

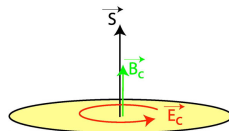
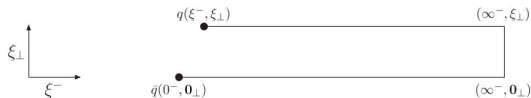
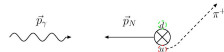
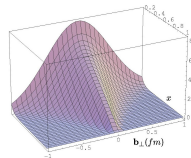
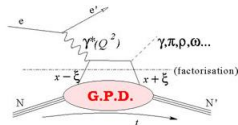
# Transverse Force Femtography

Matthias Burkardt

\*New Mexico State University

May 17, 2018

- GPDs  $\rightarrow$  3D imaging of the nucleon
  - twist-3 PDFs  $g_2(x) \rightarrow \perp$  force
  - $\hookrightarrow$  twist-3 GPDs  $\rightarrow \perp$  force tomography
- Motivation: why twist-3 GPDs
- twist-3 GPD  $G_2^q \rightarrow L^q$
  - twist 3 PDF  $g_2(x) \rightarrow \perp$  force
  - twist 2 GPDs  $\rightarrow \perp$  imaging (of quark densities)
  - $\hookrightarrow$  twist 3 GPDs  $\rightarrow \perp$  imaging of  $\perp$  forces
- Summary
  - Outlook



MB, PRD62, 071503 (2000)

- form factors:  $\overleftrightarrow{FT} \rho(\vec{r})$
  - $GPDs(x, \vec{\Delta})$ : form factor for quarks with momentum fraction  $x$
- ↪ suitable FT of  $GPDs$  should provide spatial distribution of quarks with momentum fraction  $x$

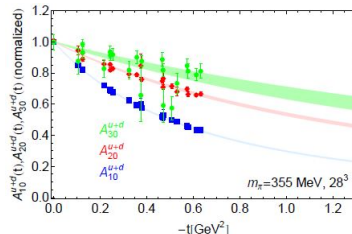
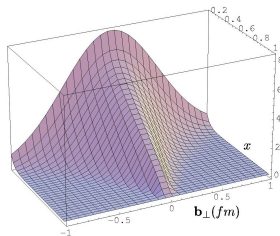
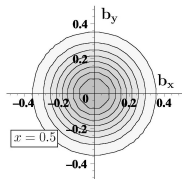
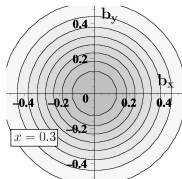
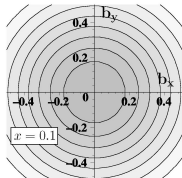
### Impact Parameter Dependent Quark Distributions

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} GPD(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

$q(x, \mathbf{b}_\perp)$  = parton distribution as a function of the separation  $\mathbf{b}_\perp$  from the transverse center of momentum  $\mathbf{R}_\perp \equiv \sum_{i \in q, g} \mathbf{r}_{\perp, i} x_i$

- probabilistic interpretation!
  - no relativistic corrections: Galilean subgroup! (MB,2000)
- ↪ corollary: interpretation of 2d-FT of  $F_1(Q^2)$  as charge density in transverse plane also free from relativistic corrections (MB,2003;G.A.Miller, 2007)

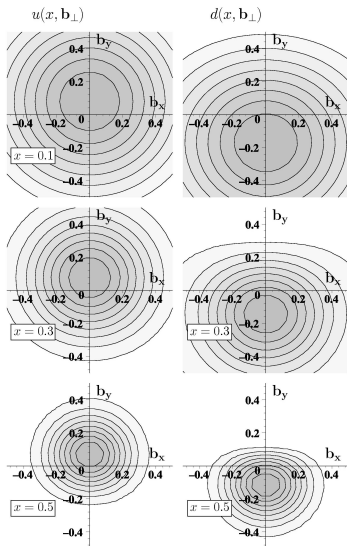
$q(x, \mathbf{b}_\perp)$  for unpol. p



### unpolarized proton

- $q(x, \mathbf{b}_\perp) = \int \frac{d^2\Delta_\perp}{(2\pi)^2} H(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$
- $F_1(-\Delta_\perp^2) = \int dx H(x, 0, -\Delta_\perp^2)$
- $x$  = momentum fraction of the quark
- $\mathbf{b}_\perp$  relative to  $\perp$  center of momentum
- small  $x$ : large 'meson cloud'
- larger  $x$ : compact 'valence core'
- $x \rightarrow 1$ : active quark becomes center of momentum

$\hookrightarrow \vec{b}_\perp \rightarrow 0$  (narrow distribution) for  $x \rightarrow 1$



proton polarized in  $+\hat{x}$  direction

no axial symmetry!

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} \\ - \frac{1}{2M} \frac{\partial}{\partial b_y} \int \frac{d^2 \Delta_\perp}{(2\pi)^2} E_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

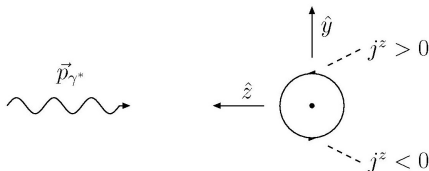
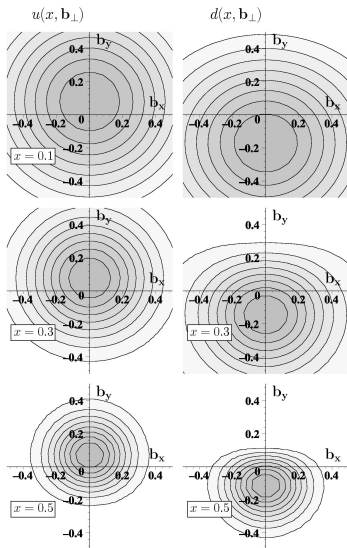
Physics: relevant density in DIS is

$j^+ \equiv j^0 + j^3$  and left-right asymmetry from  $j^3$

intuitive explanation

- moving Dirac particle with anomalous magnetic moment has electric dipole moment  $\perp$  to  $\vec{p}$  and  $\perp$  magnetic moment

$\rightarrow \gamma^*$  'sees' flavor dipole moment of oncoming nucleon

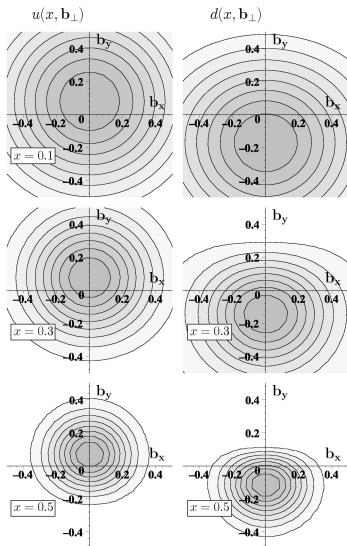


proton polarized in  $+\hat{x}$  direction

no axial symmetry!

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} \\ - \frac{1}{2M} \frac{\partial}{\partial b_y} \int \frac{d^2 \Delta_\perp}{(2\pi)^2} E_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

Physics: relevant density in DIS is  $j^+ \equiv j^0 + j^3$  and left-right asymmetry from  $j^3$



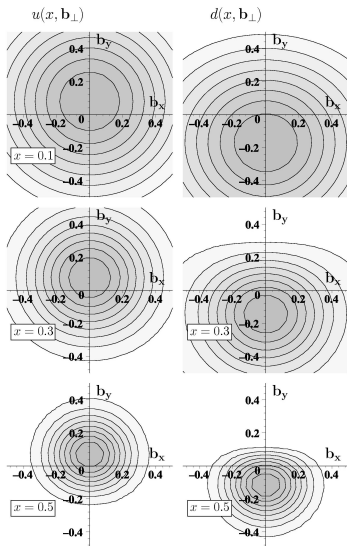
proton polarized in  $+\hat{x}$  direction

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} \\ - \frac{1}{2M} \frac{\partial}{\partial b_y} \int \frac{d^2 \Delta_\perp}{(2\pi)^2} E_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

sign & magnitude of the average shift

model-independently related to p/n  
anomalous magnetic moments:

$$\langle b_y^q \rangle \equiv \int dx \int d^2 b_\perp q(x, \mathbf{b}_\perp) b_y \\ = \frac{1}{2M} \int dx E_q(x, 0) = \frac{\kappa_q}{2M}$$



sign & magnitude of the average shift

model-independently related to p/n  
anomalous magnetic moments:

$$\begin{aligned}\langle b_y^q \rangle &\equiv \int dx \int d^2 b_\perp q(x, \mathbf{b}_\perp) b_y \\ &= \frac{1}{2M} \int dx E_q(x, 0) = \frac{\kappa_q}{2M}\end{aligned}$$

$$\kappa^p = 1.913 = \frac{2}{3}\kappa_u^p - \frac{1}{3}\kappa_d^p + \dots$$

- $u$ -quarks:  $\kappa_u^p = 2\kappa_p + \kappa_n = 1.673$

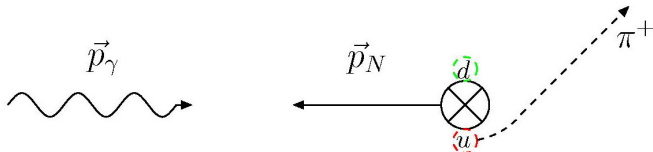
→ shift in  $+\hat{y}$  direction

- $d$ -quarks:  $\kappa_u^p = 2\kappa_n + \kappa_p = -2.033$

→ shift in  $-\hat{y}$  direction

- $\langle b_y^q \rangle = \mathcal{O}(\pm 0.2 fm)$  !!!!



example:  $\gamma p \rightarrow \pi X$ 

- $u, d$  distributions in  $\perp$  polarized proton have left-right asymmetry in  $\perp$  position space (T-even!); sign “determined” by  $\kappa_u$  &  $\kappa_d$
- attractive final state interaction (FSI) deflects active quark towards the center of momentum
- FSI translates position space distortion (before the quark is knocked out) in  $+\hat{y}$ -direction into momentum asymmetry that favors  $-\hat{y}$  direction → **chromodynamic lensing**

⇒

 $\kappa_p, \kappa_n \longleftrightarrow$  sign of SSA!!!!!!! (MB,2004)

- confirmed by HERMES & COMPASS data

$d_2 \leftrightarrow$  average  $\perp$  force on quark in DIS from  $\perp$  pol target

polarized DIS:

$$\bullet \sigma_{LL} \propto g_1 - \frac{2Mx}{\nu} g_2 \qquad \bullet \sigma_{LT} \propto g_T \equiv g_1 + g_2$$

$\hookrightarrow$  'clean' separation between  $g_2$  and  $\frac{1}{Q^2}$  corrections to  $g_1$

$$\bullet g_2 = g_2^{WW} + \bar{g}_2 \text{ with } g_2^{WW}(x) \equiv -g_1(x) + \int_x^1 \frac{dy}{y} g_1(y)$$

$$d_2 \equiv 3 \int dx x^2 \bar{g}_2(x) = \frac{1}{2M P^{+2} S_x} \langle P, S | \bar{q}(0) \gamma^+ g F^{+y}(0) q(0) | P, S \rangle$$

color Lorentz Force on ejected quark (MB, PRD 88 (2013) 114502)

$$\sqrt{2} F^{+y} = F^{0y} + F^{zy} = -E^y + B^x = -\left(\vec{E} + \vec{v} \times \vec{B}\right)^y \text{ for } \vec{v} = (0, 0, -1)$$

matrix element defining  $d_2 \leftrightarrow 1^{st}$  integration point in QS-integral

$d_2 \Rightarrow \perp$  force  $\leftrightarrow$  QS-integral  $\Rightarrow \perp$  impulse

sign of  $d_2$

$$\bullet \perp \text{ deformation of } q(x, \mathbf{b}_\perp)$$

$\hookrightarrow$  sign of  $d_2^q$ : opposite Sivers

magnitude of  $d_2$

$$\bullet \langle F^y \rangle = -2M^2 d_2 = -10 \frac{\text{GeV}}{fm} d_2$$

$$\bullet |\langle F^y \rangle| \ll \sigma \approx 1 \frac{\text{GeV}}{fm} \Rightarrow d_2 = \mathcal{O}(0.01)$$

$d_2 \leftrightarrow$  average  $\perp$  force on quark in DIS from  $\perp$  pol target

polarized DIS:

$$\bullet \sigma_{LL} \propto g_1 - \frac{2Mx}{\nu} g_2 \qquad \bullet \sigma_{LT} \propto g_T \equiv g_1 + g_2$$

$\hookrightarrow$  'clean' separation between  $g_2$  and  $\frac{1}{Q^2}$  corrections to  $g_1$

$$\bullet g_2 = g_2^{WW} + \bar{g}_2 \text{ with } g_2^{WW}(x) \equiv -g_1(x) + \int_x^1 \frac{dy}{y} g_1(y)$$

$$d_2 \equiv 3 \int dx x^2 \bar{g}_2(x) = \frac{1}{2M P^{+2} S_x} \langle P, S | \bar{q}(0) \gamma^+ g F^{+y}(0) q(0) | P, S \rangle$$

color Lorentz Force on ejected quark (MB, PRD 88 (2013) 114502)

$$\sqrt{2} F^{+y} = F^{0y} + F^{zy} = -E^y + B^x = -\left(\vec{E} + \vec{v} \times \vec{B}\right)^y \text{ for } \vec{v} = (0, 0, -1)$$

sign of  $d_2$

$$\bullet \perp \text{ deformation of } q(x, \mathbf{b}_\perp)$$

$\hookrightarrow$  sign of  $d_2^q$ : opposite Sivers

magnitude of  $d_2$

$$\bullet \langle F^y \rangle = -2M^2 d_2 = -10 \frac{\text{GeV}}{fm} d_2$$

$$\bullet |\langle F^y \rangle| \ll \sigma \approx 1 \frac{\text{GeV}}{fm} \Rightarrow d_2 = \mathcal{O}(0.01)$$

consistent with experiment (JLab, SLAC), model calculations (Weiss), and lattice QCD calculations (Göckeler et al., 2005)

chirally even spin-dependent twist-3 PDF  $g_2(x)$ 

MB, PRD 88 (2013) 114502

- $\int dx x^2 g_2(x) \Rightarrow \perp$  force on unpolarized quark in  $\perp$  polarized target  
 $\hookrightarrow$  ‘Sivers force’

scalar twist-3 PDF  $e(x)$ 

MB, PRD 88 (2013) 114502

- $\int dx x^2 e(x) \Rightarrow \perp$  force on  $\perp$  polarized quark in unpolarized target  
 $\hookrightarrow$  ‘Boer-Mulders force’

chirally odd spin-dependent twist-3 PDF  $h_2(x)$ 

M.Abdallah &amp; MB, PRD94 (2016) 094040

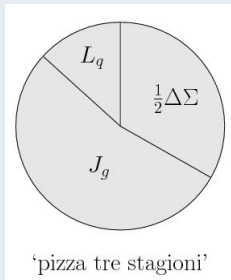
- $\int dx x^2 h_2(x) = 0$   
 $\hookrightarrow \perp$  force on  $\perp$  pol. quark in long. pol. target vanishes due to parity
- $\int dx x^3 h_2(x) \Rightarrow$  long. gradient of  $\perp$  force on  $\perp$  polarized quark in long. polarized target  
 $\hookrightarrow$  chirally odd ‘wormgear force’

## force distributions

F.Aslan & MB: *work in progress*

- use FT of twist-3 GPDs to map these forces in the  $\perp$  plane

## Ji decomposition

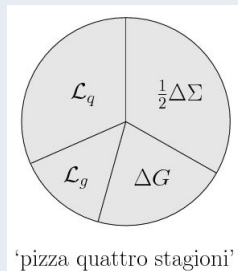


$$\frac{1}{2} = \sum_q \left( \frac{1}{2}\Delta q + L_q \right) + J_g$$

$$\vec{L}_q = \vec{r} \times (\vec{p} - g\vec{A})$$

- manifestly gauge inv. & local
- DVCS  $\longrightarrow$  GPDs  $\longrightarrow L^q$

## Jaffe-Manohar decomposition



$$\frac{1}{2} = \sum_q \left( \frac{1}{2}\Delta q + \mathcal{L}_q \right) + \Delta G + \mathcal{L}_g$$

$$\vec{\mathcal{L}}_q = \vec{r} \times \vec{p}$$

- manifestly gauge inv.  $\rightarrow$  nonlocal
- $\overleftrightarrow{p} \overleftrightarrow{p} \longrightarrow \Delta G \longrightarrow \mathcal{L} \equiv \sum_{i \in q,g} \mathcal{L}^i$

How large is difference  $\mathcal{L}_q - L_q$  in QCD and what does it represent?

Ji

$$\frac{1}{2} = \sum_q \frac{1}{2} \Delta q + \mathbf{L}_q + J_g$$

$$\mathbf{L}_q = \vec{r} \times (\vec{p} - g\vec{A})$$

Jaffe-Manohar

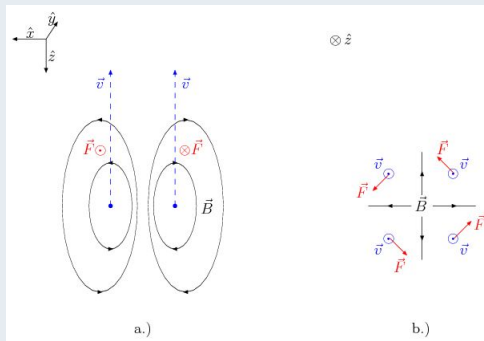
$$\frac{1}{2} = \sum_q \frac{1}{2} \Delta q + \mathcal{L}_q + \Delta G + \mathcal{L}_g$$

$$\vec{\mathcal{L}}^q = \vec{r} \times \vec{p}$$

difference  $\mathcal{L}^q - L^q$  (MB, PRD 88 (2013) 014014)

$\mathcal{L}^q - L^q = \Delta L_{FSI}^q =$  change in OAM as quark leaves nucleon

example: torque in magnetic dipole field



difference  $\mathcal{L}^q - L^q$

$\mathcal{L}_{JM}^q - L_{Ji}^q = \Delta L_{FSI}^q = \text{change in OAM as quark leaves nucleon}$

$$\mathcal{L}_{JM}^q - L_{Ji}^q = -g \int d^3x \langle P, S | \bar{q}(\vec{x}) \gamma^+ [\vec{x} \times \int_{x^-}^{\infty} dr^- F^{+\perp}(r^-, \mathbf{x}_{\perp})]^z q(\vec{x}) | P, S \rangle$$

$e^+$  moving through dipole field of  $e^-$

- consider  $e^-$  polarized in  $+\hat{z}$  direction

→  $\vec{\mu}$  in  $-\hat{z}$  direction (Figure)

- $e^+$  moves in  $-\hat{z}$  direction

→ net torque **negative**

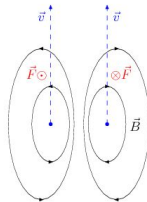
sign of  $\mathcal{L}^q - L^q$  in QCD

- color electric force between two  $q$  in nucleon attractive

→ same as in positronium

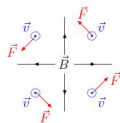
- spectator spins positively correlated with nucleon spin

→ expect  $\mathcal{L}^q - L^q < 0$  in nucleon



a.)

$\otimes \hat{z}$



b.)

difference  $\mathcal{L}^q - L^q$  $\mathcal{L}_{JM}^q - L_{Ji}^q = \Delta L_{FSI}^q =$  change in OAM as quark leaves nucleon

$$\mathcal{L}_{JM}^q - L_{Ji}^q = -g \int d^3x \langle P, S | \bar{q}(\vec{x}) \gamma^+ [\vec{x} \times \int_{x^-}^{\infty} dr^- F^{+\perp}(r^-, \mathbf{x}_{\perp})]^z q(\vec{x}) | P, S \rangle$$

 $e^+$  moving through dipole field of  $e^-$ 

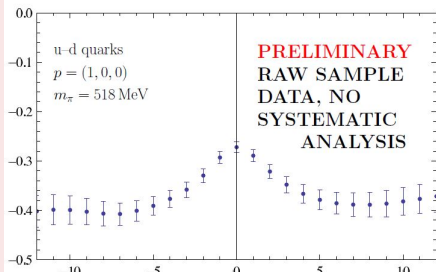
- consider  $e^-$  polarized in  $+\hat{z}$  direction
- $\rightarrow \vec{\mu}$  in  $-\hat{z}$  direction (Figure)
- $e^+$  moves in  $-\hat{z}$  direction
- $\rightarrow$  net torque **negative**

sign of  $\mathcal{L}^q - L^q$  in QCD

- color electric force between two  $q$  in nucleon attractive
- $\rightarrow$  same as in positronium
- spectator spins positively correlated with nucleon spin
- $\rightarrow$  expect  $\mathcal{L}^q - L^q < 0$  in nucleon

lattice QCD (M.Engelhardt)

- $L_{staple}$  vs. staple length
- $\rightarrow L_{Ji}^q$  for length = 0
- $\rightarrow \mathcal{L}_{JM}^q$  for length  $\rightarrow \infty$



- shown  $L_{staple}^u - L_{staple}^d$
- similar result for each  $\Delta L_{FSI}^q$



twist-3 GPDs (Meissner, Metz, Schlegel) – example  $\Gamma = \gamma^x \gamma_5$

$$\int dz^- e^{ixz^-} \bar{p}^+ \langle P' | \bar{q}(z^-/2) \gamma^x \gamma_5 q(-z^-/2) | p \rangle = \frac{-i}{2(P^+)^2} \times$$

$$\bar{u}(p') \left[ i\sigma^{+y} H'_{2T} + \frac{\gamma^+ \Delta^i - \Delta^+ \gamma^y}{2M} E'_{2T} + \frac{P^+ \Delta^y - \Delta^+ P^y}{M^2} \tilde{H}'_{2T} + \frac{\gamma^+ P^y - P^+ \gamma^y}{2M} \tilde{E}'_{2T} \right] u(p)$$

- $H'_{2T}(x, \xi, t)$  etc. all twist 3 GPDs
- similar for  $\gamma = \mathbb{1}, \gamma_5, \gamma^\perp, \sigma^{ij} \gamma_5, \sigma^{+-} \gamma_5$  (16 GPDs)

physics of twist-3 GPDs

- $x^2$ -moment  $\Rightarrow \langle p' | \bar{q}(0) \gamma^x \gamma_5 \overset{\leftrightarrow}{D}_- q(0) | p \rangle \rightsquigarrow \langle p' | \bar{q}(0) \gamma^+ F^{+y} q(0) | p \rangle +$   
twist-2
- $\hookrightarrow x^2$ -moment of twist-3 GPDs  $\rightsquigarrow$  **matrix elements of 'force operator'**  $\bar{q} \gamma^+ F^{+y} q$
- $\hookrightarrow \perp$  momentum transfer  $\rightsquigarrow$  **spatial resolution**
- $\hookrightarrow$  **transverse force femtography** ( $\perp$  vectorfields of  $\perp$  forces)
- $\Gamma = \gamma^x \gamma_5$ : force in  $\hat{y}$  direction for unpolarized quarks

twist-3 GPDs (Meissner, Metz, Schlegel) – example  $\Gamma = \gamma^\perp \gamma_5$

$$\bar{u}(p') \left[ i\sigma^{+y} H'_{2T} + \frac{\gamma^+ \Delta^i - \Delta^+ \gamma^y}{2M} E'_{2T} + \frac{P^+ \Delta^y - \Delta^+ P^y}{M^2} \tilde{H}'_{2T} + \frac{\gamma^+ P^y - P^+ \gamma^y}{2M} \tilde{E}'_{2T} \right] u(p)$$

physics of twist-3 GPDs: transverse force femtography

- $x^2$ -moment  $\Rightarrow \langle p' | \bar{q}(0) \gamma^j \gamma_5 \overset{\leftrightarrow 2}{D}_- q(0) | p \rangle \rightsquigarrow \langle p' | \bar{q}(0) \gamma^+ F^{+i} q(0) | p \rangle + \text{twist-2}$
- What can we learn from

$$\langle p' | \bar{q}(0) \gamma^+ F^{+i} q(0) | p \rangle \leftrightarrow \bar{u}(p') \left[ i\sigma^{+y} H'_{2T} + \frac{\gamma^+ \Delta^y - \Delta^+ \gamma^y}{2M} E'_{2T} + \frac{P^+ \Delta^y - \Delta^+ P^y}{M^2} \tilde{H}'_{2T} + \frac{\gamma^+ P^y - P^+ \gamma^y}{2M} \tilde{E}'_{2T} \right] u(p)$$

l.h.s.

- similar to twist-2, impact parameter interpretation only for  $\Delta^+ = 0$
  - Dirac matrix  $\gamma^+$
- $\hookrightarrow$  no sensitivity to quark pol.

r.h.s.

- different terms depend on nucleon polarization

## physics of twist-3 GPDs: transverse force femtography

$$\langle p' | \bar{q}(0) \gamma^+ F^{+i} q(0) | p \rangle \leftrightarrow$$

$$\bar{u}(p') \left[ i\sigma^{+y} H'_{2T} + \frac{\gamma^+ \Delta^y}{2M} E'_{2T} + \frac{P^+ \Delta^y}{M^2} \tilde{H}'_{2T} - \frac{P^+ \gamma^y}{2M} \tilde{E}'_{2T} \right] u(p)$$

$$i\sigma^{+y} H'_{2T}$$

- nucleon polarized in  $\hat{x}$  direction
- magnitude axially symmetric
- force in  $\hat{y}$  direction

↪ forward limit: Sivers  $\Rightarrow H'_{2T} \perp$  position dependence of Sivers

## unpolarized target

- contributing terms:  $E'_{2T}$  &  $\tilde{H}'_{2T}$
- Gordon identity:  $2P^+ \longrightarrow 2M\gamma^+ + i\sigma^{+j}\Delta^j$

↪ unpolarized target:  $\frac{\gamma^+}{2M} \Delta^y \left( E'_{2T} + 2\tilde{H}'_{2T} \right)$

↪  $\perp$  force on unpolarized quarks in unpolarized target described by  $\perp$  gradient of FT of  $E'_{2T} + 2\tilde{H}'_{2T}$

## physics of twist-3 GPDs: transverse force femtography

$$\langle p' | \bar{q}(0) \gamma^+ F^{+i} q(0) | p \rangle \leftrightarrow$$

$$\bar{u}(p') \left[ i\sigma^{+y} H'_{2T} + \frac{\gamma^+ \Delta^y}{2M} E'_{2T} + \frac{P^+ \Delta^y}{M^2} \tilde{H}'_{2T} - \frac{P^+ \gamma^y}{2M} \tilde{E}'_{2T} \right] u(p)$$

### unpolarized target

- contributing terms:  $E'_{2T}$  &  $\tilde{H}'_{2T}$
- Gordon identity:  $2P^+ \longrightarrow 2M\gamma^+ + i\sigma^{+j}\Delta^j$
- $\hookrightarrow$  unpolarized target:  $\frac{\gamma^+}{2M}\Delta^y \left( E'_{2T} + 2\tilde{H}'_{2T} \right)$
- $\hookrightarrow$   $\perp$  force on unpolarized quarks in unpolarized target described by  $\perp$  gradient of FT of  $E'_{2T} + 2\tilde{H}'_{2T}$

### $\tilde{H}'_{2T}$

- additional  $\Delta^y \frac{\sigma^{+j}\Delta^j}{2M^2} \tilde{H}'_{2T}$  from Gordon identity
- $\hookrightarrow \nabla^y \left( \vec{S}_\perp \times \vec{\nabla}_\perp \right) \text{FT} \left[ \tilde{H}'_{2T} \right]$  force

## physics of twist-3 GPDs: transverse force femtography

$$\langle p' | \bar{q}(0) \gamma^+ F^{+i} q(0) | p \rangle \leftrightarrow$$

$$\bar{u}(p') \left[ i \sigma^{+y} H'_{2T} + \frac{\gamma^+ \Delta^y}{2M} E'_{2T} + \frac{P^+ \Delta^y}{M^2} \tilde{H}'_{2T} - \frac{P^+ \gamma^y}{2M} \tilde{E}'_{2T} \right] u(p)$$

$\tilde{H}'_{2T}$

- additional  $\Delta^y \frac{\sigma^{+j} \Delta^j}{2M^2} \tilde{H}'_{2T}$  from Gordon identity
- $\hookrightarrow \nabla^y \left( \vec{S}_\perp \times \vec{\nabla}_\perp \right) \text{FT} \left[ \tilde{H}'_{2T} \right]$  force

$\tilde{E}'_{2T}$

- Gordon:  $\sigma^{yx} \Delta^x \longrightarrow S_z \nabla^x \text{FT} \left[ \tilde{E}'_{2T} \right]$
- $\hookrightarrow$  'wormgear' force ( $g_{1L}$ )

## physics of twist-3 GPDs: transverse force femtography

$$\langle p' | \bar{q}(0) \gamma^+ F^{+i} q(0) | p \rangle \leftrightarrow$$

$$\bar{u}(p') \left[ i \sigma^{+y} H'_{2T} + \frac{\gamma^+ \Delta^y}{2M} E'_{2T} + \frac{P^+ \Delta^y}{M^2} \tilde{H}'_{2T} - \frac{P^+ \gamma^y}{2M} \tilde{E}'_{2T} \right] u(p)$$

$\tilde{H}'_{2T}$

- additional  $\Delta^y \frac{\sigma^{+j} \Delta^j}{2M^2} \tilde{H}'_{2T}$  from Gordon identity
- $\hookrightarrow \nabla^y \left( \vec{S}_\perp \times \vec{\nabla}_\perp \right) \text{FT} \left[ \tilde{H}'_{2T} \right]$  force

$\tilde{E}'_{2T}$

- Gordon:  $\sigma^{yx} \Delta^x \longrightarrow S_z \nabla^x \text{FT} \left[ \tilde{E}'_{2T} \right]$
- $\hookrightarrow$  'wormgear' force ( $g_{1L}$ )

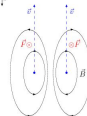
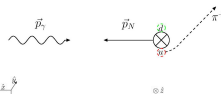
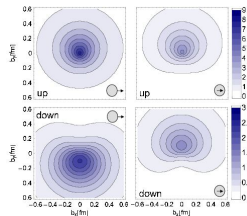
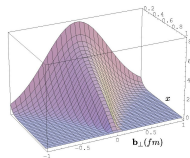
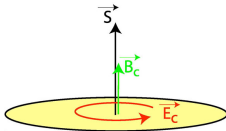
- sigimer for  $\gamma^\perp \gamma_5 \longrightarrow \gamma^\perp, \mathbb{1}, \dots$
- $x^2$  moment of each twist-3 GPD (after subtracting WW and other twist-2) yields femtoimage of force

## TMDs

| QUARKS | <i>unpolarized</i> | <i>chiral</i> | <i>transverse</i>      |
|--------|--------------------|---------------|------------------------|
| U      | $f_1$              |               | $h_1^\perp$            |
| L      |                    | $g_{1L}$      | $h_{1L}^\perp$         |
| T      | $f_{1T}^\perp$     | $g_{1T}$      | $h_{1T}, h_{1T}^\perp$ |

- each TMD affected by final state interactions (force)!
- for each case, twist-3 GPD provides insight about FSI contribution
- Sivers/Boer Mulders ( $f_{1T}^\perp/h_1^\perp$ ) only FSI
- twist-3 GPDs provide local information

- GPDs  $\xrightarrow{FT} q(x, \mathbf{b}_\perp)$  '3d imaging'
  - $x^2$  moment of twist-3 PDFs  $\rightarrow$  force
  - $x^2$  moment of twist-3 GPDs
- $\hookrightarrow \bar{q}\gamma^+ F^{+\perp} q$  distribution
- $\hookrightarrow \perp$  force femtography





- GPDs  $\xrightarrow{FT} q(x, \mathbf{b}_\perp)$  '3d imaging'
  - $x^2$  moment of twist-3 PDFs  $\rightarrow$  force
  - $x^2$  moment of twist-3 GPDs
- $\hookrightarrow \bar{q}\gamma^+ F^{+\perp} q$  distribution
- $\hookrightarrow \perp$  force femtography

