

CLAS Collaboration Meeting
Deep Processes Working Group

Deep Virtual Production of Pion Pairs

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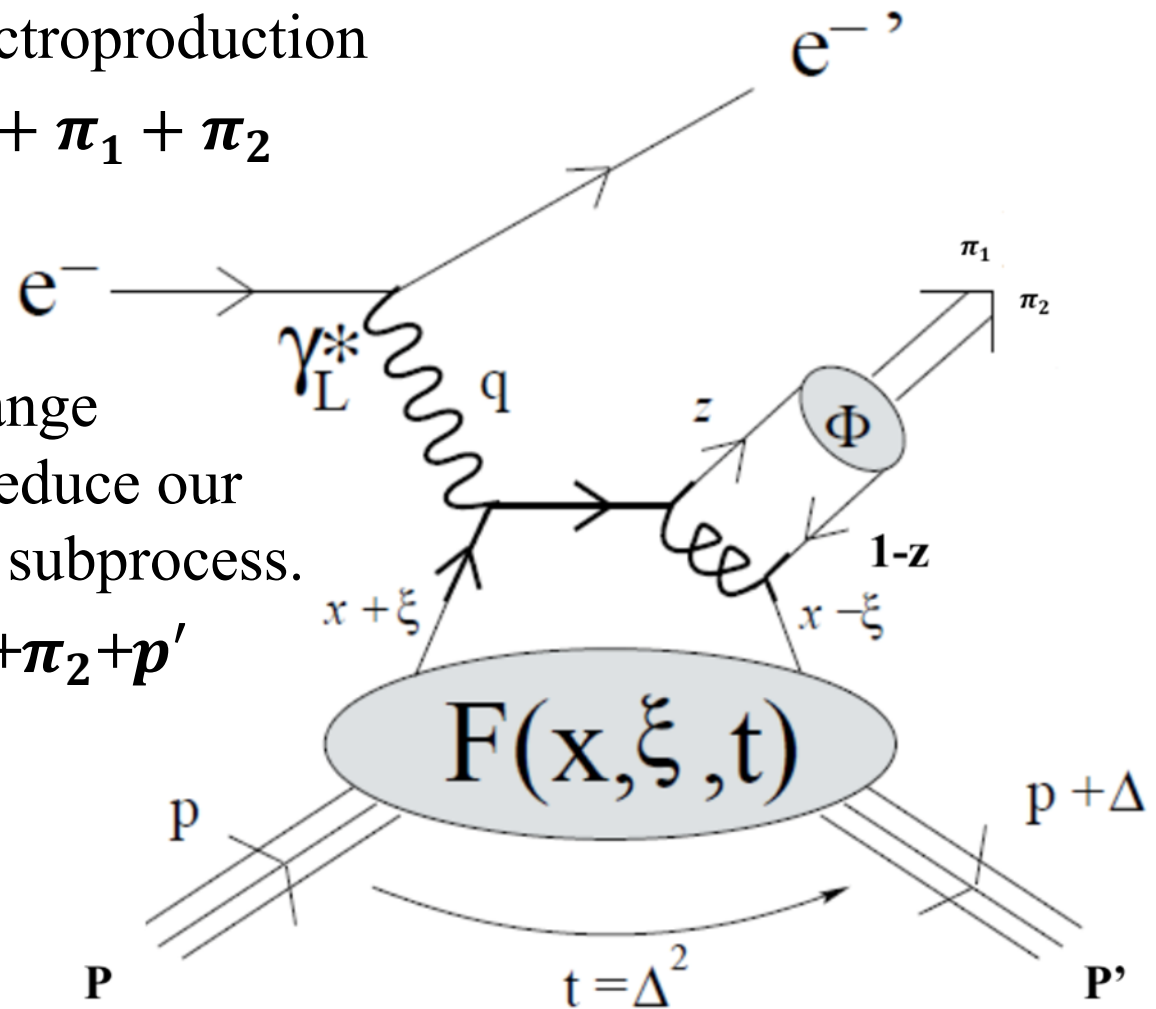
- We are mainly considering two reactions, Charged and Neutral Pion Pairs
 - $ep \rightarrow e'p' \pi^+ \pi^-$
 - Isospin $I=1$, angular momentum $J=1$
 - $\rho(770)$
 - Isospin $I=0$, angular momentum $J=0$
 - $f_0(500) = \sigma, f_0(980)$
 - $ep \rightarrow e'p' \pi^0 \pi^0$
 - Isospin zero, spin zero channel ($I:J=0:0$)
 - $f_0(500) = \sigma, f_0(980)$

- Exclusive two-pion electroproduction

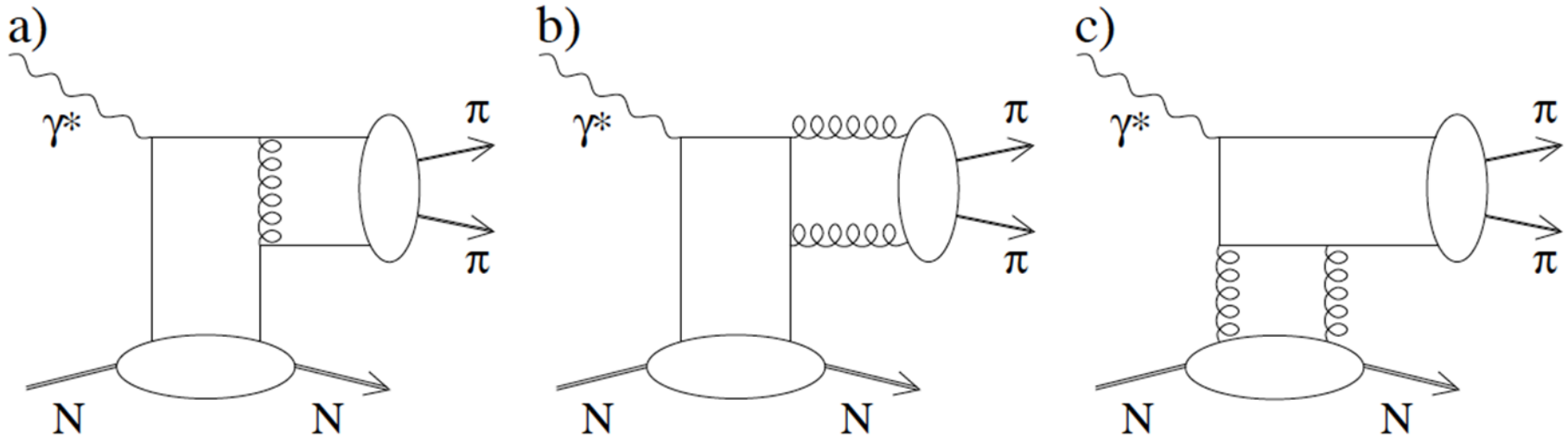
$$e + p \rightarrow e' + p' + \pi_1 + \pi_2$$

- In the one-photon exchange approximation we can reduce our analysis to the hadronic subprocess.

$$\gamma^* + p \rightarrow \pi_1 + \pi_2 + p'$$



- Leading order diagrams for exclusive deep virtual production of two pions



B. Lehmann-Dronke *et al.*, Phys Lett B **475** (2000) 147

B. Lehmann-Dronke *et al.*, Phys Rev D, **63** (2001) 114001

Neutral mesonic final state: $\pi^+\pi^-$ or $\pi^0\pi^0$

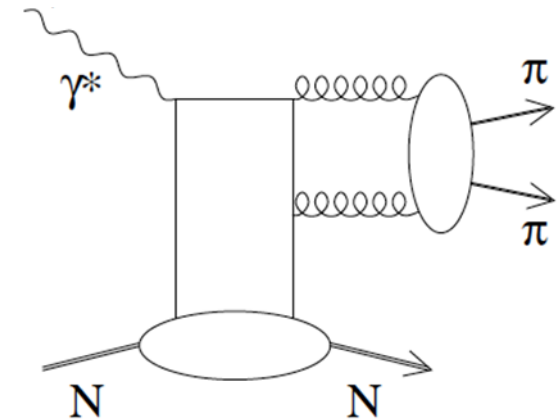
- a) [Flavor-Diagonal quark-GPD] \otimes [$q\bar{q}$ -Two-Pion Distribution Amplitude (DA)]
- b) [Flavor-Diagonal quark-GPD] \otimes [gluon-Two-Pion Distribution Amplitude(DA)]
- c) [Gluon-GPD] \otimes [$q\bar{q}$ -Two-Pion Distribution Amplitude (DA)]

- σ -meson Asymptotic Distribution Amplitudes:

- $\Phi_{\text{gluon}} = 2 \Phi_{\text{qq}}$

- σ -meson: $f_0(500)$ well established.

- ***Pole = $(450 \pm 20) \text{ MeV} - i(275 \pm 12) \text{ MeV}$***



- Microscopic structure of $f_0(500)$ not well understood.

- $q\bar{q} : {}^3P_0$

- Tetraquark

- $\pi\pi$ -molecule

- Glueball

- Superposition of all of the above

- Deep sigma-production offers intriguing evidence for gluonic content of $f_0(500)$.

- **Deep Virtual $\pi\pi$ Production Amplitude**

$$\mathcal{M} = \sum_{\substack{I \\ \lambda_N, \lambda_\pi \in (q\bar{q}, g)}} \int d\tau dz \text{GPD}_{\lambda_N}(\tau, \xi, t) \odot S_{\lambda_N, \lambda_\pi}(\tau, z, \xi) \odot \text{DA}_{\lambda_\pi}^I(z, \zeta; m_{\pi\pi}; \theta^*)$$

$$\mathcal{M} = \sum_{\substack{J^{\pi}; I \\ \lambda_N, \lambda_\pi \in (q\bar{q}, g)}} \int d\tau dz \text{GPD}_{\lambda_N}(\tau, \xi, t) \odot S_{\lambda_N, \lambda_\pi}(\tau, z, \xi) \odot \text{DA}_{\lambda_\pi}^I(z, \zeta) P_J(\cos(\theta^*)) \Omega_{J; I}(m_{\pi\pi})$$

- **Kinematics**

$$\xi \sim \frac{x_B}{2 - x_B}$$

$$t = (q - p_{\pi\pi})^2 = (P'_p - P_p)^2$$

$$\zeta, (1 - \zeta) = \frac{1}{2} [1 \pm \beta^* \cos \theta^*] = \text{pion lightcone momentum fractions}$$

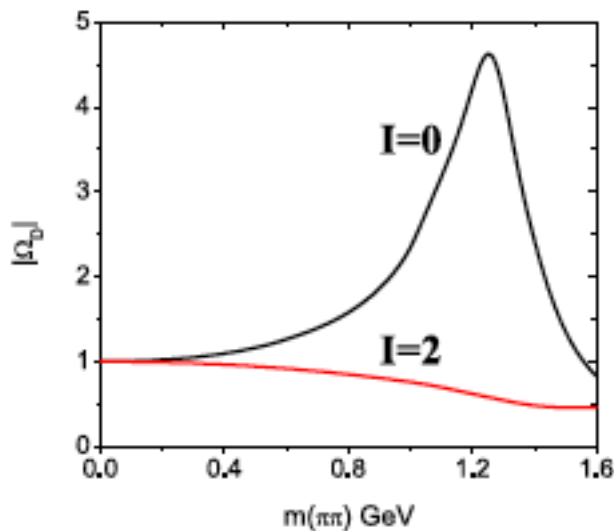
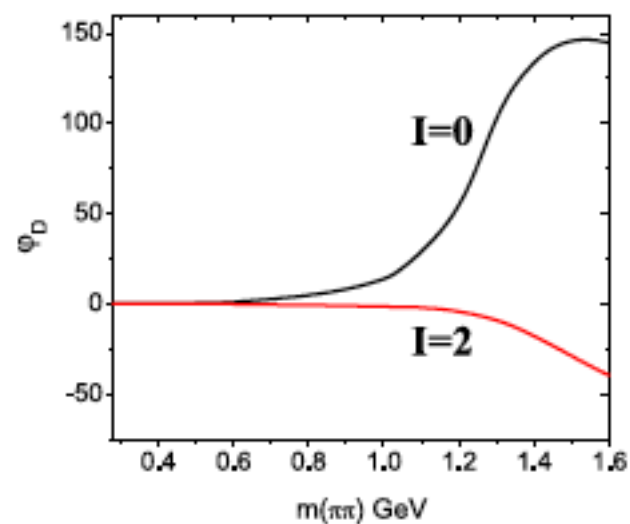
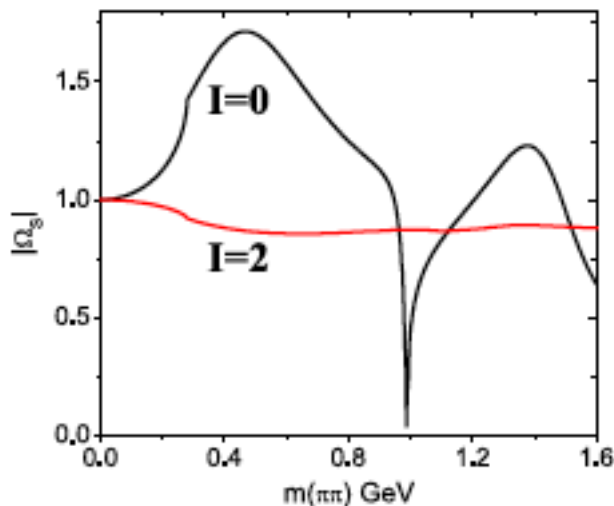
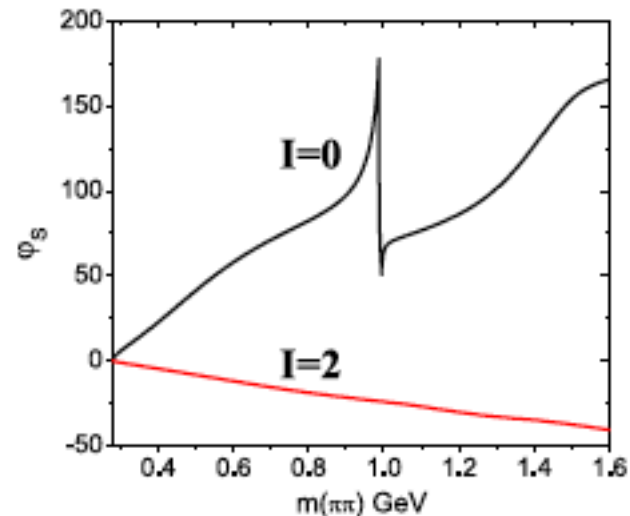
$$\beta^* = \text{pion velocity in } \pi\pi \text{ rest frame}$$

$$\theta^* = \text{pion polar angle in } \pi\pi \text{ rest frame}$$

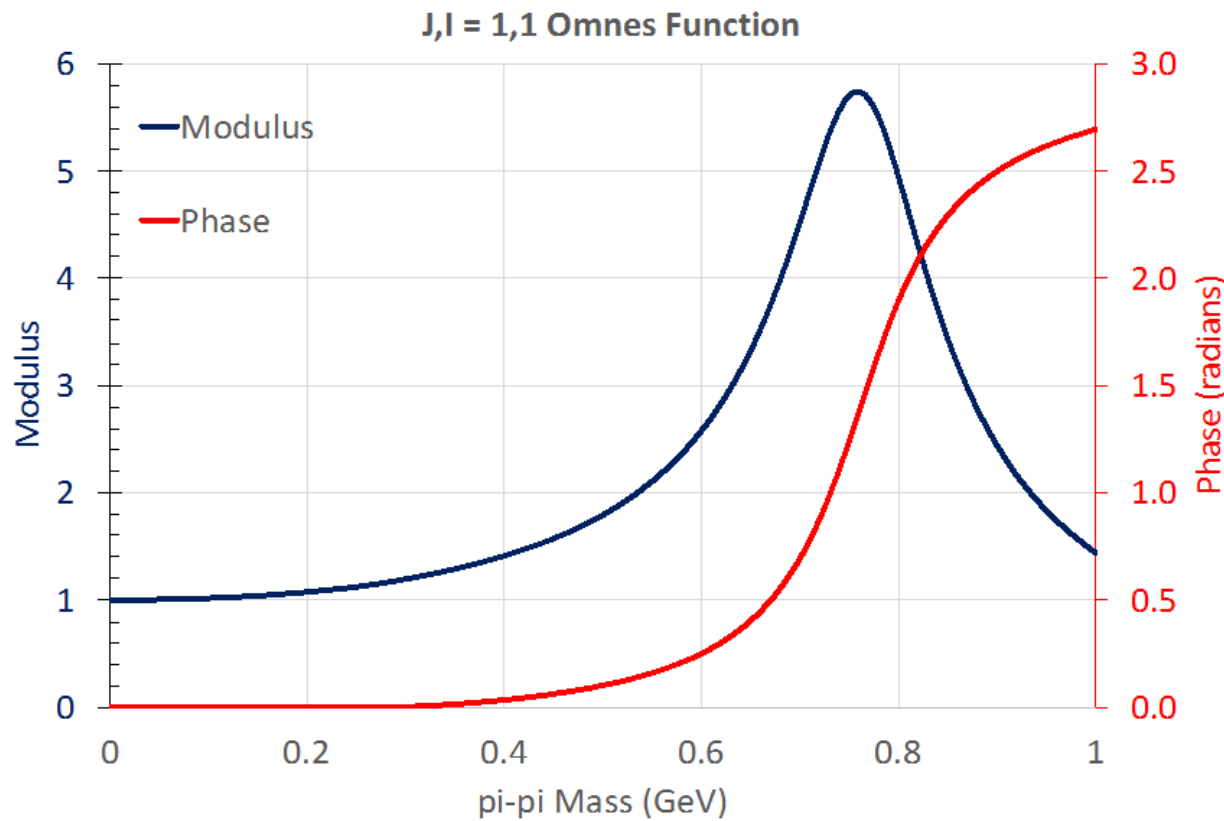
- **Dynamics**

- $S(\tau, z; \xi)$ = Hard scattering amplitude (quark-gluon propagators)
- $\Omega_{J; I}$ = Omnès-function, derived from $\pi\pi$ phase shifts
- τ = average momentum fraction of parton in nucleon
- z = momentum fraction of parton in $\pi\pi$ DA

$\pi\pi$ Mass Distribution (Omnès F'n)



- L.Dai, M.Pennington, Phys Rev D **90** 036004 (2014)
- L=0
 - $f_0(500)$
 - $f_0(980)$
 - *Small I=2 non-resonant*
- L=2
 - $f_2(1270)$
 - *Small I=2 non-resonant*



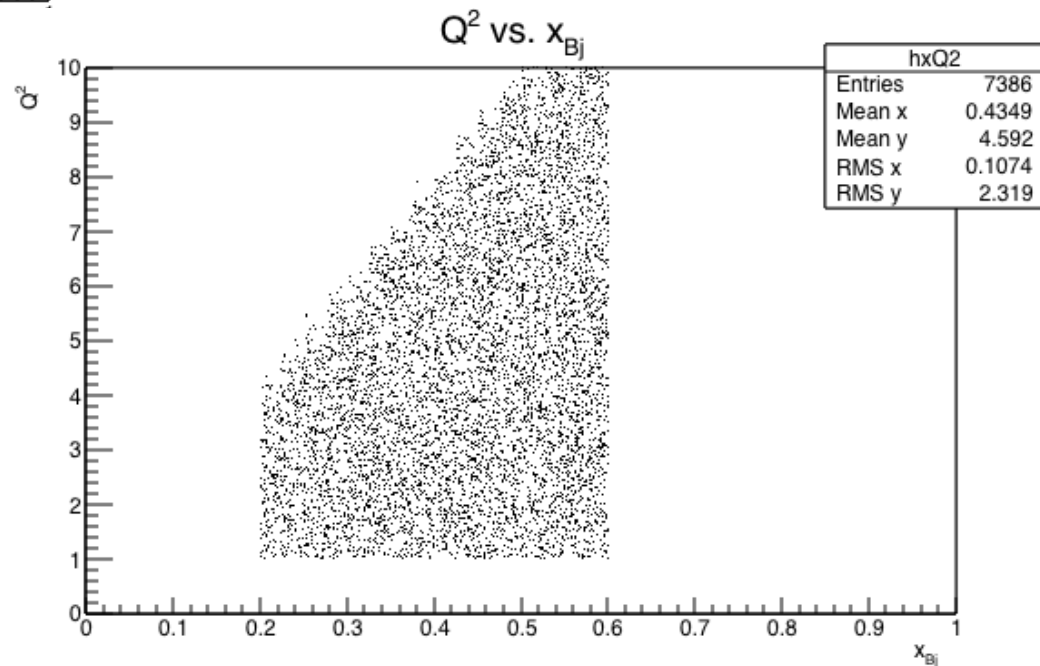
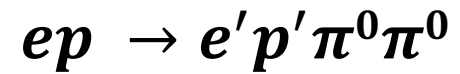
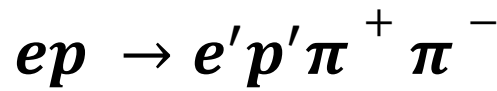
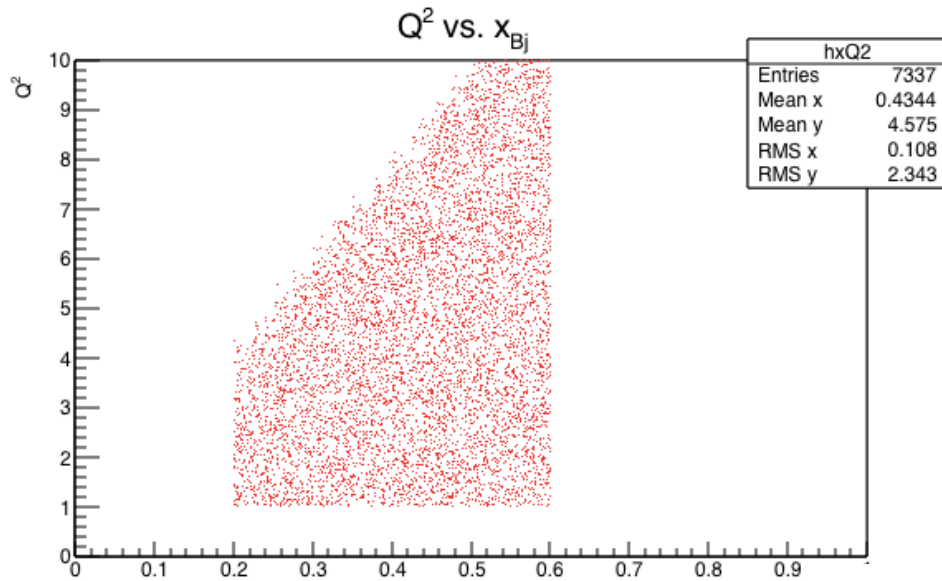
• L.Dai, M.Pennington, Phys Rev D **90** 036004 (2014)

$$\Omega_l^I(m_{\pi\pi}) = \exp \left\{ i\delta_l^I(m_{\pi\pi}) + \frac{m_{\pi\pi}^2}{\pi} \Re \left[\int_{4m_\pi^2}^{\infty} ds \frac{\delta_l^I(s)}{s(s - m_{\pi\pi}^2 - i\epsilon)} \right] \right\}$$

- **Monte-Carlo Generation of Phase Space Variables**
 - There are eight independent kinematic variables in the final state of the $ep \rightarrow e'p'\pi\pi$ reaction.

Total kinematic variables in final state (four 4-vectors)	16
Mass constraint of the four final state particles	-4
Four-Momentum Conservation, initial to final state	-4
Total number of independent variables in final state	8

- These are,
 - $Q^2, x_B, \phi_e, M_{1,2}^2, t, \phi_{1,2}^*, \cos\theta_{\sigma_Rest}, \phi_{\sigma_Rest}$



1. First consider the reaction $e + p \rightarrow e' + p' + \pi^+ + \pi^-$
 - Four Particles in final state
2. Secondly consider the reaction $e + p \rightarrow e' + p' + \pi^0 + \pi^0$,
its primary mode of decay is $\pi^0 \rightarrow \gamma \gamma$

6 particles in final state

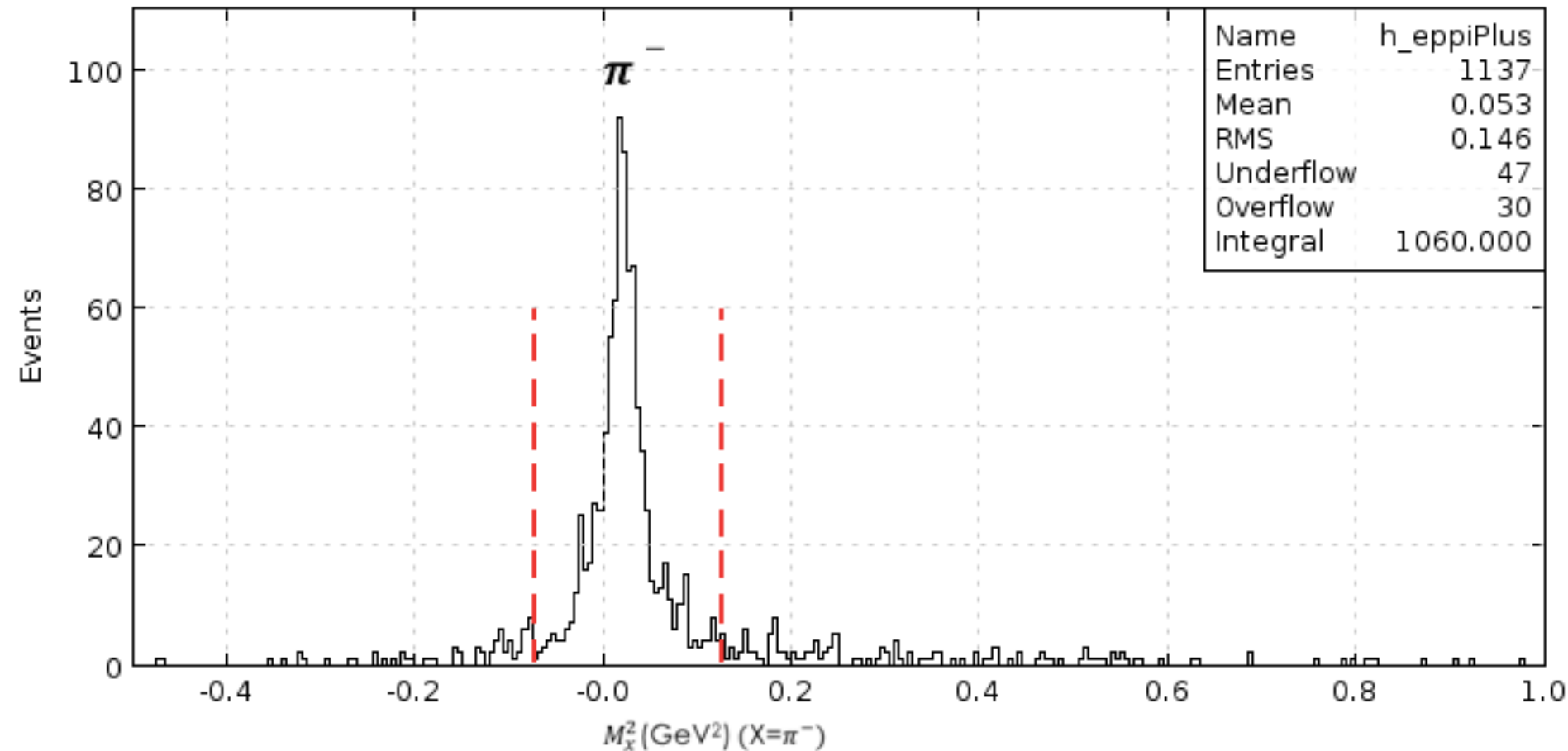
 - Scattered electron
 - Recoil Proton
 - Two π^0 s \Rightarrow Four gamma-rays

- For my simulation and reconstruction, I used
GEMC version 4a.2.1
COATJAVA version 4a.8.2

Steps :

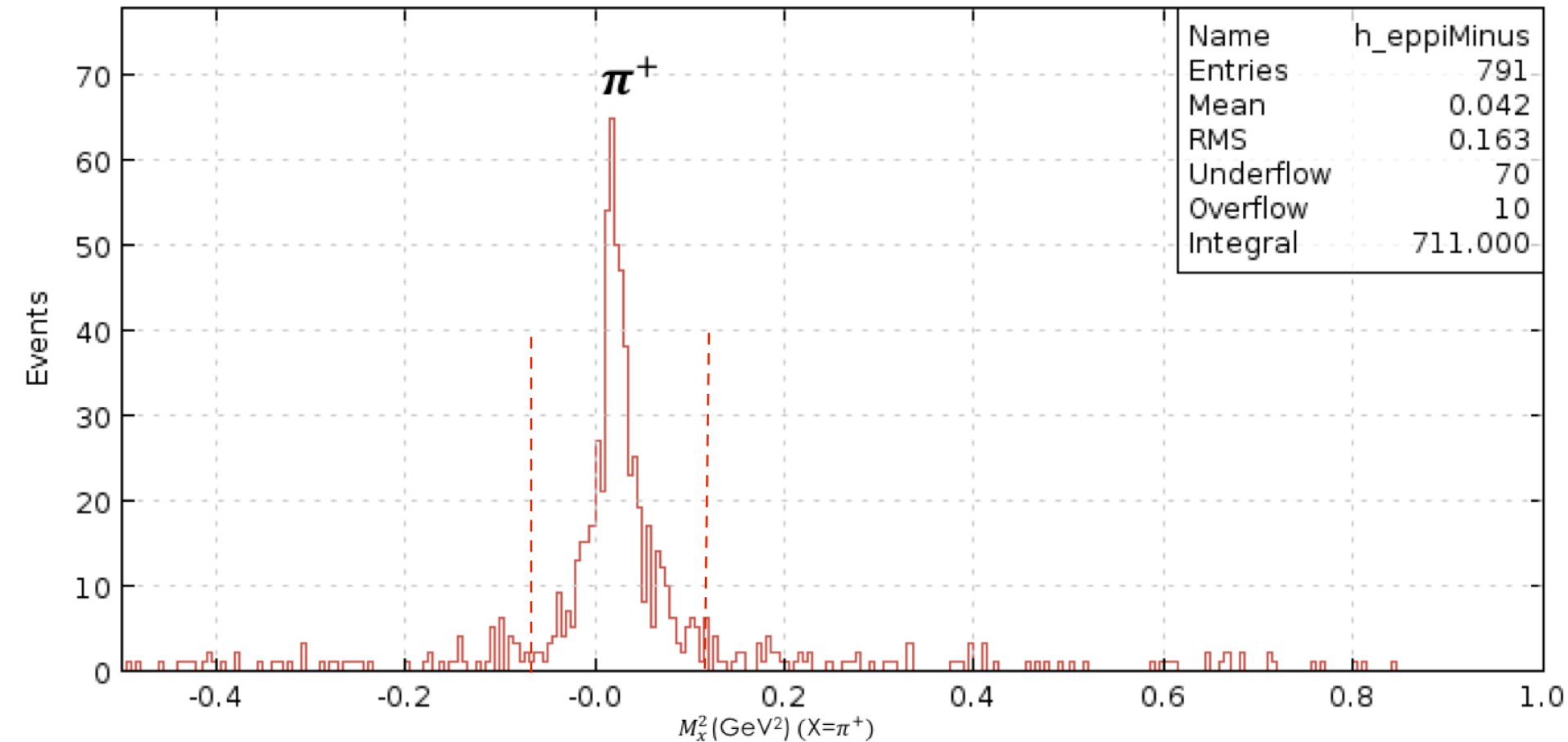
- After generation monte-carlo data is passed through the GEMC in the form of LUND format.
- Reconstruction is done with coatjava.
- CLAS12 analyses are done with **groovy** scripts (java).
- This method ties well with the coatjava framework and provides standard tools for reading EVIO files and reconstructed banks.

- Missing mass square of $H(e, e' p \pi^+) X$, $X = \pi^-$



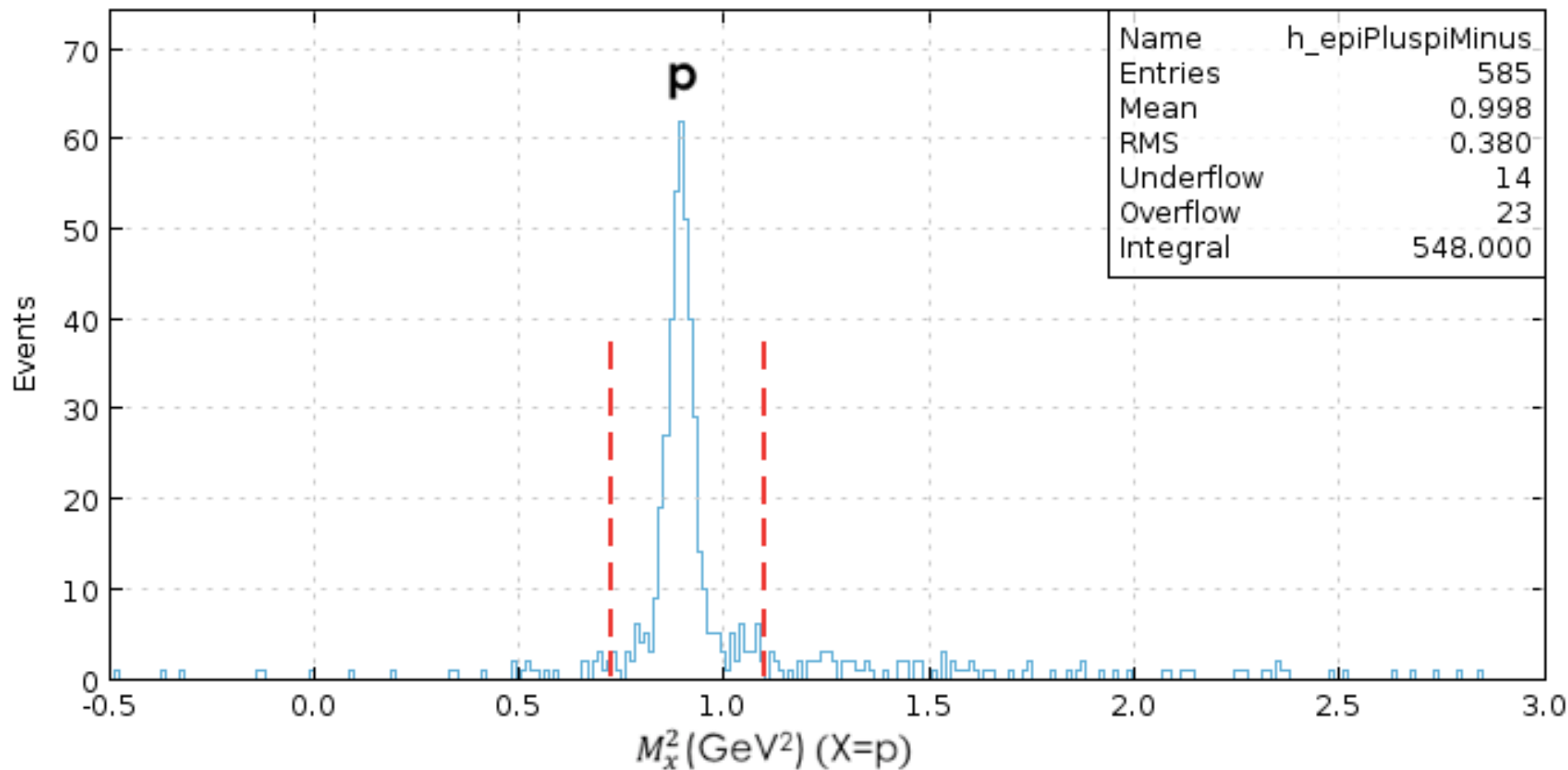
**Detection \otimes reconstruction
efficiency $\approx 14\%$**

- Missing mass square of π^+



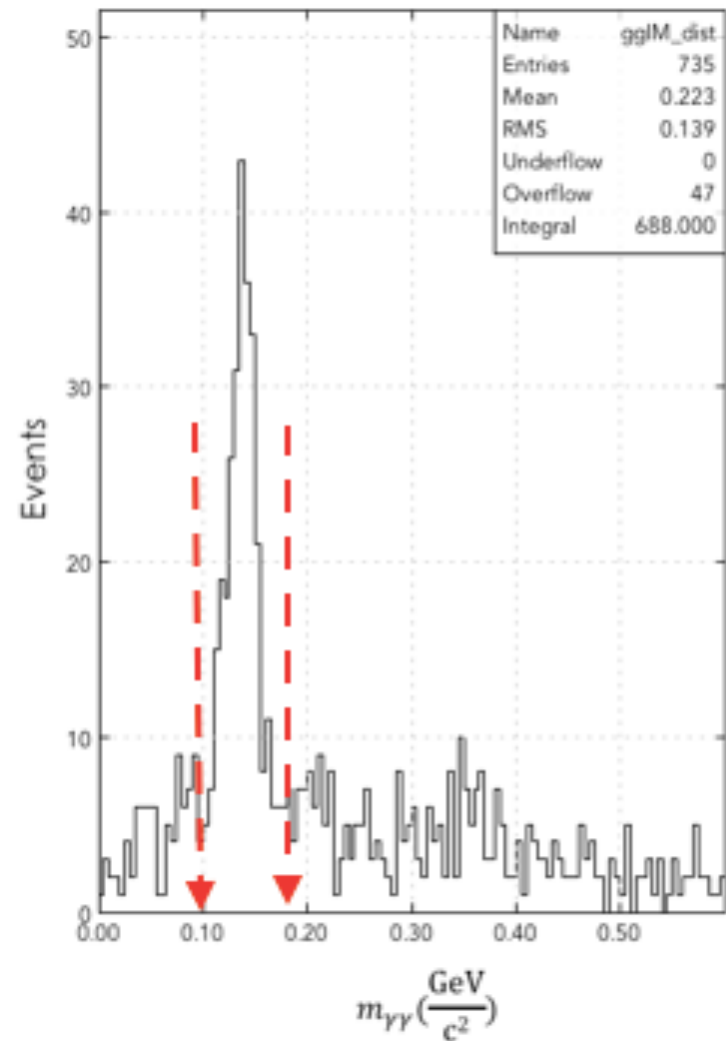
**Detection \otimes reconstruction
efficiency $\approx 11\%$**

- Missing mass square of *proton*

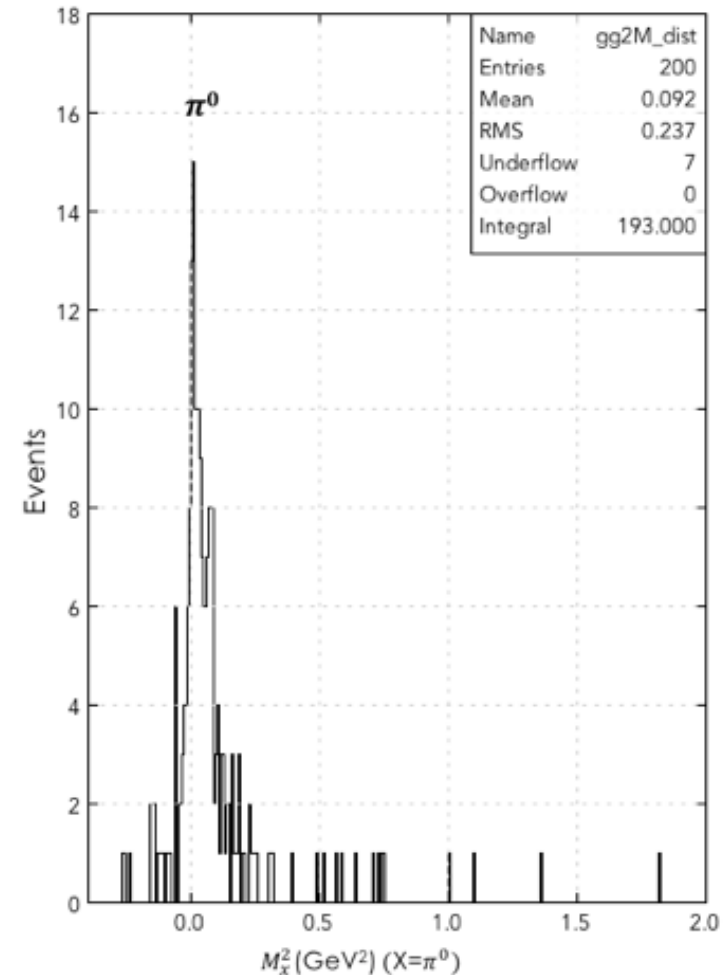


**Detection \otimes reconstruction
efficiency \approx 8%**

- Secondly, consider the reaction, $ep \rightarrow e' p' \pi^0 \pi^0$, and π^0 decays into two gammas ($\pi^0 \rightarrow \gamma\gamma$).
- Expected two photon invariant mass peak

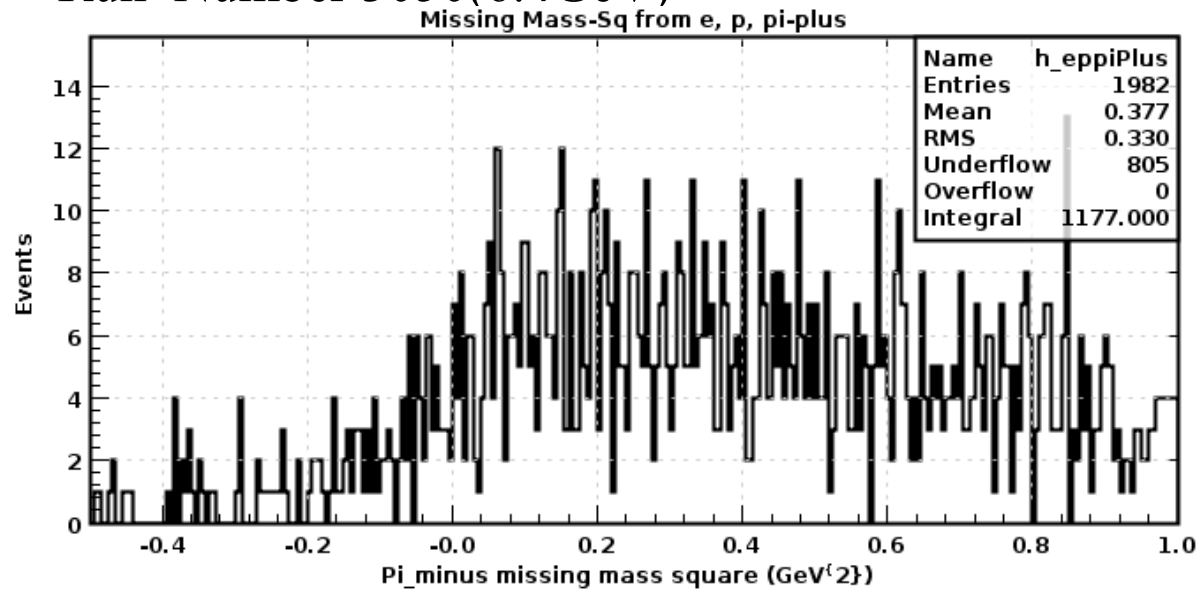


- Applied a cut on invariant mass :
around $0.10 < m_{\gamma\gamma}^2 < 0.17 \text{ GeV}^2$
- *Peak around 0.02 GeV^2*

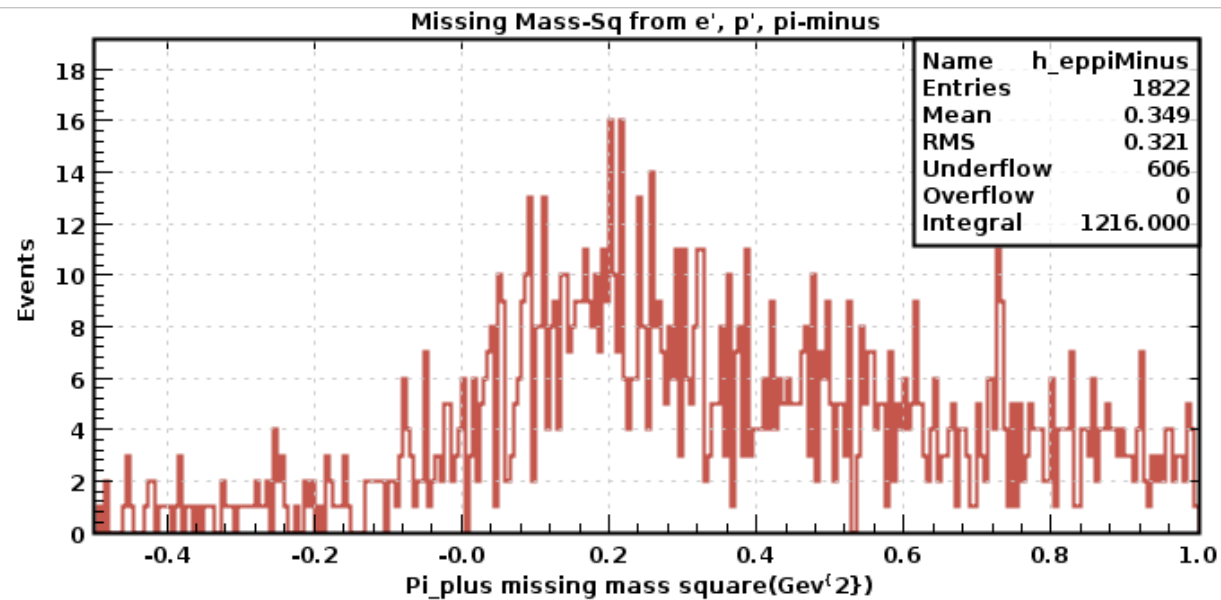


**Detection \otimes reconstruction
efficiency $\approx 2\%$**

- Run Number 3050(6.4GeV)

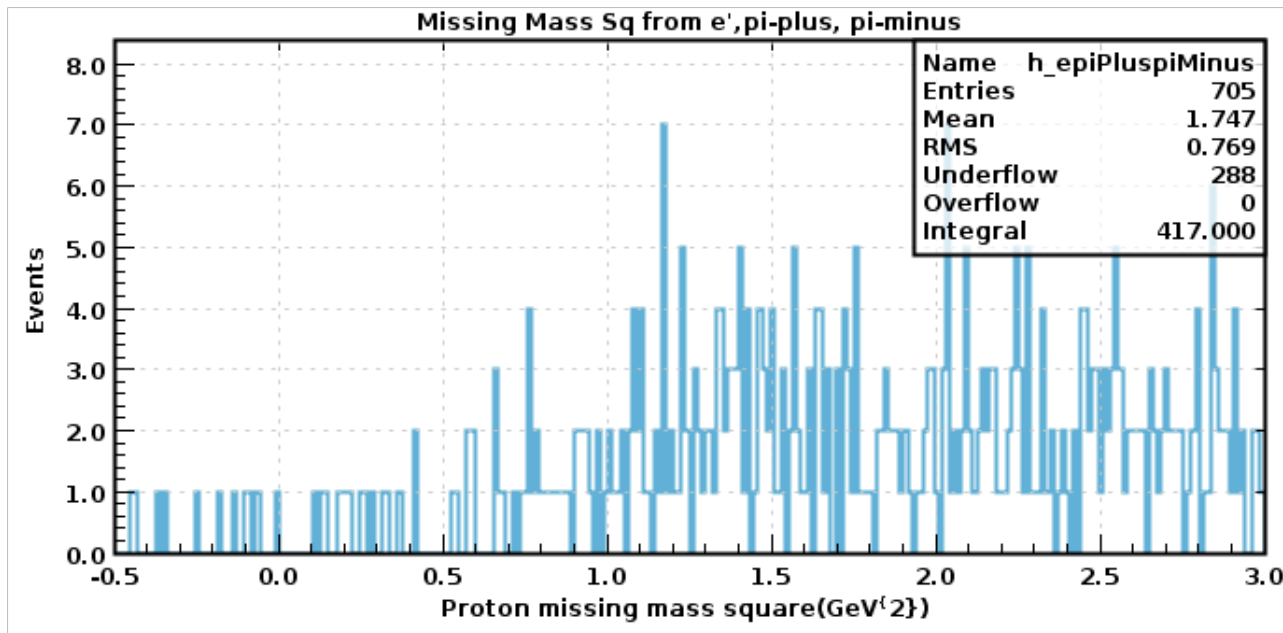


- Missing mass square of π^-



- Missing mass square of π^+

- Run_Number 3050(6.4GeV)



- Missing mass square of Proton

Thresholds look physical, but no sign of an exclusive peak in any of the three channels

- Particle ID?
- Momentum Calibrations?

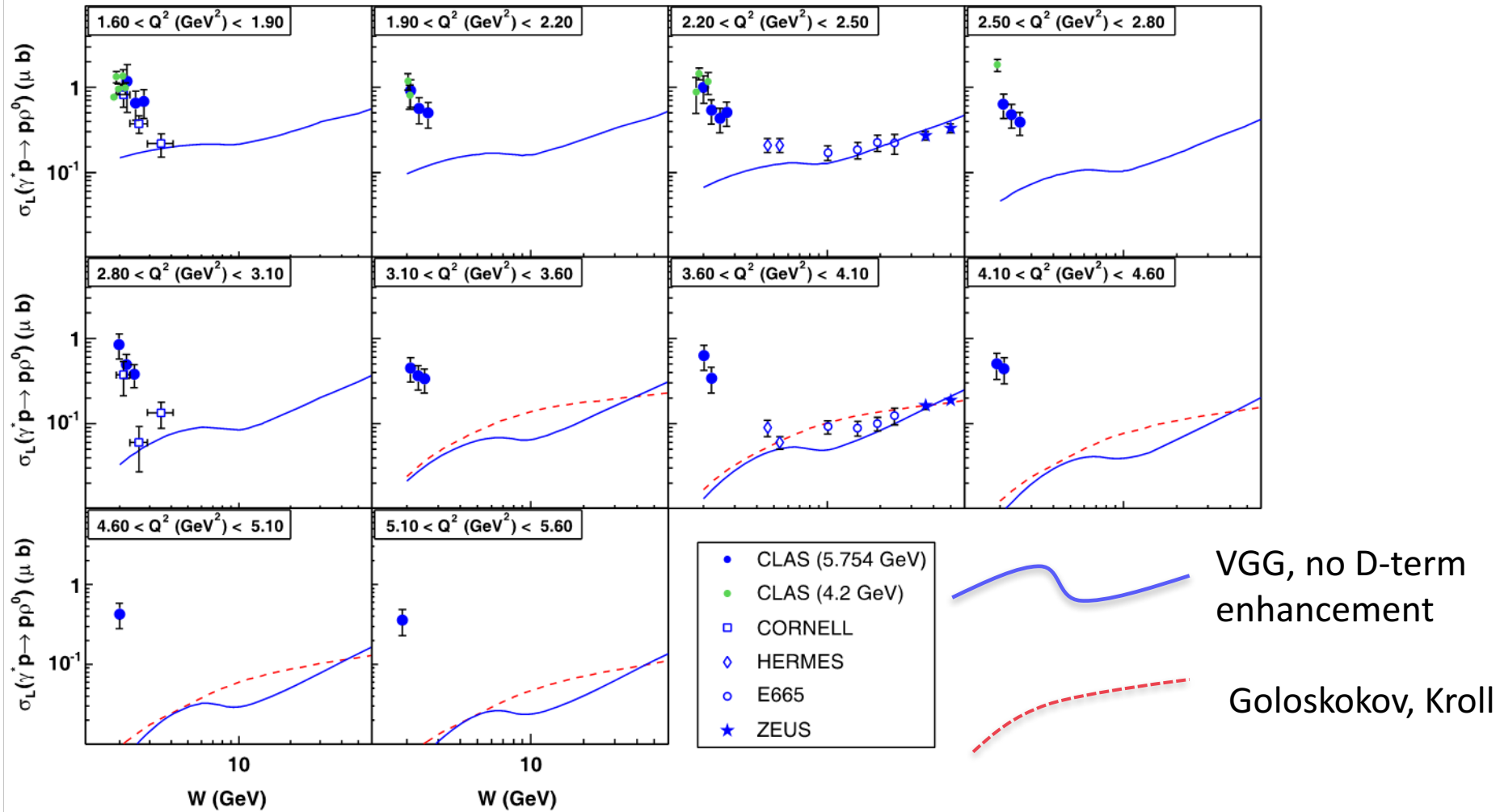
- **Preparing a run group proposal**
 - **Working on full cross section model**
- **Continue with 10.6 GeV data analysis**

Back Up Slides

- Deep ρ is a background to deep σ in $\pi^+\pi^-$ channel
 - Theory work on deep ρ
 - G-K Transversity
 - C.Weiss: Instanton dynamics study in progress.
- Detecting deep σ in the $\pi^0\pi^0$ channel is challenging in CLAS12.

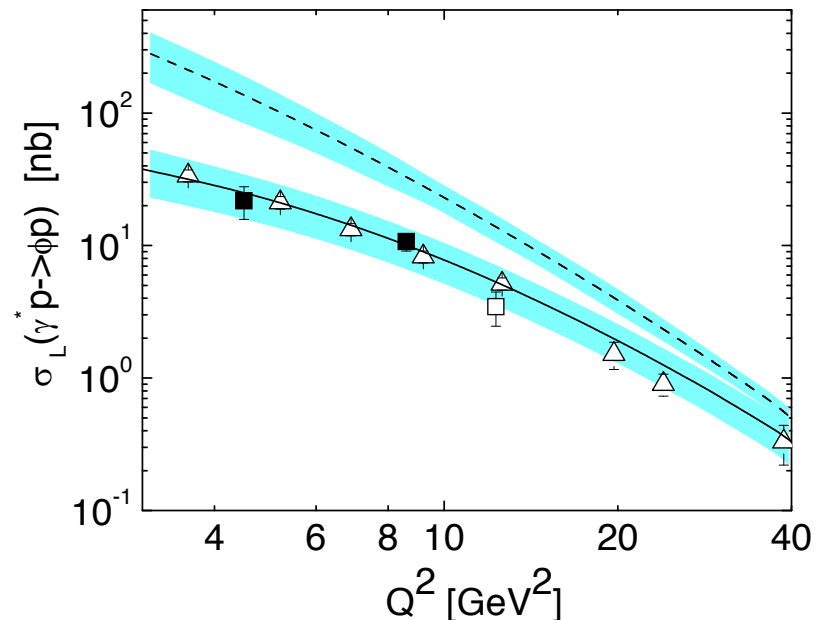
Deep ρ meson Problem

- S-channel helicity conservation violated
- Cross section is anomalously large at low W

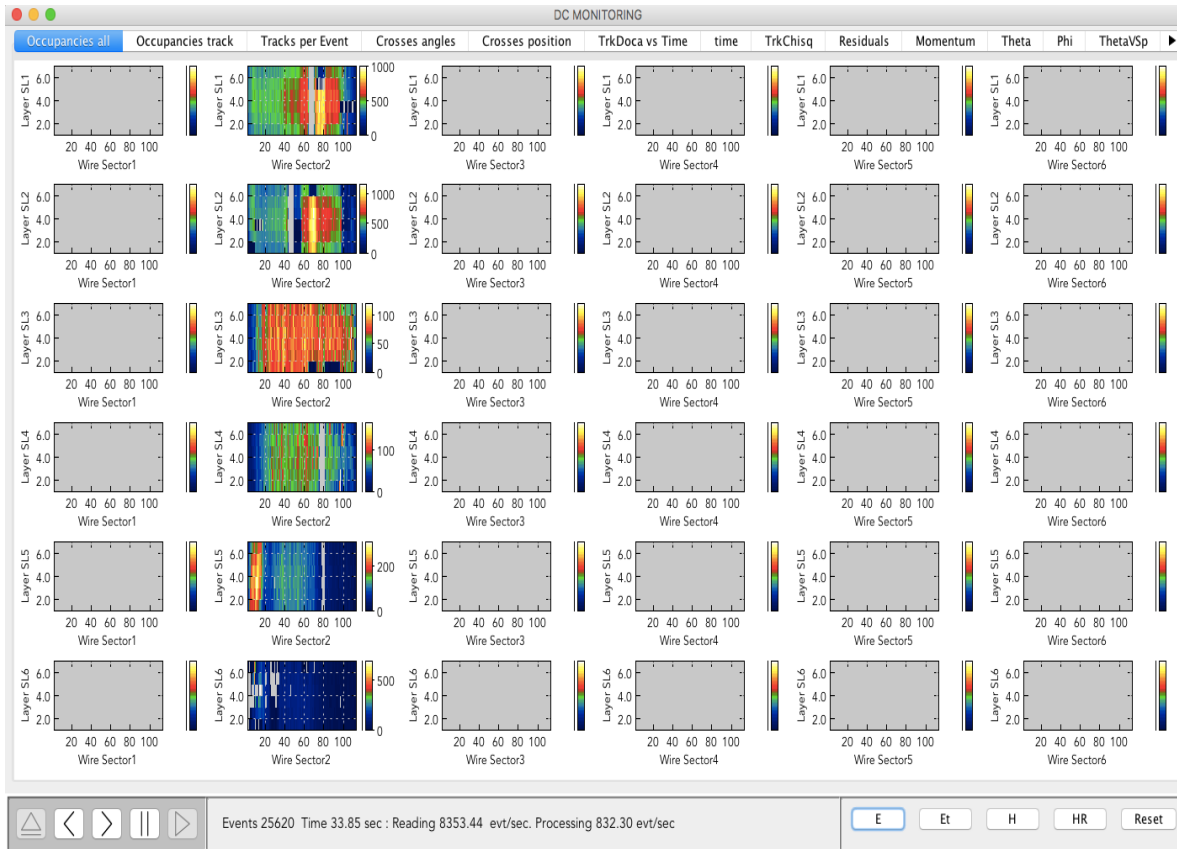


The Deep ϕ -meson

- Corrections up to factor of 10 to leading-order factorization at Jlab kinematics
- Successful phenomenology with finite-size/ χ SB in $\gamma \rightarrow$ meson amplitude and kinematic higher twist in proton GPD.
 - Deep π^0, η : χ SB Twist-3 $DA \otimes GPD_T$
 - $d\sigma_T \gg d\sigma_L$
 - (Recent Hall A and CLAS results)
- Deep ϕ : Sudakov form factor (finite-size) suppression:
 - CLAS/HERMES/HERA data \rightarrow



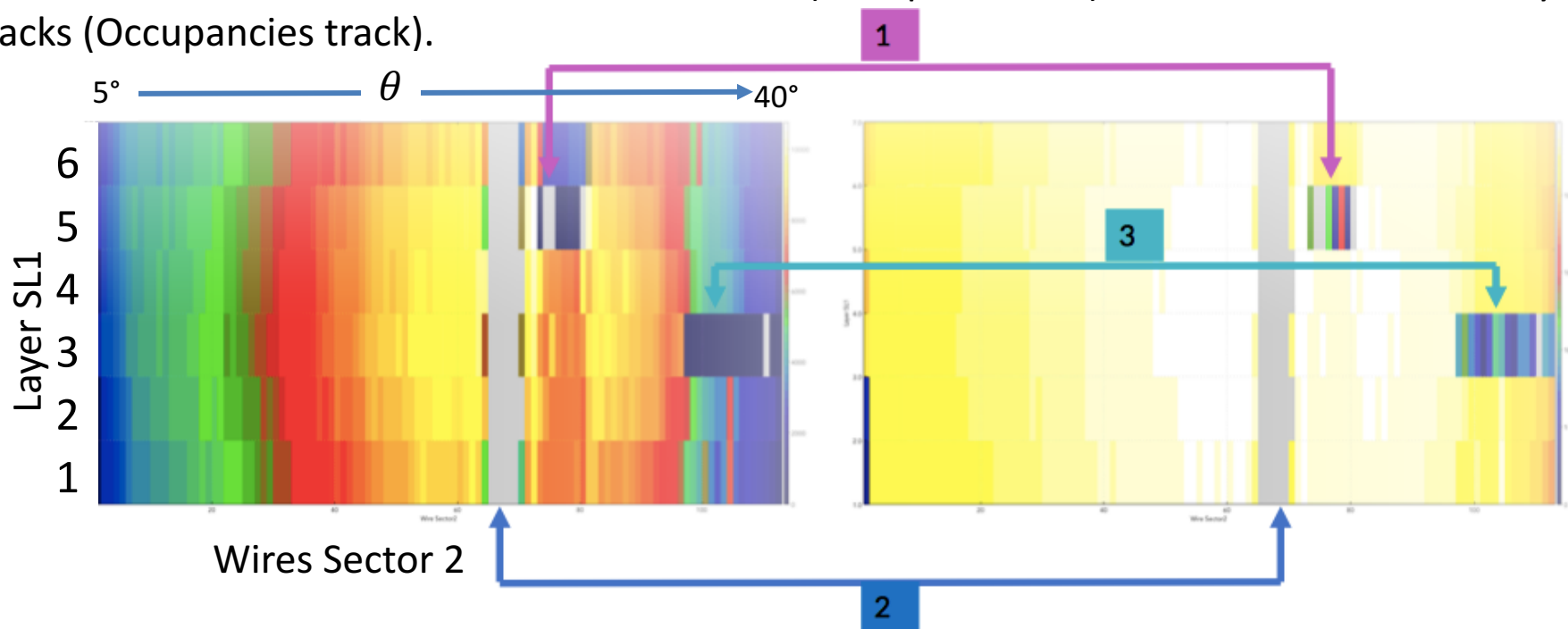
Service Project : DC Monitoring



- Occupancy Plots
- Track per Event
- Time
- Track DOCA vs Time
- Residual
- Δt Plot
- Hits per track
- Residual vs Track DOCA
- Crosses angles
- Crosses Position

Service Project : Occupancy Plots

Occupancy Plots: it displays plots of layer versus wire for all the hits (Occupancy Raw) in the DC::TDC bank, for all the hits in the TBHits bank (Occupancies all) and for the hits used only in tracks (Occupancies track).



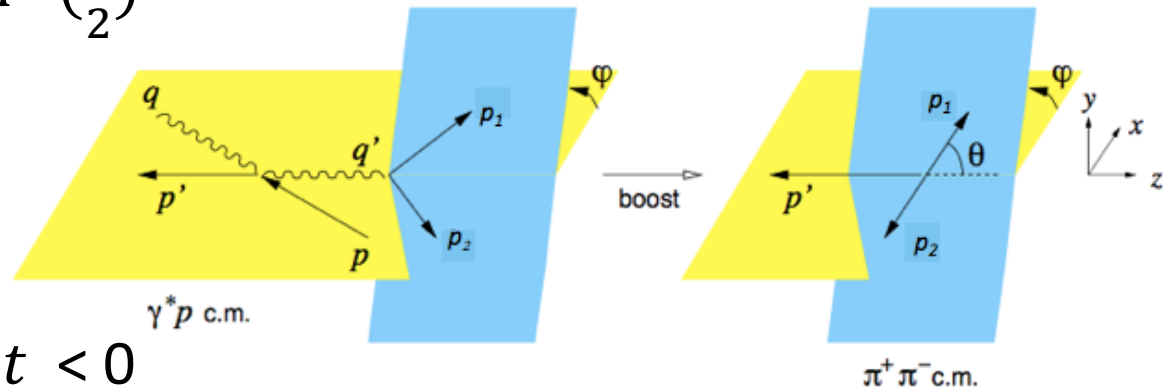
Anomaly	How to spot it
Individual dead wire	A single hole
Individual hot wire	A single hot spot
Unplugged Signal connector	16 wires: Vertical stripe, 6 layers high by 2 or 3 wires wide
Blown lv fuse	32 wires: Vertical stripe, 6 layers high by 5-6 wires wide
Unplugged hv pin, sense wire	Horizontal stripe, 1 layer high by 8 wires wide (or 16 if past wire 80)
2 Unplugged hv pin, sense wire	Horizontal stripe, 2 layer high by 8 wires wide (or 16 if past wire 80) almost no counts
Unplugged hv pin, field wire	Horizontal stripe, 2 layer high by 8 wires wide (or 16 if past wire 80) only depleted
Tripped hv supply channel or rarely a malfunction signal board	Rectangular hole: 6 layer high by 8 wires wide (wire 1-48); 16 wires wide (wire 49 - 80); and 32 wires wide (wires 81 - 112)

Basic Kinematics and Observables

- Here are the exclusive two-pion electroproduction kinematics on a proton using the following momentum variables:

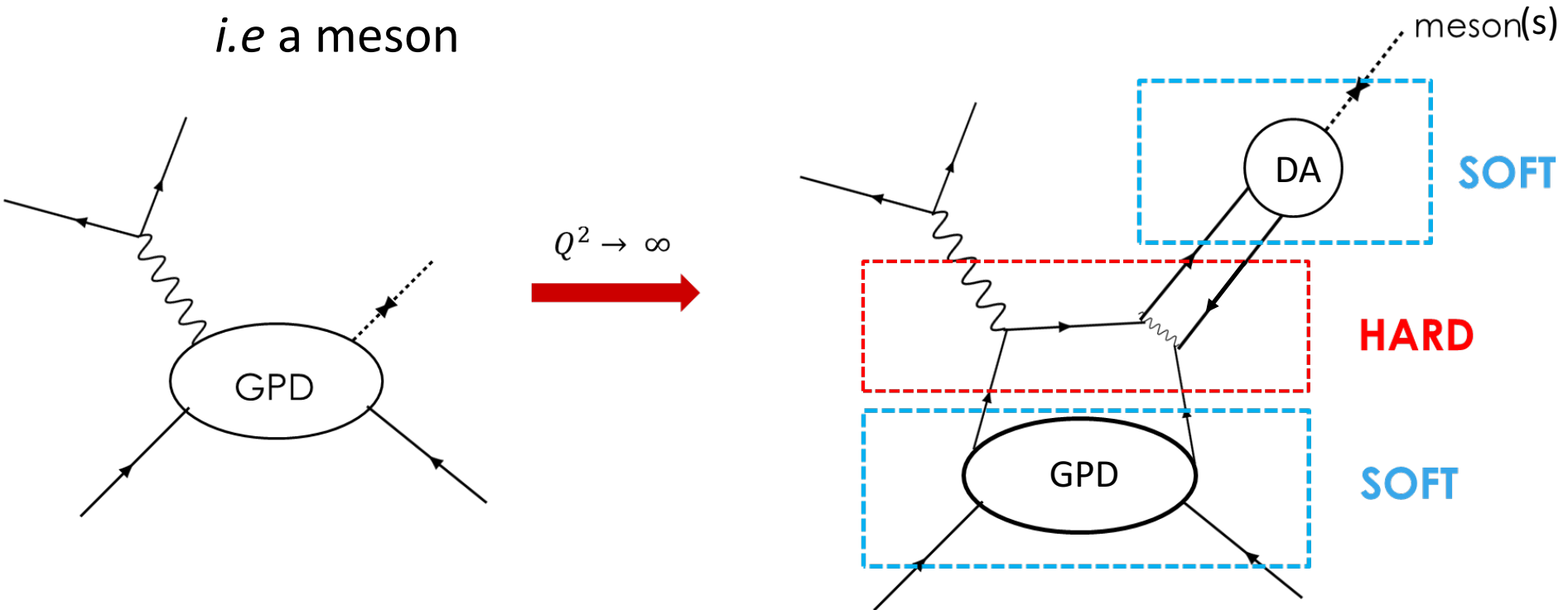
$$e(\mathbf{k}) + P(P) \rightarrow e(\mathbf{k}') + \pi_1(\mathbf{p}_1) + \pi_2(\mathbf{p}_2) + P(P').$$

- $q = k - k'$
- $q^2 = -Q^2 = 4EE' \sin^2\left(\frac{\theta}{2}\right)$
- $\nu = E - E'$
- $W^2 = (P + q)^2$
- $x_B = \frac{Q^2}{2P \cdot q}$
- $\Delta = P' - P$ and $\Delta^2 = t < 0$
- $(p_1 + p_2)^2 = m_{\pi\pi}^2$
- $q' = p_1 + p_2$ (e.g. σ or ρ meson)



Deep Virtual Exclusive Scattering (DVES)

- $ep \rightarrow e'p'h$ where h is the hadronic system
i.e a meson



- The interaction of the scattered electron with a parton (HARD), calculable through perturbative QCD, and the parton interaction with the proton (SOFT), described in terms of GPDs and another soft part describes the meson production.

Traditionally, **elastic form factors** and **parton distribution functions** (PDFs) were considered totally unrelated:

Elastic form factors give information on the charge and magnetization distributions in the transverse plane;

PDFs describe the distribution of partons in the longitudinal direction.

Generalized Parton Distributions encode information on the distribution of partons both in the transverse plane and in the longitudinal direction

→ natural extension and continuation of studies of the hadron structure.

Leading Order Amplitudes

$$\begin{aligned}
 & \langle B_2(p'), \pi^a \pi^b(q') | J^{e.m.} \cdot \varepsilon_L | B_1(p) \rangle \\
 &= -(e4\pi\alpha_s) \frac{N_c^2 - 1}{N_c^2} \frac{1}{4Q} \int_{-1}^1 d\tau \int_0^1 dz \\
 & \times \left[\sum_{f,f'} F_{ff'}(\tau, \xi) \Phi_{ff'}^{ab}(z, \zeta, m_{\pi\pi}) \right. \\
 & \times \left\{ \frac{e_{f'}}{z(\tau + \xi) - i0} + \frac{e_f}{(1-z)(\tau - \xi) + i0} \right\} \\
 & - \frac{2N_c}{N_c^2 - 1} \sum_f e_f F_{ff}(\tau, \xi) \Phi_G^{ab}(z, \zeta, m_{\pi\pi}) \\
 & \times \frac{1}{z(1-z)} \left\{ \frac{1}{(\tau + \xi) - i0} - \frac{1}{(\tau - \xi) + i0} \right\} \\
 & + \frac{4N_c}{N_c^2 - 1} \sum_f e_f \tau F_G(\tau, \xi) \Phi_{ff}^{ab}(z, \zeta, m_{\pi\pi}) \\
 & \left. \times \frac{1}{z(1-z)} \frac{1}{[\tau + \xi - i0][\tau - \xi + i0]} \right]. \quad (4)
 \end{aligned}$$

e_f is the charge of a quark of flavor $f=u,d,s$ in units of the proton charge.

ξ is the skewness parameter

ζ is the longitudinal momentum fraction carried by the pion.

The function $F_{ff'}(\tau, \xi)$ is a skewed quark distribution defined as:

$$\begin{aligned}
 F_{ff'}(\tau, \xi) = \int \frac{d\lambda}{2\pi} e^{i\lambda\tau} \langle B_2(p') | T \{ \bar{\psi}_{f'}(-\lambda n/2) \\
 \times \hat{n} \psi_f(\lambda n/2) \} | B_1(p) \rangle, \quad (6)
 \end{aligned}$$

The quark two-pion distribution amplitude $\Phi_{ff'}^{ab}(z, \zeta)$ is defined as:

$$\Phi_{ff'}^{ab}(z, \zeta) = \int \frac{d\lambda}{2\pi} e^{-i\lambda z(q' \cdot n^*)} \langle \pi^a \pi^b | T \{ \bar{\psi}_f(\lambda n^*)$$

and the two-pion gluon distribution amplitude $\Phi_G^{ab}(z, \zeta)$ as:

$$\begin{aligned}
 \Phi_G^{ab}(z, \zeta, m_{\pi\pi}^2) = \frac{1}{n^* \cdot q'} \int \frac{d\lambda}{2\pi} e^{-i\lambda z(q' \cdot n^*)} \\
 \times \langle \pi^a \pi^b | \{ n^* \cdot \mu n^{*\nu} G_{\alpha\mu}^A(\lambda n^*) \\
 \times G_\nu^{A\alpha}(0) \} | 0 \rangle, \quad (9)
 \end{aligned}$$

- $\pi^+ \pi^-$ production:

$$\begin{aligned}\Phi_{\pi^+ \pi^-}^{f' f} (z, \zeta, m_{\pi\pi}) &= \delta^{f' f} \Phi^{I=0} (z, \zeta, m_{\pi\pi}) \\ &+ \tau_3^{f' f} \Phi^{I=1} (z, \zeta, m_{\pi\pi}),\end{aligned}\quad (13)$$

$$\Phi_G^{\pi^+ \pi^-} (z, \zeta, m_{\pi\pi}) = \Phi^G (z, \zeta, m_{\pi\pi}),\quad (14)$$

- $\pi^0 \pi^0$ production:

$$\Phi_{\pi^0 \pi^0}^{f' f} (z, \zeta, m_{\pi\pi}) = \delta^{f' f} \Phi^{I=0} (z, \zeta, m_{\pi\pi}),\quad (15)$$

$$\Phi_G^{\pi^0 \pi^0} (z, \zeta, m_{\pi\pi}) = \Phi^G (z, \zeta, m_{\pi\pi}),\quad (16)$$

CLAS12 Detector

HTCC- Detection with charged pions with momentum above 5 GeV/c

LTCC- Charged pion identification for $p > 3 \text{ GeV}/c$

FTOF- for precise time-of-flight measurements for charged particle identification

FT-measures the small angle scattered electrons using a lead tungsten inner calorimeter

CTOF- particle identification is achieved with the central time of flight scintillator array, CTOF array allows for particle identification in momentum range upto 1.2 GeV/c, 0.65 GeV/c for π/p and π/k separation

EC-To detect showering particles

IC-(Inner Calorimeter) High energy resolution photon detector

Cherenkov Counter- e/π separation

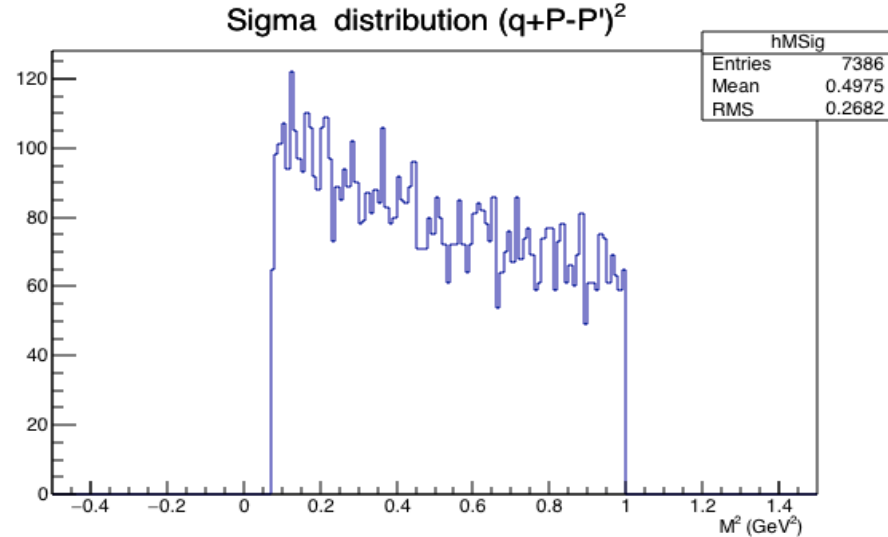
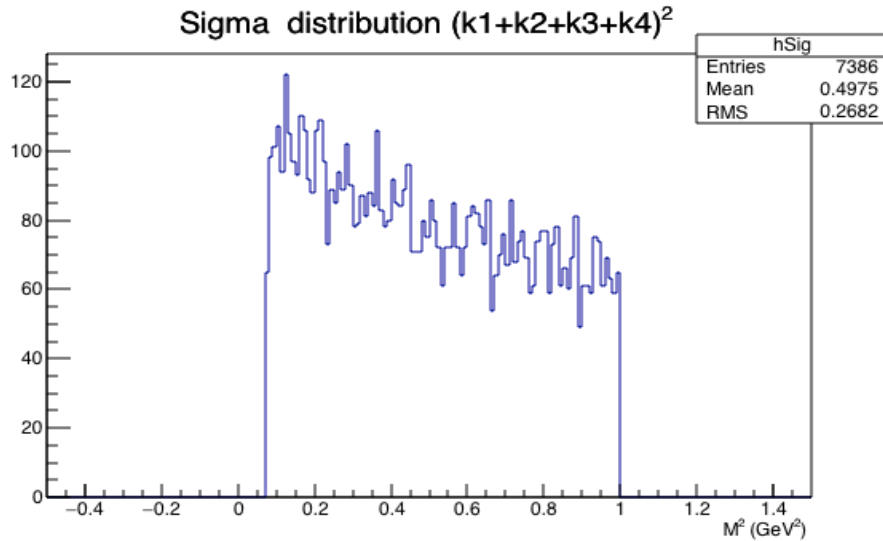
RICH- π/K and p/K separation

Micromegas-moderate magnetic field

SVT-very little thing, around target. take tracking information. Silicon tracker is great for locating a vertex

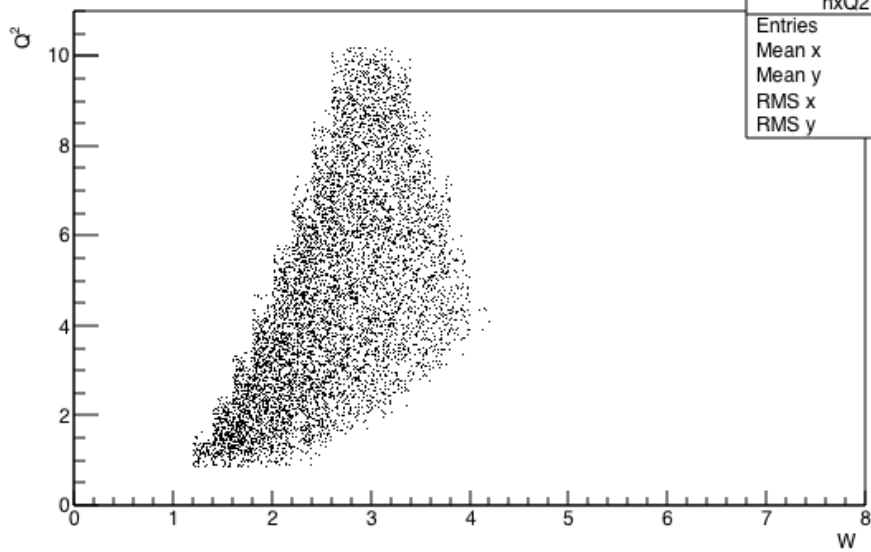
Event Generation

- From $ep \rightarrow ep'\pi^0 \pi^0$, π^0 decays into two gammas ($\pi^0 \rightarrow \gamma \gamma$)
- Now we have 4 gammas. We can check Sigma Distribution in two ways
 - $(k_1 + k_2 + k_3 + k_4)^2$
 - $(q + p - p')^2$



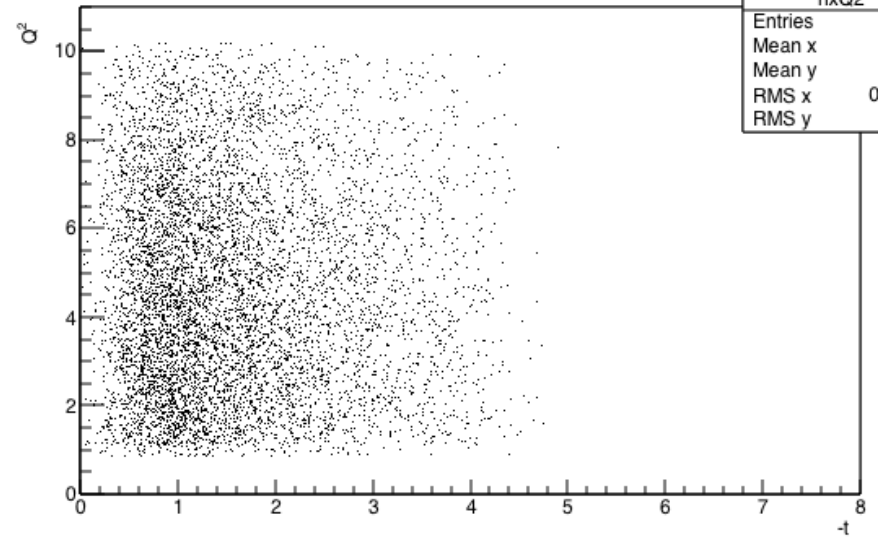
Event Generation

Q^2 vs. W



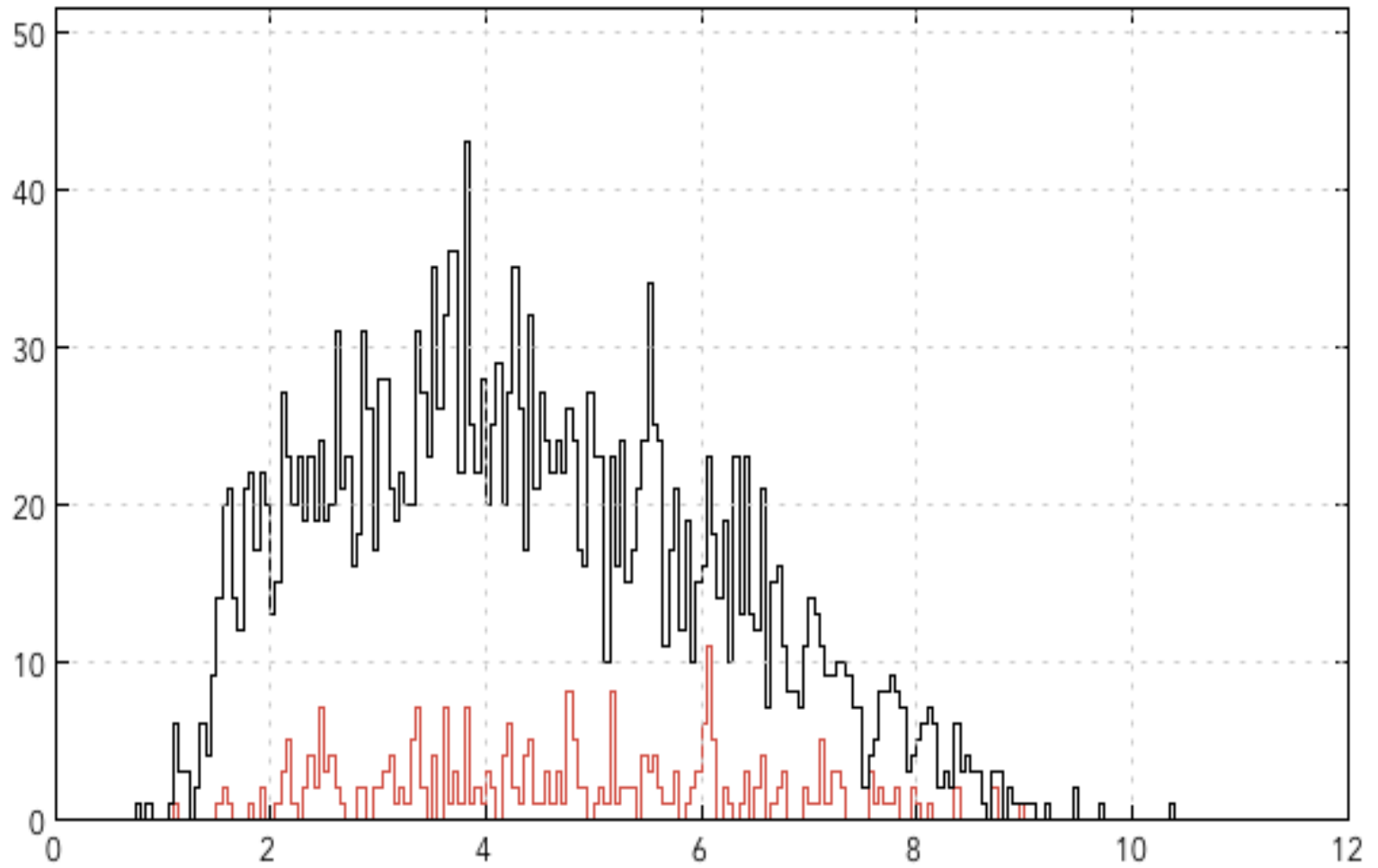
hxQ2	
Entries	7386
Mean x	2.579
Mean y	4.592
RMS x	0.605
RMS y	2.319

Q^2 vs. $-t$



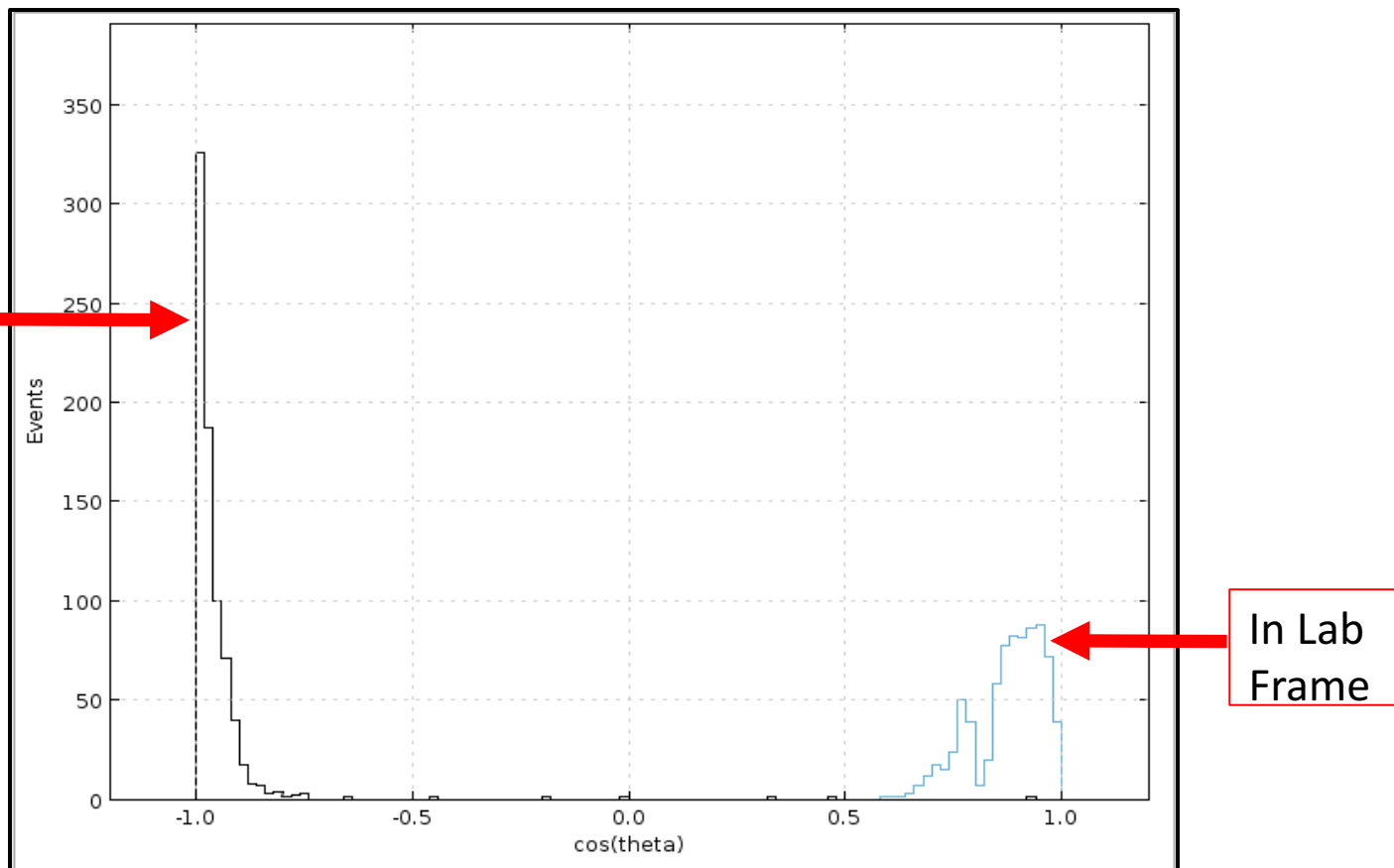
hxQ2	
Entries	7386
Mean x	1.571
Mean y	4.592
RMS x	0.9321
RMS y	2.319

Analysis



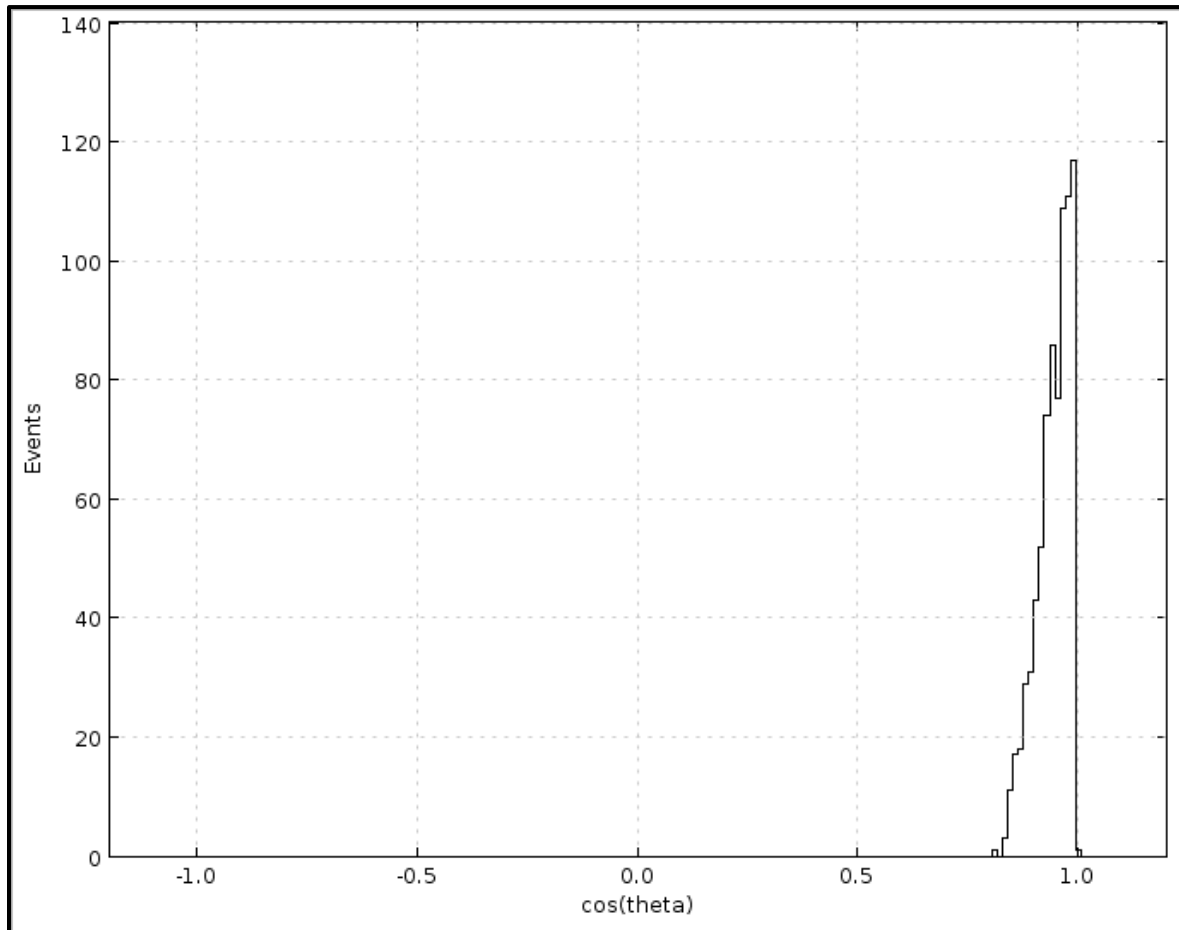
Analysis

- Here is the proton cosine distribution in both lab frame and CM frame
- Protons are forward direction in the Lab frame
- But in the CM frame they are 100% backward

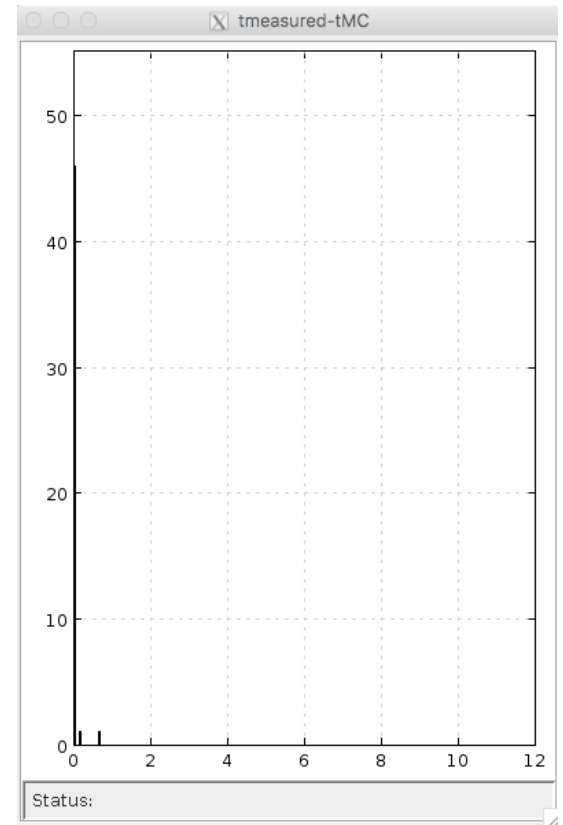
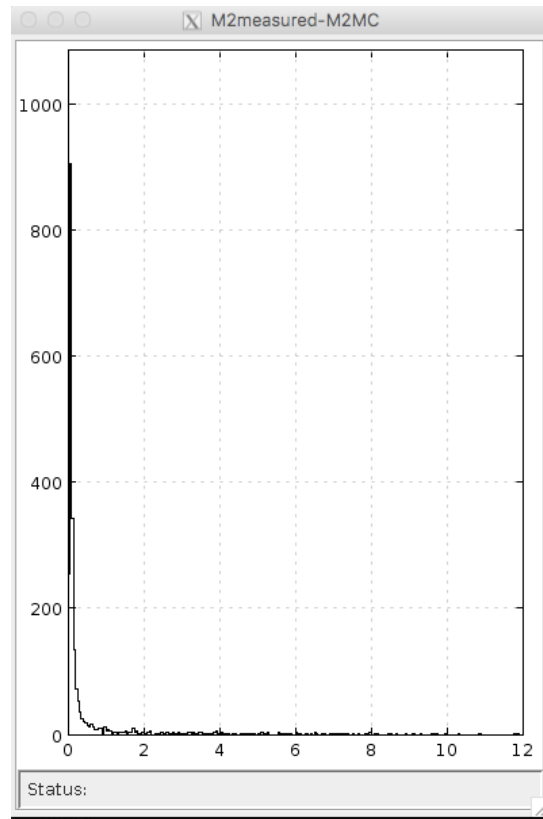
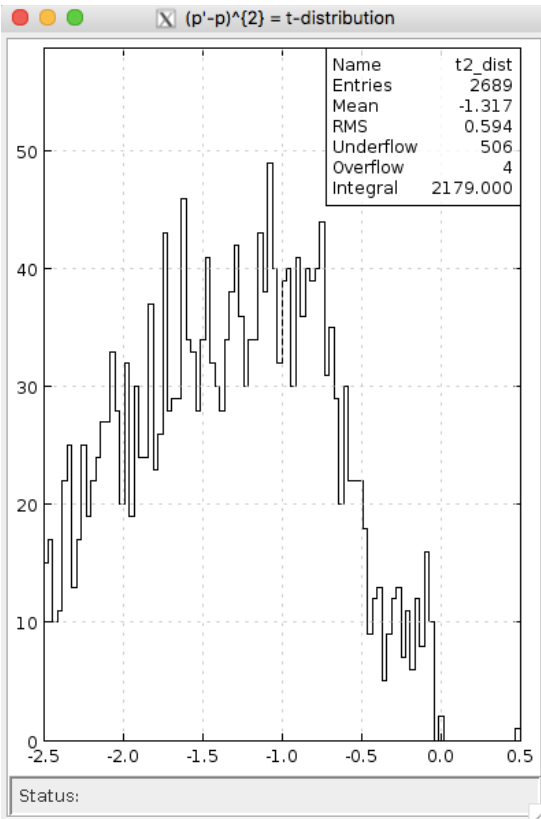


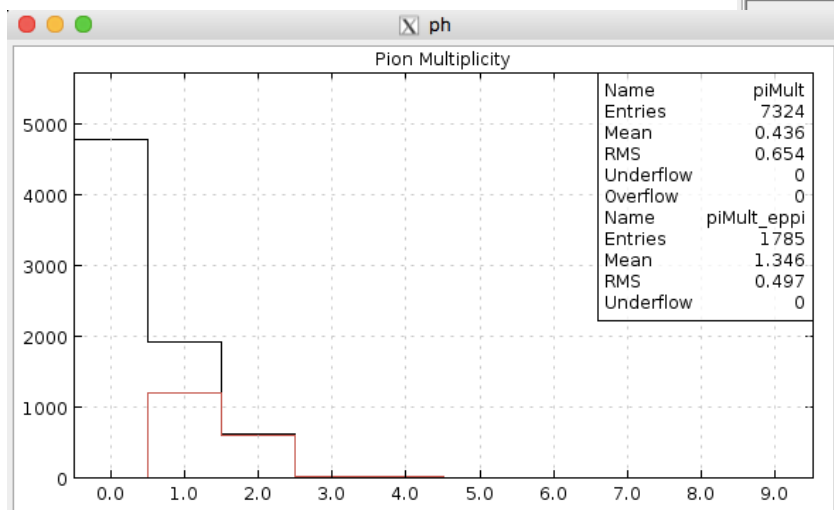
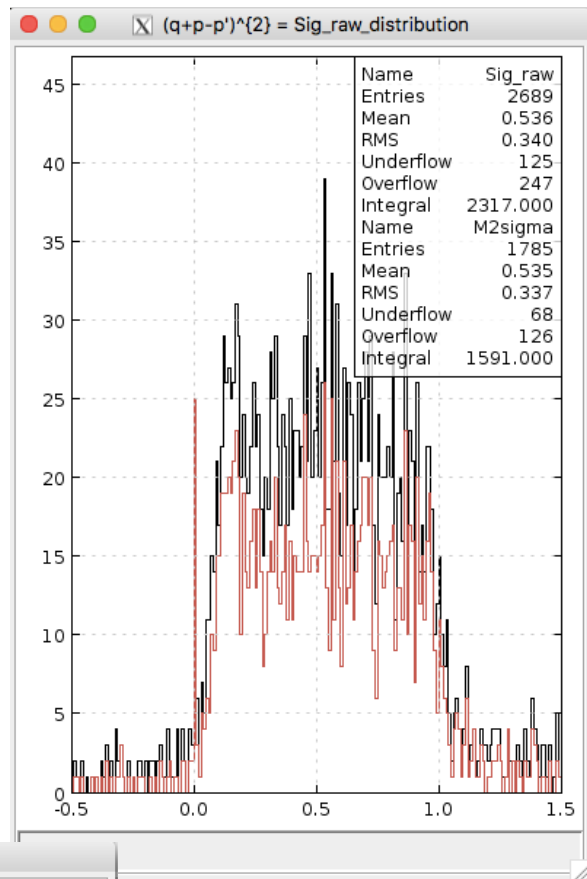
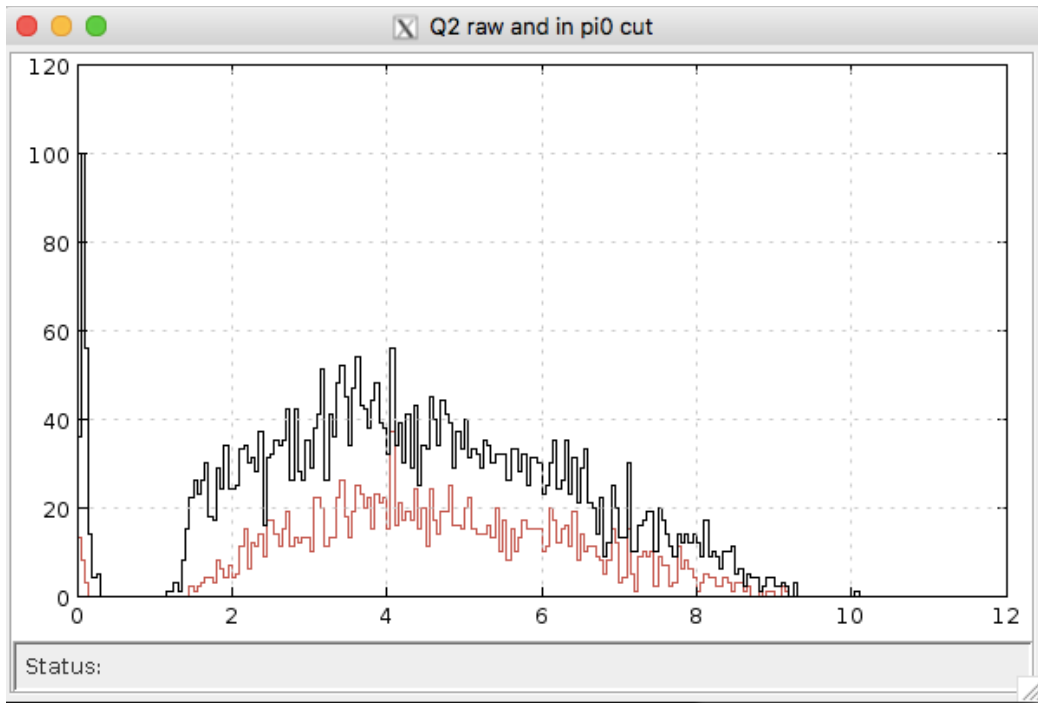
Analysis

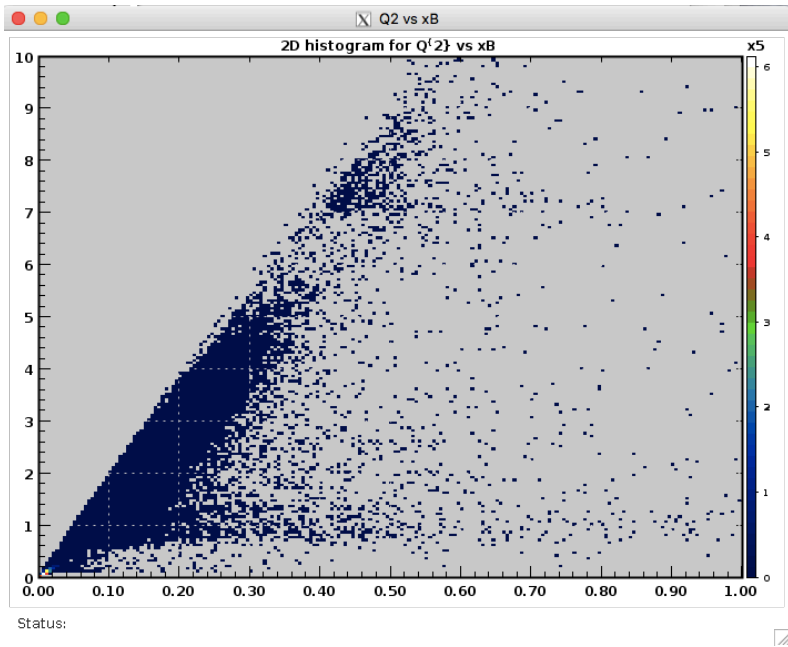
- Treat pi-minus as "missing" even if detected
- Here is the cosine distribution in rest frame
- This said that, pi-plus is always forward.



Analysis







After cut
 $Q^2 > 2$
log plot

