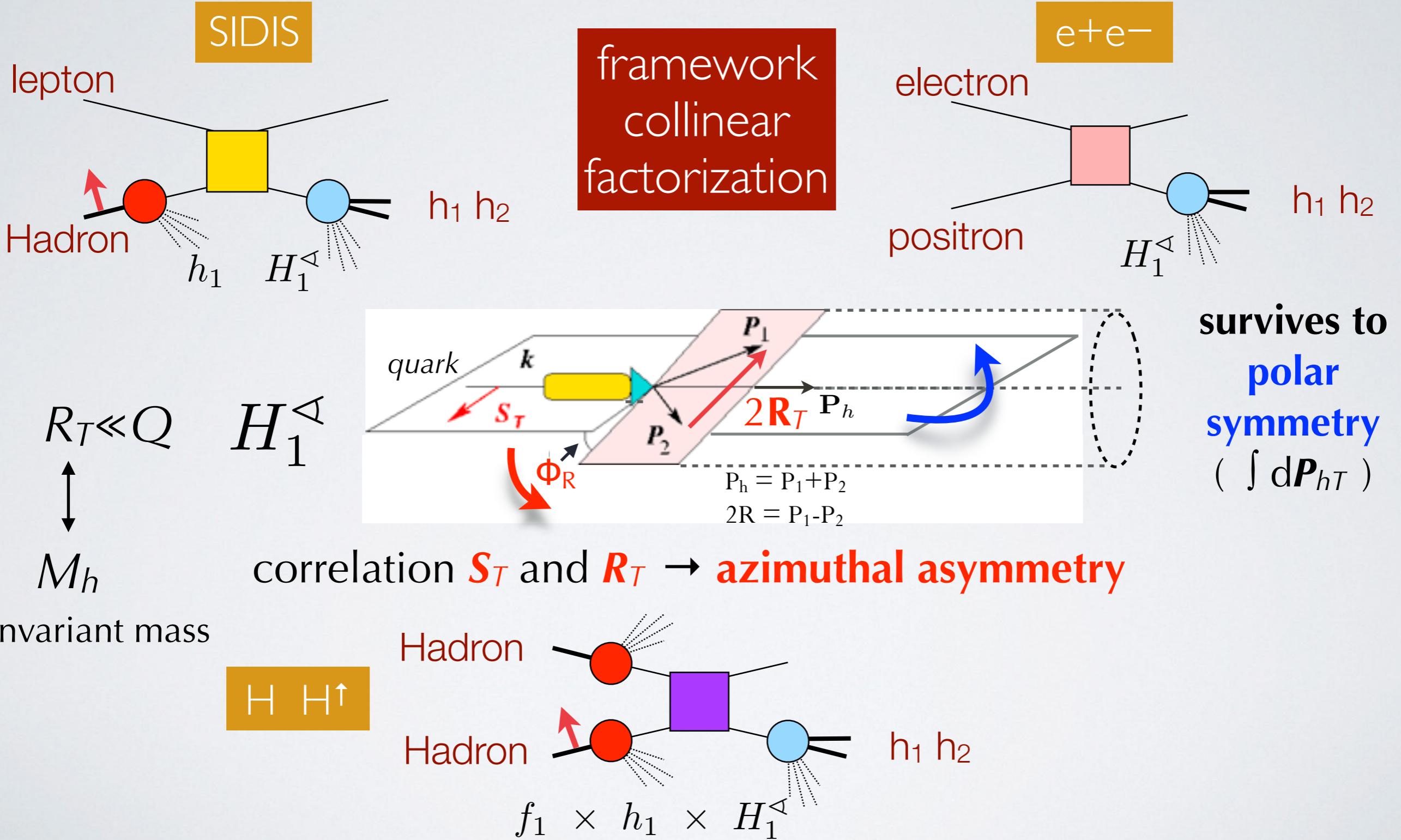


survives to  
polar  
symmetry  
(  $\int dP_{hT}$  )

correlation  $s_T$  and  $R_T \rightarrow$  azimuthal asymmetry

invariant mass

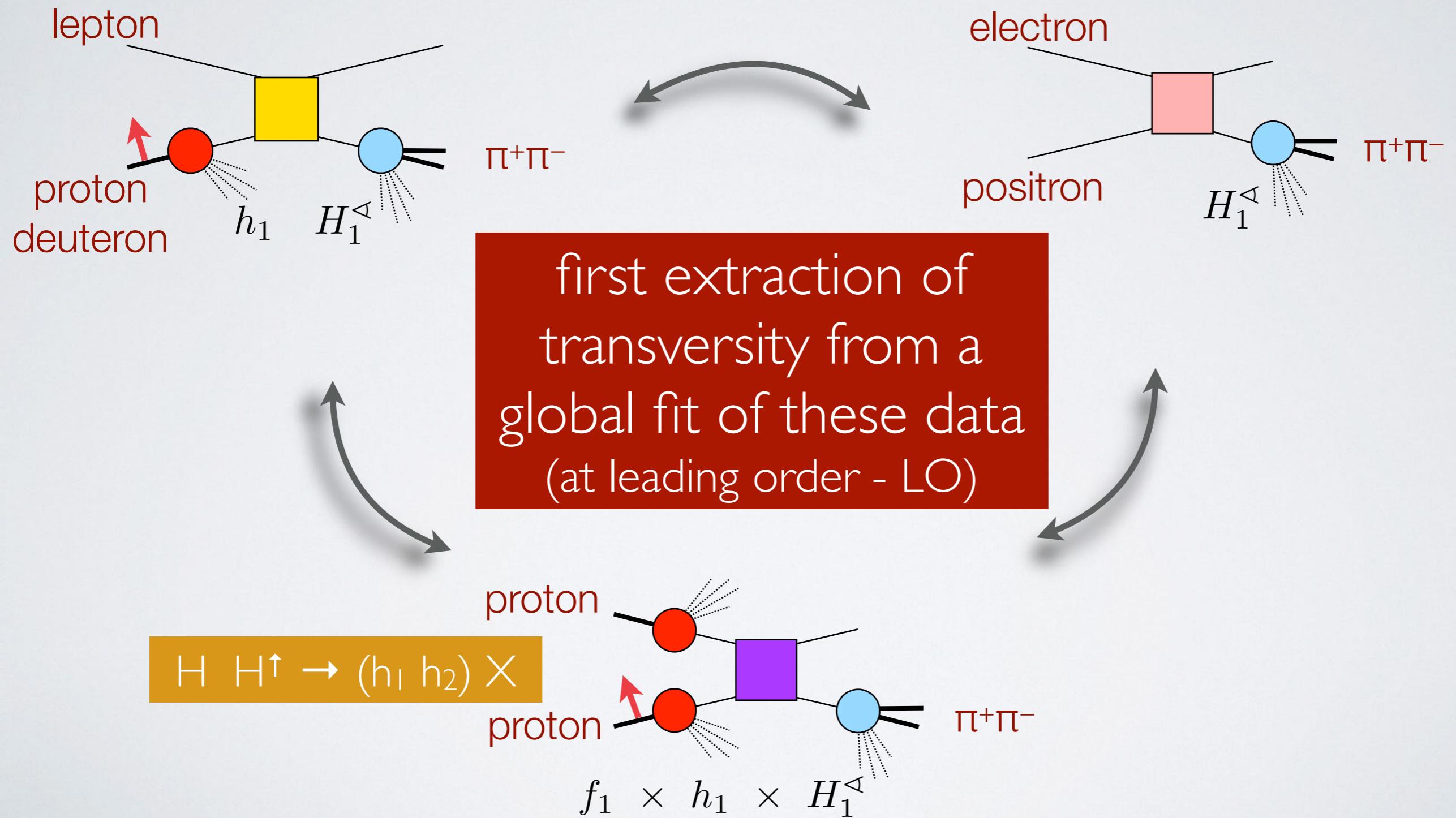
# 2-hadron-inclusive production



# take-away message

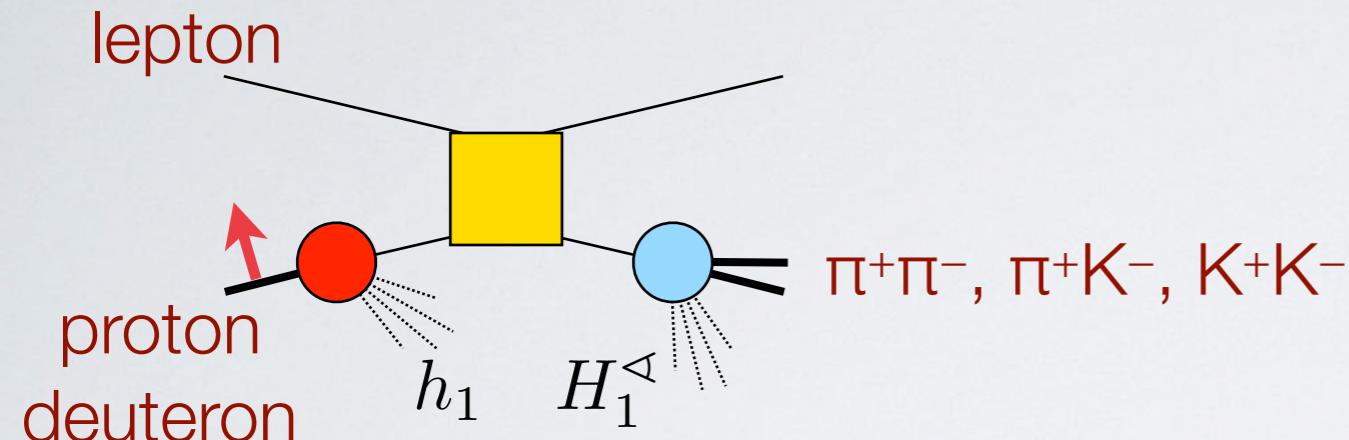
SIDIS  $\ell^- H^\uparrow \rightarrow \ell' (h_1 h_2) X$

$e^+e^- \rightarrow (h_1 h_2) X$



# exp. data for 2-hadron-inclusive production

SIDIS  $\ell^- H^\uparrow \rightarrow \ell^+ (h_1 h_2) X$

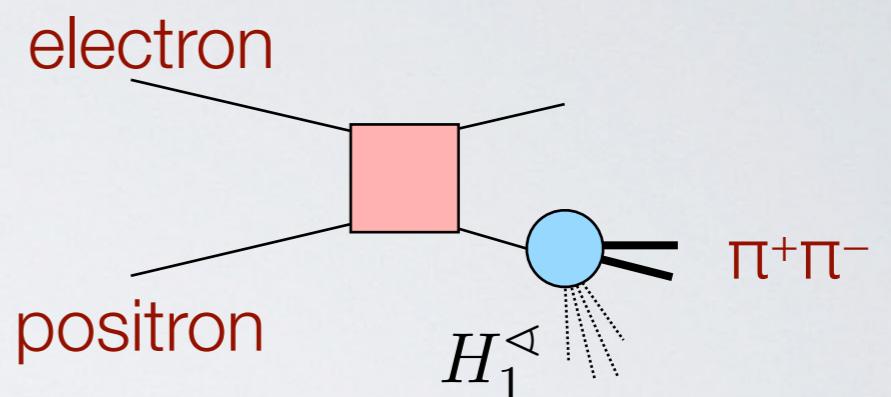


Airapetian et al.,  
*JHEP* **0806** (08) 017



Adolph et al., *P.L.* **B713** (12)  
Braun et al., *E.P.J. Web Conf.* **85** (15) 02018

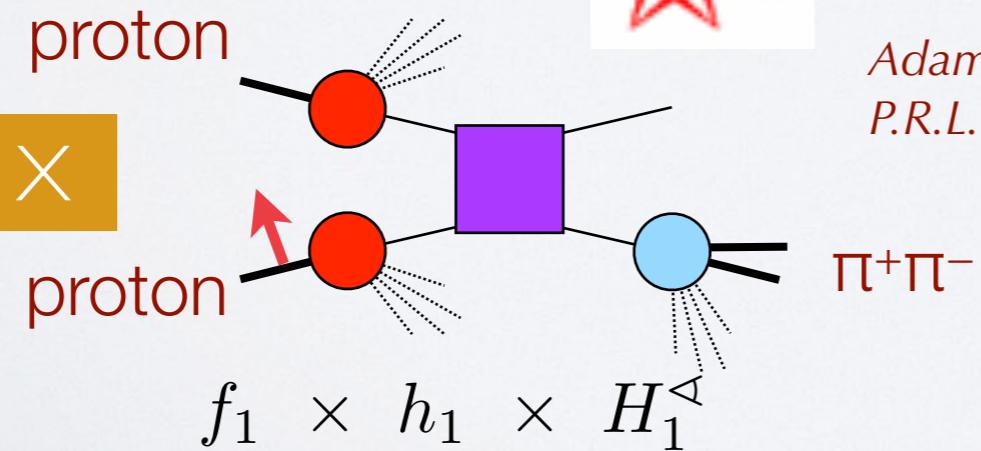
$e^+e^- \rightarrow (h_1 h_2) X$



Vossen et al., *P.R.L.* **107** (11) 072004

$D_1$  Seidl et al., *P.R.* **D96** (17) 032005

$H^- H^\uparrow \rightarrow (h_1 h_2) X$



run 2006 ( $s=200$ )

Adamczyk et al. (STAR),  
*P.R.L.* **115** (2015) 242501



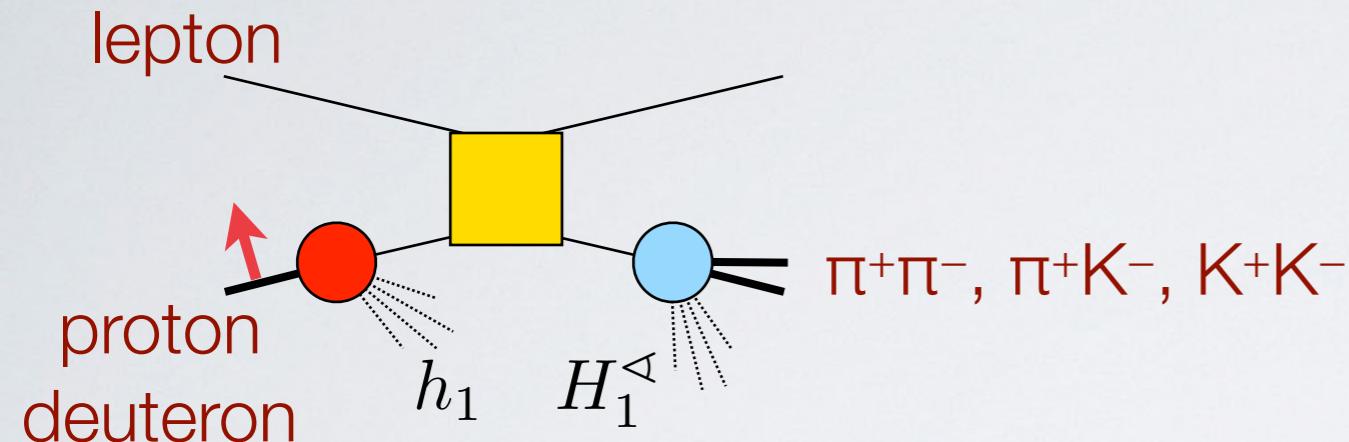
run 2011 ( $s=500$ )

Adamczyk et al. (STAR),  
*P.L.* **B780** (18) 332

$A_{UT}(\eta, M_h, P_T)$

# exp. data for 2-hadron-inclusive production

SIDIS  $\ell^- H^\uparrow \rightarrow \ell' (h_1 h_2) X$

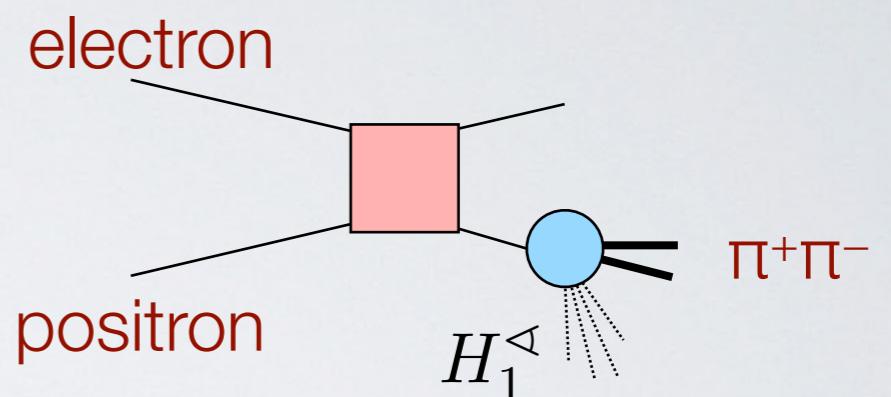


Airapetian et al.,  
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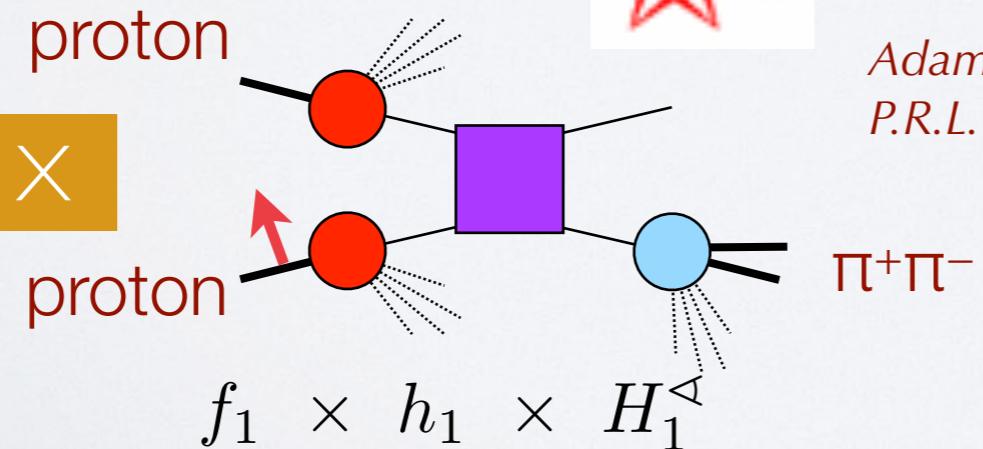
$e^+e^- \rightarrow (h_1 h_2) X$



Vossen et al., *P.R.L.* **107** (11) 072004

$D_1$  Seidl et al., *P.R.* **D96** (17) 032005  
from Montecarlo

$H^- H^\uparrow \rightarrow (h_1 h_2) X$



run 2006



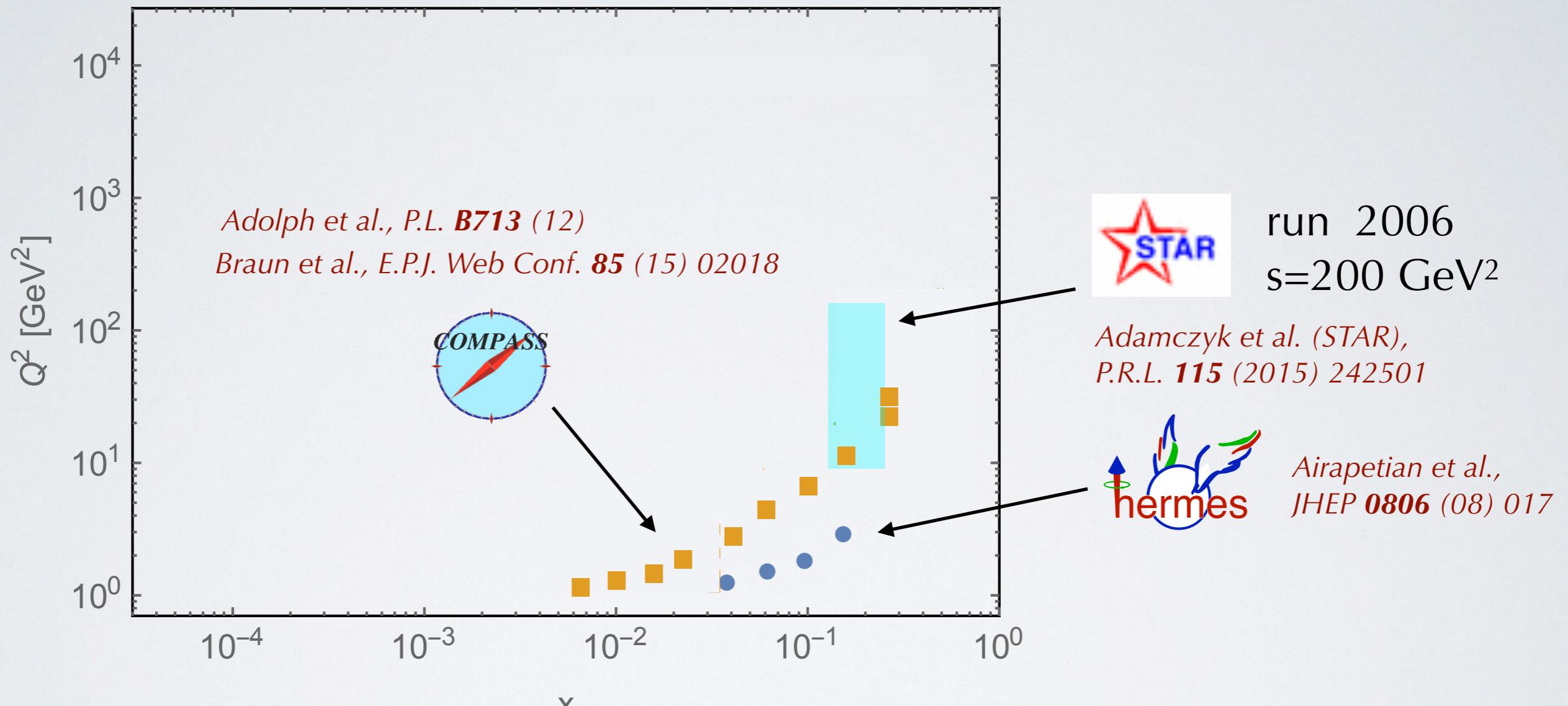
Adamczyk et al. (STAR),  
*P.R.L.* **115** (2015) 242501

run 2011

Adamczyk et al. (STAR),  
*P.L.* **B780** (18) 332

$A_{UT}(\eta, M_h, P_T)$

# the kinematics



explore only valence quarks

# choice of functional form

different funct. form whose Mellin transform can be computed analytically  
but keep main feature: comply with Soffer Bound at any x and scale  $Q^2$

$$h_1^{q_v}(x; Q_0^2) = F^{q_v}(x) \left[ \text{SB}^q(x) + \overline{\text{SB}}^{\bar{q}}(x) \right]$$

↓  
Soffer Bound

$$2|h_1^q(x, Q^2)| \leq 2 \text{ SB}^q(x, Q^2) = |f_1^q(x, Q^2) + g_1^q(x, Q^2)|$$

MSTW08      DSSV

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MSTW08      DSSV

$$F^{q_v}(x) = \frac{N_{q_v}}{\max_x [|F^{q_v}(x)|]} x^{A_{q_v}} [1 + B_{q_v} \text{Ceb}_1(x) + C_{q_v} \text{Ceb}_2(x) + D_{q_v} \text{Ceb}_3(x)]$$

$$|N_{q_v}| \leq 1 \Rightarrow |F^{q_v}(x)| \leq 1$$

Ceb<sub>n</sub>(x) Cebyshev polynomial

10 fitting parameters

Soffer Bound satisfied at any  $Q^2$

# choice of functional form

different funct. form whose Mellin transform can be computed analytically  
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$$h_1^{q_v}(x; Q_0^2) = F^{q_v}(x) \left[ \text{SB}^q(x) + \overline{\text{SB}}^{\bar{q}}(x) \right]$$

## Soffer Bound

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MSTW08	DSSV
--------	------

$$F^{q_v}(x) = \frac{N_{q_v}}{\max_x |F^{q_v}(x)|} x^{A_{q_v}} [1 + B_{q_v} \text{Ceb}_1(x) + C_{q_v} \text{Ceb}_2(x) + D_{q_v} \text{Ceb}_3(x)]$$

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## Ceb<sub>n</sub>(x) Cebyshев polynomial

## 10 fitting parameters

# Soffer Bound satisfied at any $Q^2$

if  $\lim_{x \rightarrow 0} x \text{SB}^q(x) \propto x^{a_q}$  then  $h_1^q(x) \stackrel{x \rightarrow 0}{\approx} x^{A_q + a_q - 1}$

$A_q + a_q > \frac{1}{3}$  grants that tensor charge  $\delta q(Q^2) = \int_{x_{\min}}^1 dx h_1^{q-\bar{q}}(x, Q^2)$  is finite and error  $O(1\%)$

# theoretical uncertainties

quark  $D_1^q$  is well constrained by  $e^+e^-$  (Montecarlo) but

we don't know anything about the gluon  $D_1^g$  ( $e^+e^-$  doesn't help..)

Single-Spin Asymmetry  
in  $p-p^\uparrow$  collisions

$$A_{UT}(\eta, M_h, P_T) = \frac{d\sigma_{UT}}{d\sigma_0}$$

chiral-odd  
not important

typical cross section for  $a+b \rightarrow c+d$  process

$$d\sigma_0 \propto \sum_{a,b,c,d} \int \frac{dx_a dx_b}{8\pi^2 \bar{z}} f_1^a(x_a) f_1^b(x_b) \frac{d\hat{\sigma}_{ab \rightarrow cd}}{d\hat{t}} D_1^c(\bar{z}, M_h)$$

# theoretical uncertainties

quark  $D_{1q}$  is well constrained by  $e^+e^-$  (Montecarlo) but

we don't know anything about the gluon  $D_{1g}$  ( $e^+e^-$  doesn't help..)

Single-Spin Asymmetry  
in  $p-p^\uparrow$  collisions

$$A_{UT}(\eta, M_h, P_T) = \frac{d\sigma_{UT}}{d\sigma_0}$$

chiral-odd  
not important  
important !

typical cross section for  $a+b \rightarrow c+d$  process

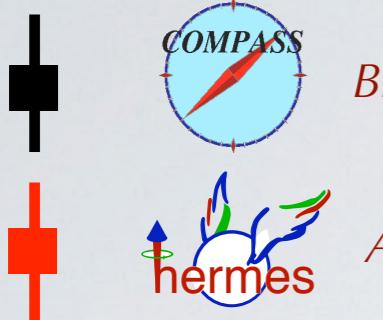
$$d\sigma_0 \propto \sum_{a,b,c,d} \int \frac{dx_a dx_b}{8\pi^2 \bar{z}} f_1^a(x_a) f_1^b(x_b) \frac{d\hat{\sigma}_{ab \rightarrow cd}}{d\hat{t}} D_1^c(\bar{z}, M_h)$$

our choice: compute  $d\sigma_0$  with  $D_{1g}(Q_0) = \begin{cases} 0 \\ D_{1u}(Q_0) / 4 \\ D_{1u}(Q_0) \end{cases}$

deteriorates our  $e^+e^-$  fit as  $\chi^2/\text{dof} = \begin{cases} 1.69 & 1.28 \\ 1.81 & 1.37 \\ 2.96 & 2.01 \end{cases}$

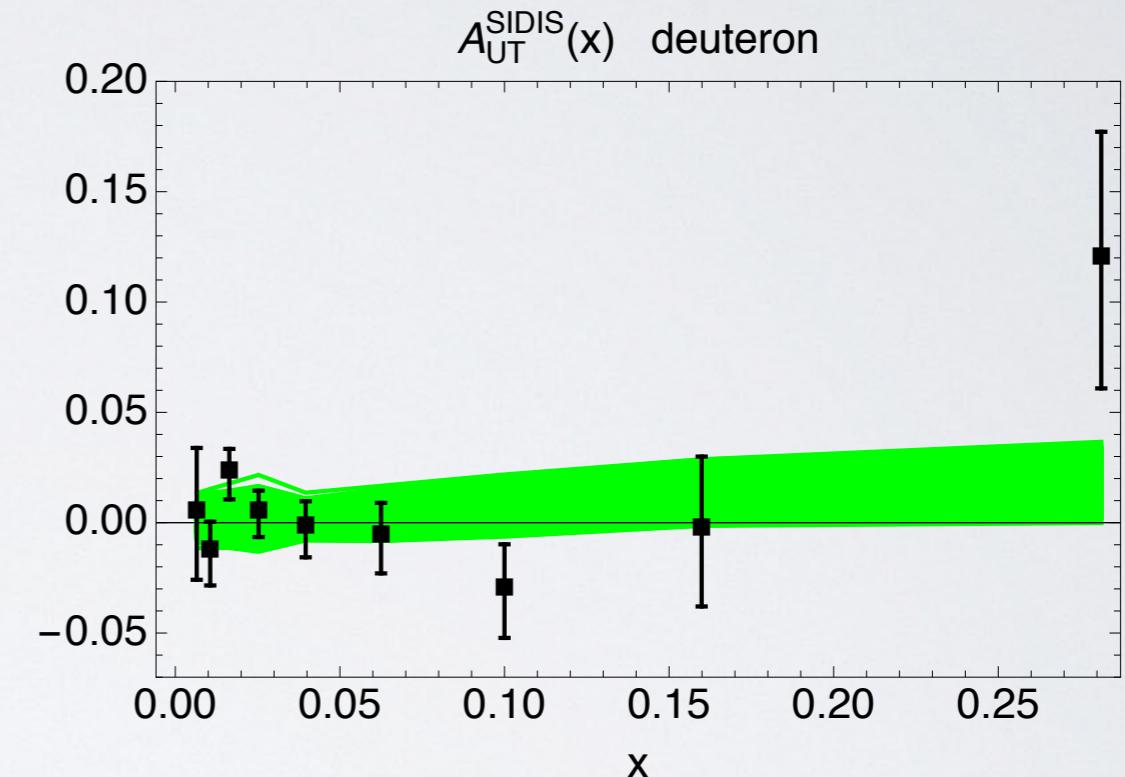
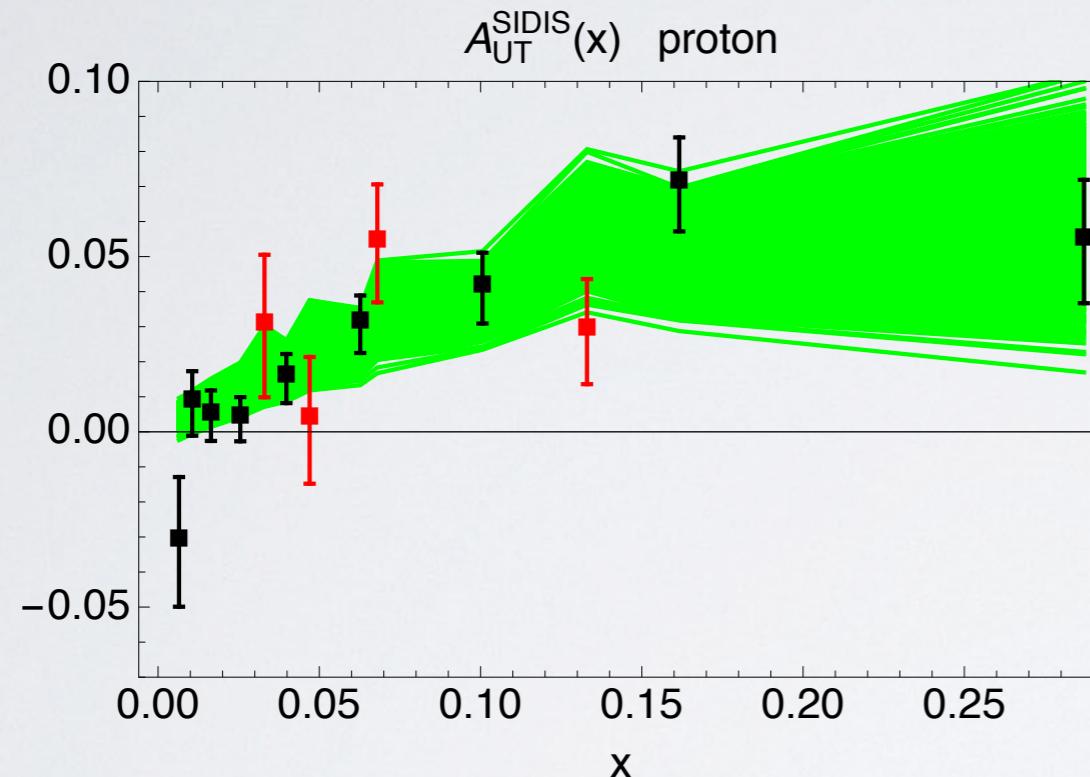
background       $\rho$       channels

# statistical uncertainty: the bootstrap method



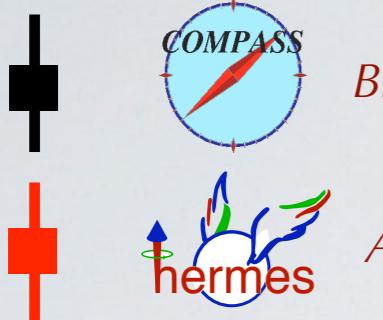
Braun et al., E.P.J. Web Conf. **85** (15) 02018

Airapetian et al., JHEP **0806** (08) 017



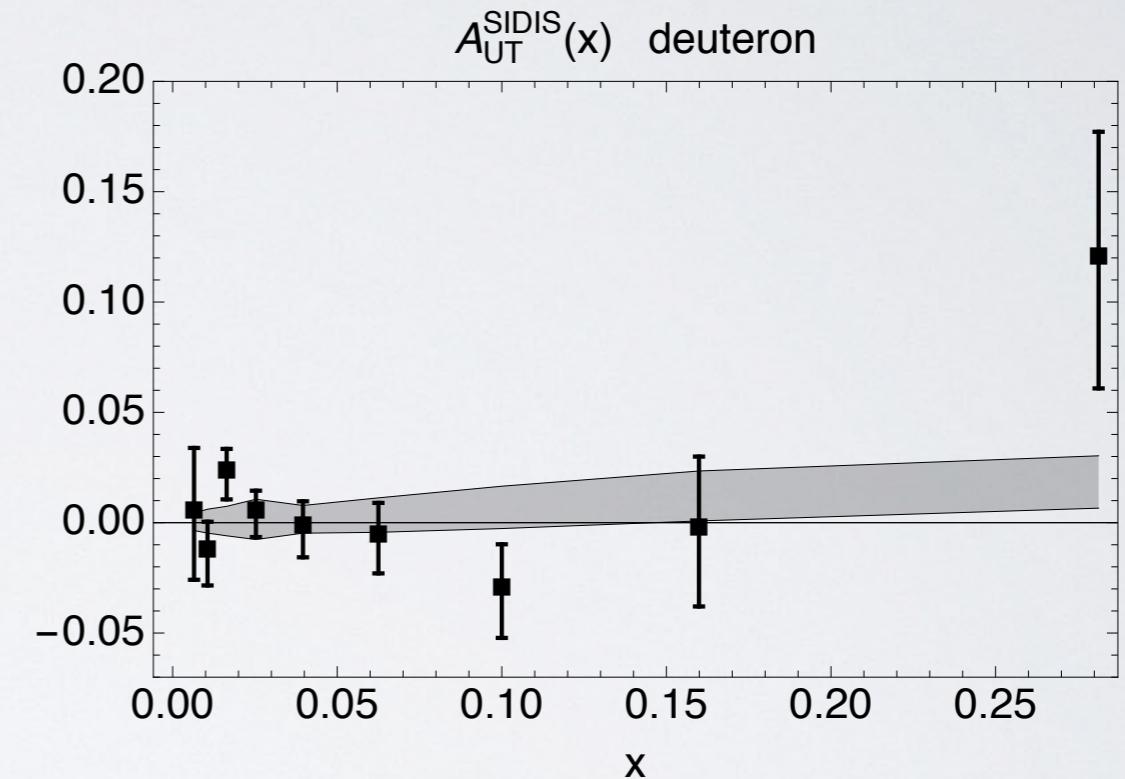
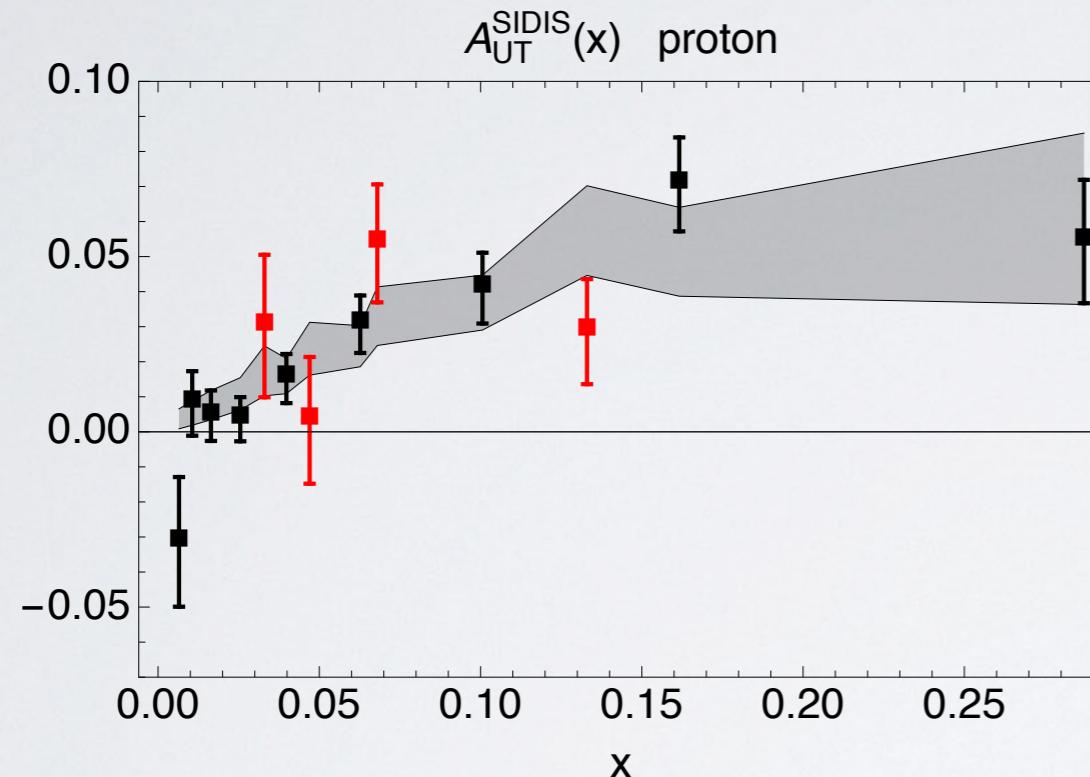
all 600 replicas

# statistical uncertainty: the bootstrap method



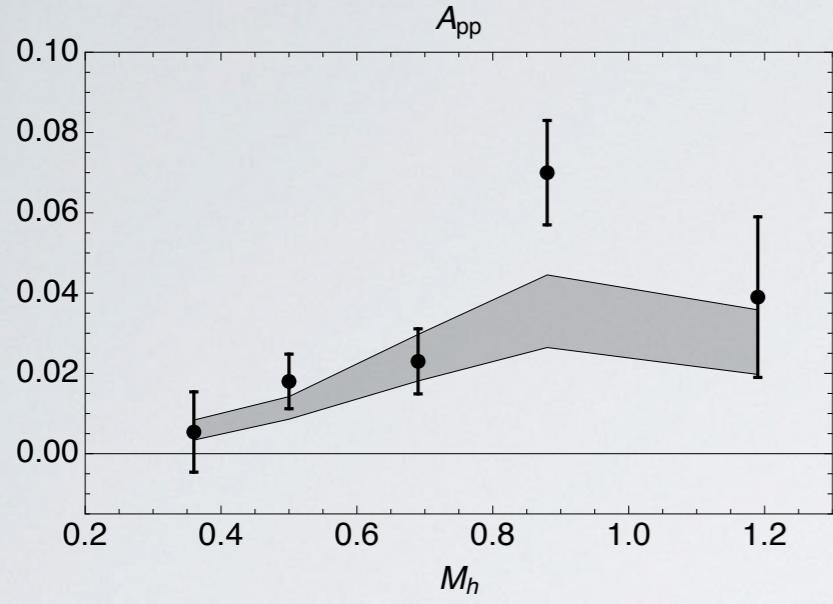
Braun et al., E.P.J. Web Conf. **85** (15) 02018

Airapetian et al., JHEP **0806** (08) 017

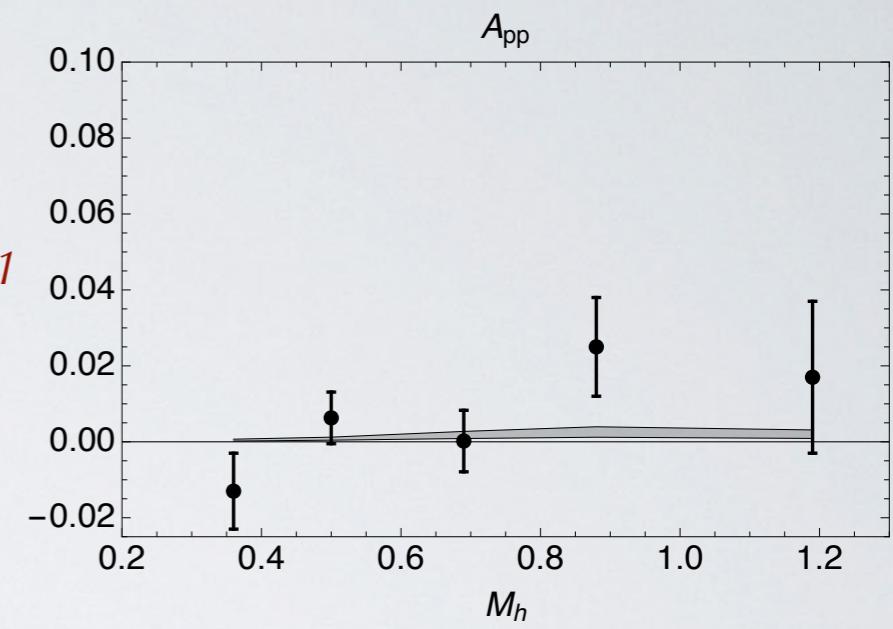


90% replicas

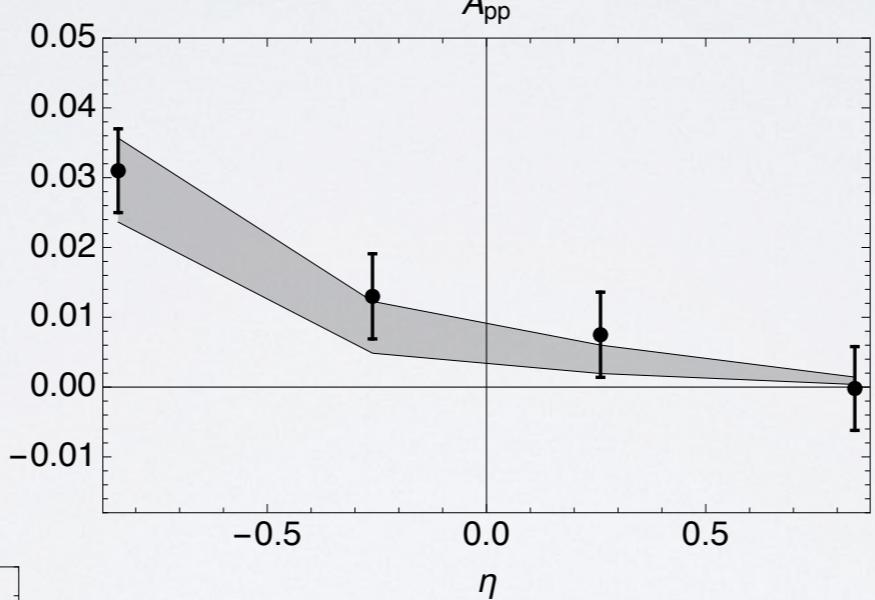
# fit STAR asymmetry



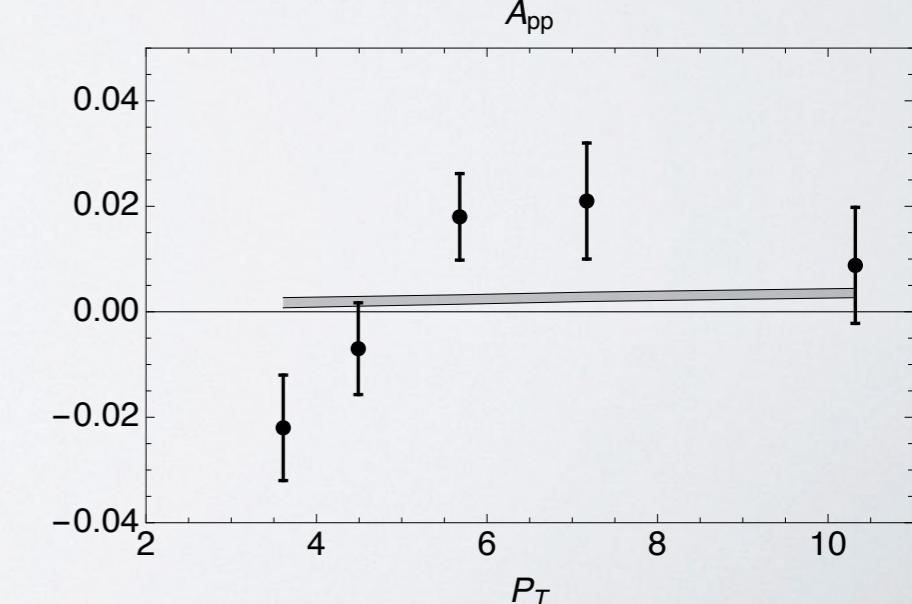
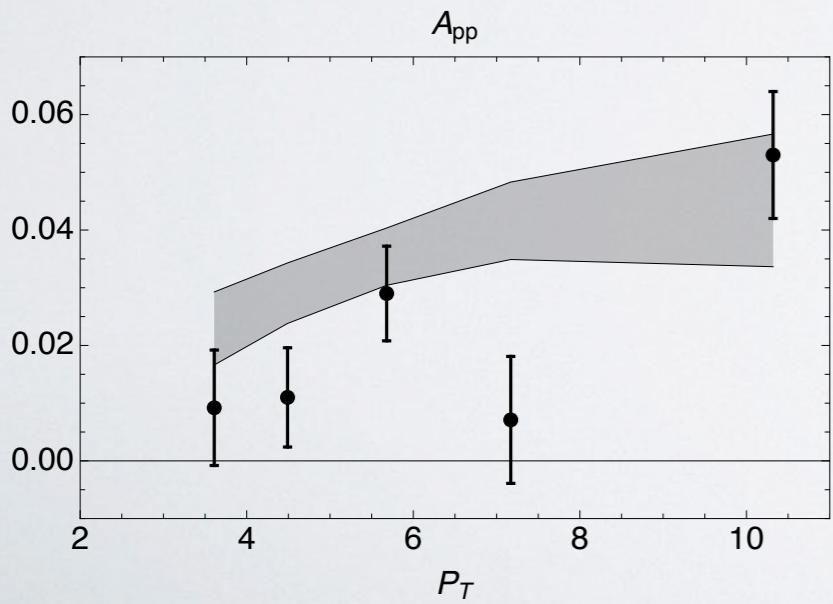
Adamczyk et al. (STAR),  
P.R.L. 115 (2015) 242501



$\eta > 0$



90% uncertainty band



# $\chi^2$ of the fit

**46** data points, **10** parameters  
global  $\chi^2/\text{dof} = 2.08 \pm 0.09$

$\approx 38\%$

$\approx 62\%$

**SIDIS**

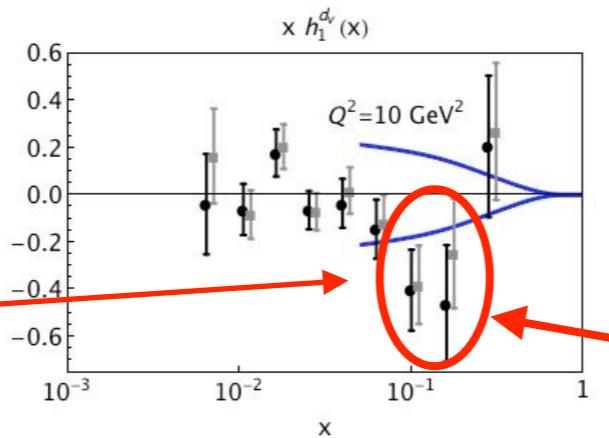
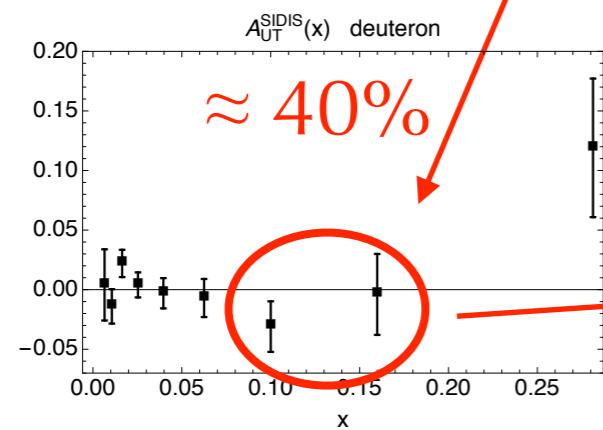


$\approx 24\%$



$\approx 76\%$

$\approx 60\%$   
rest



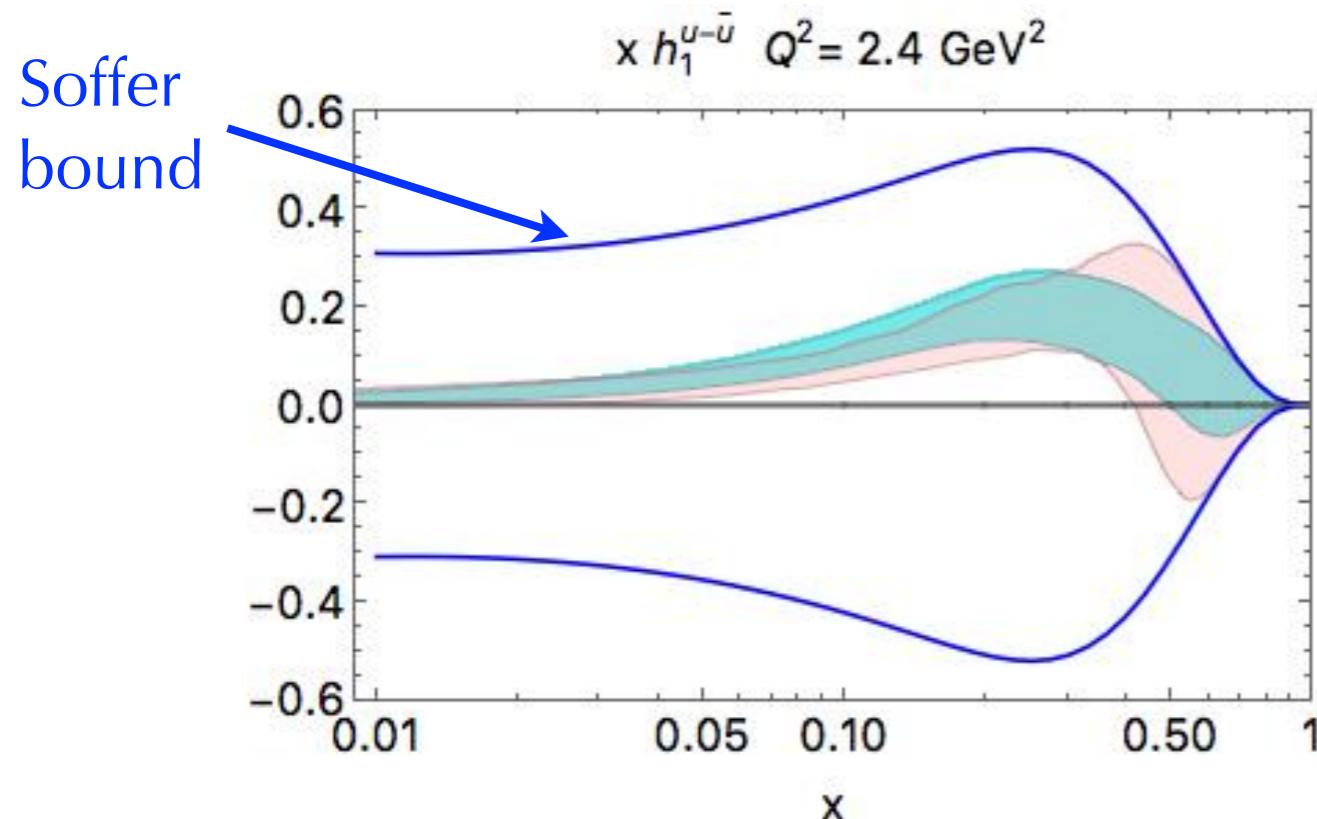
**STAR**

$\rightarrow P_T$  bins  $\approx 70\%$

$\rightarrow M_h$  bins  $\approx 28\%$

$\rightarrow \eta$  bins  $\approx 2\%$

# comparison with previous fit



*Radici & Bacchetta,  
P.R.L. **120** (18) 192001*

global fit

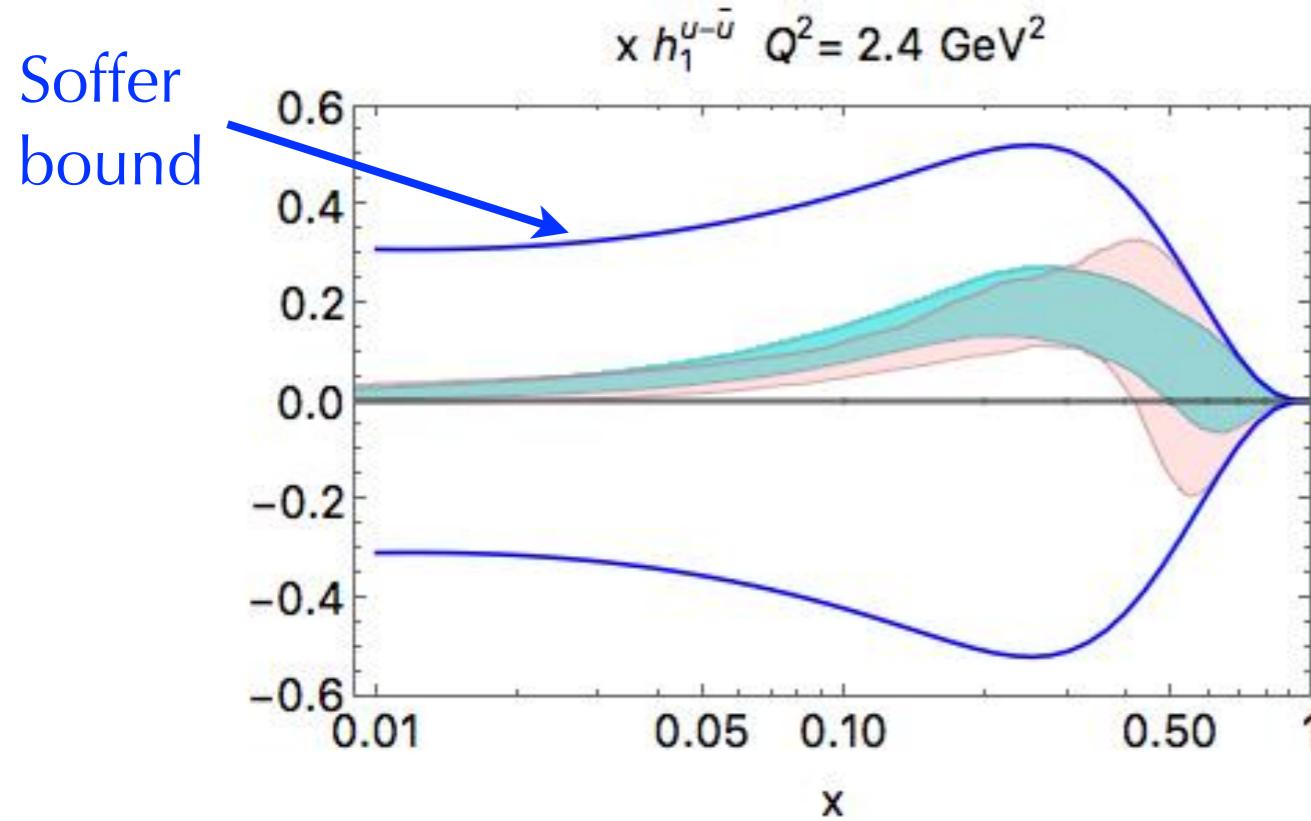
up

higher  
precision

old fit (only SIDIS data)

*Radici et al.,  
JHEP **1505** (15) 123*

# comparison with previous fit



*Radici & Bacchetta,  
P.R.L. **120** (18) 192001*

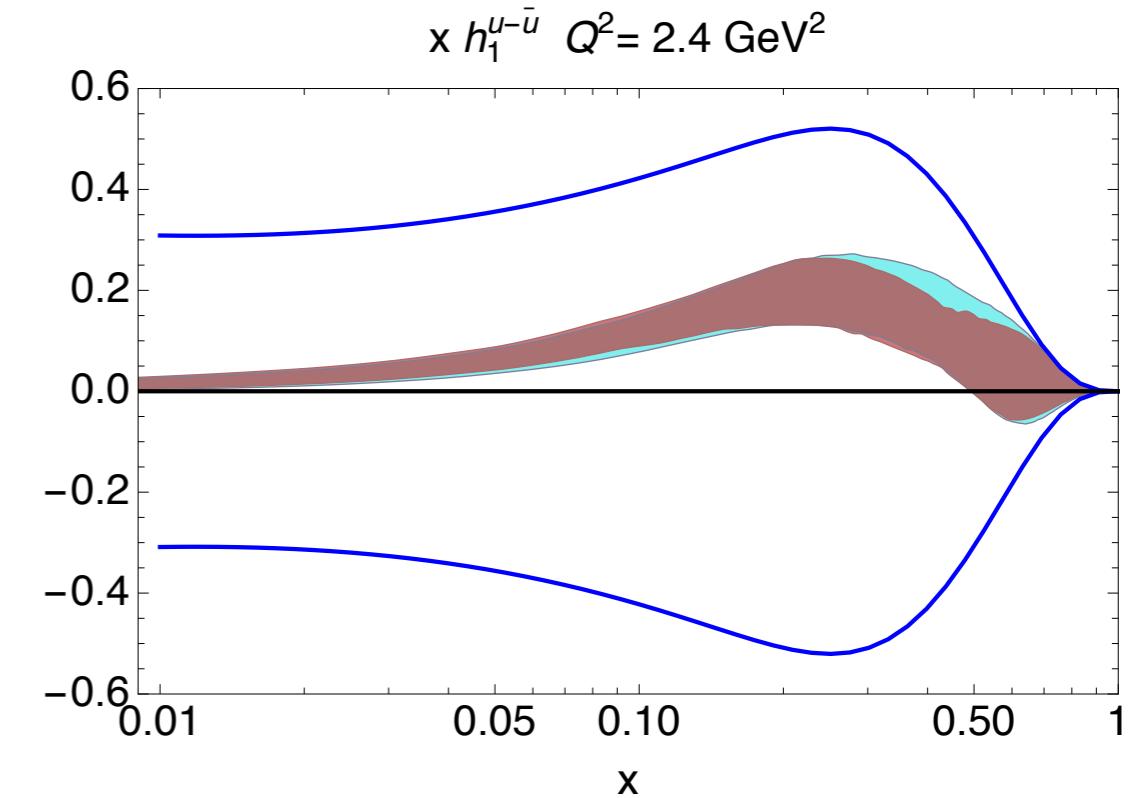
up

higher precision

up

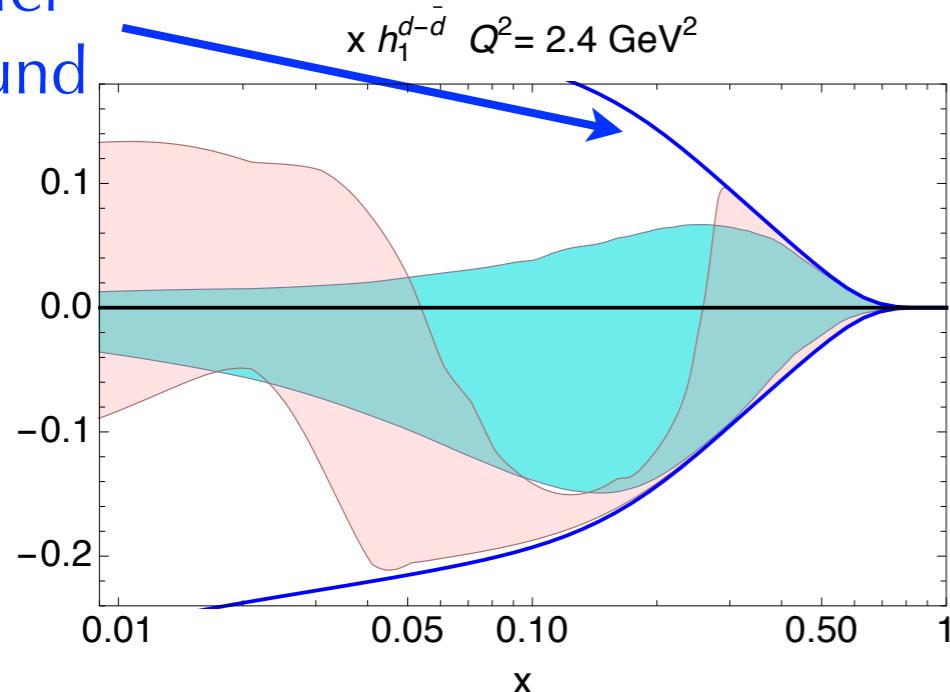
insensitive to  
uncertainty on  
gluon  $D_1$

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_{1u}/4 \\ D_{1u} \end{cases}$$



# comparison with previous fit

Soffer  
bound



*Radici & Bacchetta,  
P.R.L. **120** (18) 192001*

global fit

down

old fit

*Radici et al.,  
JHEP **1505** (15) 123*

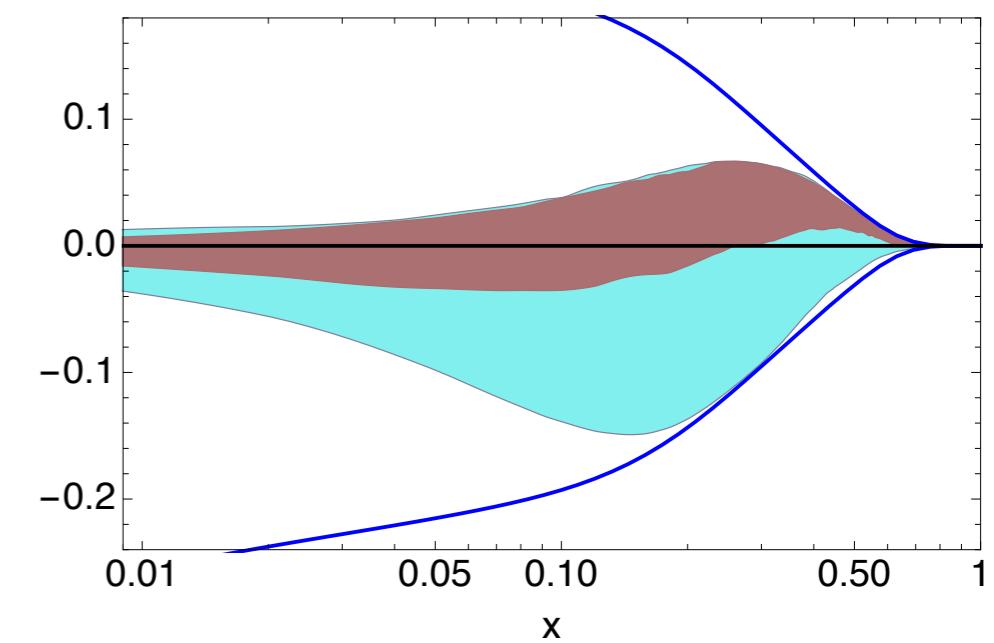
down

sensitive to  
uncertainty on  
gluon  $D_1$

$$D_1 g(Q_0) = 0$$

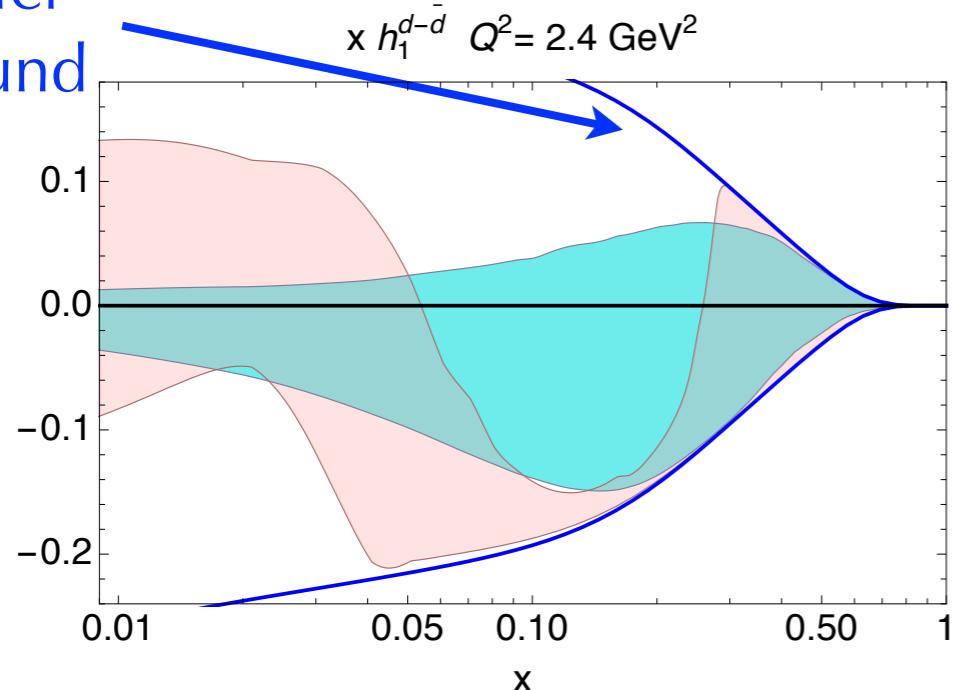
$$D_1 g(Q_0) = \begin{cases} 0 \\ D_1^u / 4 \\ D_1^u \end{cases}$$

$x h_1^{d-\bar{d}} Q^2 = 2.4 \text{ GeV}^2$



# comparison with previous fit

Soffer  
bound



*Radici & Bacchetta,  
P.R.L. **120** (18) 192001*

global fit

old fit

*Radici et al.,  
JHEP **1505** (15) 123*

down

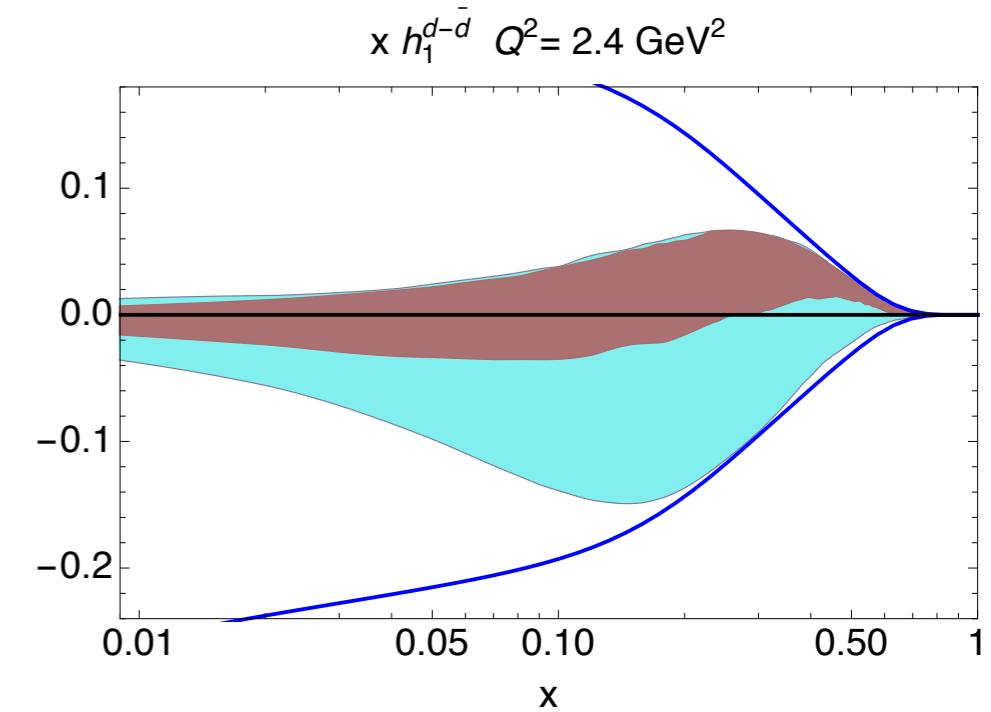
need dihadron multiplicities  
from RHIC

down

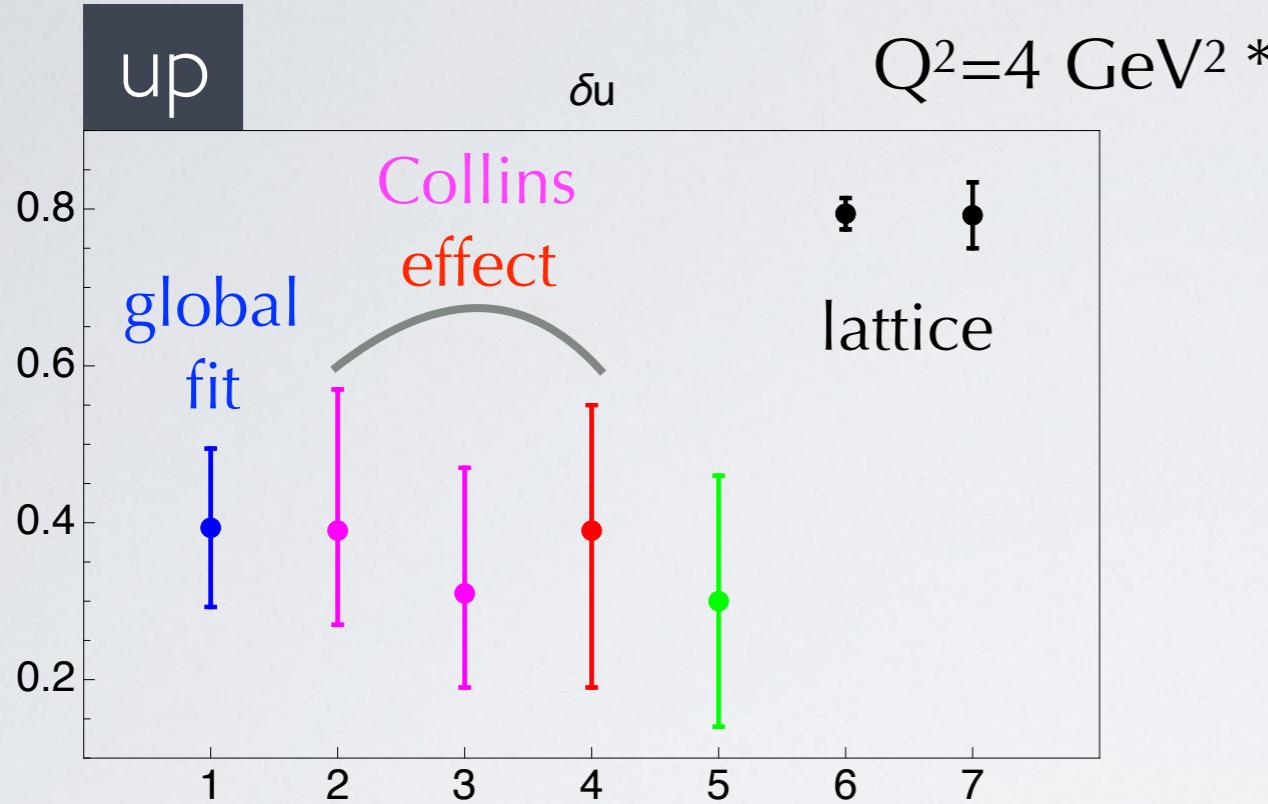
sensitive to  
uncertainty on  
gluon  $D_1$

$$D_{1g}(Q_0) = 0$$

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_{1u}/4 \\ D_{1u} \end{cases}$$



**tensor charge**  $\delta q(Q^2) = \int dx h_1^{q\bar{q}}(x, Q^2)$



- 1- global fit** *Radici & Bacchetta, P.R.L. 120 (18) 192001*

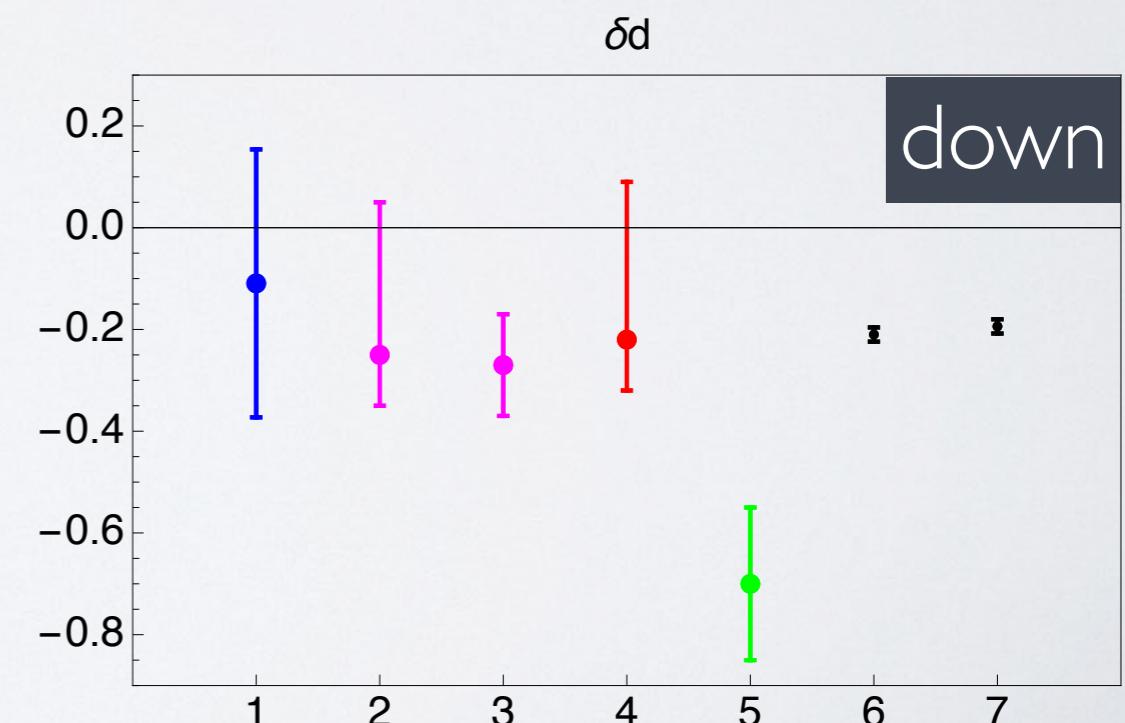
**2,3- Torino** *Anselmino et al., P.R. D87 (13) 094019* \*  $Q^2=1$

**4- TMD fit** *Kang et al., P.R. D93 (16) 014009* \*  $Q^2=10$

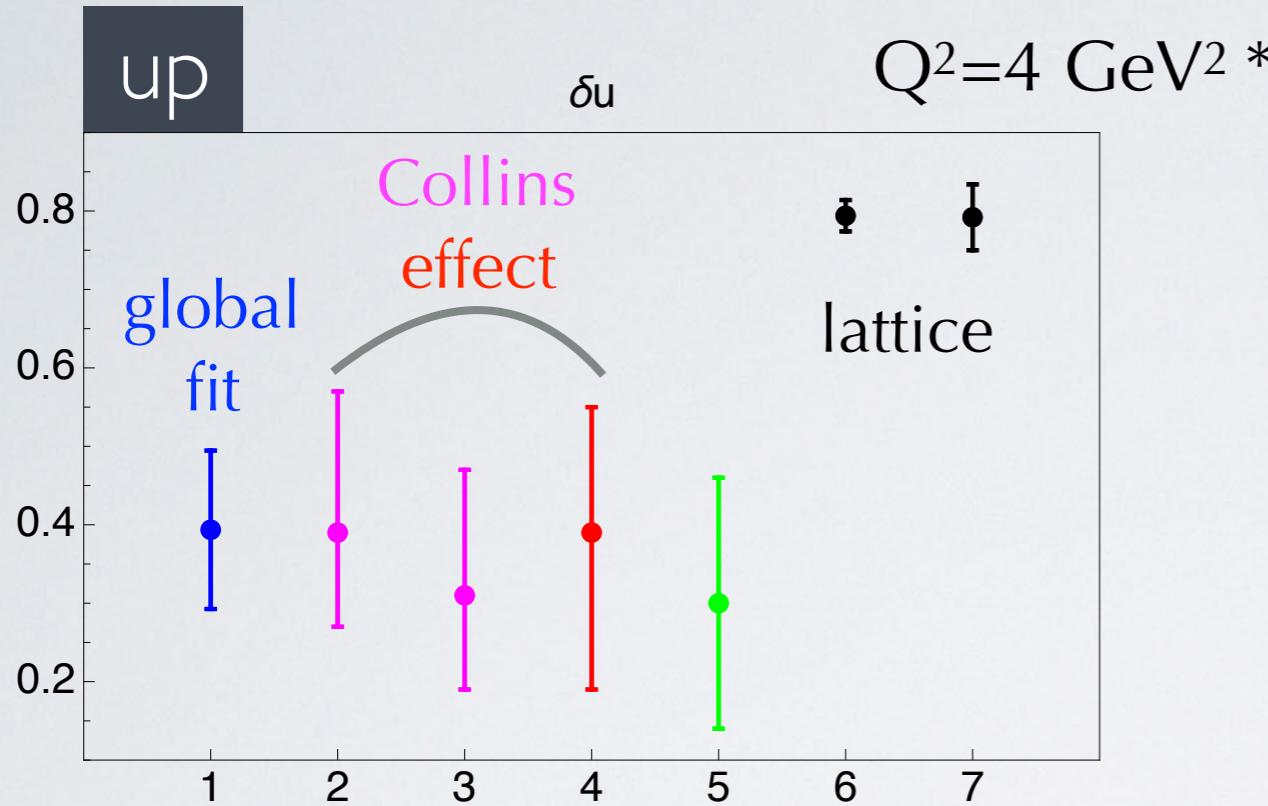
**5- JAM fit** *Lin et al., P.R.L. 120 (18) 152502* {**Collins effect + lattice  $g_T = \delta u - \delta d$** } \*  $Q_0^2=2$

**6- ETMC17** *Alexandrou et al., P.R. D95 (17) 114514;  
E P.R. D96 (17) 099906*

**7- PNDME16** *Bhattacharya et al., P.R. D94 (16) 054508*



$$\text{tensor charge } \delta q(Q^2) = \int dx h_1 q\bar{q} (x, Q^2)$$



incompatibility for up  
compatible for down  
but with large errors  
(except JAM)

**1- global fit** *Radici & Bacchetta, P.R.L. 120 (18) 192001*

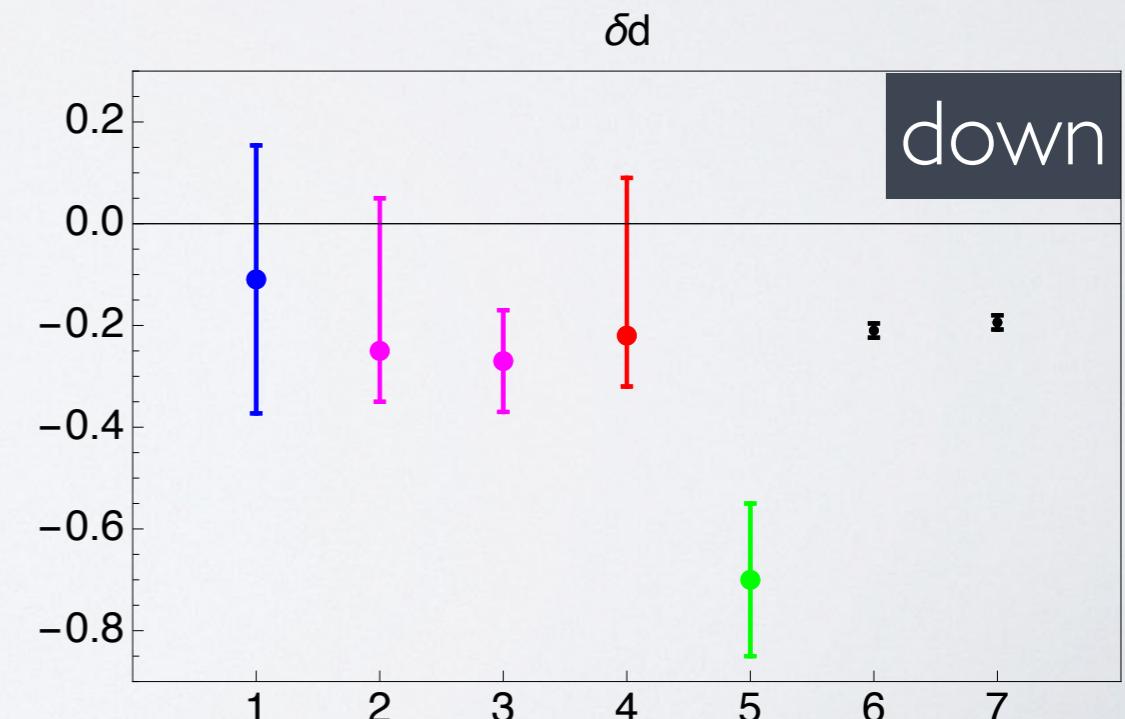
**2,3- Torino** *Anselmino et al., P.R.D 87 (13) 094019* \*  $Q^2=1$

**4- TMD fit** *Kang et al., P.R.D 93 (16) 014009* \*  $Q^2=10$

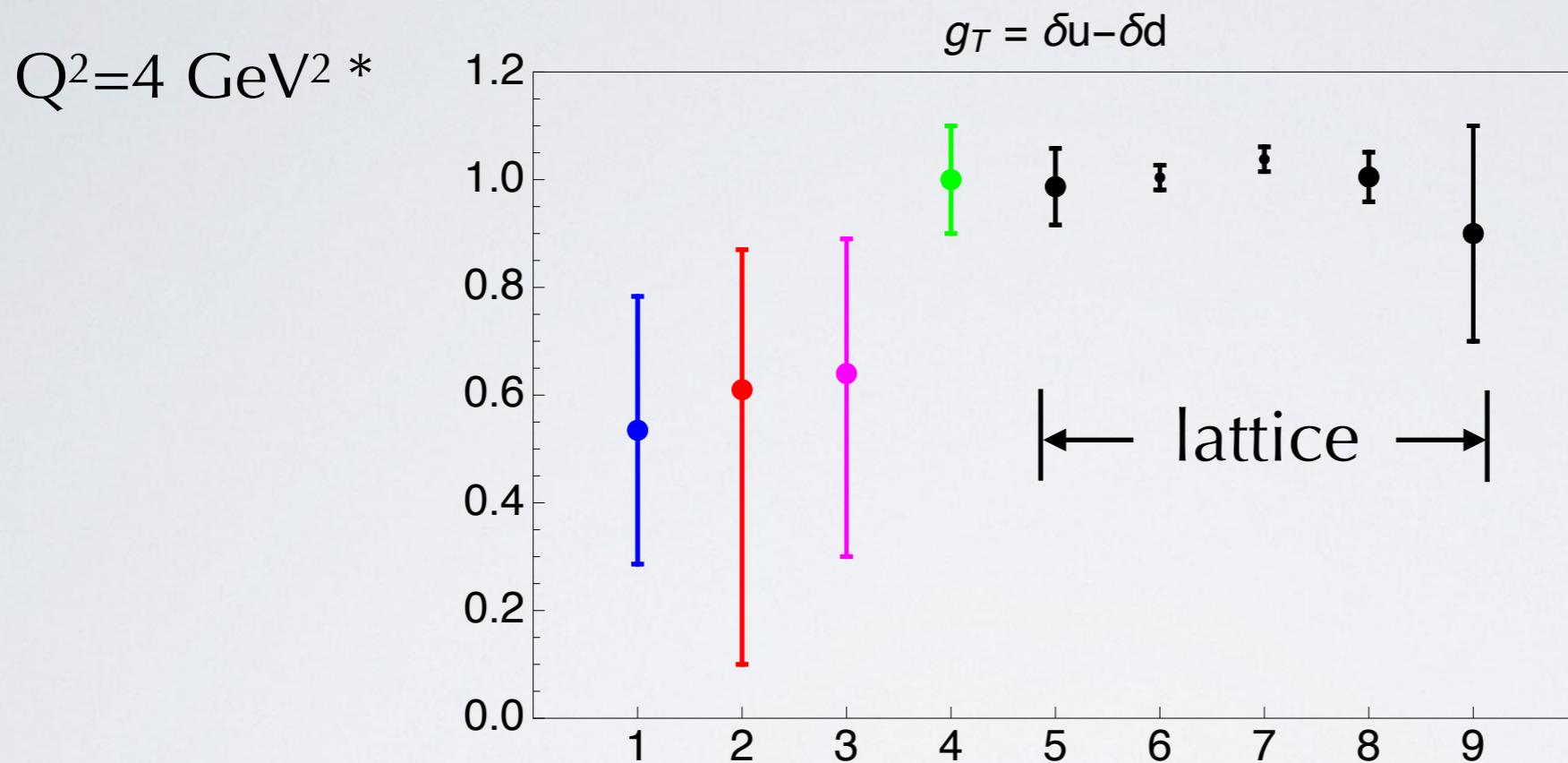
**5- JAM fit** *Lin et al., P.R.L. 120 (18) 152502* {Collins effect + lattice  $g_T = \delta u - \delta d$  \*  $Q_0^2=2$

**6- ETMC17** *Alexandrou et al., P.R.D 95 (17) 114514; E P.R.D 96 (17) 099906*

**7- PNNDME16** *Bhattacharya et al., P.R.D 94 (16) 054508*



# isovector tensor charge $g_T = \delta u - \delta d$



Radici & Bacchetta,  
P.R.L. 120 (18) 192001

Kang et al., P.R. D93 (16) 014009

Anselmino et al., P.R. D87 (13) 094019

Lin et al., P.R.L. 120 (18) 152502

1) **global fit '17**

2) **"TMD fit" \*  $Q^2=10$**

3) **Torino fit \*  $Q^2=1$**

4) **JAM fit '17 \*  $Q_0^2=2$**

5) PNDME '16

6) ETMC '17

7) LHPC '12

8) RQCD '14

9) RBC-UKQCD

Bhattacharya et al., P.R. D94 (16) 054508

Alexandrou et al., P.R. D95 (17) 114514;  
E P.R. D96 (17) 099906

Green et al., P.R. D86 (12)

Bali et al., P.R. D91 (15)

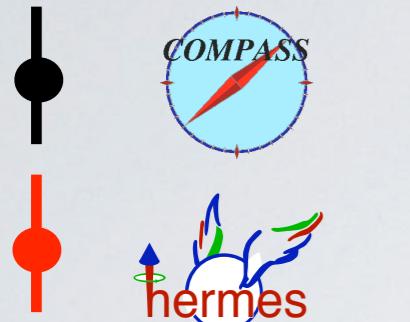
Aoki et al., P.R. D82 (10)

# “transverse-spin puzzle” ?

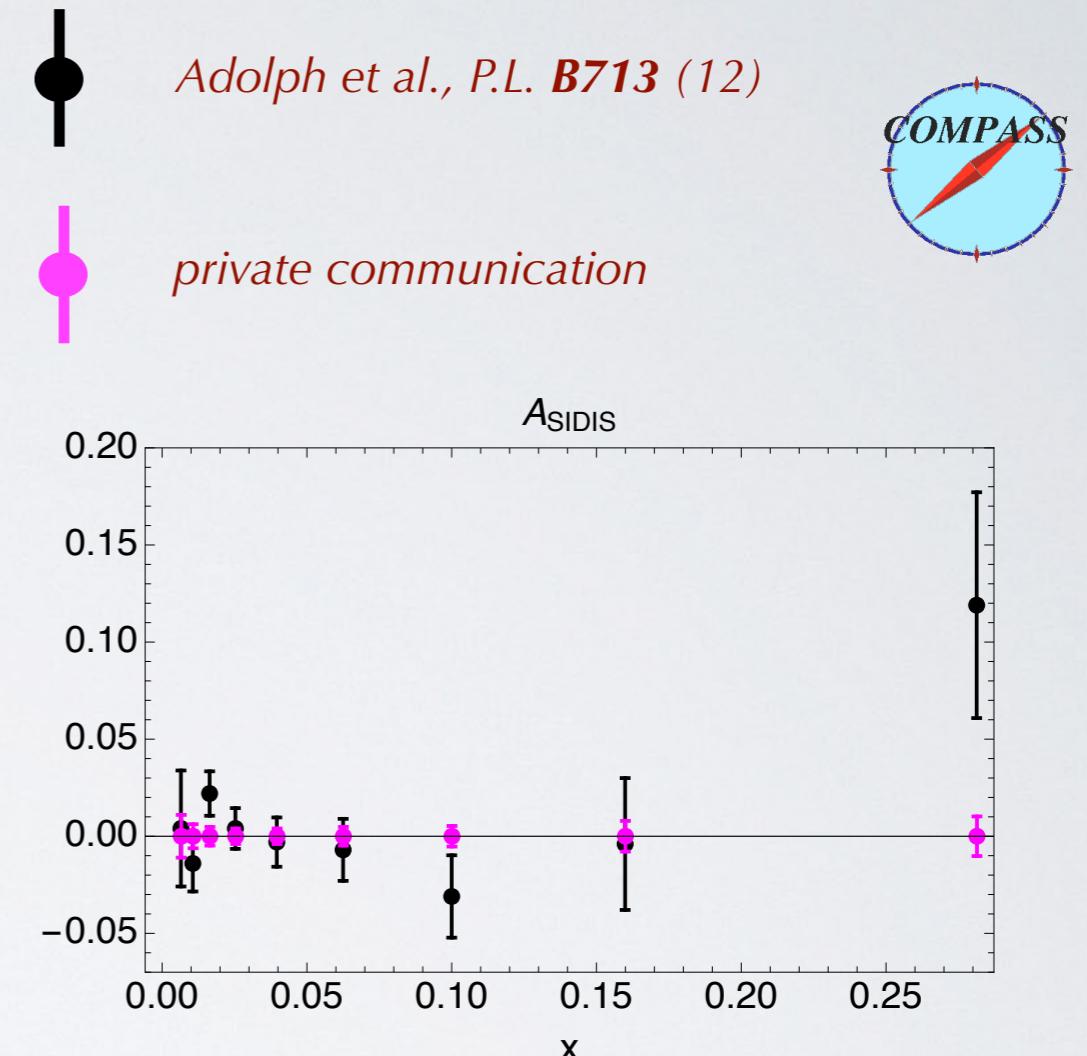
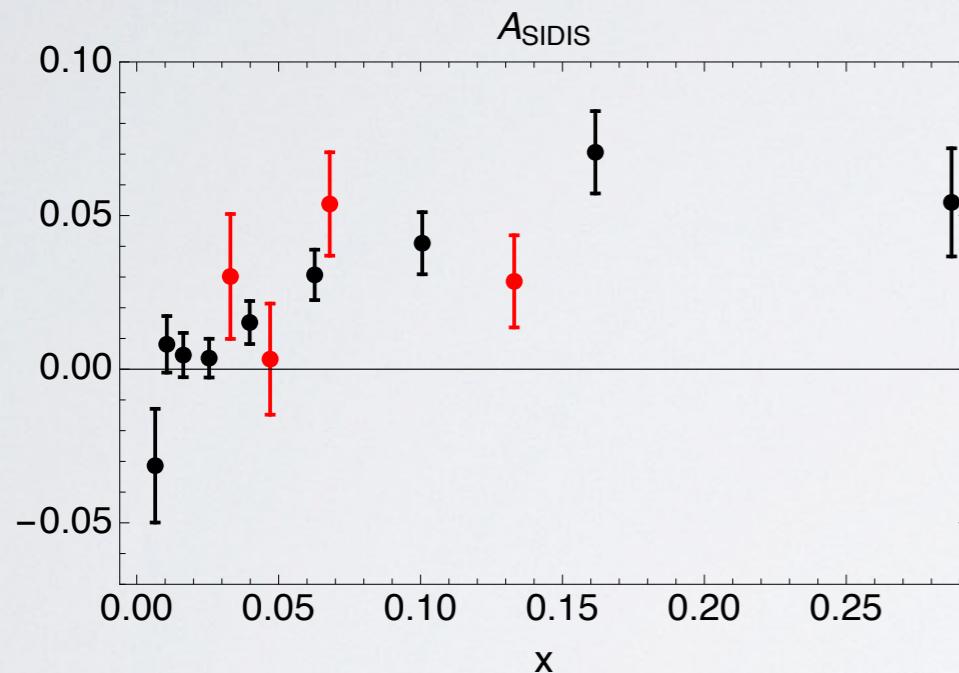
there seems to be no simultaneous compatibility  
about  $\delta u$ ,  $\delta d$ ,  $g_T = \delta u - \delta d$   
between lattice and  
phenomenological extractions  
of transversity

**BUT**

# add Compass pseudodata for future deuteron run



proton



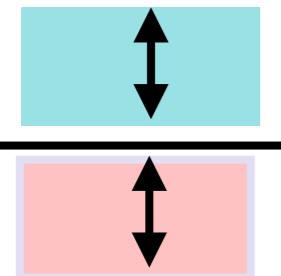
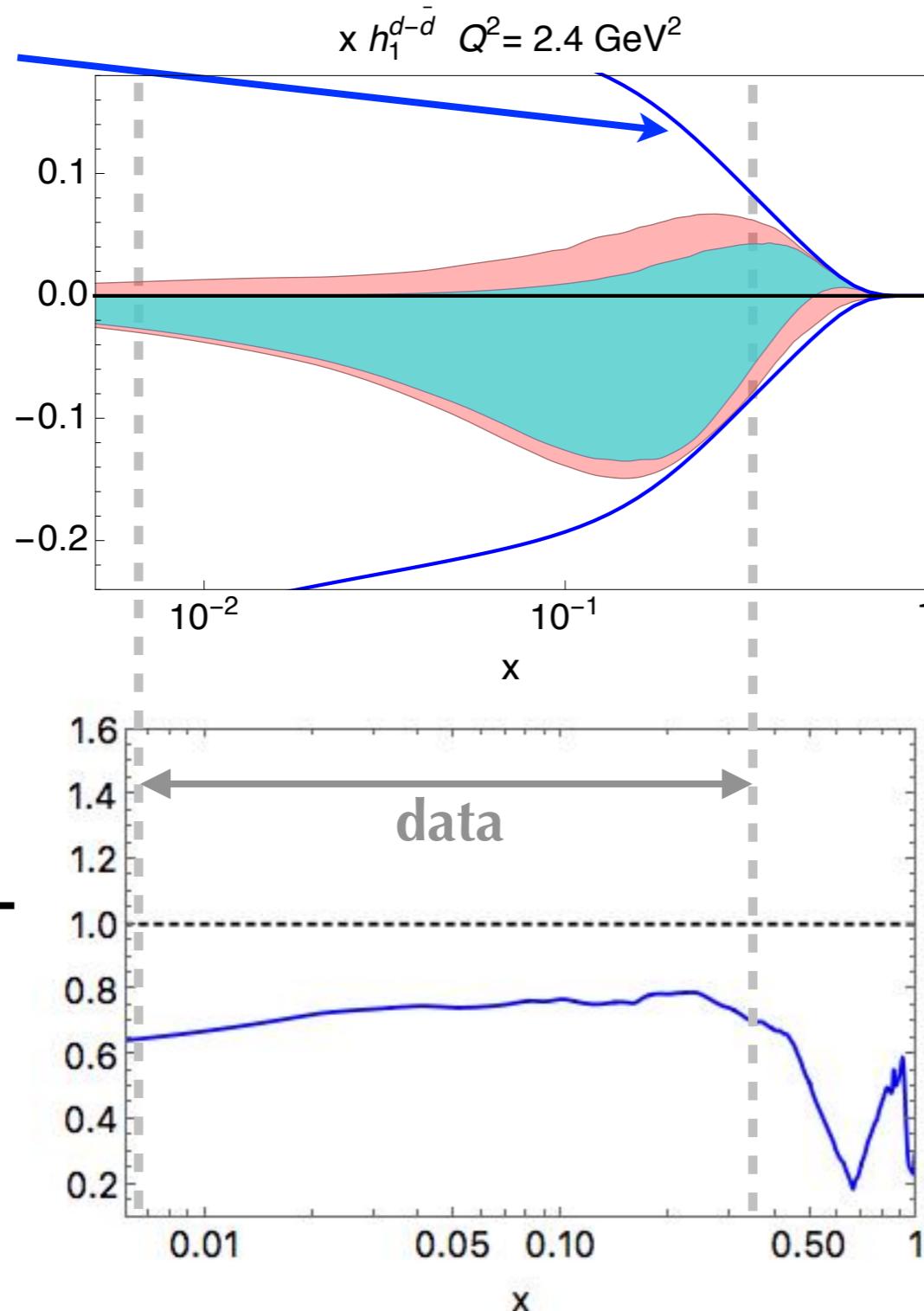
$$A_{\text{SIDIS}} \sim 4h_1^{u_v} - h_1^{d_v}$$

$$A_{\text{SIDIS}} \sim h_1^{u_v} + h_1^{d_v}$$

# pseudodata impact on down

Soffer  
bound

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_1^u/4 \\ D_1^u \end{cases}$$



ratio of  
widths

global fit + pseudodata

global fit

Radici & Bacchetta,  
*PRL* **120** (18) 192001

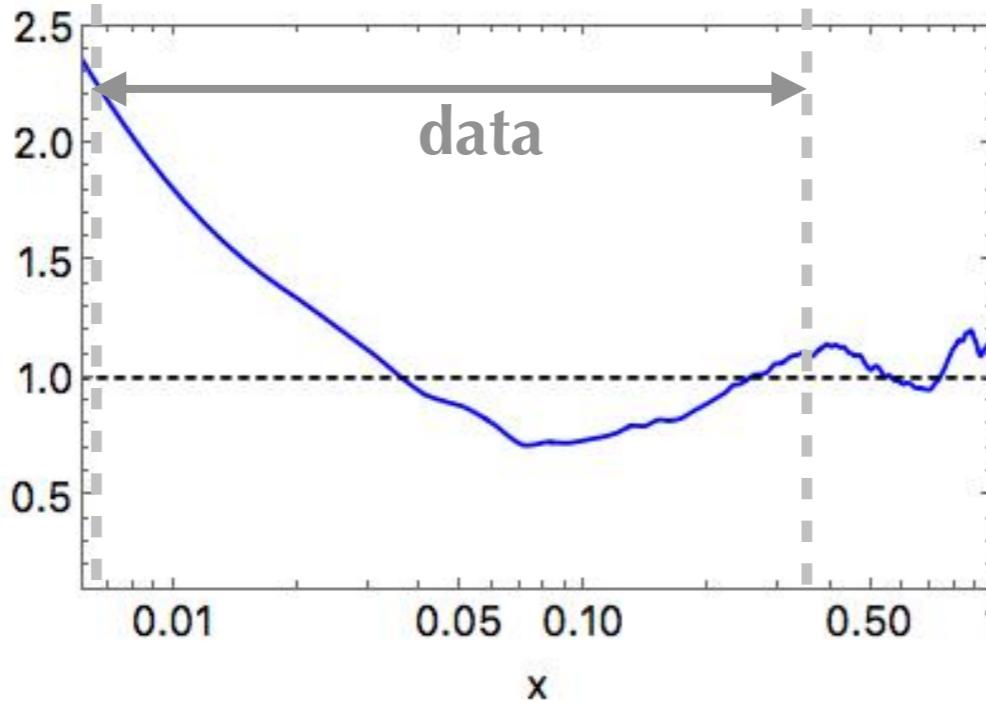
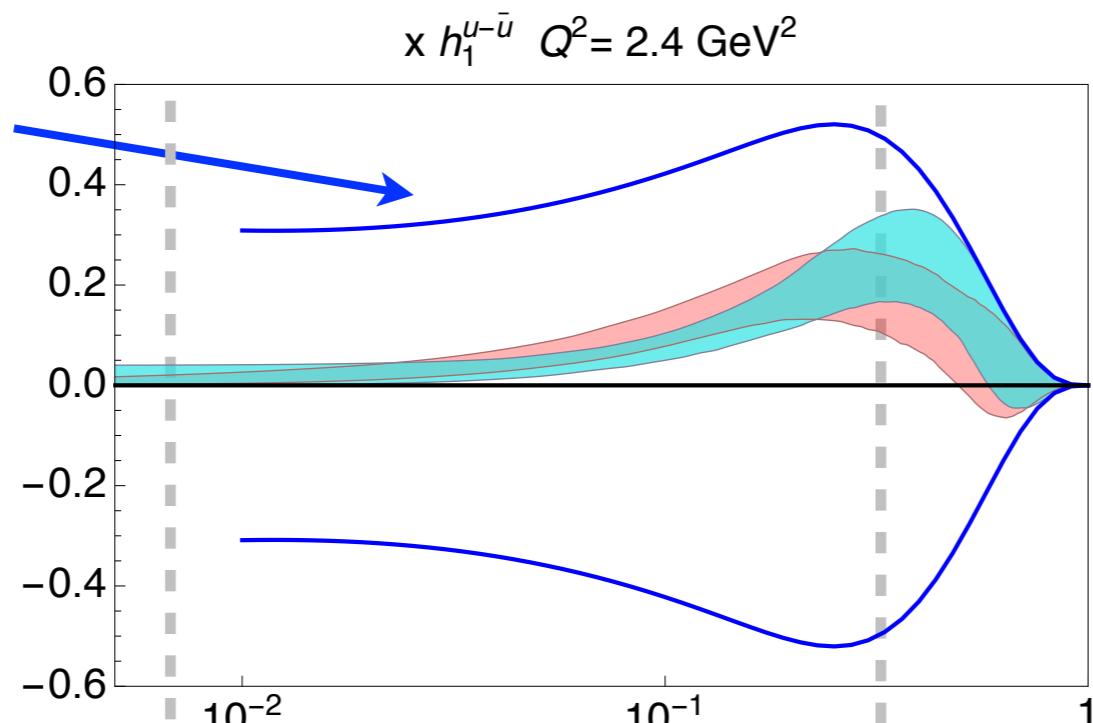
down

< 20% >  
increase in  
precision

# pseudodata impact on up

Soffer  
bound

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_1^u/4 \\ D_1^u \end{cases}$$



ratio of  
widths

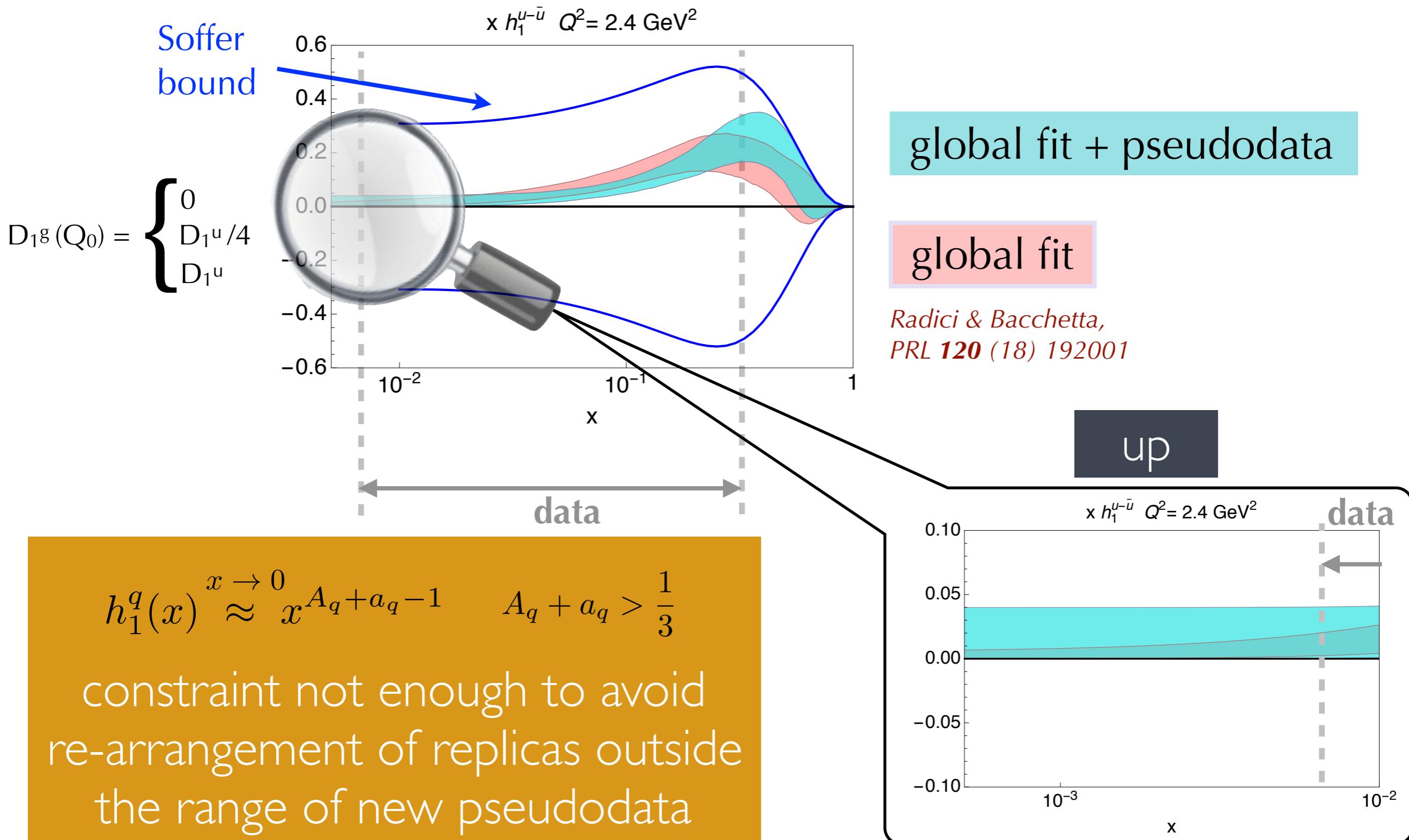
global fit + pseudodata

global fit

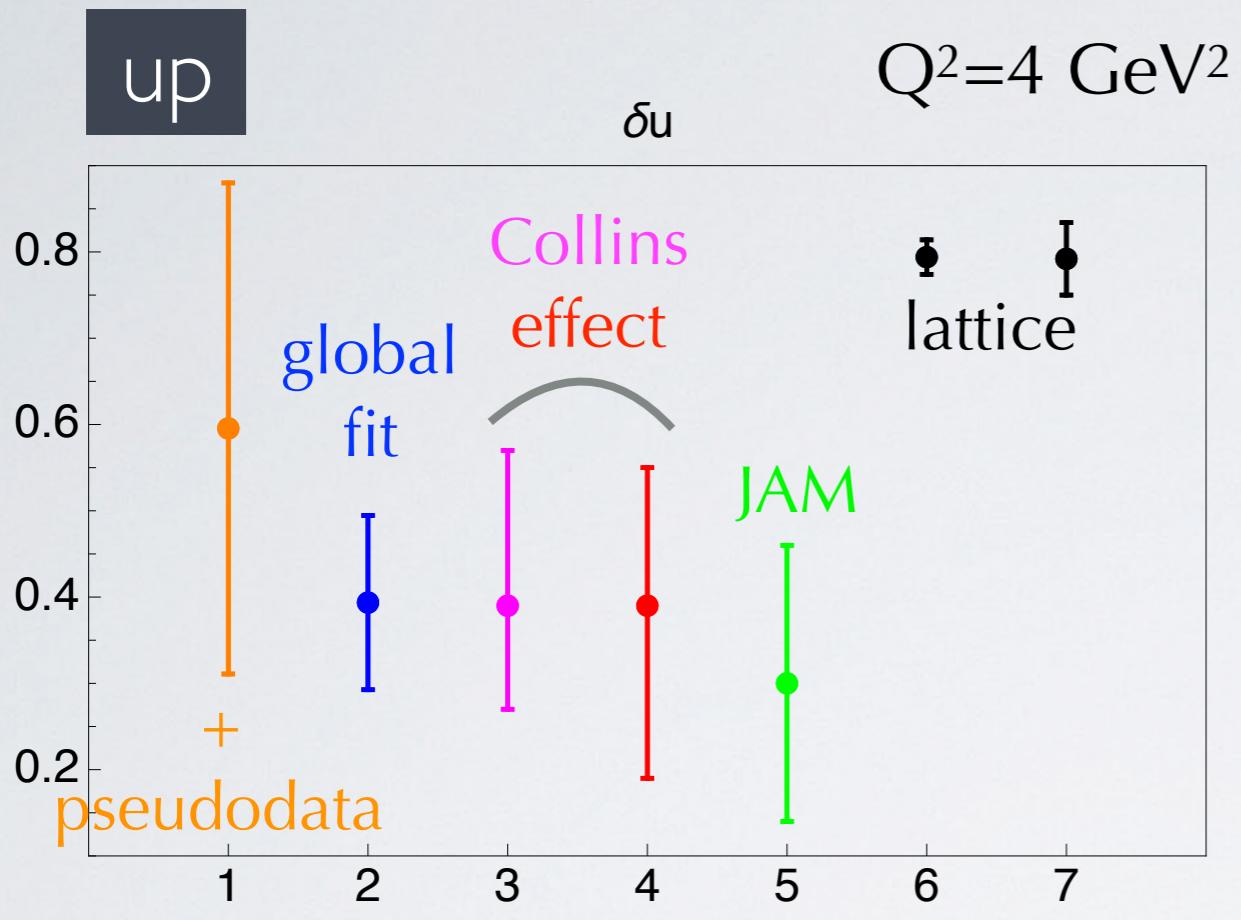
Radici & Bacchetta,  
*PRL* **120** (18) 192001

up

# pseudodata impact on up

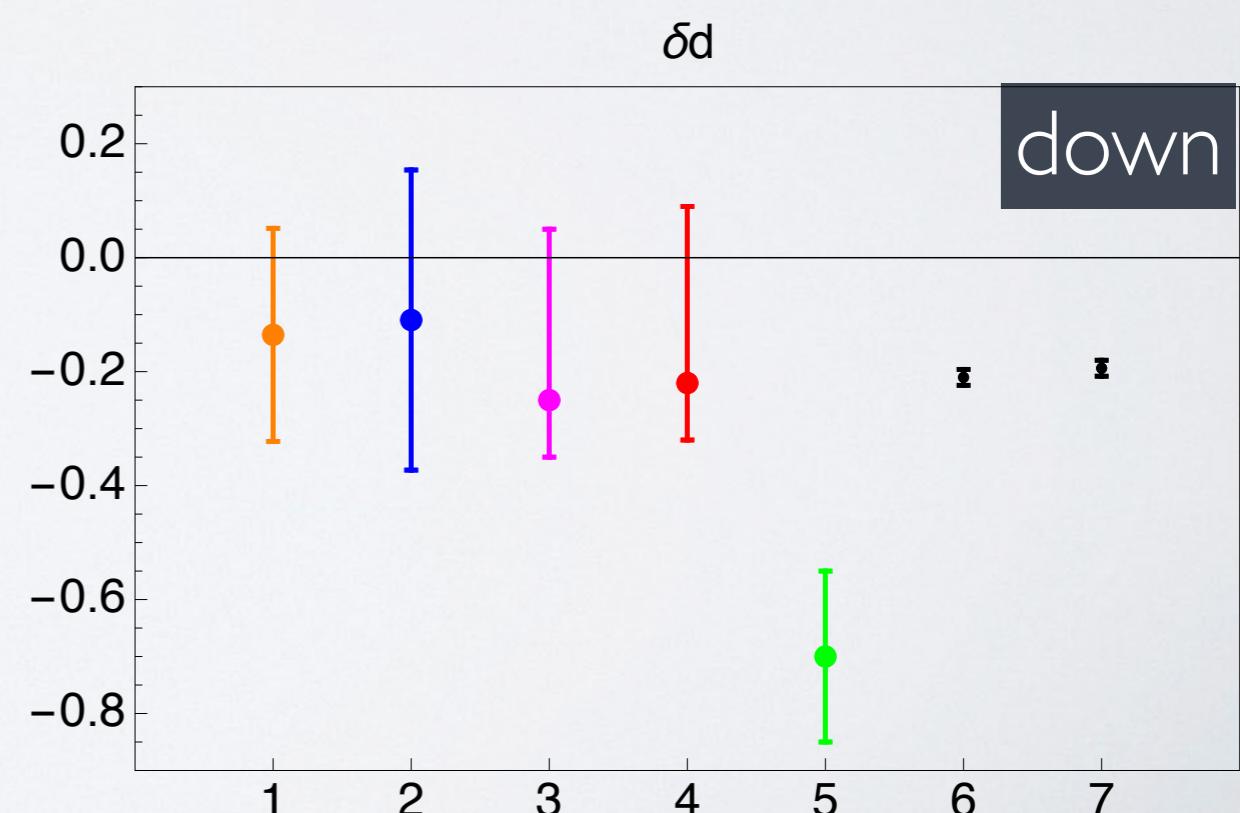


# pseudodata impact on tensor charge

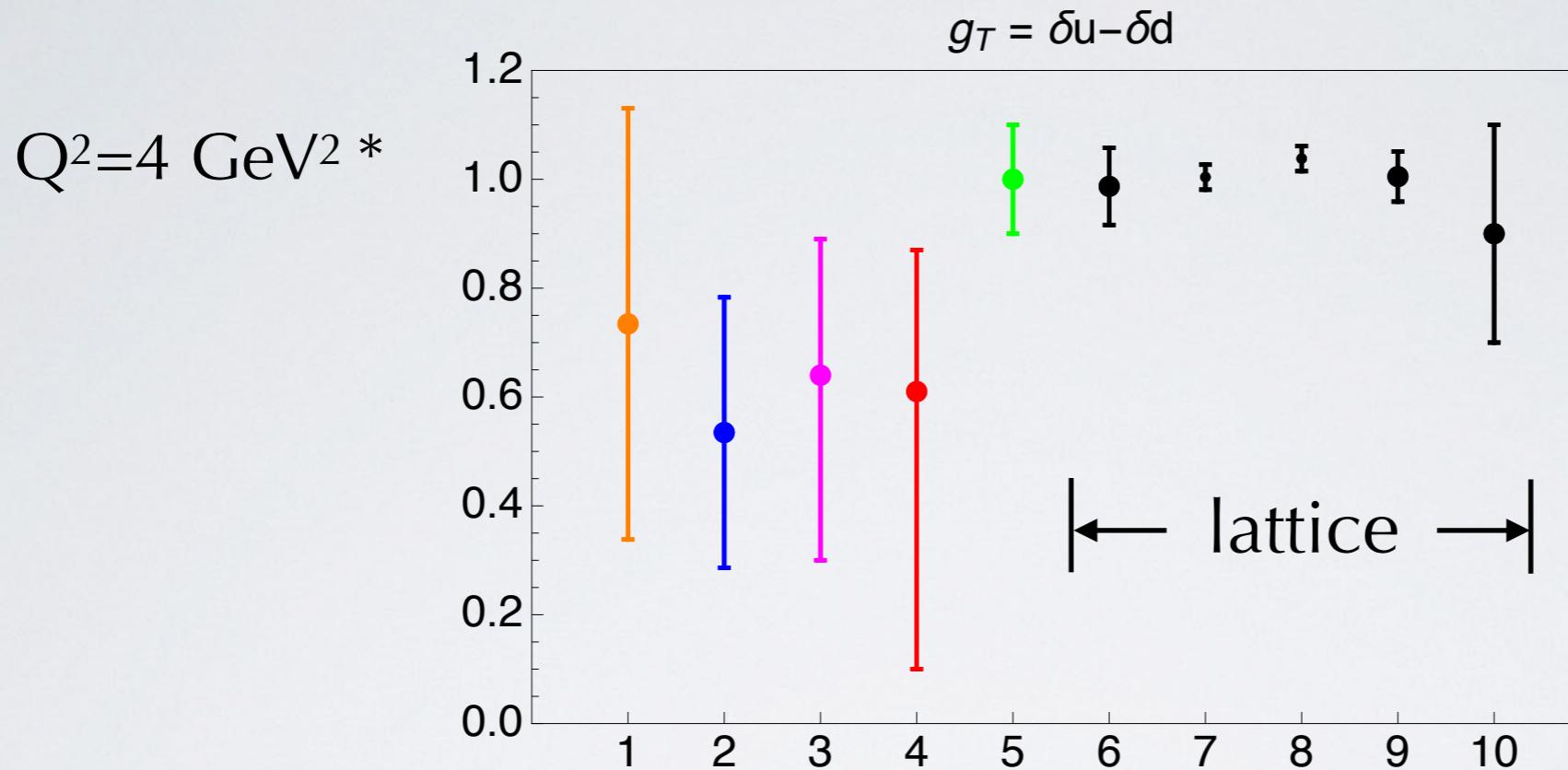


better precision on down  
larger uncertainty on up  
(“reversed role” of flavors..)

but global fit re-establishes  
full compatibility  
with lattice



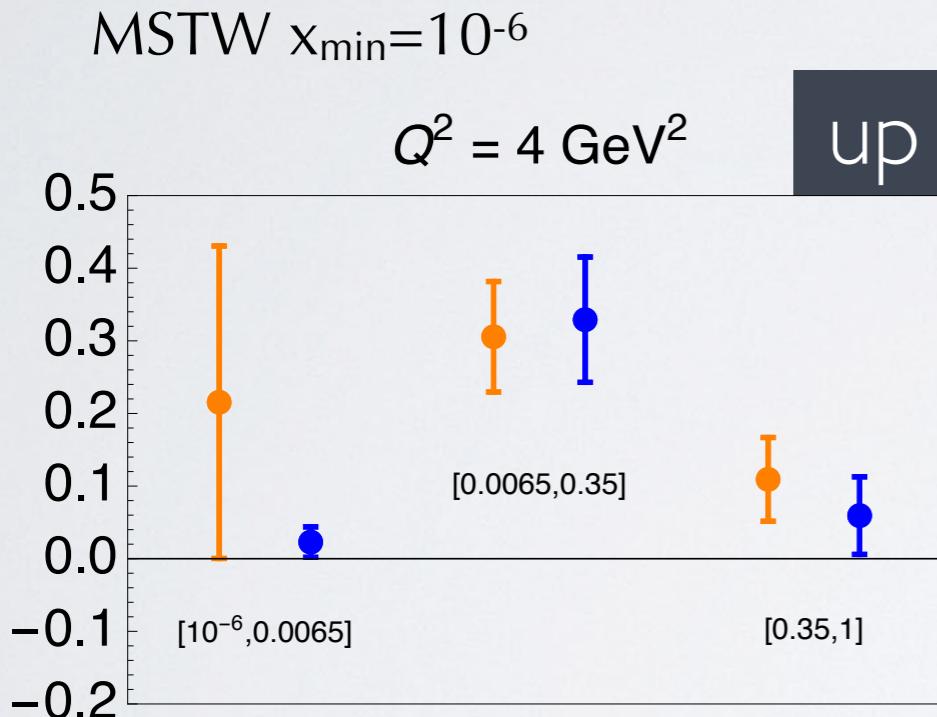
# pseudodata impact on isovector tensor charge



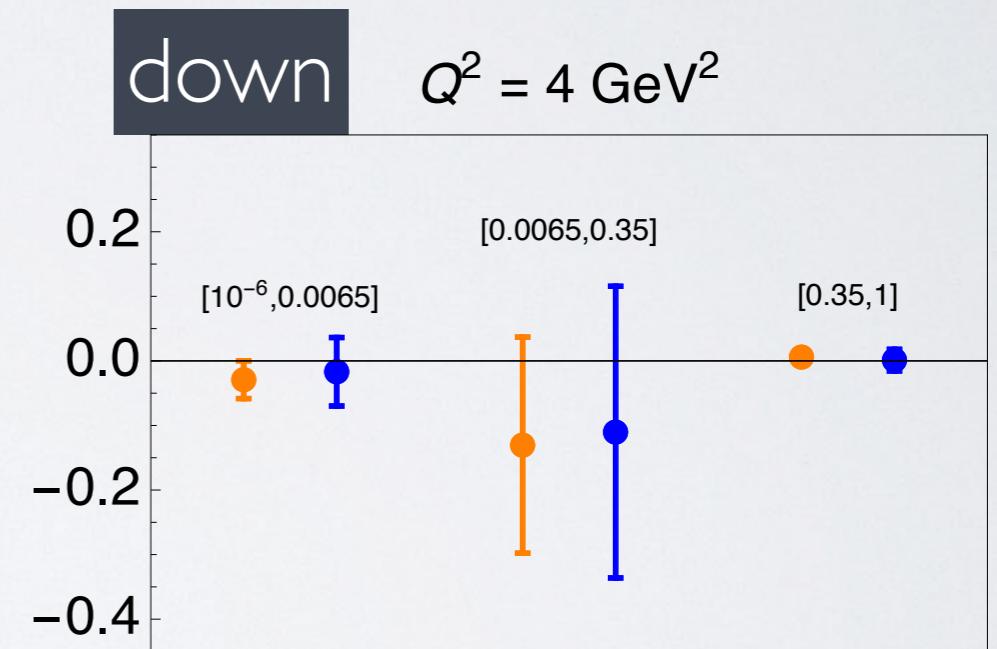
**we can have** simultaneous compatibility about  $\delta u$ ,  $\delta d$ ,  $g_T$  between lattice and our global fit, but because of large uncertainties coming from extrapolation outside the x-range of data (mainly at low x)

# impact of extrapolation outside data

$$\delta q(Q^2) = \int_{x_{\min}}^1 dx h_1^{q-\bar{q}}(x, Q^2)$$

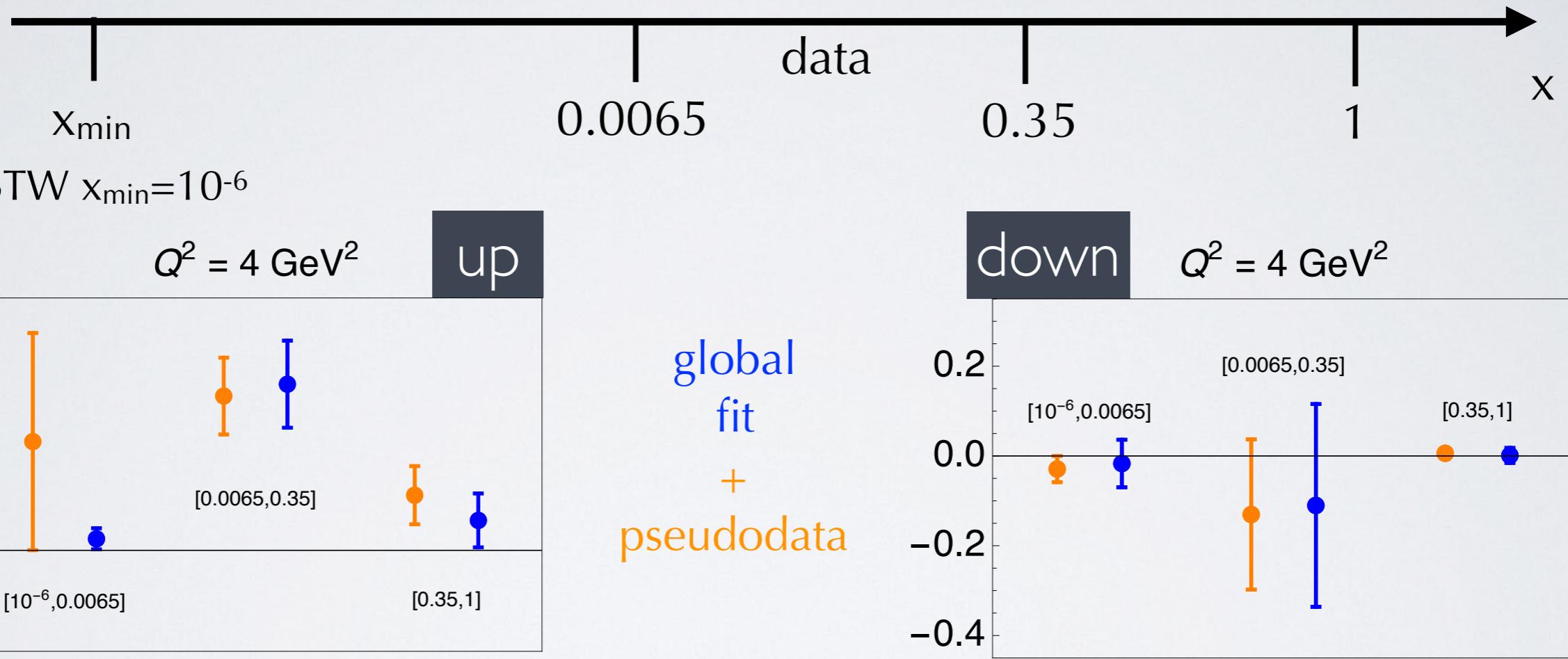


global  
fit  
+  
pseudodata



# impact of extrapolation outside data

$$\delta q(Q^2) = \int_{x_{\min}}^1 dx h_1^{q-\bar{q}}(x, Q^2)$$

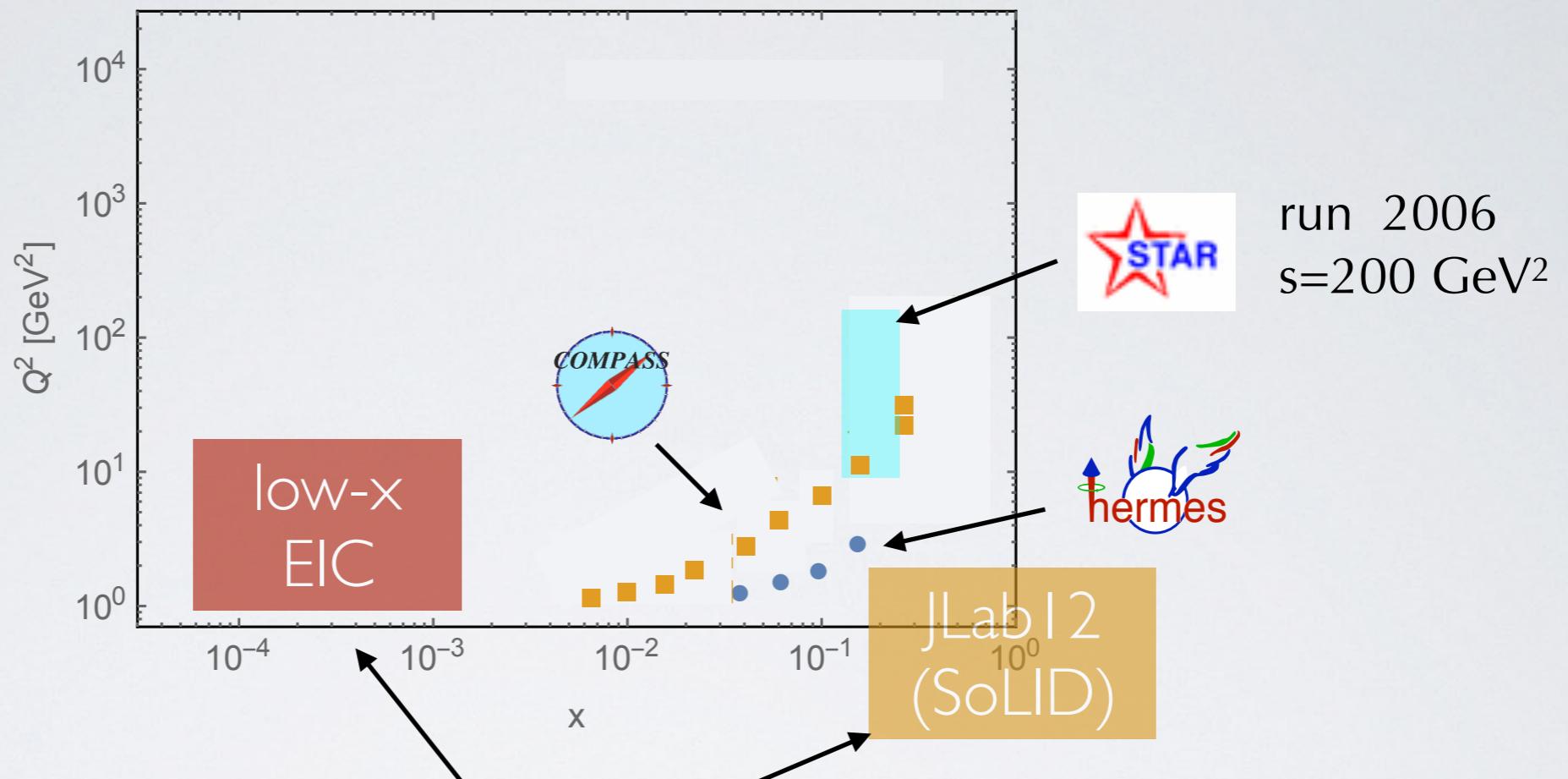


impact of pseudodata

for down: better precision everywhere

for up: large uncertainties in extrapolation at low  $x$

# more constraints on extrapolation



- of course, need more data
- theoretical constraints from low-x behavior in dipole picture  
(generalize work on helicity)  $\Delta q^S(x, Q^2) \approx \left(\frac{1}{x}\right)^{\alpha_h}$   $\alpha_h = \frac{4}{\sqrt{3}} \sqrt{\frac{\alpha_s N_c}{2\pi}}$  by  
*Kovchegov et al., P.L. B772 (17) 136*
- theoretical constraints from Burkardt-Cottingham sum rule  
*Accardi & Bacchetta, P.L. B773 (17) 632*

$$h_1^q(x) \approx x^{A_q + a_q - 1}$$

$$A_q + a_q > \frac{1}{3}$$

$$A_q + a_q > 1$$

# Conclusions

- first global fit of di-hadron inclusive data leading to extraction of transversity as a PDF in collinear framework
- inclusion of STAR p-p<sup>↑</sup> data increases precision of up channel; large uncertainty on down due to unconstrained gluon unpolarized di-hadron fragmentation function
- no apparent simultaneous compatibility with lattice for tensor charge of up, down, and isovector
- adding Compass pseudodata for deuteron reverses the scenario: better down but larger uncertainties on up from extrapolation (mainly at low x); reach overall compatibility with lattice
- need data spanning larger x range; meantime, look for other theoretical constraints on extrapolation (mostly, at low x)

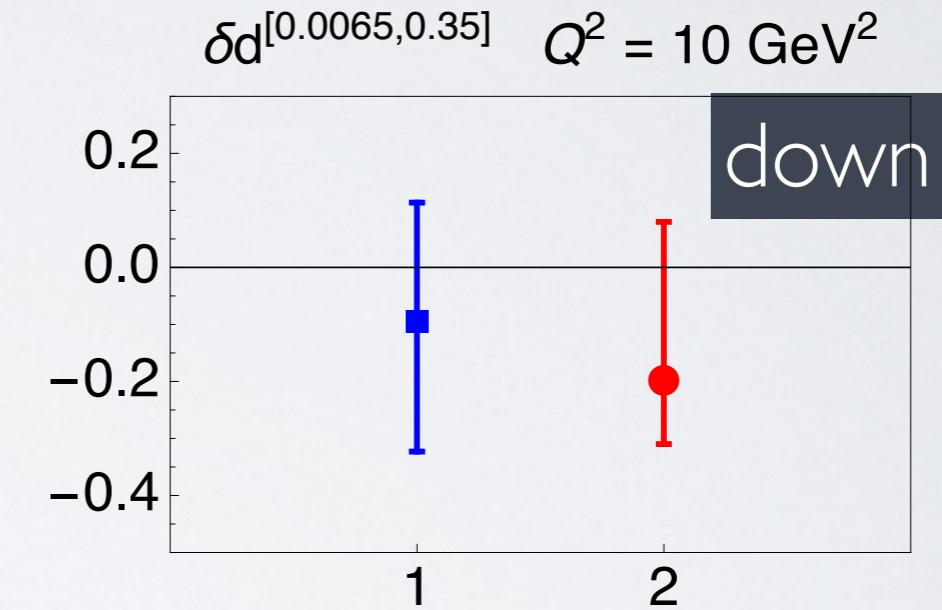
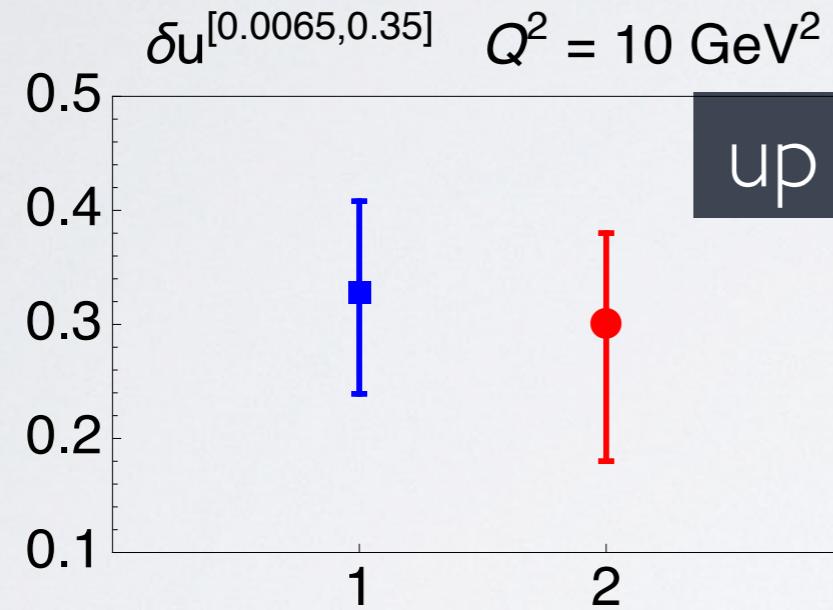
**THANK YOU**

# Back-up

tensor charge  $\delta q(Q^2) = \int dx h_1 q\bar{q} (x, Q^2)$

truncated

$$\delta q^{[0.0065, 0.35]} \quad Q^2 = 10$$



global fit

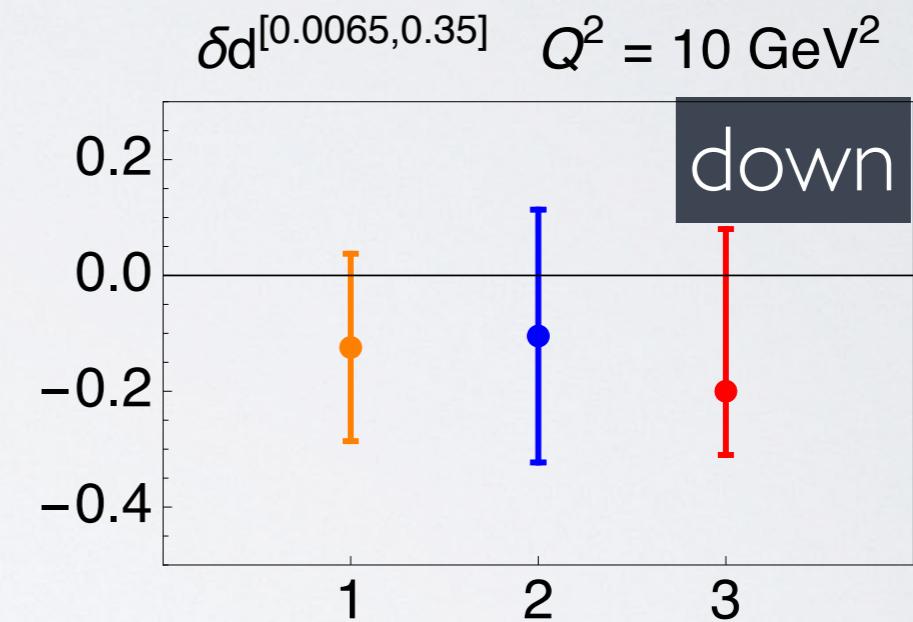
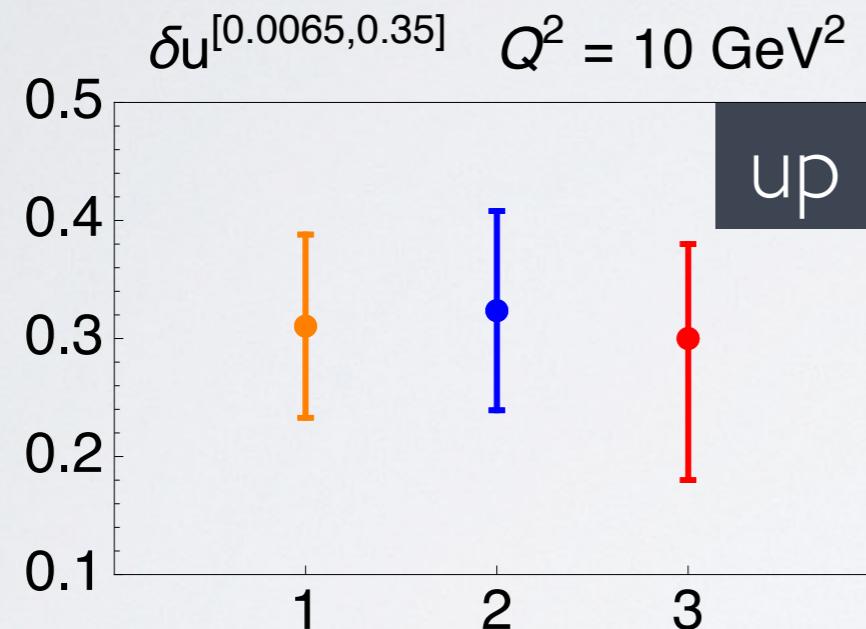
*Radici & Bacchetta,  
P.R.L. **120** (18) 192001*

TMD fit

*Kang et al.,  
P.R. D93 (16) 014009*

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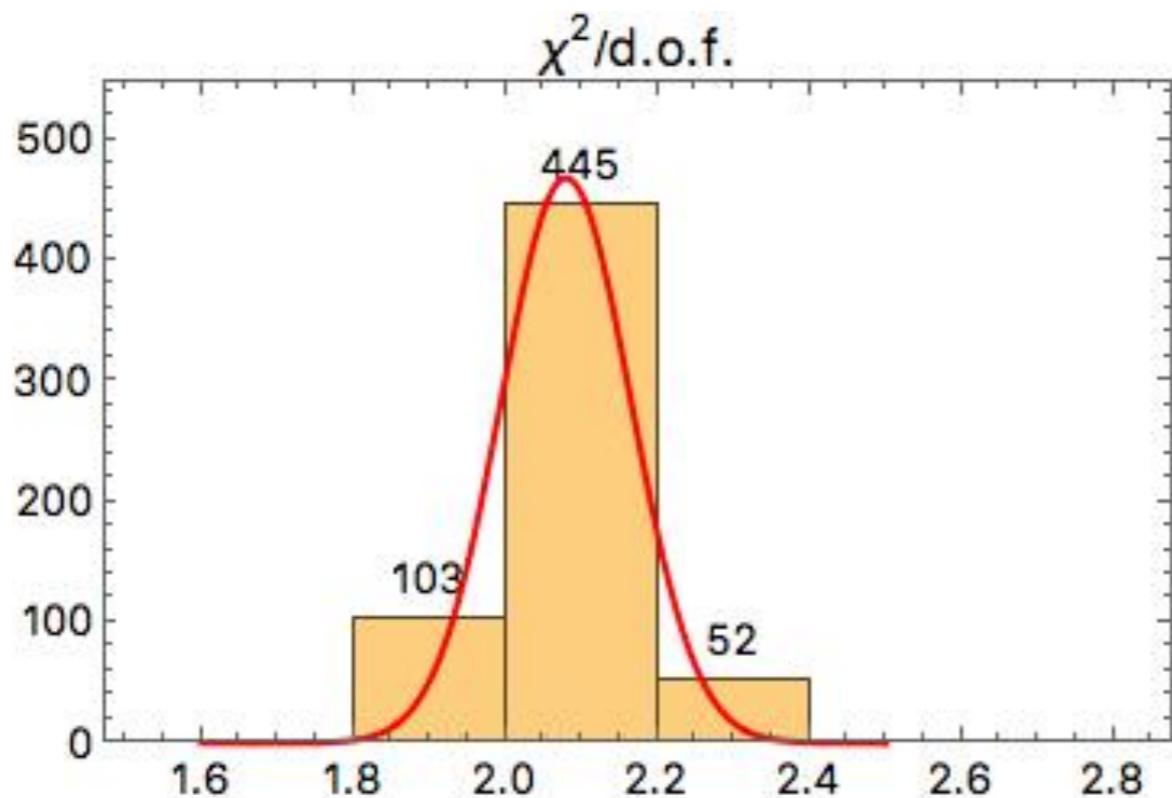


+  
pseudodata

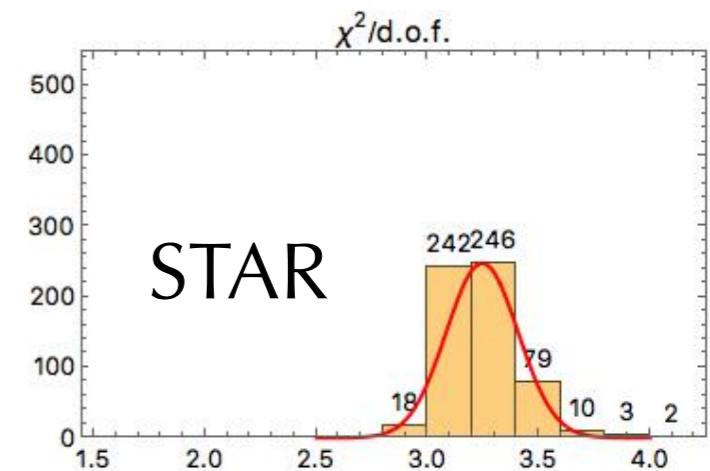
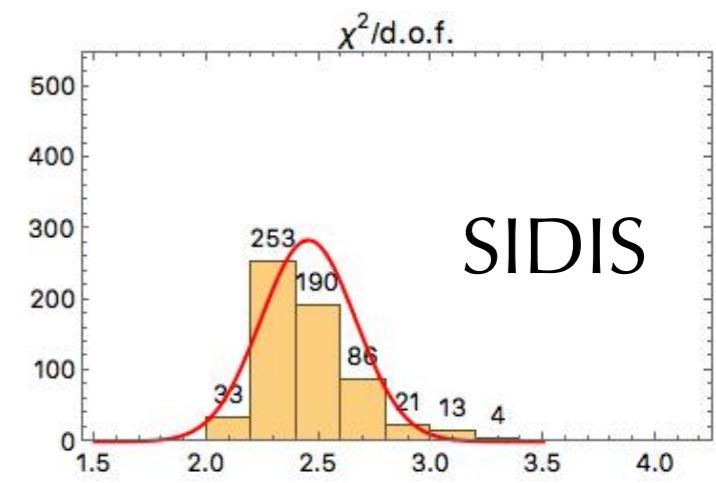
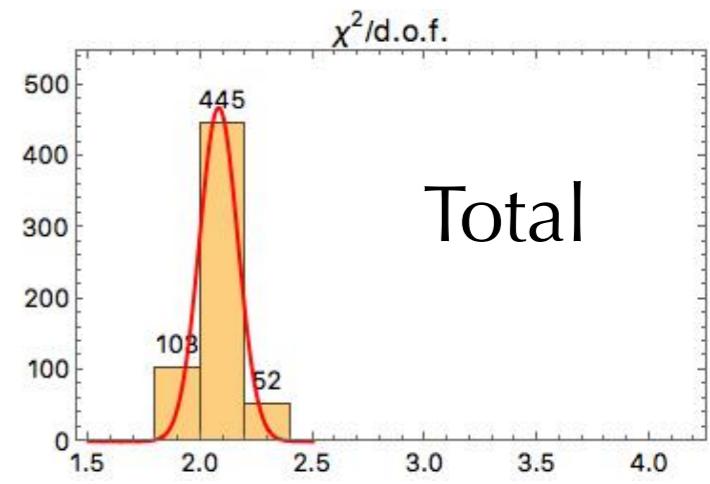
global fit  
*Radici & Bacchetta,  
P.R.L. **120** (18) 192001*

TMD fit  
*Kang et al.,  
P.R. D93 (16) 014009*

# $\chi^2$ of the fit

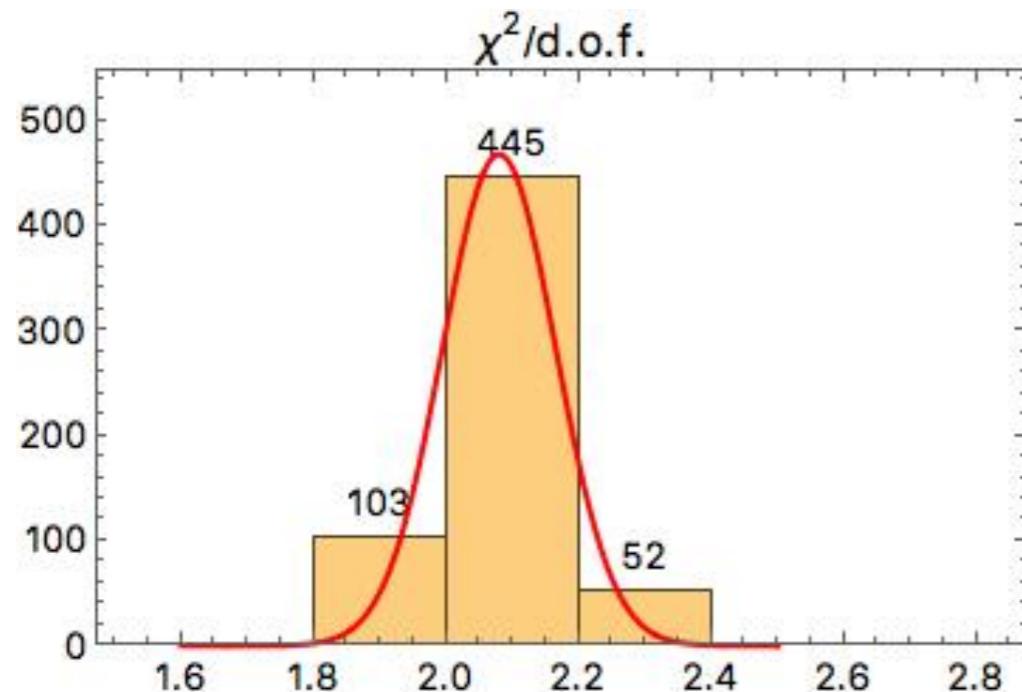


$$\chi^2/\text{dof} = 2.08 \pm 0.09$$



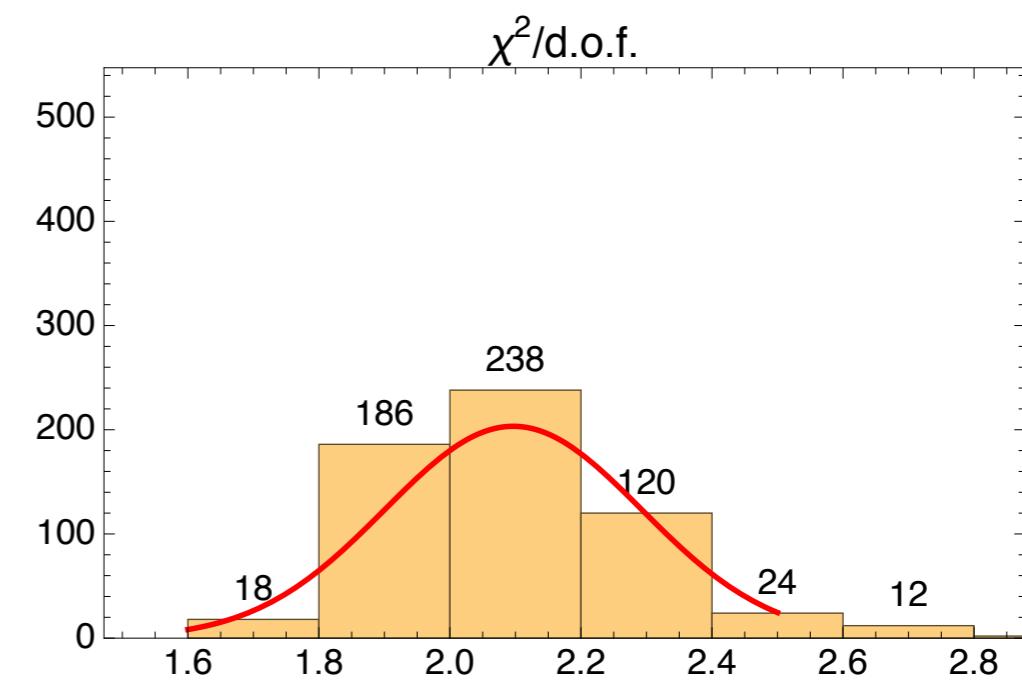
# $\chi^2$ of the fit

global fit



$$\chi^2/\text{dof} = 2.08 \pm 0.09$$

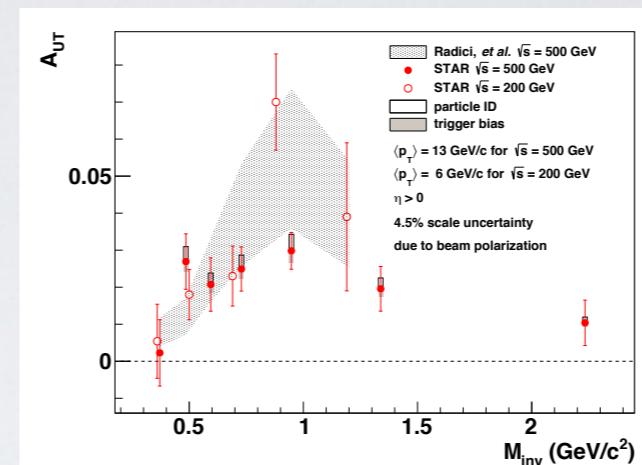
+ pseudodata



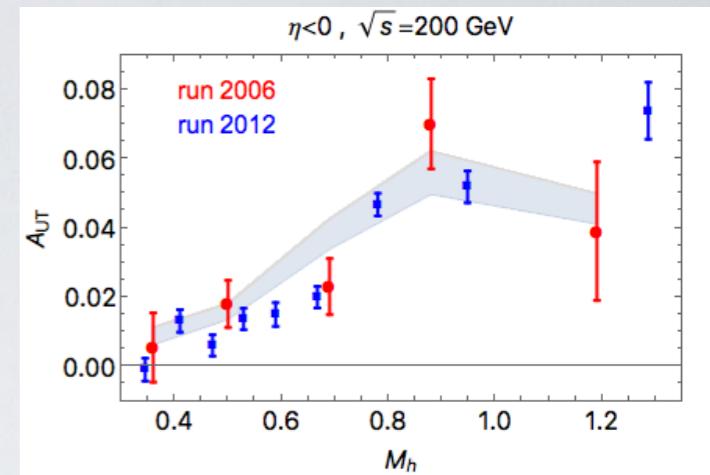
$$\chi^2/\text{dof} = 2.10 \pm 0.20$$

# To do list

- use also other (multi-dimensional) data from STAR run 2011 ( $s=500$ ) and (later) run 2012 ( $s=200$ )



Adamczyk et al. (STAR), P.L. **B780** (18) 332



Radici et al., P.R. **D94** (16) 034012

- need data on  $p+p \rightarrow (\pi\pi) X$  constrains gluon  $D_{1g}$
- refit di-hadron fragmentation functions using new data:  
 $e^+e^- \rightarrow (\pi\pi) X$  constrains  $D_{1q}$   
 (currently only by Montecarlo)
- use COMPASS data on  $\pi K$  and  $KK$  channels, and from  $\Lambda^\uparrow$  fragmentation:  
 constrain strange contribution ?
- explore other channels, like inclusive DIS via Jet fragm. funct.'s



Seidl et al.,  
 P.R. **D96** (17) 032005