Collinear Distributions from Monte Carlo Global QCD Analyses

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WILLIAM & MARY





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Motivation

• Want to obtain reliable information of nonperturbative dynamics associated with hadron structure and hadronization

<u>Factorization</u> \rightarrow separation of short and long distance physics in pQCD expressions of experimental observables, e.g.

Unpolarized deep inelastic
scattering (DIS) observable
$$\ell + (p,d) \rightarrow \ell' + X$$

 $d\sigma(x,Q^2) \simeq \sum_f \int_x^1 \frac{d\xi}{\xi} f\left(\frac{x}{\xi},Q^2\right) \frac{d\hat{\sigma}_f(\xi,Q^2)}{Hard \ scattering \ coefficient}$

- Collinear factorization → distributions depend on some fraction of longitudinal momentum
- Nonperturbative distributions are typically determined empirically through global QCD analyses
 - → Objects are parameterized $xf(x) = Nx^a(1-x)^b(1+c\sqrt{x}+dx)$
 - \rightarrow Parameters are optimized with a least-squares fit

$$\chi^2 = \sum_{e}^{N_{exp}} \sum_{i}^{N_{data}} \frac{(D_i^e - T_i)^2}{(\sigma_i^e)^2}$$

Motivation

• However, many observables can depend on more than one type of distribution

Polarized semi-inclusive
DIS observable
$$\begin{array}{l}
A_{1}^{h} \sim \frac{g_{1}^{h}}{F_{1}^{h}} \sim \frac{\sum_{f,f'} \Delta f \otimes \Delta C_{ff'} \otimes D_{f'}^{h}}{\sum_{f,f'} f \otimes C_{ff'} \otimes D_{f'}^{h}}
\end{array}$$

→ Unpolarized (Polarized) PDFs: describe nucleon's momentum (spin) structure

→ Fragmentation functions (FFs) D_f : describe parton-to-hadron fragmentation

...and there are many issues with performing single chi-squared minimizations

→ Uncertainties computed by Hessian or Lagrange multiplier method introduce tolerance criteria (uncertainties inflated by arbitrary factor)

 \rightarrow Parameters difficult to constrain (flat eigendirections) are typically fixed

 \rightarrow Highly non-linear chi-squared function means many local minima that a single fit can be trapped in

• With a consistent theoretical framework and rigorous fitting procedure we can (more effectively): 1. Test universality

2. Explore the limits of collinear factorization

3. Study power suppressed corrections

JAM Collaboration Efforts

• <u>Recent efforts by the JAM collaboration:</u>

→ JAM15: Iterative Monte Carlo Analysis of spin PDFs (DIS only) – studies impact of high precision Jefferson Lab data on proton spin structure N. Sato *et. al.* Phys. Rev. D93 074005 (2016)

→ JAM16: First Monte Carlo analysis of FFs (SIA only) – preformed to obtain reliable determination of FFs and their uncertainties N. Sato *et. al.* Phys. Rev. D94 114004 (2016)

		JAM15	JAM16	JAM17	JAM18
Process	DIS SIA SIDIS DY			\bigvee	\bigvee
Function	$\begin{array}{c} f \\ \Delta f \\ D_f^h \end{array}$				

→ <u>JAM17</u>: First combined Monte Carlo analysis of polarized DIS, polarized SIDIS, and SIA data – studies impact of SIDIS on sea quark helicity distributions

JE, N. Sato, W. Melnitchouk PRL 119 132001 (2017)

 \rightarrow <u>JAM18</u>: Universal extraction of all nonperturbative input (in progress)

N. Sato, JE, C. Andrés, W. Melnitchouk, et. al. (2018)

• Other JAM projects:

→ Monte Carlo extraction of transversity distribution with lattice QCD constraints H.-W. Lin, W. Melnitchouk, A. Prokudin, N. Sato, H. Shows, PRL 120 152052 (2018)

 \rightarrow Monte Carlo analysis of pion PDFs (see P. Barry's talk Tuesday p.m.)

P. C. Barry, N. Sato, W. Melnitchouk, C.-R. Ji, arXiv:1804.01965 (2018)

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• Based on Bayesian statistical methods – robust determination of "observables" *O* (PDFs,FFs,etc.) and their uncertainties

$$E\left[\mathcal{O}\right] = \int d^{n} a \mathcal{P}(\vec{a}|data) \mathcal{O}(\vec{a})$$
$$V\left[\mathcal{O}\right] = \int d^{n} a \mathcal{P}(\vec{a}|data) \left[\mathcal{O}(\vec{a}) - E[\mathcal{O}]\right]^{2}$$

• Bayes' theorem defines probability ${\cal P}$ as

$$\mathcal{P}(\vec{a}|data) = \frac{1}{Z}\mathcal{L}(data|\vec{a})\pi(\vec{a})$$

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$$\underbrace{\mathbf{Likelihood function}}_{\mathbf{Likelihood function}}$$

$$\mathcal{L} = \exp\left(-\frac{1}{2}\chi^{2}(\vec{a})\right) \Rightarrow \text{Gaussian form in data with } \chi^{2} = \sum_{e}^{N_{exp}} \sum_{i}^{N_{data}} \frac{(D_{i}^{e} - T_{i})^{2}}{(\sigma_{i}^{e})^{2}}$$

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$$\stackrel{\text{Priors}}{\uparrow}$$

$$\stackrel{\text{Evidence"}}{=} Z = \int d^{n}a\mathcal{L}(data|\vec{a})\pi(\vec{a})$$

• Based on Bayesian statistical methods – robust determination of "observables" *O* (PDFs,FFs,etc.) and their uncertainties

$$E\left[\mathcal{O}\right] = \int d^{n}a\mathcal{P}(\vec{a}|data)\mathcal{O}(\vec{a})$$
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• Monte Carlo technique is used to evaluate expectation value and variance integrals

 \rightarrow samples parameter space and assigns weights w_k to each parameter a_k such that

$$E[\mathcal{O}(\vec{a})] = \sum_{k} w_k \mathcal{O}(\vec{a}_k) \qquad V[\mathcal{O}(\vec{a})] = \sum_{k} w_k (\mathcal{O}(\vec{a}_k) - E[\mathcal{O}])^2$$

Iterative Monte Carlo (IMC)

(JAM17 Analysis)



- → Samples wide region of parameter space
- → Data is partitioned for cross-validation – training set is fitted via chi-square minimization
- → Posteriors used to construct sampler (multi-dimensional Guassian, kernel density estimation, etc) – where parameters are chosen for the next iteration
- \rightarrow Procedure iterated until converged

$$E[\mathcal{O}] = \frac{1}{n} \sum_{k}^{n} \mathcal{O}(\vec{a}_{k})$$
$$V[\mathcal{O}] = \frac{1}{n} \sum_{k}^{n} (\mathcal{O}(\vec{a}_{k}) - E[\mathcal{O}])^{2}$$

Nested Sampling

• Statistical mapping of multidimensional integral to 1-D

$$Z = \int d^n a \mathcal{L}(data | \vec{a}) \pi(\vec{a}) = \int_0^1 dX \mathcal{L}(X)$$

where the *prior volume* $dX = \pi(\vec{a})d^n a$



• Algorithm:

→ Initialize $X_0 = 1, L = 0$ and choose N active points $X_1, X_2, ..., X_N$ from prior

→ For each iteration, sample new point and remove lowest L_i , replacing with point such that L is monotonically increasing

 \rightarrow Repeat until entire parameter space has been explored

Proton Spin Structure from DIS

• Typically measure longitudinal and transverse spin asymmetries

$$A_{\parallel} = \frac{\sigma^{\uparrow \Downarrow} - \sigma^{\uparrow \Uparrow}}{\sigma^{\uparrow \Downarrow} + \sigma^{\uparrow \Uparrow}} = D\left(A_1 + \eta A_2\right) \qquad A_{\perp} = \frac{\sigma^{\uparrow \Rightarrow} - \sigma^{\uparrow \Leftarrow}}{\sigma^{\uparrow \Rightarrow} + \sigma^{\uparrow \Leftarrow}} = d\left(A_2 + \zeta A_1\right)$$

→ Virtual photoproduction asymmetries: $A_1 = \frac{(g_1 - \gamma^2 g_2)}{F_1}$ $A_2 = \gamma \frac{(g_1 + g_2)}{F_1}$ $\gamma^2 = \frac{4M^2 x^2}{Q^2}$

• Leading contribution to polarized structure function g₁:

$$g_1(x,Q^2) = \frac{1}{2} \sum_q e_q^2 \left[(\Delta C_q \otimes \Delta q^+)(x,Q^2) + (\Delta C_g \otimes \Delta g)(x,Q^2) \right] + \mathcal{O}\left(\frac{1}{Q}\right)$$

• First moment of polarized structure function g₁:

$$\int_0^1 dx g_1^p(x, Q^2) = \frac{1}{36} \left[8\Delta\Sigma + 3g_A + a_8 \right] \left(1 - \frac{\alpha_s}{\pi} + \mathcal{O}(\alpha_s^2) \right) + \mathcal{O}\left(\frac{1}{Q}\right)$$

 \rightarrow DIS requires assumptions about triplet and octet axial charges to extract $\Delta\Sigma$

• Assuming exact $SU(2)_f$ and $SU(3)_f$ values from weak baryon decays

$$\int dx \left(\Delta u^+ - \Delta d^+ \right) = g_A \sim 1.269 \qquad \int dx \left(\Delta u^+ + \Delta d^+ - 2\Delta s^+ \right) = a_8 \sim 0.586$$
$$\Delta \Sigma_{[10^{-3}, 0.8]} \sim 0.3$$

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Proton Spin Structure from SIDIS

• Measured via longitudinal double spin asymmetries

$$A_1^h(x, z, Q^2) = \frac{g_1^h(x, z, Q^2)}{F_1^h(x, z, Q^2)}$$



• Polarized structure function at NLO defined in terms of 2-D convolution

$$g_1^h(x, z, Q^2) = \frac{1}{2} \sum_q e_q^2 \left\{ \Delta q(x, Q^2) D_q^h(z, Q^2) + \frac{\alpha_s(Q^2)}{2\pi} \times \left(\Delta q \otimes \Delta C_{qq} \otimes D_q^h + \Delta q \otimes \Delta C_{gq} \otimes D_g^h + \Delta g \otimes \Delta C_{qg} \otimes D_q^h \right) \right\}$$

• To include SIDIS observables in the JAM global analyses, fragmentation functions (FFs) must be known

→ Choice of FF parameterizations available (HKNS & DSS) differed significantly in kaon sector – strongly impacts Δs^+ extraction

• JAM17 Analysis: first to fit simultaneously polarized PDFs + FFs and release SU(3) constraints

JAM17 Polarized PDF Distributions



- Isoscalar sea distribution consistent with zero
- Isovector sea slightly prefers positive shape at low *x*
 - \rightarrow Non-zero asymmetry given by small contributions from SIDIS asymmetries

JE, N. Sato, W. Melnitchouk PRL 119 132001 (2017)

- Δu^+ consistent with previous analysis
- Δd^+ slightly larger in magnitude
 - → Anti-correlation with Δs^+ , which is less negative than JAM15 at $x \sim 0.2$



JAM17 – Resolution of the Strange Polarization

JAM17 + SU(3)

DSSV09

JAM15



Why does DIS+SU(3) give large negative Δs^+ ?

- Low *x* DIS deuterium data from COMPASS prefers small negative Δs^+
- Negative polarization shifted to intermediate region to satisfy SU(3) constraint
- b parameter for Δs^+ typically fixed to values ~6-10, producing a peak at *x*~0.1

JE, N. Sato, W. Melnitchouk PRL 119 132001 (2017)

- Δs^+ distribution consistent with zero, slightly positive in intermediate *x* range
- Primarily influenced by HERMES K⁻ data from deuterium target





N. Sato, JE, C. Andrés, W. Melnitchouk, et. al. (2018)

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JAM18 Data vs Theory (Preliminary)

SIDIS

N. Sato, JE, C. Andrés, W. Melnitchouk, et. al. (2018)



• Difficulty fitting low- Q^2 data \rightarrow only $Q^2 > 5$ GeV² included

JAM18 Unpolarized PDFs (Preliminary)

PDFs



- Central value and uncertainties from maximum likelihood + data resampling method
- Distributions mostly consistent with previous analyses
 - \rightarrow Light sea asymmetry differs at large-*x*
- SIDIS supports suppression of strange distribution





JAM18 Impact of SIDIS



- Decrease in central value and uncertainty of strange PDF with SIDIS
- Large effect on the gluon distribution
 - → Correlation with strange PDF (momentum sum rule)



Summary and Outlook

• Monte Carlo statistical methods are important for robust extractions of nonperturbative functions and their uncertainties

→ Necessary for future global QCD studies that contain large data sets and have many fit parameters (TMDs, GPDs)

- New approaches being developed:
 - → Likelihood sampling methods (Nested Sampling)
 - → Generalization of Gaussian likelihood (systematic treatment of incompatible data sets)
- First universal analysis of unpolarized + polarized measurements underway
 - \rightarrow Simultaneous extraction of all nonperturbative input
 - \rightarrow Strict test of universality
 - \rightarrow Can separate individual aligned/anti-aligned helicity distributions
- Longer term: extracting transverse momentum dependent (TMD) PDFs and FFs