# Beam Spin Asymmetries for Positively Charged Kaons

David Riser, University of Connecticut in collaboration with Kyungseon Joo, Nick Markov



- Motivation
- Measurement of BSA
- Analysis of BSA Moments
- Results, Conclusion, Outlook



Currently there is big interest in 3-D nucleon structure in the form of GPDs and TMDs



"Leading-Twist" TMD Quark Distributions

The **SIDIS** cross section can be expressed in terms of model independent structure functions

$$\frac{d\sigma}{dx_B \, dQ^2 \, dz \, d\phi_h \, dp_{h\perp}^2} = K(x, y, Q^2) \Big\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \Big\}$$
(without polarized target)

Mulders, Tangerman (1995) Complete tree-level result for polarized deep-inelastic lepto-production.



TMD and Fragmentation Functions

$$C[\omega fD] = x \sum_{a} e_{a}^{2} \int d^{2} \vec{p}_{\perp} d^{2} \vec{k}_{\perp} \delta^{(2)} \left( \vec{p}_{\perp} - \vec{k}_{\perp} - \vec{P}_{h\perp}/z \right) \omega(\vec{p}_{\perp}, \vec{k}_{\perp}) f^{a}(x, p_{\perp}^{2}) D^{a}(z, k_{\perp}^{2})$$

$$F_{LU}^{\sin\phi} = \frac{2M}{Q} \mathcal{C} \left( -\frac{\hat{\mathbf{h}} \cdot \mathbf{k_T}}{M_h} \left( x e H_1^{\perp} + \frac{M_h}{M} f_1 \frac{\tilde{G}^{\perp}}{z} \right) + \frac{\hat{\mathbf{h}} \cdot \mathbf{p_T}}{M} \left( x g^{\perp} D_1 + \frac{M_h}{M} h_1^{\perp} \frac{\tilde{E}}{z} \right) \right)$$



**TMD** and Fragmentation Functions

$$C[\omega fD] = x \sum_{a} e_{a}^{2} \int d^{2} \vec{p}_{\perp} d^{2} \vec{k}_{\perp} \delta^{(2)} \left( \vec{p}_{\perp} - \vec{k}_{\perp} - \vec{P}_{h\perp}/z \right) \omega(\vec{p}_{\perp}, \vec{k}_{\perp}) f^{a}(x, p_{\perp}^{2}) D^{a}(z, k_{\perp}^{2})$$



Structure function has a **twist-3** piece in every term



Beam Spin Asymmetry (BSA) measurements are a good tool for extracting "moments".

$$BSA = \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-} = \frac{A_{LU}^{\sin\phi}\sin\phi}{1 + A_{UU}^{\cos\phi}\cos\phi + A_{UU}^{\cos(2\phi)}\cos(2\phi)}$$

Experimentally, they are measured by selecting your event of interest and calculating the quantity defined below.

$$BSA_{i} = \frac{1}{P_{e}} \frac{N_{i}^{+} - N_{i}^{-}}{N_{i}^{+} + N_{i}^{-}}$$

If acceptance corrections don't depend on helicity, the correction cancels.



# **Experimental Details**

# Run period: E1-F

- Beam energy 5.5 GeV
- Torus field run at 60% (2250 A)
- Negatively charged particles Inbending
- Over 1B event triggers
- LH\_2 target, 5 cm







A brief overview of electron identification.

- Sampling fraction cut (momentum dependent)
- Geometrical fiducial cuts (DC, EC, CC)
- Cherenkov PMT matching cuts (theta, phi)
- EC energy deposition cut
- Z-vertex cut
- Vertex and kinematic corrections done after ID.



# A brief overview of kaon identification.



Drift chamber region 1 fiducial cut  $10^{4}$ Electron vertex difference cut Likelihood ratio maximization condition  $\bullet$ 0.8  $10^{3}$ • Significance level cut 0.6  $rac{\mathcal{L}_h}{\mathcal{L}_\pi + \mathcal{L}_K + \mathcal{L}_p}$  $\sim$ 10<sup>2</sup> 0.4 10 0.2  $\alpha = 1 - \int_{\mu - \beta_{obs}}^{\mu + \beta_{obs}} P(\beta; p, h) d\beta$  $0^{1}_{0}_{0}_{0.5}_{0.5}_{1}_{1.5}_{1.5}_{2}_{2.5}_{2.5}_{2.5}_{3}$ <u>...</u> 3.5 4.5 p (GeV/c)

1.2







 $Q^2 > 1, W > 2.0$ 

 $M_X(ep \to eK^+X) > 1.27$ 

12 bins in phi 6 bins in other axes



### More kinematic distributions





David Riser University of Connecticut

#### Result of measurement with statistical uncertainties



Shown above: Different x bins.



#### Result of measurement with statistical uncertainties





Systematic uncertainty - Result of uncertainty in estimate of systematic effects arising from selection bias, backgrounds, calibrations etc.

Source	Cut type	Source	Size
Kaon ID SL	Min	Beam polarization	3%
CC theta match	Min/Max	Kaon momentum TBD In the case that none of the parameters are correlated, $\rho_{ij} = \delta_{ij}$ $(\delta \mathcal{O})^2 = \sum_{i=1}^{N} (\mathcal{O}(\theta_i + \delta \theta_i/2) - \mathcal{O}(\theta_i - \delta \theta_i/2))$	ТРП
DC R1 Fid	Min		
DC R3 Fid	Min		
Sampling frac.	Min/Max		
EC Edep	Min		
EC-U	Min/Max		s
EC-V	Max		$o_{ij}$
EC-W	Max		$(2) - \mathcal{O}(\theta_i - \delta \theta_i/2))^2$
z-vertex	Min/Max	$i{=}1$	



x axis





z axis





Momentum range of accepted kaon is varied to check impact of contamination at higher momentum. Difference is included as systematic.







Shown above: Different x bins.





Shown above: Different z bins.



Extraction of Moments



Recall what we are measuring

$$BSA = \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-} = \frac{A_{LU}^{\sin\phi}\sin\phi}{1 + A_{UU}^{\cos\phi}\cos\phi + A_{UU}^{\cos(2\phi)}\cos(2\phi)}$$

We minimize the total chi-2 (with respect to the parameters a) for each kinematic bin by gradient descent

$$\chi^2 = \sum_{i=1}^{n_{\phi}} \frac{(A_{obs}(\phi_i) - A_{pred}(\phi_i, \vec{a}))^2}{\sigma_{stat}^2 + \sigma_{sys}^2}$$

One would be very happy to extract all 3 moments, in some cases it may be possible.







(Sine, Full) Results almost identical.





Points of conclusion

- Measured BSA for 4 kinematic variables on a sample of SIDIS kaons (+), statistical and systematic uncertainty is calculated for each point.
  - Most systematic uncertainty comes from kaon ID at high momentum, but it is still reasonably small.
- An analysis of moments has been performed for each kinematic point.
- Theoretical predictions are in progress for the kinematic dependence of this measurement based on TMD PDF/FF functions



Thank you for your time.

Thanks to the following people for helping with this analysis: Kyungseon Joo, Harut Avakian, Nick Markov, Andrey Kim, Nobuo Sato, Nathan Harrison, Kemal Tezgin, Frank Cao, and Brandon Clary

For various parts of the analysis I thank previous workers: Wes Gohn, Marco Mirazita, Nathan Harrison



Extra Slides



Parameter Estimation

Single fit - Minimize the chi-2 between your data points and the predictions from your model.

$$\chi^2 = \sum_{i=1}^{n_{\phi}} \frac{(A_{obs}(\phi_i) - A_{pred}(\phi_i, \vec{a}))^2}{\sigma_{stat}^2 + \sigma_{sys}^2}$$

This is done using gardient descent. The errors and correlations on/between the fit parameters a are calculated by using the topology of the chi-2 function around the parameter minima. More explicitly the inverse Hessian matrix.

$$\sigma_{ij} = H_{ij}^{-1} = C_{ij}$$



### Parameter Estimation (2)

Give a probabilistic interpretation of our measurement and give a probabilistic interpretation to the parameters of our model for the result (using Bayes Theorem).

$$P(\vec{a}|A_i) = \frac{P(A_i|\vec{a})P(\vec{a})}{P(A_i)}$$

Use Markov chain MC (MCMC) to explore the parameter space {a} (Metropolis algorithm).

$$P_{Accept} = \frac{P(A_i | \vec{a}_{n+1}) P(\vec{a}_{n+1})}{P(A_i | \vec{a}_n) P(\vec{a}_n)}$$



Max likelihood ratio method for charged hadrons

$$\mathcal{L}_h = \prod_{i=1}^N P_i(x_i; p, h) \qquad \qquad \frac{\mathcal{L}_h}{\mathcal{L}_\pi + \mathcal{L}_K + \mathcal{L}_p}$$

Simple case of one distribution used

$$P(\beta; p, h) = \frac{1}{\sqrt{2\pi}\sigma_{\beta}(p, h)} exp\left\{-\frac{1}{2}\left(\frac{\beta - \mu_{\beta}(p, h)}{\sigma_{\beta}(p, h)}\right)^{2}\right\}$$

"Significance", use to tighten/loosen cut on beta, discard low likelihood events (positron example).

$$\alpha = 1 - \int_{\mu - \beta_{obs}}^{\mu + \beta_{obs}} P(\beta; p, h) d\beta$$



Measurement of the beam spin asymmetry is done experimentally by recording events with different electron helicity states and counting the ratio below. Helicity flipping occurs at high enough frequency that acceptance effects are expected to cancel.

$$A_{LU}^{\sin\phi} = \frac{1}{P_e} \frac{N^+ - N^-}{N^+ + N^-}$$

The **average beam polarization** was determined to be (75 +/- 3) %, and the wave-plate position was determined as a function of run number by analyzing the sine phi moment for positive pions.





Treatment of systematic uncertainties

Analysis depends on some set of parameters, cut values, calibration values, etc.

 $\vec{\theta} = (\theta_i, \theta_2, ..., \theta_N)$ 

The values of these parameters impact the outcome of the measurement of your observable. The standard formula for "error propagation" is,

$$(\delta \mathcal{O})^2 = \sum_{i=1}^N \sum_{j=1}^N \frac{\partial \mathcal{O}}{\partial \theta_i} \frac{\partial \mathcal{O}}{\partial \theta_j} \rho_{ij} \delta \theta_i \delta \theta_j$$

However it is not analytically possible to find in most cases the relationship between the parameter and the observable.

$$\frac{\partial(BSA)}{\partial(ECU\ Cut\ Value)} = ?$$



Numerically find the derivative around the nominal parameter value





$$\frac{\partial \mathcal{O}}{\partial \theta_i} \frac{\partial \mathcal{O}}{\partial \theta_j} \rho_{ij} \delta \theta_i \delta \theta_j \approx \left( \mathcal{O}(\theta_i + \delta \theta_i/2) - \mathcal{O}(\theta_i - \delta \theta_i/2) \right) \\ \times \left( \mathcal{O}(\theta_j + \delta \theta_j/2) - \mathcal{O}(\theta_j - \delta \theta_j/2) \right) \rho_{ij}$$

In the case that none of the parameters are correlated,  $ho_{ij}=\delta_{ij}$ 

$$(\delta \mathcal{O})^2 = \sum_{i=1}^{N} (\mathcal{O}(\theta_i + \delta \theta_i/2) - \mathcal{O}(\theta_i - \delta \theta_i/2))^2$$





#### Electron Momentum Corrections Before and After













