



University of New Hampshire
Nuclear & Particle Physics Group

The g_2p Experiment: A Measurement of the Proton's Spin Structure Functions 2018 Status Update

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Hall A Collaboration Meeting

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Most Slides & Figures by Ryan Zielinski

Essential Quantities in ep Scattering

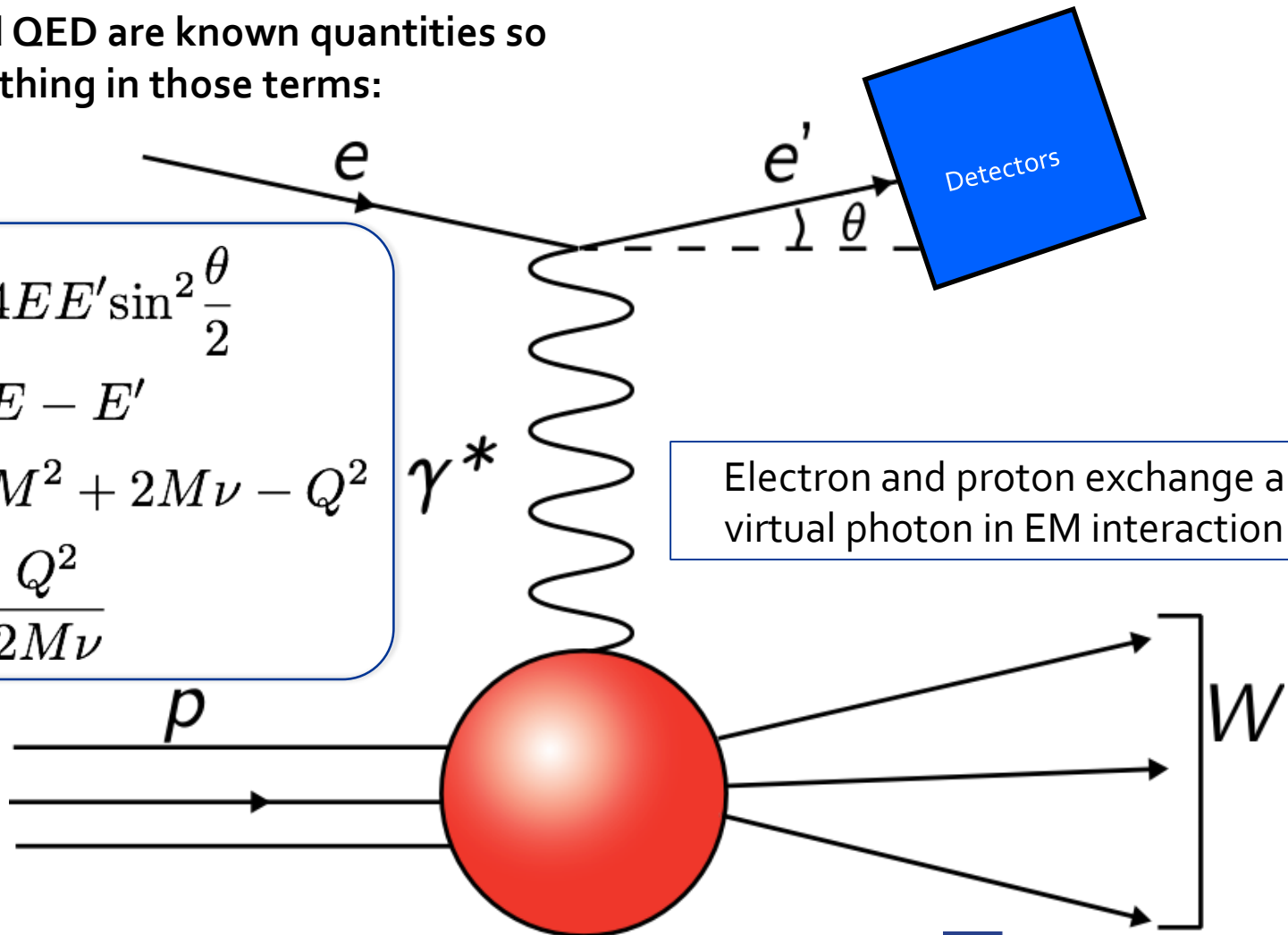
Electron and QED are known quantities so define everything in those terms:

$$Q^2 = 4EE' \sin^2 \frac{\theta}{2}$$

$$\nu = E - E'$$

$$W^2 = M^2 + 2M\nu - Q^2$$

$$x = \frac{Q^2}{2M\nu}$$



Inclusive ep Scattering Cross Sections describe normalized interaction rate

Elastic scattering: target remains in the ground state after interaction

$$E'_{\text{elas}} = \frac{E}{1 + \frac{2E}{M} \sin^2 \frac{\theta}{2}}$$

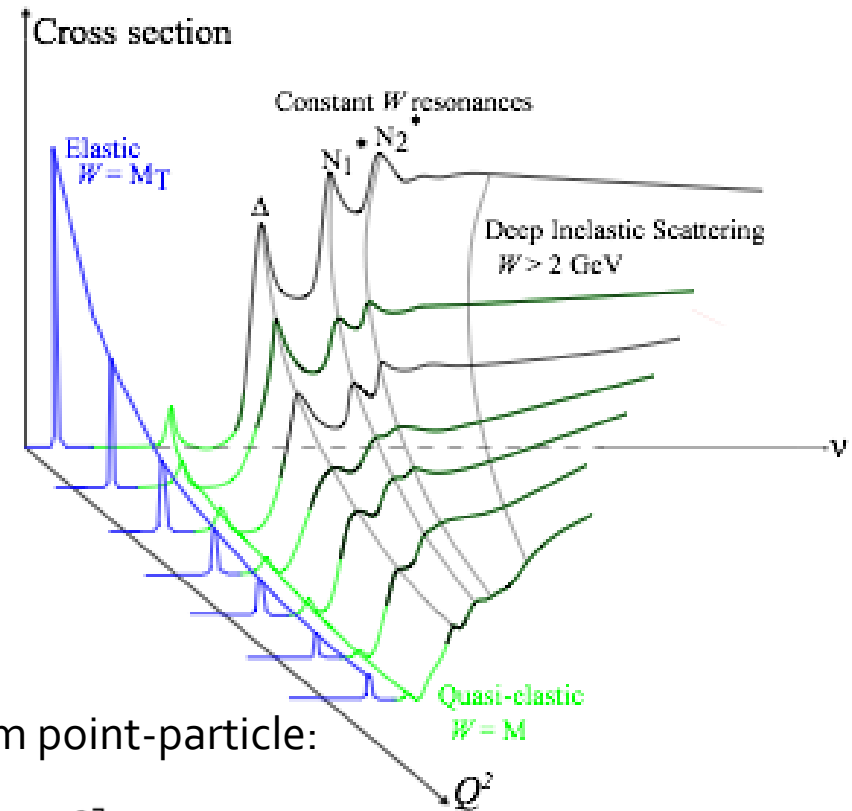
Mott cross section describes scattering from point-particle:

$$\left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} = \frac{\alpha^2}{4E^2 \sin^4 \theta/2} \cos^2 \frac{\theta}{2}$$

Rosenbluth cross section describes deviation from point-particle:

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} \left[\frac{G_E^2(Q^2) + \tau G_M^2(Q^2)}{1 + \tau} + 2\tau G_M^2(Q^2) \tan^2 \frac{\theta}{2} \right]$$

G_E and G_M related to charge and current distributions



Inclusive ep Scattering Cross Sections describe normalized interaction rate

Inelastic scattering: Target is in excited state after interaction

Structure Functions:

Inclusive *unpolarized* cross sections

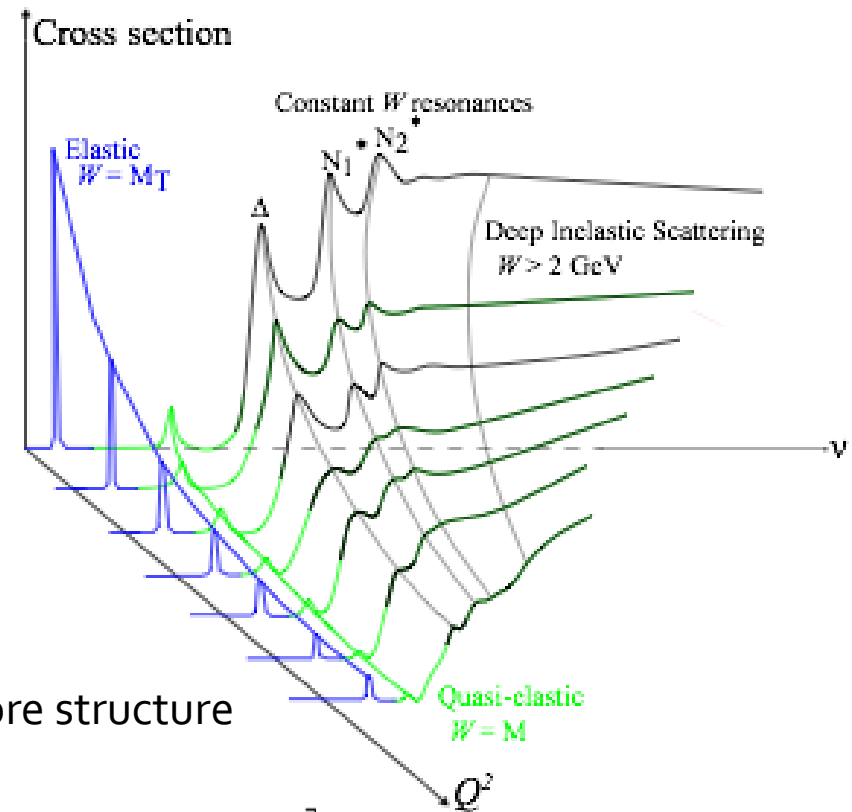
$$\frac{d^2\sigma}{d\Omega dE'} = \sigma_{\text{Mott}} \left[\frac{1}{\nu} F_2(x, Q^2) + \frac{2}{M} F_1(x, Q^2) \tan^2 \frac{\theta}{2} \right]$$

F_1 and F_2 related to quark/gluon distribution

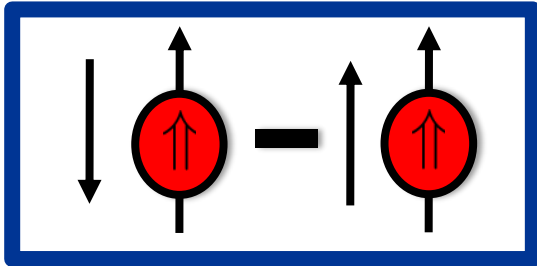
Adding a *polarized* beam and target adds two more structure functions

$$\frac{d^2\sigma^\pm}{d\Omega dE'} = \sigma_{\text{Mott}} \left[\alpha F_1(x, Q^2) + \beta F_2(x, Q^2) \pm \gamma g_1(x, Q^2) \pm \delta g_2(x, Q^2) \right]$$

g_1 and g_2 related to spin distribution



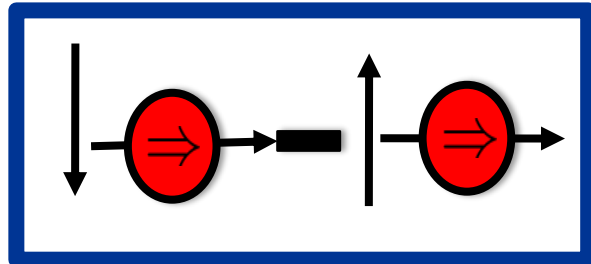
Extracting Spin Structure by Looking at Cross Section Differences



Parallel

Inclusive *polarized* cross sections

$$\frac{d^2\sigma^{\uparrow\uparrow}}{dE'd\Omega} - \frac{d^2\sigma^{\downarrow\uparrow}}{dE'd\Omega} = \frac{4\alpha^2}{M\nu Q^2} \frac{E'}{E} \left[g_1(x, Q^2) \{E + E' \cos\theta\} - \frac{Q^2}{\nu} g_2(\nu, Q^2) \right]$$



Perpendicular

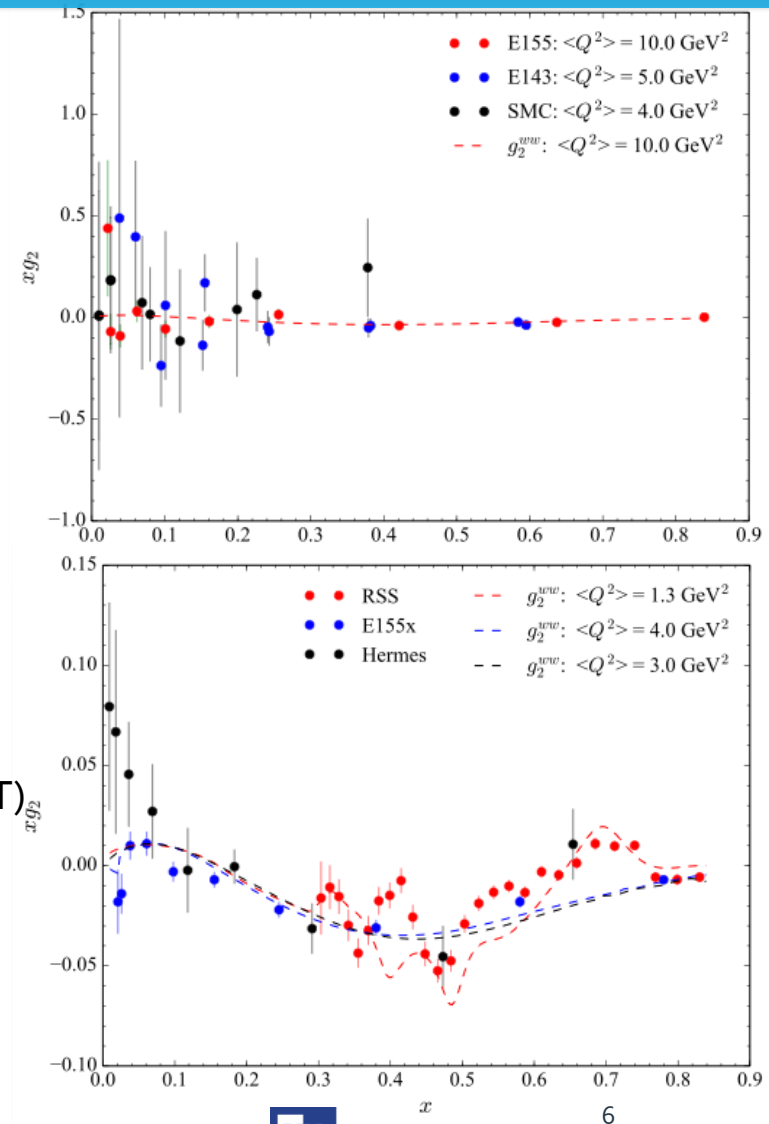
$$\frac{d^2\sigma^{\uparrow\Rightarrow}}{dE'd\Omega} - \frac{d^2\sigma^{\downarrow\Rightarrow}}{dE'd\Omega} = \frac{4\alpha^2}{M\nu Q^2} \frac{E'^2}{E} \sin\theta \left[\nu g_1(x, Q^2) + 2E g_2(\nu, Q^2) \right]$$

Two equations, two unknowns...

Motivation:

Measure a fundamental spin observable (g_2) in the region $0.02 < Q^2 < 0.20 \text{ GeV}^2$ for the first time

- Measurements at Jefferson Lab:
 - RSS – medium Q^2 (1-2 GeV^2) (published)
 - SANE – high Q^2 (2-6 GeV^2) (analysis)
 - **g_2p – low Q^2 (0.02-0.20 GeV^2) (analysis)**
- Low Q^2 is difficult:
 - Electrons strongly influenced by target field
 - Strong kinematic dependence on observables
- Low Q^2 is useful:
 - Test predictions of Chiral Perturbation Theory (χPT)
 - Test sum rules and measure moments of g_2
 - Study finite size effects of the proton
- **g_2p experiment ran spring 2012 in Hall A**

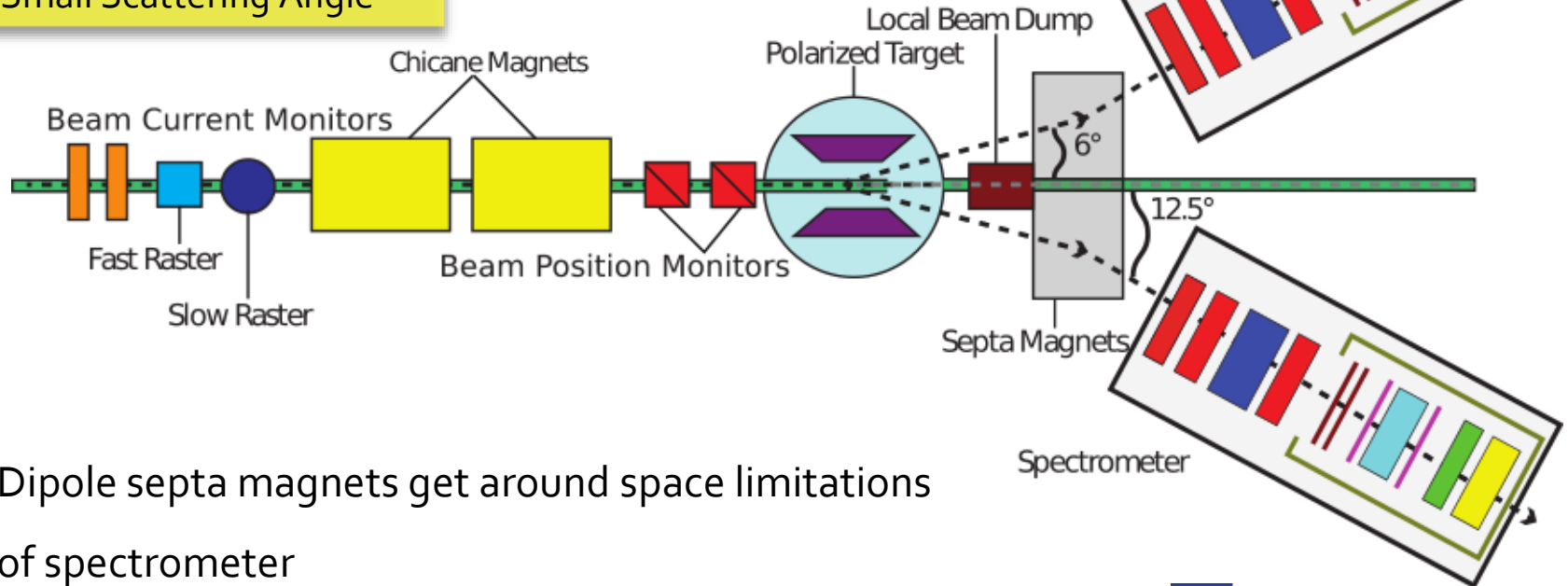


Hall A Experimental Setup:

Measuring g_2^p

- Electron Beam
- Polarized Proton Target
- Spectrometer/Detectors
- Small Scattering Angle

- Transverse polarized NH_3 target (2.5/5.0T)
- Dipole chicane magnets help compensate for target field bending of beam



- Dipole septa magnets get around space limitations of spectrometer

Polarized Protons Created with Dynamic Nuclear Polarization (DNP)

Creating initial polarization:

- Align spins in large B and low T

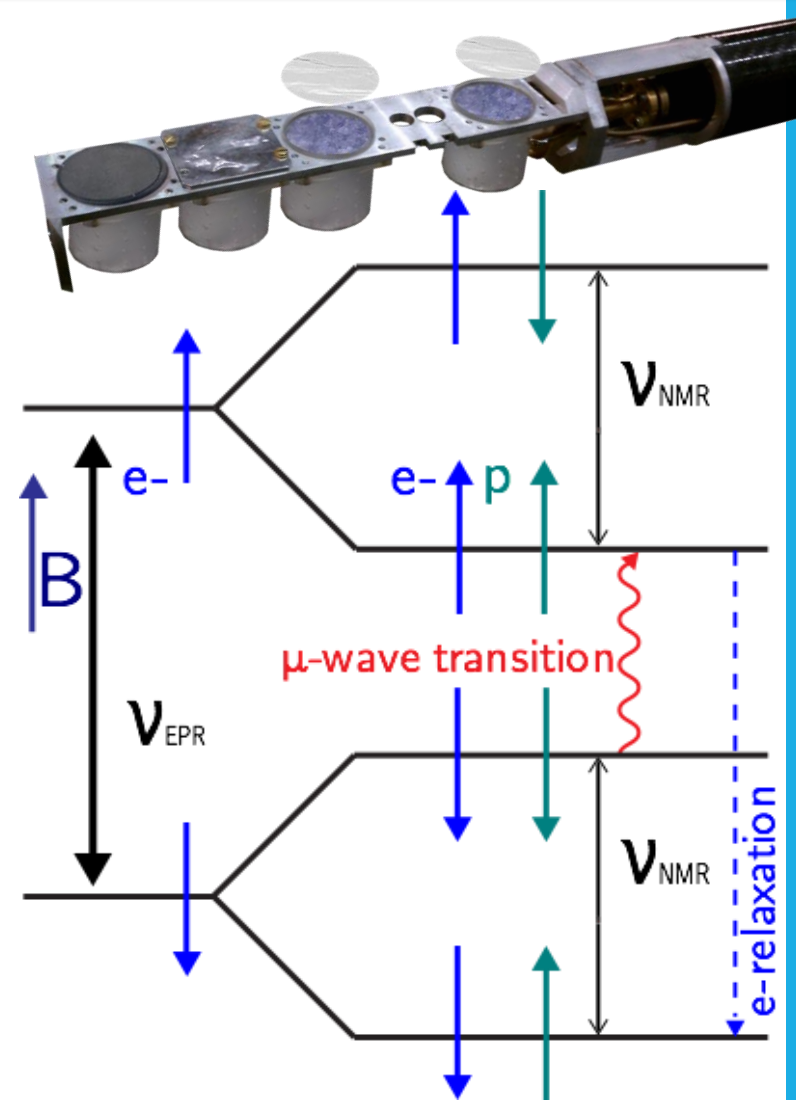
- 5.0 T / 2.5 T @ 1 K

$$P_{TE} = \frac{N^\uparrow - N^\downarrow}{N^\uparrow + N^\downarrow} = \frac{e^{\frac{\mu B}{kT}} - e^{\frac{-\mu B}{kT}}}{e^{\frac{\mu B}{kT}} + e^{\frac{-\mu B}{kT}}}$$

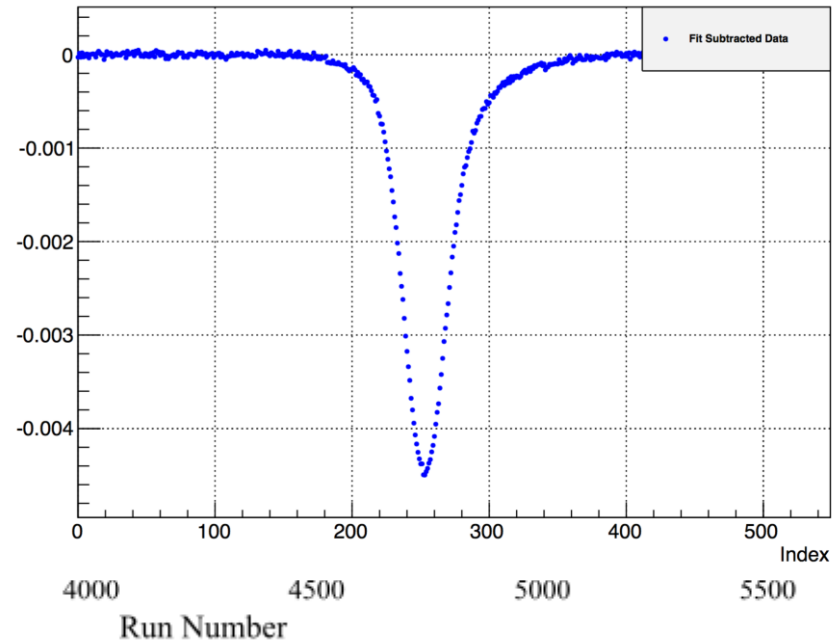
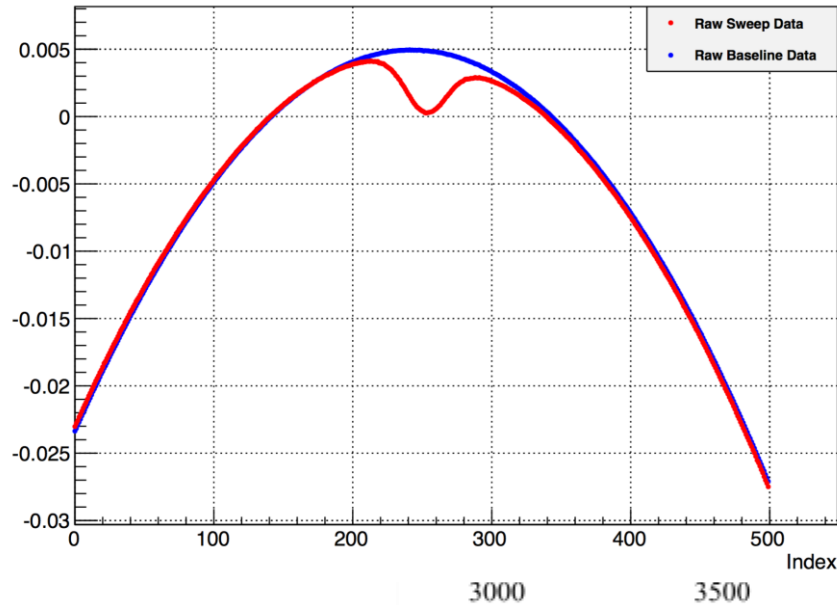
- Large μ_e ($\sim 660\mu_p$) creates large electron polarization ($\sim 99\%$ at 5T/1K)

Enhancing initial polarization:

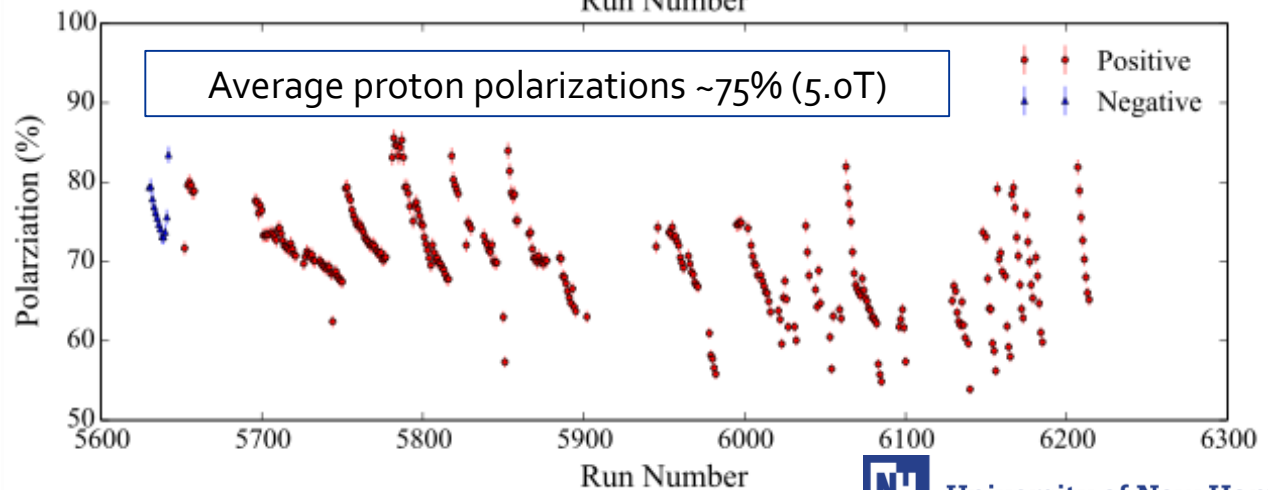
- Proton pol. much smaller ($\sim 0.5\%$ 5T) at TE
- ep spin coupling and microwaves drive pol.
- Electrons relax much quicker than protons so polarization is sustained



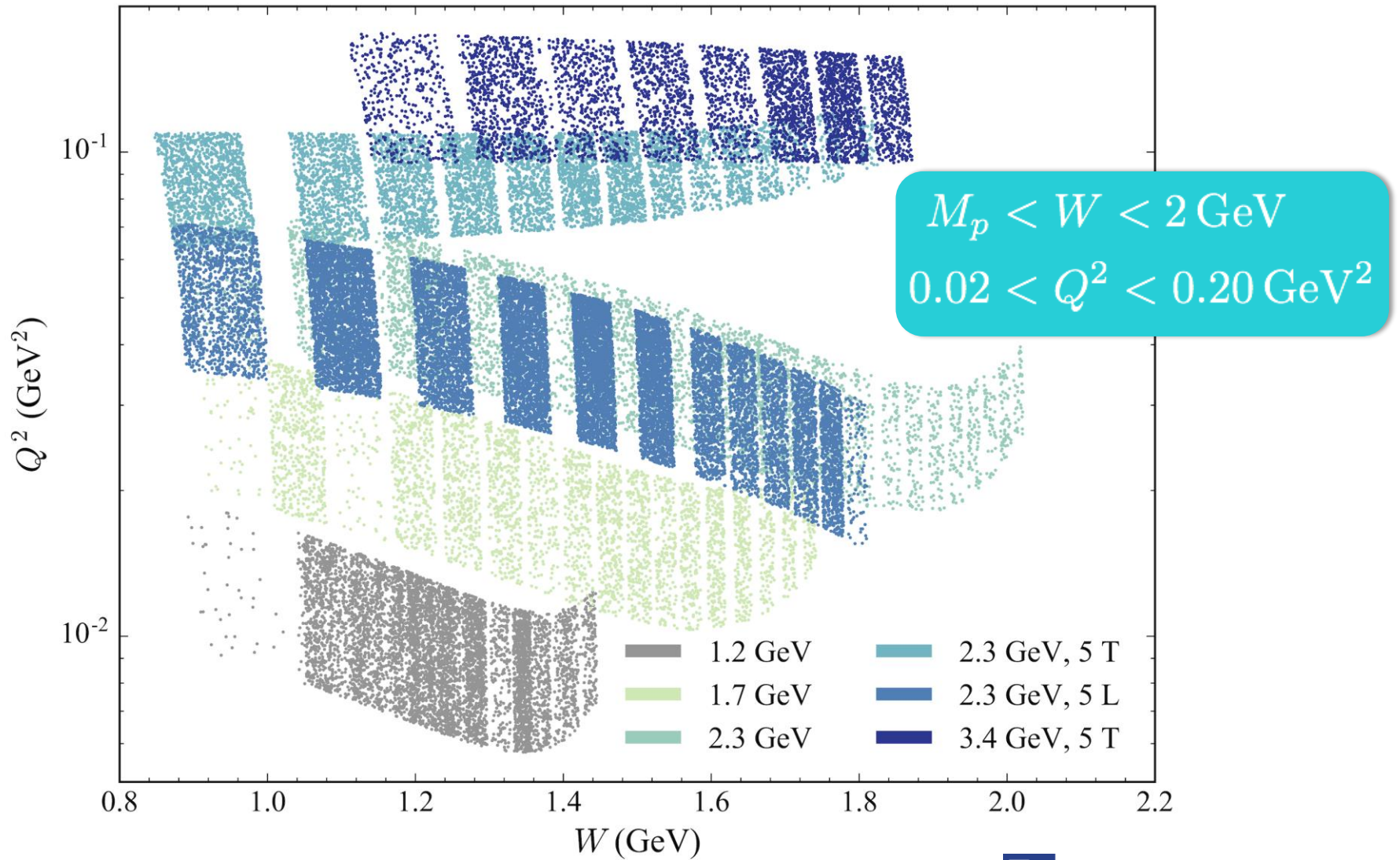
Proton Polarization Measured with Q-Meter



- LRC circuit where proton spin's couple with and change inductance



g_2p Kinematic Coverage



MEASURING $g_{1,2}$ from data

What can we measure?

1. Helicity dependent asymmetries
2. Unpolarized cross sections
3. Polarized cross sections

1. $A_{\perp} = \frac{\frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow - \uparrow\Rightarrow)}{\frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow + \uparrow\Rightarrow)}$

2. $\sigma_0 = \frac{1}{2} \frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow + \uparrow\Rightarrow)$

3. $\Delta\sigma_{\perp} = \frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow - \uparrow\Rightarrow) = 2 \cdot A_{\perp} \sigma_0$

Similar equation for parallel polarized cross section

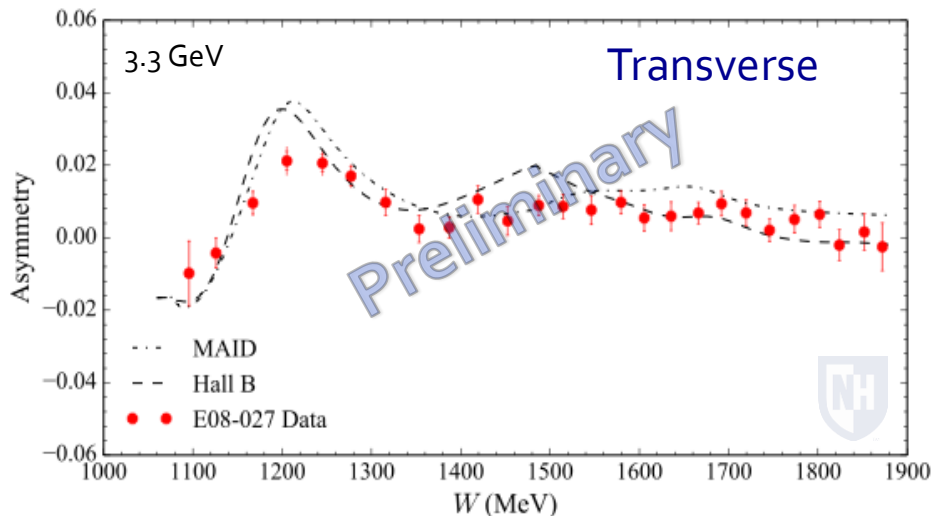
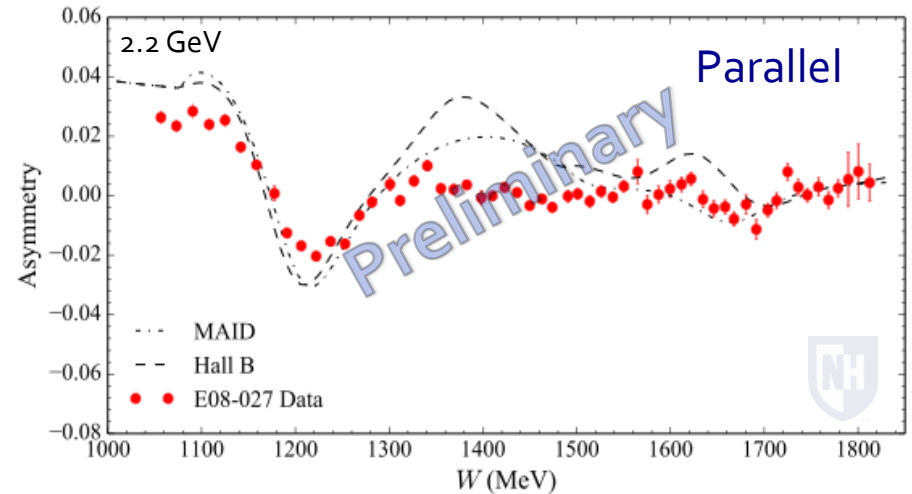
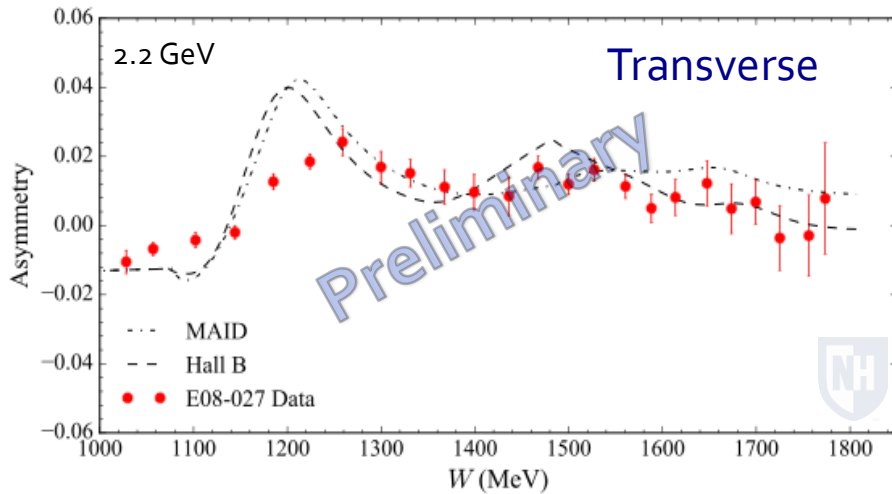
Why do it this way?

- Asymmetries are easy to measure
- Lots of data on unpolarized cross sections so models are a possibility

Need to be mindful of contributions from scattering from anything other than protons



5T Proton Asymmetries



Raw Counts:

$$Y_{\pm} = \frac{N_{\pm}}{LT_{\pm}Q_{\pm}}$$

Measured Asymmetries:

$$A^{\text{raw}} = \frac{Y_+ - Y_-}{Y_+ + Y_-}, \quad \text{Combine both HRS for best statistics!}$$

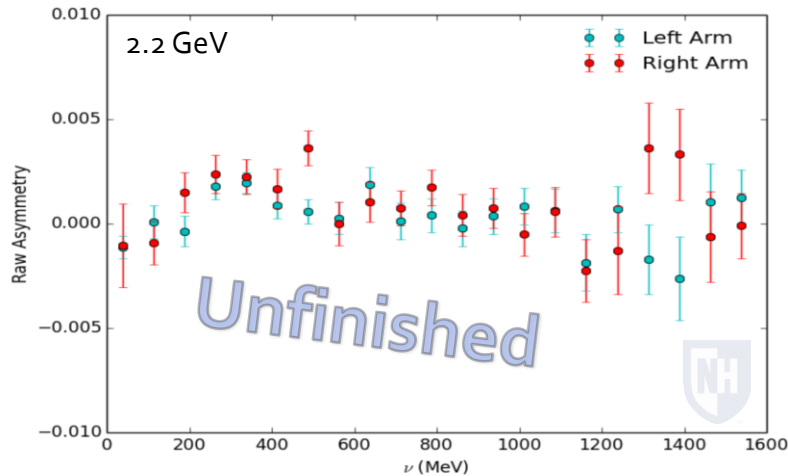
$$A^{\text{exp}} = \frac{1}{f \cdot P_t \cdot P_b} A^{\text{raw}}$$

dilution factor

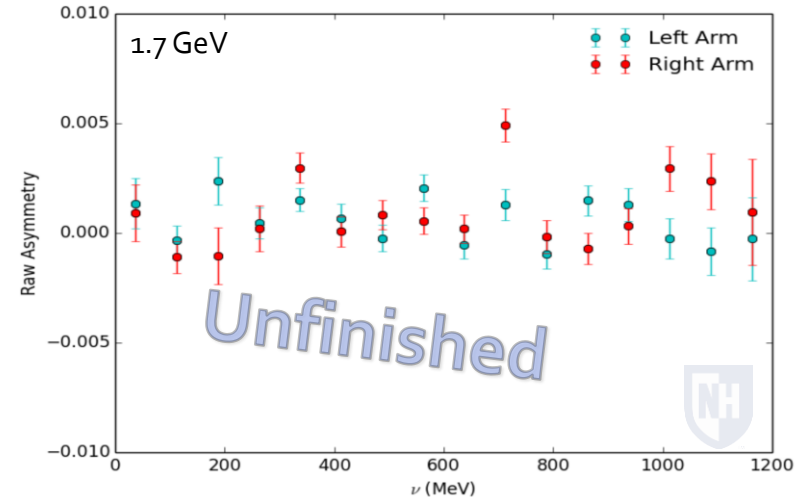
beam/target pol



2.5T Proton Asymmetries

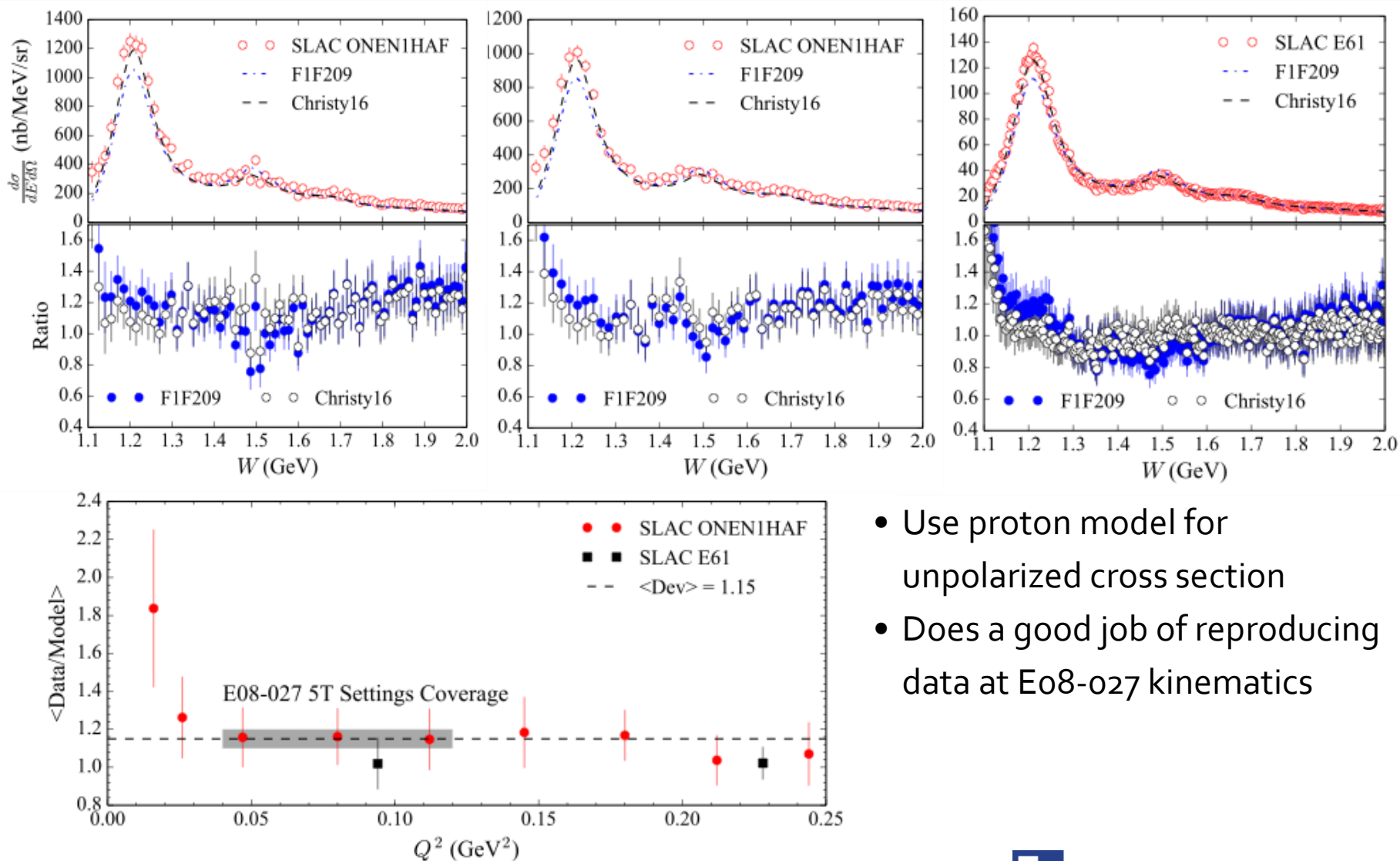


1.1 GeV data has large systematics that make it hard to work with.



2.5T Asymmetries show systematic shift in many places, working on resolving before data is usable

Model Cross Section



- Use proton model for unpolarized cross section
- Does a good job of reproducing data at E08-027 kinematics



Acceptance Study

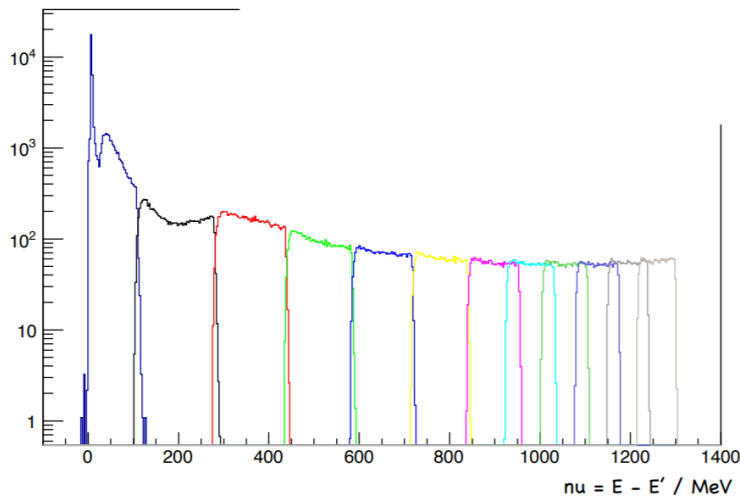
Acceptance study is on-going, we hope to be able to use data instead of model cross sections soon

$$\sigma_0 = \frac{P_S N}{\frac{Q}{e}(\rho\Delta Z)T_L\epsilon_{\text{det}}} \frac{1}{\Delta\Omega\Delta E' A}$$

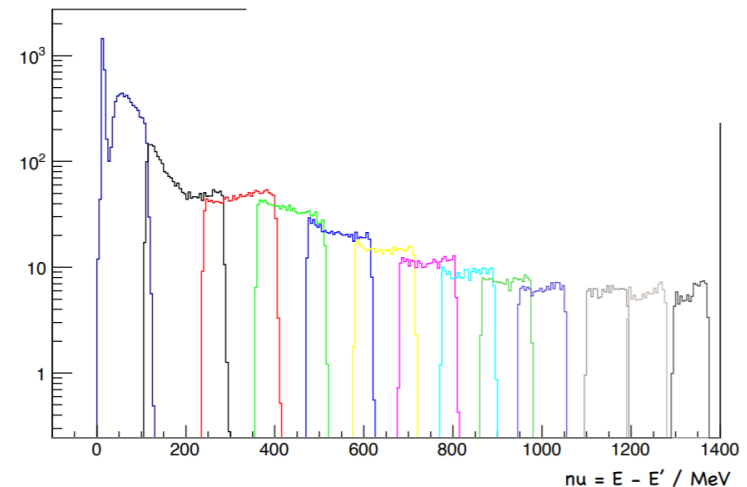
initial angle and momentum coverage in simulation \nearrow \nwarrow ratio of accepted events and total events

Continuity on longitudinal spectrum for most recent acceptance is good, not as much on transverse

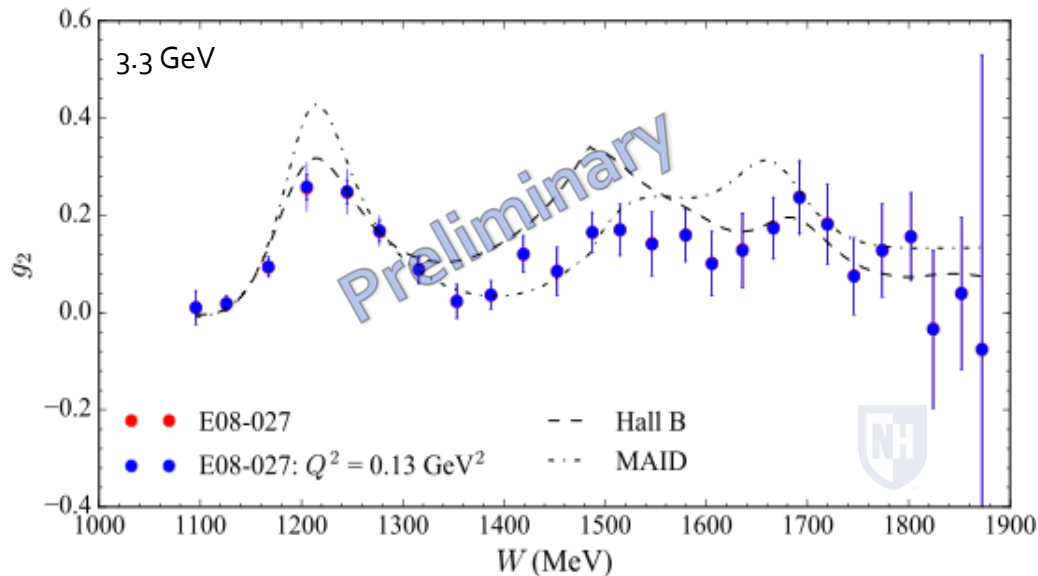
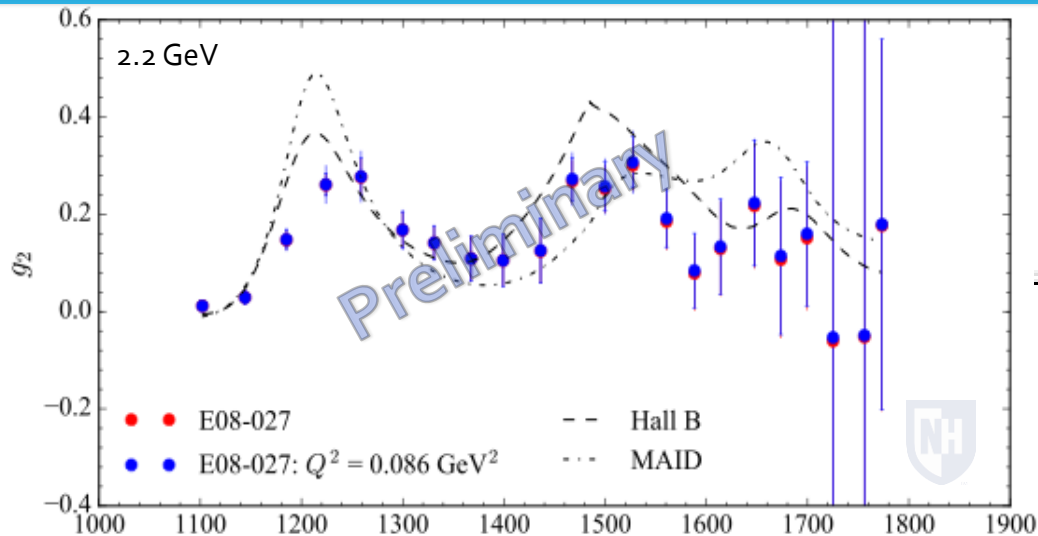
Longitudinal



Transverse



Extracting the Spin Structure Functions : g_2



Model driven procedure for unmeasured part

$$g_2(x, Q^2) = \frac{K_1 y}{2} \left[\Delta\sigma_{\perp} \left(K_2 + \tan\frac{\theta}{2} \right) \right] - \frac{g_1(x, Q^2) y}{2}$$

$$K_1 = \frac{M Q^2}{4\alpha} \frac{y}{(1-y)(2-y)}$$

$$K_2 = \frac{1 + (1-y)\cos\theta}{(1-y)\sin\theta}$$

Adjusting to a constant Q^2

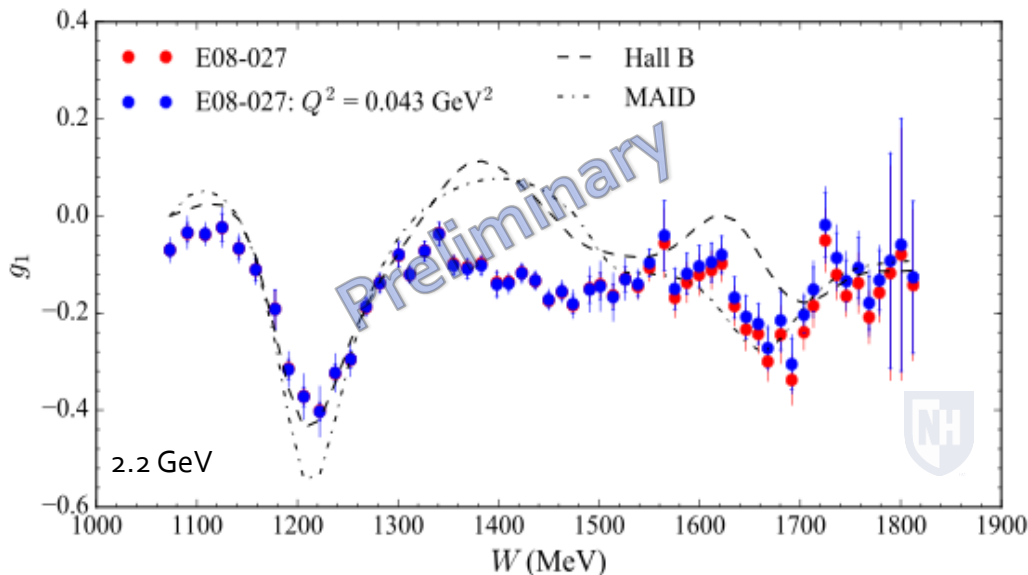
$$\delta_{\text{evolve}} = g_{1,2}^{\text{mod}}(x_{\text{data}}, Q_{\text{data}}^2) - g_{1,2}^{\text{mod}}(x_{\text{const}}, Q_{\text{const}}^2),$$

$$x_{\text{const}} = Q_{\text{const}}^2 / (W^2 - M^2 + Q_{\text{const}}^2),$$

Small effect at the transverse settings



Extracting the Spin Structure Functions: g_1

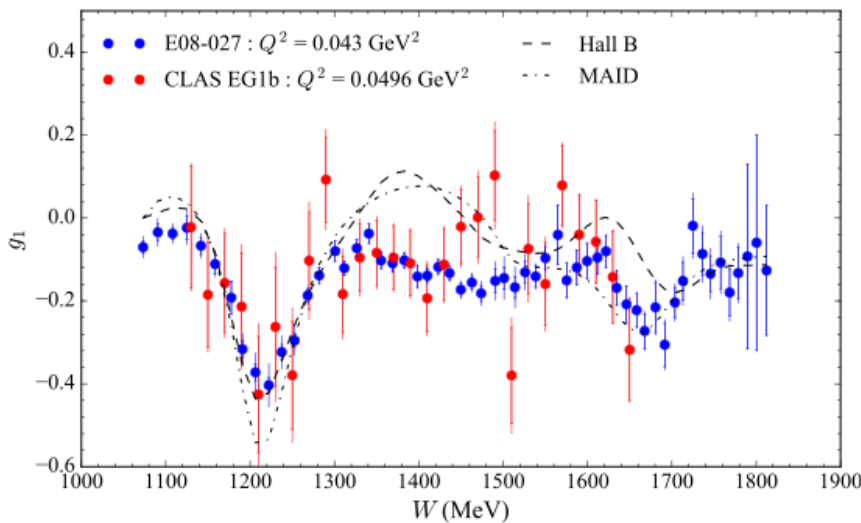


Model driven procedure for unmeasured part

$$g_1(x, Q^2) = K_1 \left[\Delta\sigma_{\parallel} \left(1 + \frac{1}{K_2} \tan \frac{\theta}{2} \right) \right] + \frac{2g_2(x, Q^2)}{K_2 y} \tan \frac{\theta}{2}$$

$$K_1 = \frac{MQ^2}{4\alpha} \frac{y}{(1-y)(2-y)}$$

$$K_2 = \frac{1 + (1-y)\cos\theta}{(1-y)\sin\theta}$$



- E08-027 data is consistent with previously published data from CLAS
- But with much better statistics!!

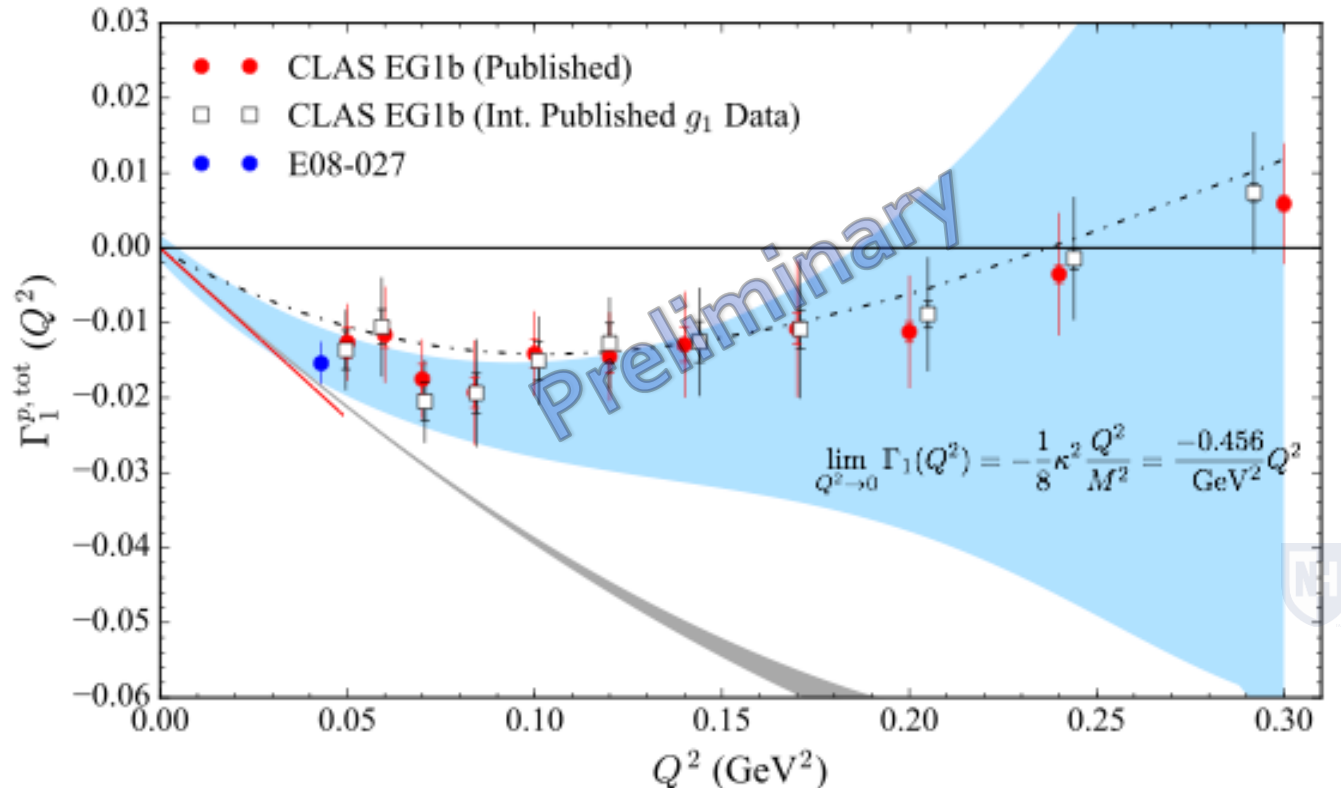
<http://clas.sinp.msu.ru/cgi-bin/jlab/db.cgi>



First Moment of $g_1(x, Q^2)$

$$\Gamma_1(Q^2) = \int_0^{x_{th}} g_1(x, Q^2) dx$$

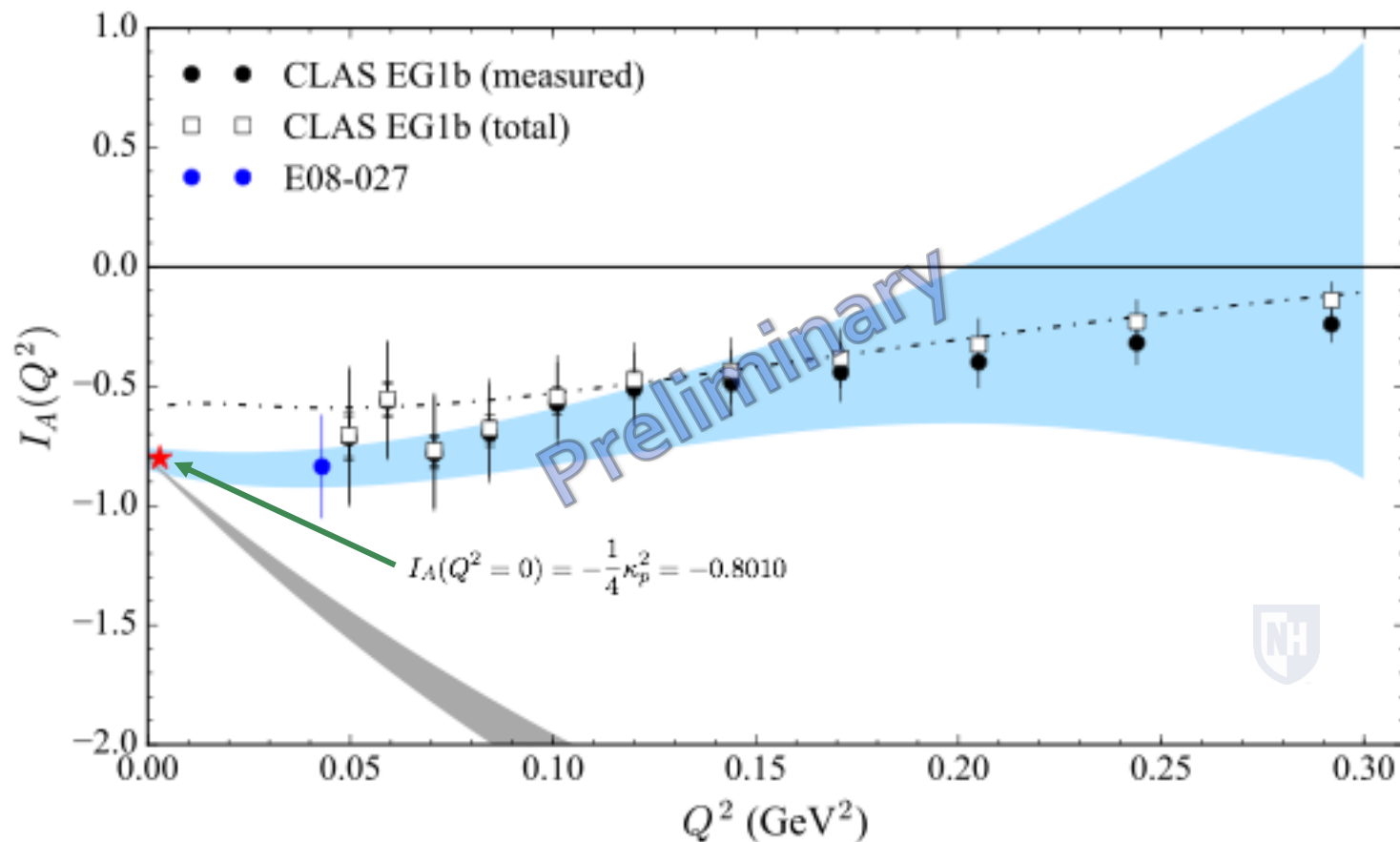
Moments provide a useful quantity that can be related back to theory predictions!



Setting	$\Gamma_1^{\text{meas.}}$	$\Gamma_1^{\text{low } x}$	$\Gamma_1^{\text{tot.}}$	$\delta_{\text{stat}}^{\text{tot.}}$	$\delta_{\text{sys}}^{\text{tot.}}$
2254 5T Long.	-0.01541	0.0003	-0.01541	0.0006	0.0028

Extended GDH Sum

$$I_A(Q^2) = \frac{2M^2}{Q^2} \int_0^{x_{th}} \left(g_1(x, Q^2) - \frac{4M^2}{Q^2} x^2 g_2(x, Q^2) \right) dx$$

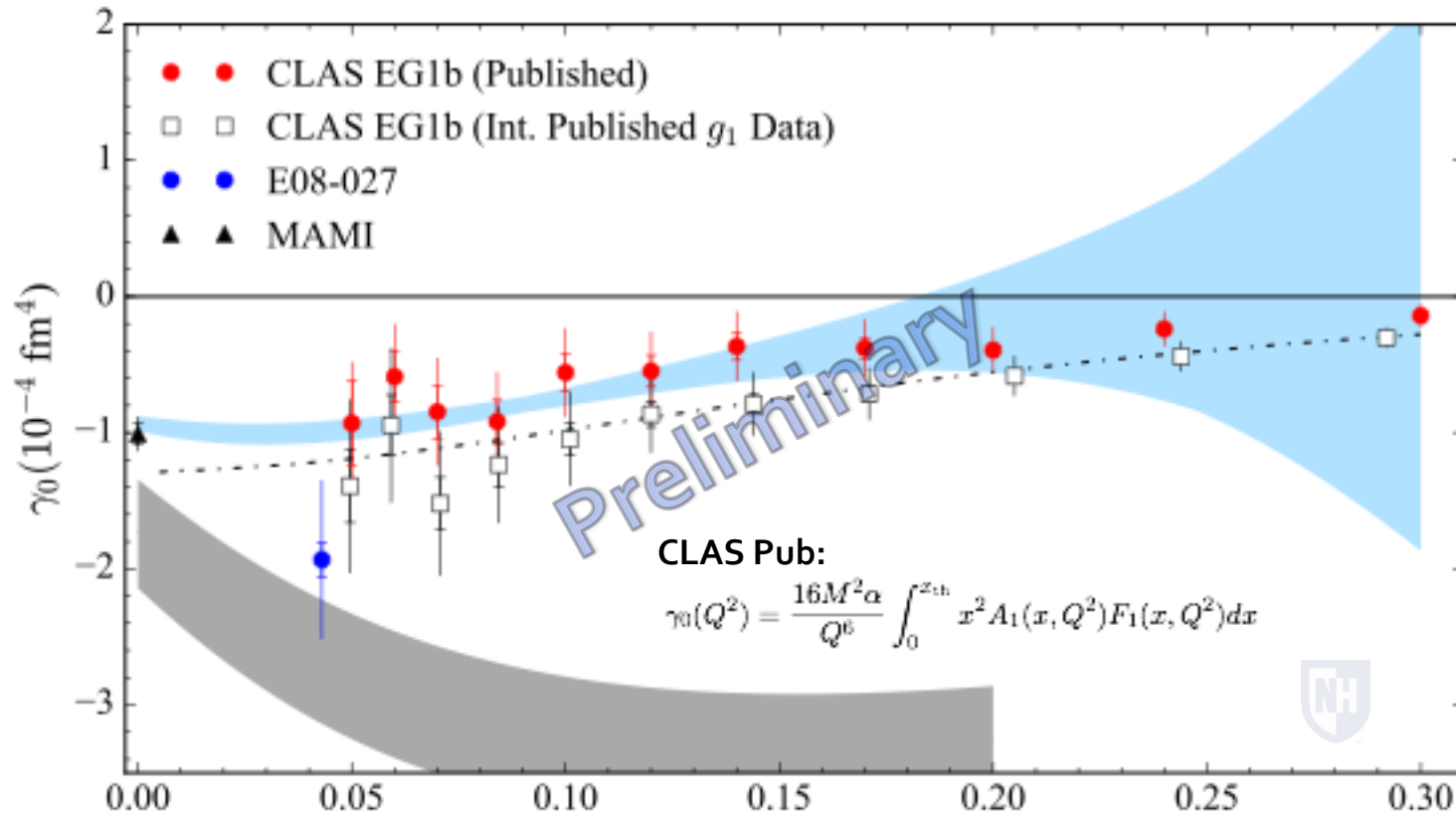


Setting	$I_A^{meas.}$	$I_A^{low z}$	$I_A^{tot.}$	$\delta_{stat}^{tot.}$	$\delta_{sys}^{tot.}$
2254 5T Long.	-0.83669	0.0119	-0.83669	0.0240	0.2155



Forward Spin Polarizability

$$\gamma_0(Q^2) = \frac{16\alpha M^2}{Q^6} \int_0^{x_{th}} x^2 \left(g_1(x, Q^2) - \frac{4M^2}{Q^2} x^2 g_2(x, Q^2) \right) dx$$



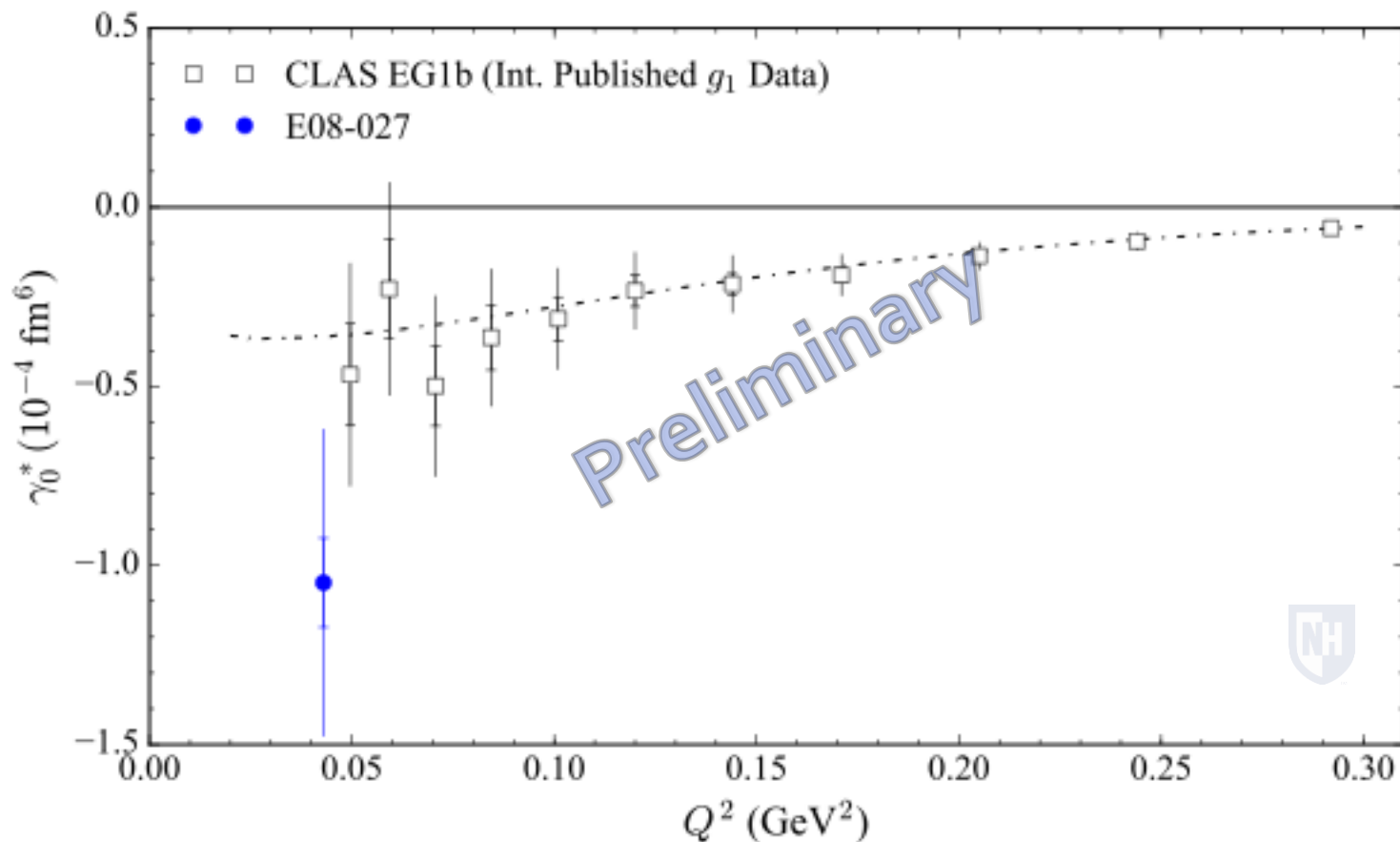
Setting	$\gamma_0^{mens.}$	Q^2 (GeV ²) $\gamma_0^{low x}$	$\gamma_0^{tot.}$	$\delta_{stat}^{tot.}$	$\delta_{sys}^{stat.}$
2254 5T Long.	-1.9352	-0.0007	-1.9352	0.1202	0.5663

$$\gamma_0(Q^2 = 0) = [-1.01 \pm 0.08 \text{ (stat)} \pm 0.10 \text{ (sys)}] \cdot 10^{-4} \text{ fm}^4$$



Higher Order Polarizability

$$\gamma_0^*(Q^2) = \frac{64\alpha M^4}{Q^{10}} \int_0^{x_{th}} x^4 \left(g_1(x, Q^2) - \frac{4M^2}{Q^2} x^2 g_2(x, Q^2) \right) dx$$



Setting	γ_0^* mens.	γ_0^* low x	γ_0^* tot.	$\delta_{\text{stat}}^{\text{tot.}}$	$\delta_{\text{sys}}^{\text{tot.}}$
2254 5T Long.	-1.0501	-2×10^{-5}	-1.0501	0.1203	0.4072



Conclusions

- Experimental measurements of proton structure are key to understanding the proton!
- The g_{2p} experiment was a precision measurement of proton g_2 in low Q^2 region **for the first time!**
- Longitudinal data agrees with previous measurements.
- 5T Data Analysis is complete, 2.5T Analysis is ongoing. We hope to have publishable data at both the 2.2 GeV and 1.7 GeV energy levels at this target field.
- Acceptance study is on-going, to allow us to use experimental unpolarized cross sections.
- Packing fraction/dilution analysis also complete for 5T data, but may need more investigation for 2.5T data.



Acknowledgements

g2p Analysis Team

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