Pseudo Distributions on the



WILLIAM & MARY

CHARTERED 1693

Lattice



Joe Karpie William & Mary / Jefferson Lab Light Cone 2018

In Collaboration with

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Introduction

- Lattice calculations moving from Hadron bulk properties (masses, charges,...) to Hadron distributions (Form Factors, Structure Functions, Distribution functions, ...)
- Project Goals
 - Long Term: Study various methods of calculating parton distributions from ab initio Lattice QCD
 - Short Term: Understand systematic effects in the simple case of iso-vector unpolarized quark PDF
- Mellin moments and OPE
 - Restricted to low moments by reduced rotational symmetry
- Hadronic Tensor Methods
 - "Light-like" separated Hadronic TensorK-F Liu et al Phys. Rev. Lett. 72 1790 (1994), Phys. Rev. D62 (2000) 074501
 - O Good lattice cross sections Y.-Q. Ma J.-W. Qiu (2014) 1404.6860 Y.-Q. Ma, J.-W. Qiu (2017) 1709.03018
- Ioffe Time Pseudo Distribution Methods
 - Quasi PDF X. Ji, Phys.Rev.Lett. 110, (2013)
 - Pseudo PDF A. Radyushkin Phys.Lett. B767 (2017)
- J.-W. Chen et.al. (2018) 1803.04393 C Alexandrou et.al. (2018) 1803.02685

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A Chambers et.al (2017) 1703.01153

K. Orginos, A Radyushkin, JK, S Zafeiropoulos (2017) 1706.05373

Ioffe Time distribution

$$\nu = p \cdot \xi$$

V. Braun, et. al Phys. Rev. D 51, 6036 (1995)

•
$$\mathcal{I}(\nu,\mu^2) = \int_{-1}^{1} dx e^{i\nu x} f(x,\mu^2) = \frac{1}{2p^+} \langle h(p) | \bar{\psi}_q(0,\xi^-,0_T) \gamma^+ W((0,\xi^-,0_T);0) \psi_q(0) | h(p) \rangle_{\mu^2}$$

- CP Even/Odd combinations
 - $\circ \quad \text{Even:} \quad q_-(x) = f(x) + f(-x) = q(x) \bar{q}(x) \equiv q_V(x)$

$$\begin{array}{l} \circ \quad \text{Odd:} \quad q_+(x) = f(x) - f(-x) = q(x) + \bar{q}(x) = q_V(x) + 2\bar{q}(x) \\ \\ \mathfrak{Re}\left[\mathcal{I}(\nu)\right] = \int_0^1 dx \cos(\nu x) q_V(x) \equiv \mathcal{I}_V(\nu) \\ \\ \mathfrak{Im}\left[\mathcal{I}(\nu)\right] = \int_0^1 dx \sin(\nu x) (q(x) + \bar{q}(x)) \end{array}$$

• Perturbative DGLAP evolution I.I. Balitsky and V.M. Braun, Nucl. Phys. B311, 541 (1988)

$$\mathcal{I}_{v}(\nu,\mu_{2}^{2}) = \mathcal{I}_{v}(\nu,\mu_{1}^{2}) - \frac{C_{F}\alpha_{s}}{2\pi}\log\frac{\mu_{2}^{2}}{\mu_{1}^{2}}\int_{0}^{1}du\left[\frac{1}{2}\delta(1-u) - (1-u) - 2[\frac{u}{1-u}]_{+}\right]\mathcal{I}_{v}(u\nu,\mu_{1}^{2})$$

Ioffe Time Pseudo Distributions

• A general matrix element of interest

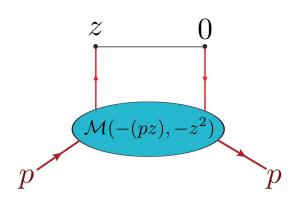
 $M^{\alpha}(z,p) = \langle h(p) | \bar{\psi}_q(z) \gamma^{\alpha} W(z;0) \psi_q(0) | h(p) \rangle$

- Lorentz decomposition
 - Use of symmetry
 - \circ Choice of p, z, and α can remove higher twist term

$$M^{\alpha}(z,p) = 2p^{\alpha}\mathcal{M}_p(\nu,z^2) + z^{\alpha}\mathcal{M}_z(\nu,z^2)$$

- Factorization Relation to ITDF
 - Perturbatively calculable Wilson coefficients for each parton

$$\mathcal{M}(\nu, -z^2) = \sum_i C_i(z^2\mu^2, \alpha_s) \otimes \mathcal{I}_i(\nu, \mu^2) + H.T.$$



A. Radyushkin (2017) 1710.08813 J.-H. Zhang (2018) 1801.03023 T. Izubuchi (2018) 1801.03917

Special Cases

Light cone PDF

 $M^{\alpha}(z,p) = \langle h(p) | \bar{\psi}_q(z) \gamma^{\alpha} W(z;0) \psi_q(0) | h(p) \rangle$ $M^{\alpha}(z,p) = 2p^{\alpha}\mathcal{M}_p(\nu,z^2) + z^{\alpha}\mathcal{M}_z(\nu,z^2)$

• Light cone PDF
$$p = (p^+, \frac{m^2}{2p^+}, 0_T)$$
 $z = (0, z^-, 0_T)$ $\alpha = + \mathcal{M}_p((p^+z^-), 0) = \int_{-1}^1 dx e^{ix(p^+z^-)} f(x)$ A. Radyushkin (2017) 1612.05170

0

A. Radyushkin (2017) 1612.05170

Straight Link "Primordial" TMD

$$p = (p^+, \frac{m^2}{2p^+}, 0_T)$$
 $z = (0, z^-, z_T)$ $\alpha = +$

$$\mathcal{M}_p((p^+z^-), -z_T^2) = \int_{-1}^1 dx e^{ix(p^+z^-)} \int d^2k_T e^{ik_T \cdot z_T} F(x, k_T^2)$$

Pseudo PDF $p = (E, 0, 0, p_3)$ $z = (0, 0, 0, z_3)$ $\alpha = 0$ -1

$$\mathcal{M}_p((-z_3 * p_3), -z_3^2) = \int_{-1}^1 dx e^{ix(-z_3 * p_3)} P(x, -z_3^2)$$
⁵

Pseudo PDF vs Quasi PDF

$0.2 GeV \approx 1 fm^{-1}$

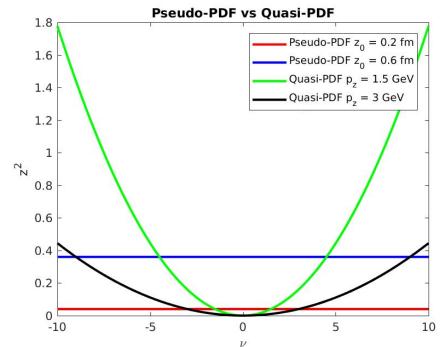
- Both are integrals of pseudo ITDF
 - Pseudo PDF has fixed invariant scale dependence

$$P(x,z_0^2) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\nu e^{-i\nu x} \mathcal{M}(\nu,z_0^2)$$

 Quasi PDF mixes invariant scales until p_z is effectively large enough

$$Q(x, p_z^2) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\nu e^{-i\nu x} \mathcal{M}(\nu, \frac{\nu^2}{p_z^2})$$

- Expedite desired limit of $z^2 \longrightarrow 0$
 - Pseudo-PDFs use reduced distributions
 - Quasi-PDFs use LaMET



Numerical Lattice Field Theory

• Importance sampling of path integral

$$\begin{array}{l} \langle O(\bar{\psi},\psi,A_{\mu})\rangle = \frac{1}{Z} \int D[\bar{\psi}] D[\psi] D[A_{\mu}] O(\bar{\psi},\psi,A_{\mu}) e^{-S(\bar{\psi},\psi,A_{\mu})} \\ \text{Correlation functions} \\ C_2(\vec{p},T) = \langle O_N(-\vec{p},T) \bar{O}_N(\vec{p},0) \rangle \end{array} \approx \frac{1}{N} \sum_{i}^{N} F_O(U_{\mu}^{(i)}) \end{array}$$

$$C_{op}(O_{op}; \vec{p}, T) = \sum_{t} \sum_{\vec{x}} \langle O_N(-\vec{p}, T) O_{op}(\vec{x}, t) \bar{O}_N(\vec{p}, 0) \rangle$$

• Feynman-Hellman matrix element extraction C. Bouchard et.al Phys. Rev. D 96, no. 1, 014504 (2017)

$$\frac{\langle N(p)|O_{op}|N(p)\rangle}{2E_{N(p)}} = \lim_{T \to \infty} \frac{1}{\tau} (R(T+\tau) - R(T)) \qquad R(T) = \frac{C_{op}(O_{op}; \vec{p}, T)}{C_2(\vec{p}, T)}$$

Renormalization and the Reduced distribution

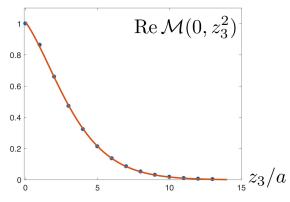
- Effective Bare matrix element $M^{eff}(T) = (R(T+1) R(T)) + O(e^{-\Delta T})$
- Vector current to cancel quark field renormalization
 - Forces matrix elements to give unit nucleon charge

$$Z_p^{-1} = M^4(0, p)$$

- Reduced distribution
 - TMD "Factorization" and suppression of polynomial corrections

$$F(x, k_T^2) = f(x)g(k_T^2) \quad \mathcal{M}(\nu, z^2) = \mathcal{M}(\nu, 0)\mathcal{M}(0, z^2)$$

 BONUS: UV corrections from Wilson line exactly cancels. No complicated renormalizations.



$$\mathfrak{M}(\nu, z^2) = \frac{\mathcal{M}(\nu, z^2)}{\mathcal{M}(0, z^2)}$$

Numerical Study

$$O_q^{\alpha}(z;T) = \sum_{\vec{x}} \bar{\psi}_q(\vec{x}+\vec{z},T)\lambda^3 \gamma^{\alpha} W((\vec{x}+\vec{z},T);(\vec{x},T))\psi_q(\vec{x},T)$$

Quenched K. Orginos, A Radyushkin, JK, S Zafeiropoulos (2017) 1706.05373

• $\beta = 6.0$ $m_{\pi} = 600$ MeV $32^3 \times 64$ a = 0.1 fm

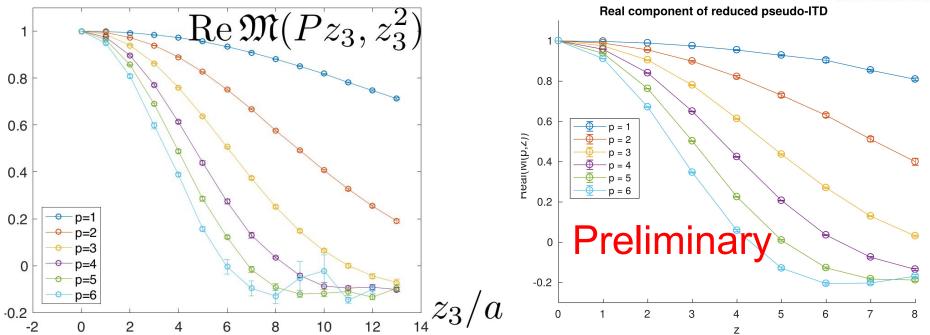
Dynamical (Preliminary) Unpublished

- a127m440: $\beta = 6.1$ $m_{\pi} = 440$ MeV $24^3 \times 64$ a = 0.127 fm
- a127m440L: $\beta = 6.1$ $m_{\pi} = 440$ MeV $32^3 \times 96$ a = 0.127 fm
- a094m400: $\beta = 6.3$ $m_{\pi} = 400$ MeV $32^3 \times 64$ a = 0.094 fm

Quenched Results

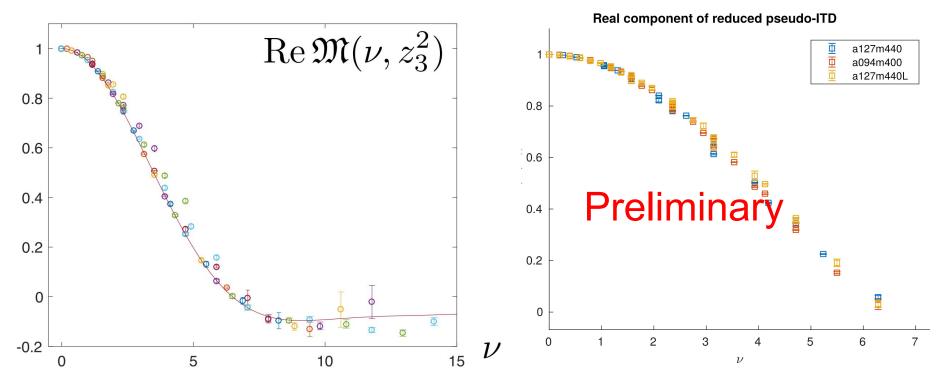
Dynamical Results

a127m440



Quenched Results

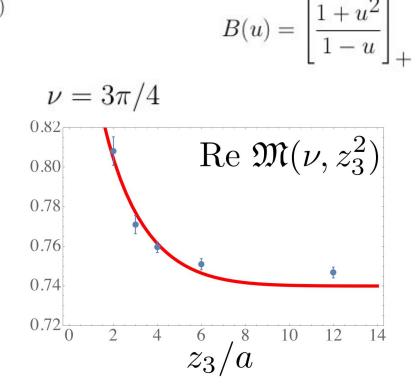
Dynamical Results



Perturbative Evolution of Lattice data $B \otimes \mathfrak{M}(\nu, z^2) = \int_0^1 du B(u) M(u\nu, z^2)$

$$\mathfrak{M}(\nu, z_0^2) = \mathfrak{M}(\nu, z^2) + \frac{C_F \alpha_s}{2\pi} \log(\frac{z^2}{z_0^2}) B \otimes \mathfrak{M}(\nu, z^2)$$

- Position space DGLAP evolution
- Improvement of Almost Universal curve
- Separation of Regimes
 - Small separation matrix elements follow log behavior expected from perturbation theory
 - Large separation matrix elements seem
 z₃ independent expected from
 cancellation of polynomial effects

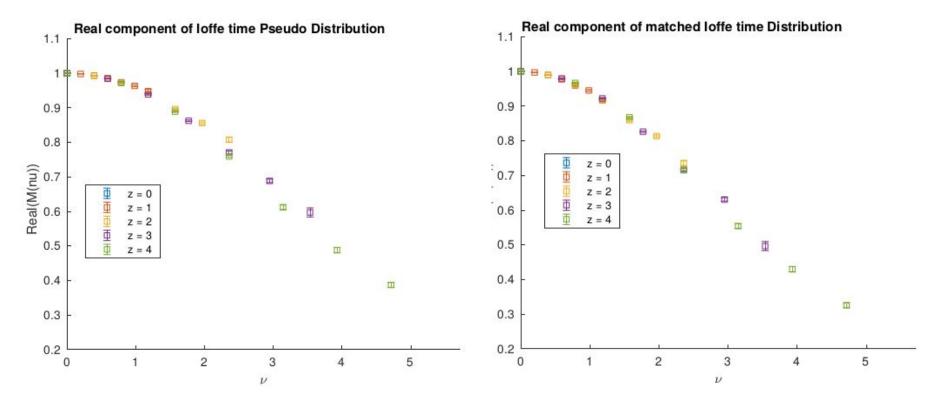


Matching Lattice data to loffe Time distribution

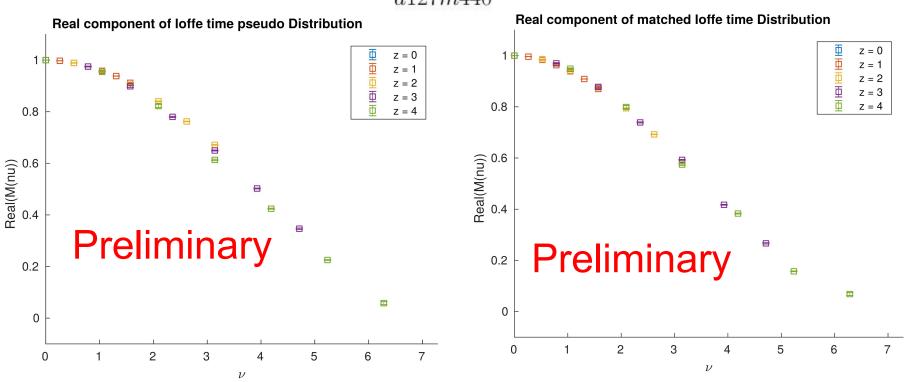
- Without proper matching, the results are only comparable to global fits up to α_s corrections
- Yet another convolution of reduced pseudo-ITD
- At 1-loop, scale evolution and matching can be simultaneous
- Results are matched to $\mu = 1$ GeV MS bar ITD

$$\begin{split} \mathcal{I}(\nu,\mu^2) &= \mathfrak{M}(\nu,z^2) + \frac{C_F \alpha_S}{2\pi} \int_0^1 du \Big(B(u) \left(\log(z^2 \mu^2 \frac{e^{2\gamma_E}}{4}) + 1 \right) \\ &+ \Big[4 \frac{\log(1-u)}{1-u} - 2(1-u) \Big]_+ \Big) M(u * \nu, z^2) \\ &= \mathfrak{M}(\nu,z^2) + \frac{C_F \alpha_S}{2\pi} \left[(\log(z^2 \mu^2 \frac{e^{2\gamma_E}}{4}) + 1) B \otimes M(\nu,z^2) + L \otimes M(\nu,z^2) \right] \end{split}$$

Perturbative Matching of Quenched Lattice data



Perturbative Matching of Dynamical Lattice data



a127m440

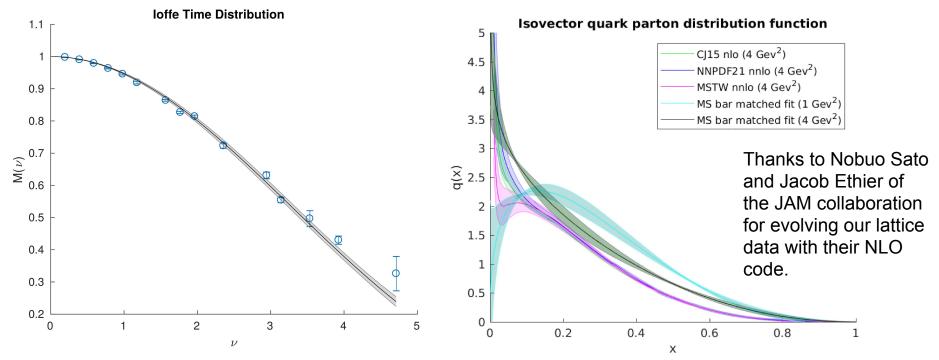
Real component and the Valence Quark distribution

• A more general model PDF

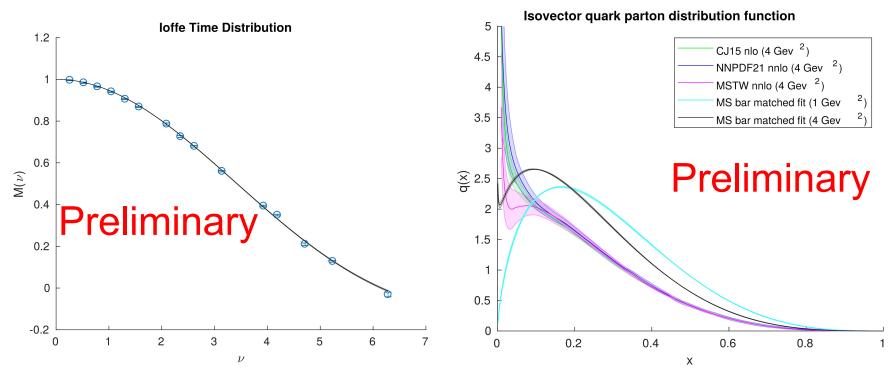
$$f_{abcd}(x) = N_{abcd} x^a (1-x)^b (1+c\sqrt{x}+dx)$$

- Limiting behaviors
 - \circ Regge a = -0.5
 - \circ Quark counting b=3
 - \circ Small Corrections $c \sim 0$ $d \sim 0$

Quenched Pseudo PDF Matched to MS bar Compared to Global fit PDFs



Dynamical Pseudo PDF Matched to MS bar Compared to Global fit PDFs



Summary

- First study of pseudo ITDF analyzed as reduced pseudo PDFs
- Quenched and Dynamical Results are in agreement with PDF fits at large x
- Treatment of z² dependence guided by data
- Application of proper matching to perturbative scheme
- Missing divergent behavior improves after scale evolution to 4 GeV²
- Systematics left to thoroughly study
 - Continuum limit
 - Control of Excited states
 - Finite Volume
 - Physical Pion mass limit

Thank you for listening