

# Old & New in Strangeness Nucl. Phys.

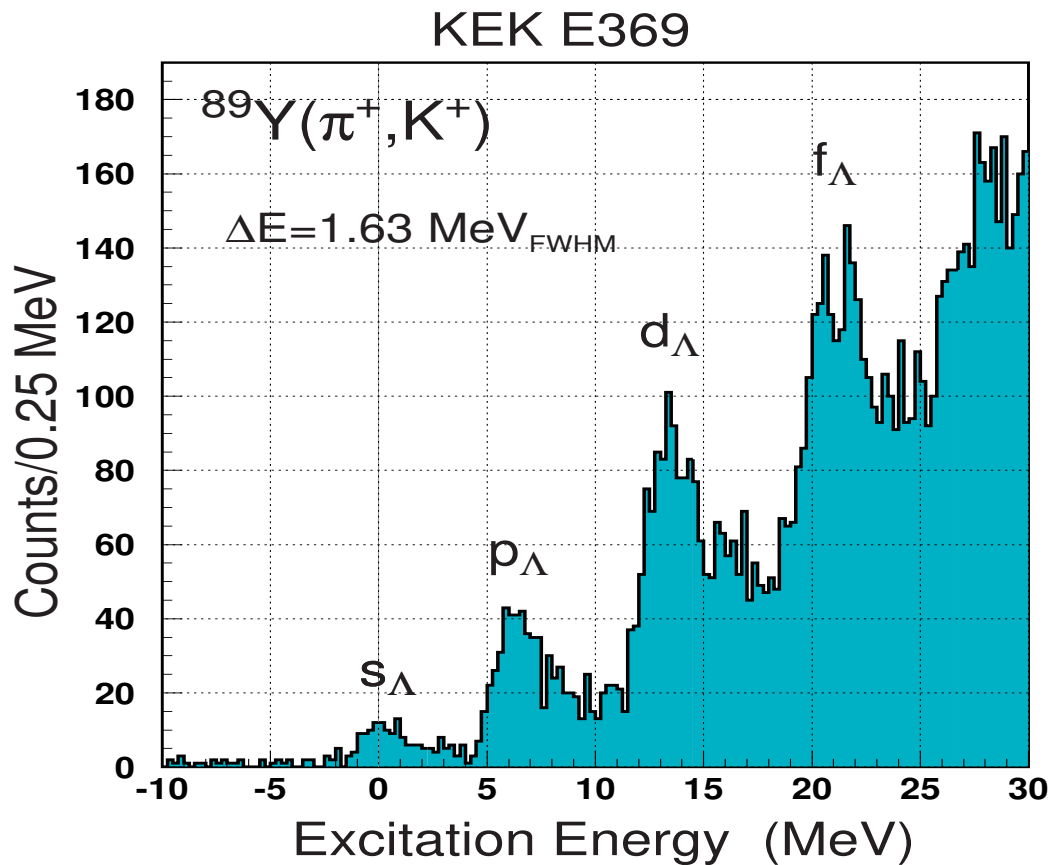
(my first 50 years in SNP)

Avraham Gal, Hebrew University, Jerusalem

- Dynamics of  $\Lambda$  hypernuclei ( ${}^A_{\Lambda}Z$ )
  - (i) s-shell few-body (ii) p-shell & beyond
- $\Lambda\Lambda$  hypernuclei: onset of  $\Lambda\Lambda$  binding?
- Hyperons ( $\Lambda, \Sigma, \Xi$ ) in nuclear matter & beyond
  - (i) neutron stars: hyperon puzzle
  - (ii) competition with  $\bar{K}$  condensation?
  - (iii)  $\Lambda^*(1405), \Sigma^*(1385)$ ?  $K^-$  nuclear clusters
- NPA Topical Issues: 881 (2012) & 954 (2016)
- Review: A.Gal E.V.Hungerford D.J.Millener  
Rev. Mod. Phys. 88 (2016) 035004

# $\Lambda$ hypernuclear dynamics

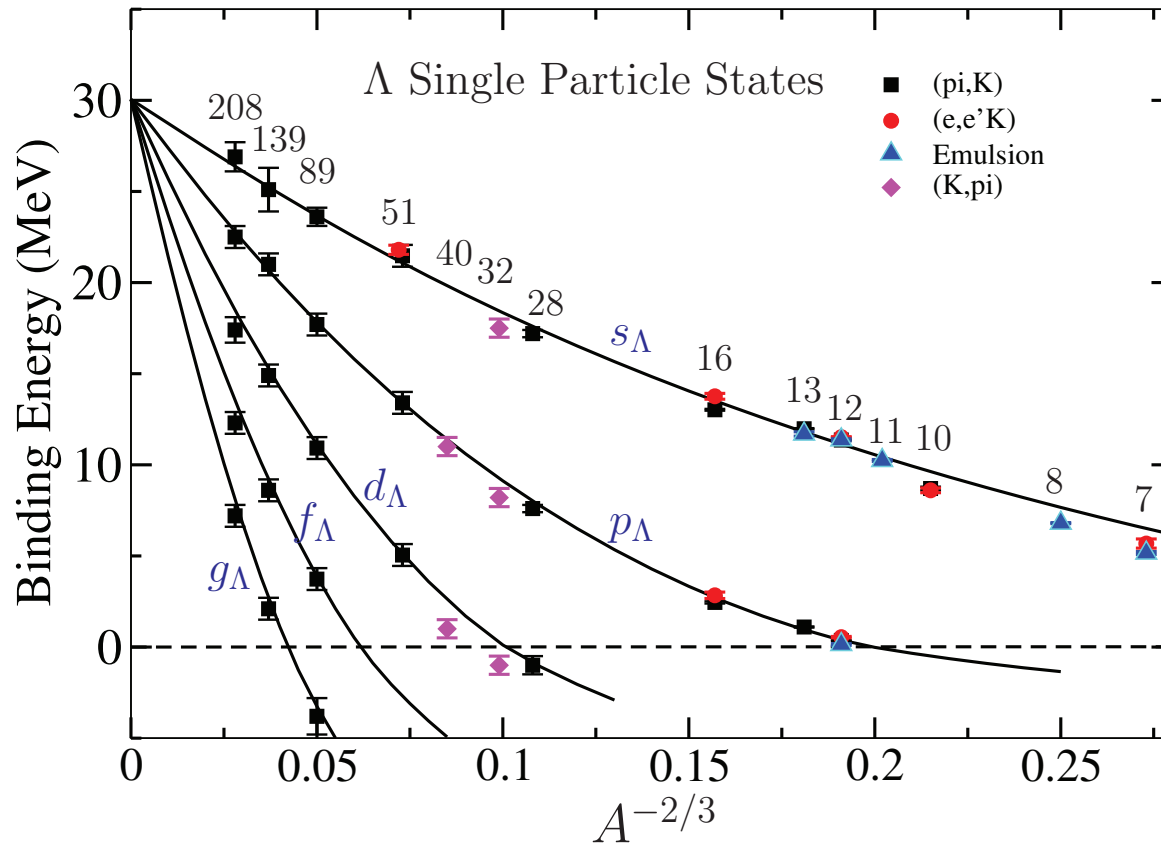
# Observation of $\Lambda$ single-particle states



Hotchi et al., PRC 64 (2001) 044302  $B_\Lambda = 23.1(1) \rightarrow 23.6(5) \text{ MeV}$

Motoba-Lanskoy-Millener-Yamamoto, NPA 804 (2008) 99:  
negligible  $\Lambda$  spin-orbit splittings, 0.2 MeV for  $1f_\Lambda$

Update: Millener, Dover, Gal PRC 38, 2700 (1988)



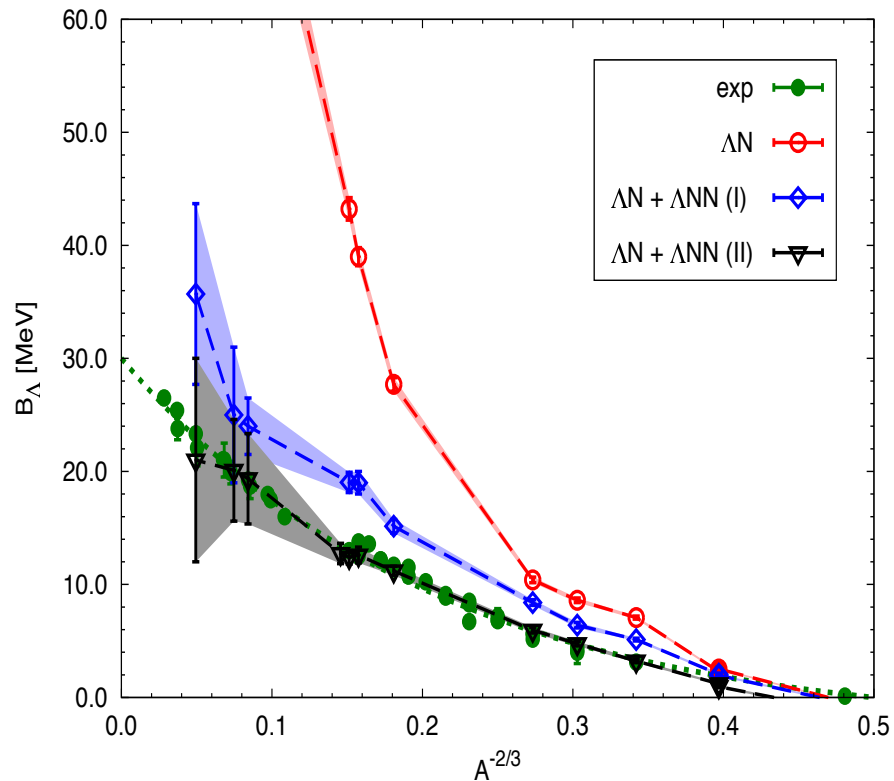
Woods-Saxon  $V = 30.05$  MeV,  $r = 1.165$  fm,  $a = 0.6$  fm

$V(\text{MeV}) = 27 \pm 3$  (1964)  $27.2 \pm 1.3$  (1965)  $27.8 \pm 0.3$  (1988)

from  $\pi^-$  decays of heavy spallation hypernuclei to  $(\pi^+, K^+)$

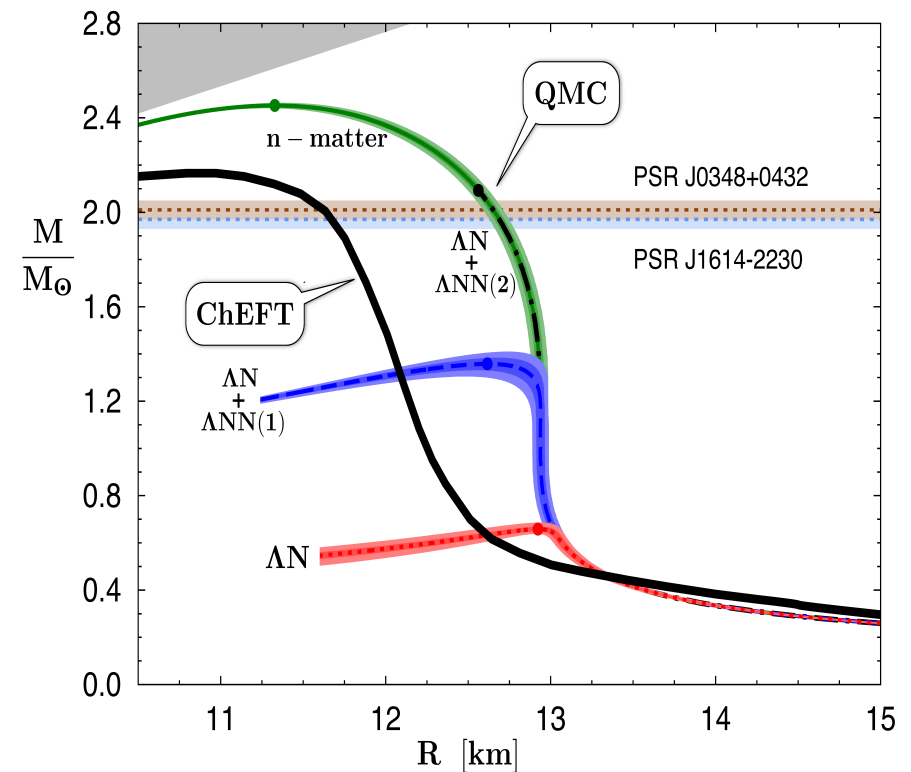
Skyrme-Hartree-Fock studies suggest  $\Lambda NN$  repulsion.

# Hyperon puzzle: QMC calculations



Lonardoni et al, PRC 89 (2014) 014314

$\Lambda NN$  effect on  $B_\Lambda$ (g.s.)



PRL 114 (2015) 092301

$\Lambda NN$  effect on neutron stars

- $\Lambda N$  overbinds, adding  $\Lambda NN$  stiffens EOS of neutron stars.
- YY add  $0.3M_\odot$  to  $M_{\max}$  (Rijken-Schulze 2016).
- Overbinding has been a problem in s-shell hypernuclei.

# The lightest, s-shell, $\Lambda$ hypernuclei

${}^A_{\Lambda}Z$	$T$	$J_{\text{g.s.}}^{\pi}$	$B_{\Lambda}$ (MeV)	$J_{\text{exc.}}^{\pi}$	$E_x$ (MeV)
${}^3_{\Lambda}\text{H}$	0	$1/2^+$	0.13(5)		
${}^4_{\Lambda}\text{H}-{}^4_{\Lambda}\text{He}$	1/2	$0^+$	<b>2.16(8)</b> – <b>2.39(3)</b>	$1^+$	<b>1.09(2)</b> – <b>1.406(3)</b>
${}^5_{\Lambda}\text{He}$	0	$1/2^+$	3.12(2)		

- **No  $\Lambda\text{N}$  and no  $\Lambda\text{nn}$  bound state are expected.**
- **$\Delta B_{\Lambda}({}^4_{\Lambda}\text{He}-{}^4_{\Lambda}\text{H})=0.23(9)$  (g.s.)  $-0.083(94)$  (exc.) in MeV.**

## Recent $A = 3, 4$ few-body calculations

- **A. Nogga, NPA 914 (2013) 140**  
Faddeev & Faddeev-Yakubovsky (chiral LO & NLO).
- **E. Hiyama et al., PRC 89 (2014) 061302(R)**  
Jacobi-coordinates Gaussian basis (Nijmegen soft-core).
- **R. Wirth et al., PRL 113 (2014) 192502**  
**D. Gazda, A. Gal, (2016) PRL 116 122501 NPA 954 161**  
ab-initio Jacobi-NCSM (chiral LO).

# Overbinding problem in s-shell

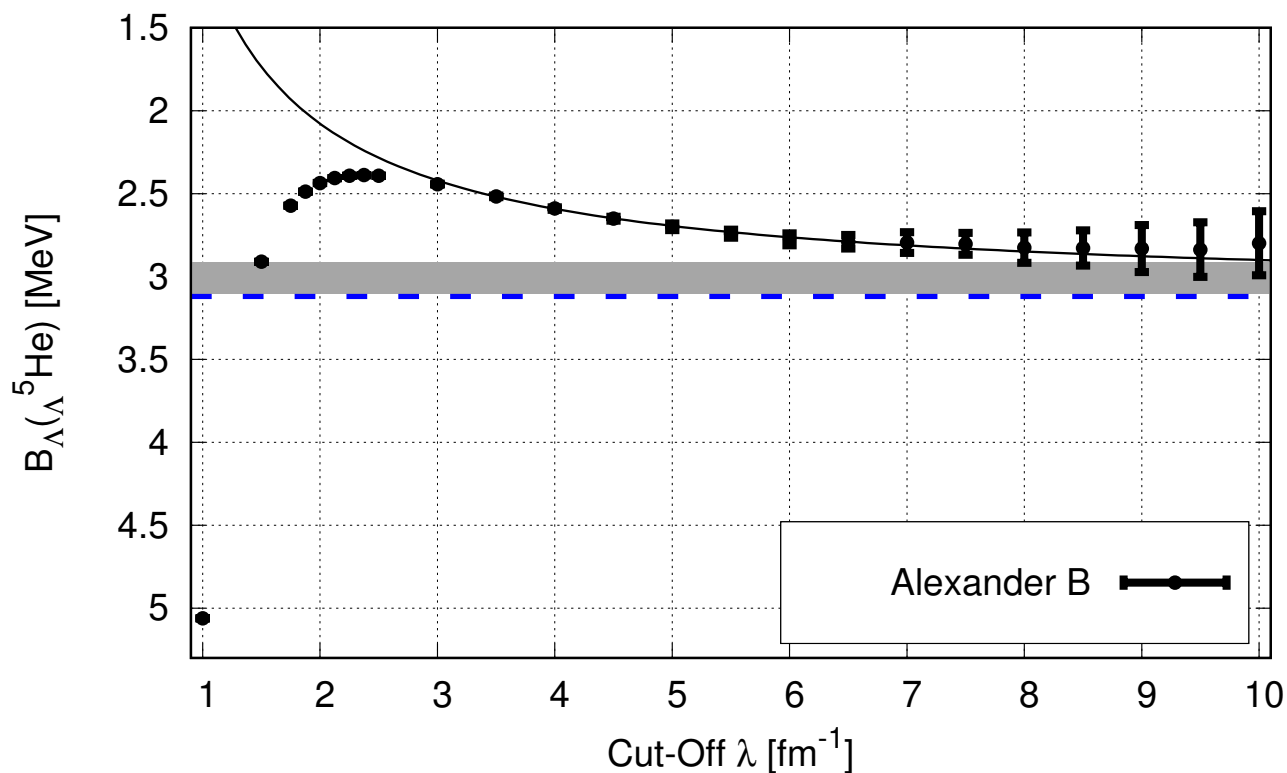
(MeV)	$B_\Lambda(^3_\Lambda\text{H})$	$B_\Lambda(^4_\Lambda\text{H}_{\text{g.s.}})$	$E_x(^4_\Lambda\text{H}_{\text{exc.}})$	$B_\Lambda(^5_\Lambda\text{He})$
Exp.	0.13(5)	2.16(8)	1.09(2)	3.12(2)
Dalitz	0.10	2.24	0.36	$\geq 5.16$
NSC97f(S)	0.18	2.16	1.53	2.10
AFDMC(I)	–	1.97(11)	–	5.1(1)
AFDMC(II)	–1.2(2)	1.07(8)	–	3.22(14)
LO $\chi$ EFT(600)	0.11(1)	2.31(3)	0.95(15)	5.82(2)
LO $\chi$ EFT(700)	–	2.13(3)	1.39(15)	4.43(2)

L. Contessi, N. Barnea, A. Gal, arXiv:1805.04302

Resolving the overbinding problem in  $\not\chi$ EFT

- Fit 2  $\Lambda N$  LECs to  $\Lambda N$  scattering lengths.
- Fit 3  $\Lambda NN$  LECs to the 3  $A=3,4$  levels.
- Calculate in SVM  $^5_\Lambda\text{He}$  binding.

# Contessi et al. Hyp. s-shell $\not\equiv$ EFT



$B_{\Lambda}(^5_{\Lambda}\text{He})$  vs. cut-off  $\lambda$  in LO  $\not\equiv$ EFT SVM calculations

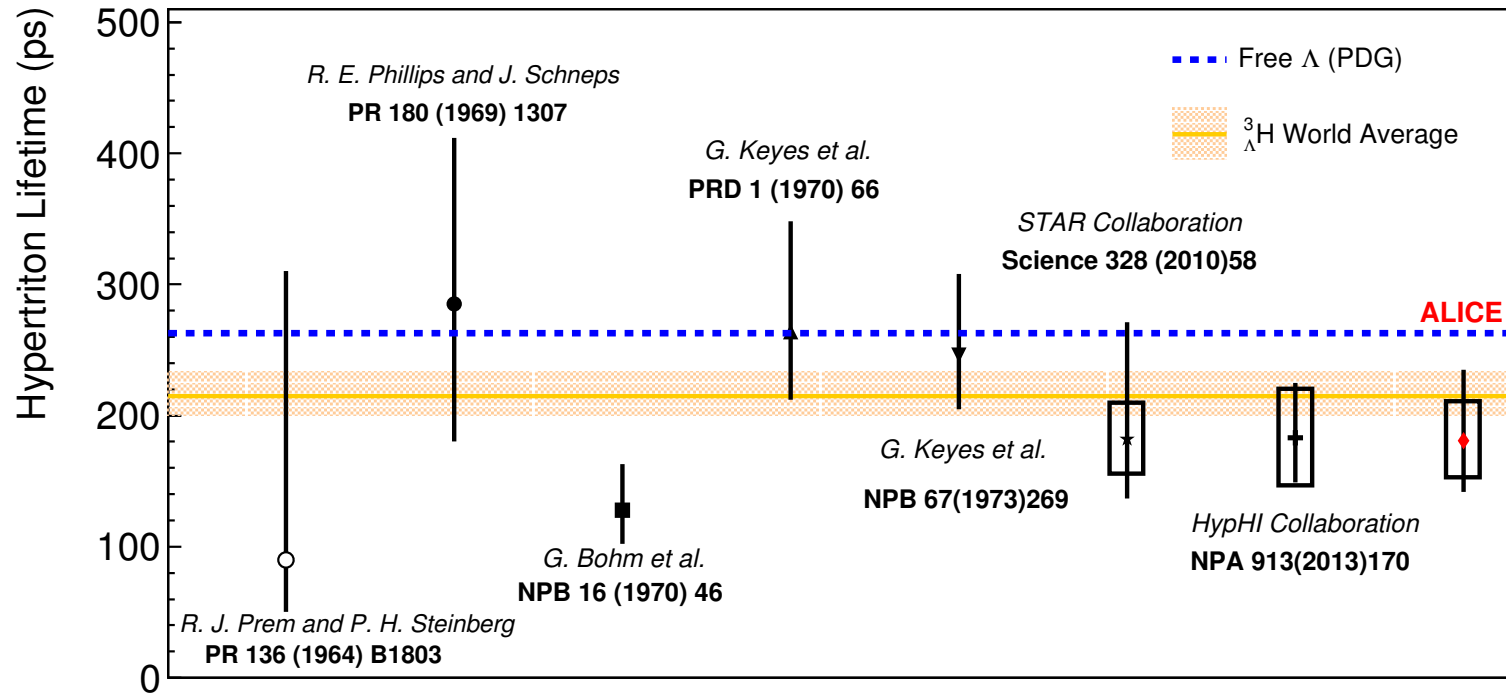
Solid line: a two-parameter fit  $a+b/\lambda$ ,  $\lambda \geq 4 \text{ fm}^{-1}$ .

Gray horizontal band:  $\lambda \rightarrow \infty$  extrapolation uncertainties.

Dashed horizontal line:  $B_{\Lambda}^{\text{exp.}}(^5_{\Lambda}\text{He}) = 3.12 \pm 0.02 \text{ MeV}$ .



# ${}^3_{\Lambda}\text{H}$ lifetime puzzle



The weakly-bound  ${}^3_{\Lambda}\text{H}$ ,  $B_{\Lambda}=0.13\pm 0.05$  MeV, expected to have lifetime within a few % of the free  $\Lambda$  lifetime. Recent heavy-ion  ${}^3_{\Lambda}\text{H}$  production experiments yield lifetimes shorter by  $\geq 30\%$ . ALICE, PLB 754 (2016) 360. STAR, PRC 97 (2018) 054909:  $\tau=142^{+24}_{-21}\pm 29$  ps ( $0.54^{+0.09}_{-0.08}\tau_{\Lambda}$ ).

# Brief review of lifetime calculations

Reference	Method	$R_3$	$\Gamma({}^3_{\Lambda}\text{H}, J=\frac{1}{2}, T=0)/\Gamma_{\Lambda}$
<b>Experiment</b>	<b>wo. av.</b>	<b><math>0.35 \pm 0.04</math></b>	<b><math>1.22 \pm 0.07</math></b>
Dalitz-Rayet (1966)	closure	–	$1.05 \pm 0.01$
Congleton (1992)	$\Lambda d$	$0.33 \pm 0.02$	1.12
Kamada et al (1998)	$\Lambda p n$ Fad.	0.379	1.03
Gal-Garcilazo (2018)	$\Lambda p n$ Fad.		$1.05 \pm 0.02$ <b>prelim.</b>

- $R_3 = \Gamma({}^3_{\Lambda}\text{H} \rightarrow \pi^- + {}^3\text{He}) / \Gamma({}^3_{\Lambda}\text{H} \rightarrow \pi^- + \text{all}) \Rightarrow J = \frac{1}{2}$ .
- Closure:  $\Gamma({}^3_{\Lambda}\text{H}, J=\frac{1}{2}, T=0) / \Gamma_{\Lambda} = 1 + 0.14\sqrt{B_{\Lambda}}$ .
- A bound, isomeric  ${}^3_{\Lambda}\text{H}(J=\frac{3}{2}, T=0)$  (unlikely) would decay **much slower** than a free  $\Lambda$ .
- A bound  ${}^3_{\Lambda}\text{H}(J=\frac{1}{2}, T=1)$ , analog of  $\Lambda_{nn}$ , would decay to  $\Lambda d$  or by  $\gamma(M1)$  to  ${}^3_{\Lambda}\text{H}(J=\frac{1}{2}, T=0)$ .

## ${}^4_{\Lambda}\text{H}$ & ${}^4_{\Lambda}\text{He}$ lifetimes

$$\Gamma({}^4_{\Lambda}\text{H})/\Gamma_{\Lambda} \approx \frac{3}{2} \times \left( \frac{2}{3} \times 0.7 + 1 \times 0.3 \right) + 0.25 = 1.40$$

$$\Gamma({}^4_{\Lambda}\text{He})/\Gamma_{\Lambda} \approx \frac{3}{2} \times \left( \frac{1}{3} \times 0.7 + 1 \times 0.3 \right) + 0.25 = 1.05$$

**Input:**  $\frac{3}{2}$  for nuclear structure,  $R_4=0.7$

$\frac{2}{3}$  &  $\frac{1}{3}$  for  $\pi^-$  or  $\pi^0$  and  ${}^4\text{He}$ ,  $\Gamma_{\text{n.m.}}/\Gamma_{\Lambda} \approx 0.25$

$\Rightarrow \tau({}^4_{\Lambda}\text{H}) \approx 190$  ps,  $\tau({}^4_{\Lambda}\text{He}) \approx 250$  ps

in rough agreement with measured lifetimes.

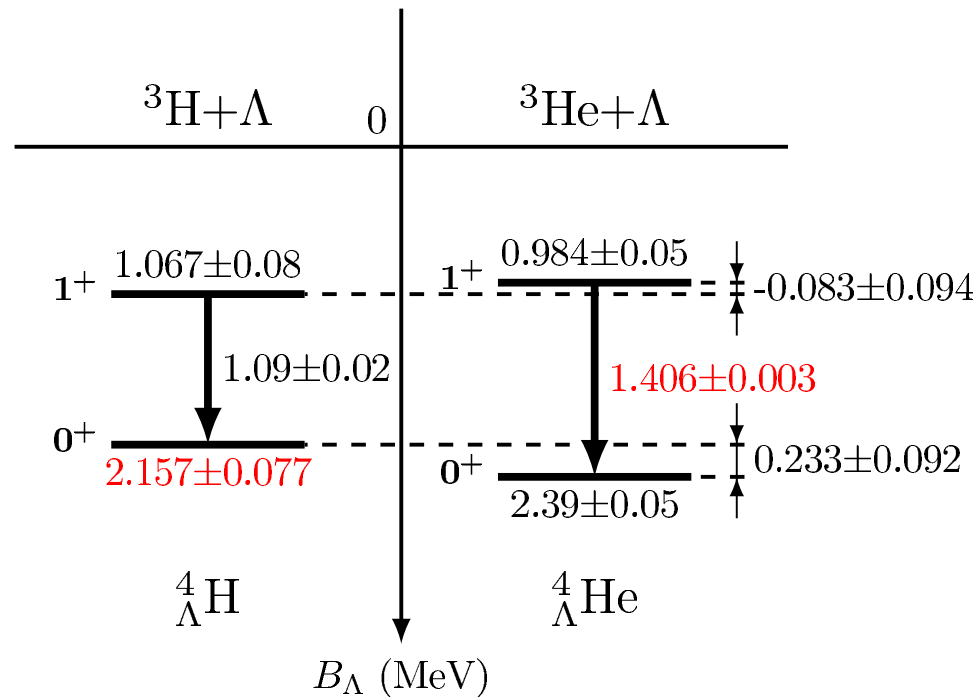
Looks like **Lifetime Puzzle** is limited to  ${}^3_{\Lambda}\text{H}$ .

- For  $A \geq 12$ ,  $\tau({}^A_{\Lambda}\text{Z}) \sim 200$  ps, from KEK and very recently from **HKS JLab E02-E017 NPA 973 (2018) 116**. Lifetime is due to  $\Lambda\text{N} \rightarrow \text{NN}$ .

# The ${}^4_{\Lambda}\text{H}-{}^4_{\Lambda}\text{He}$ complex & CSB since 2015

MAMI's A1,  ${}^4_{\Lambda}\text{H} \rightarrow {}^4\text{He} + \pi^-$ , PRL 114 (2015) 232501

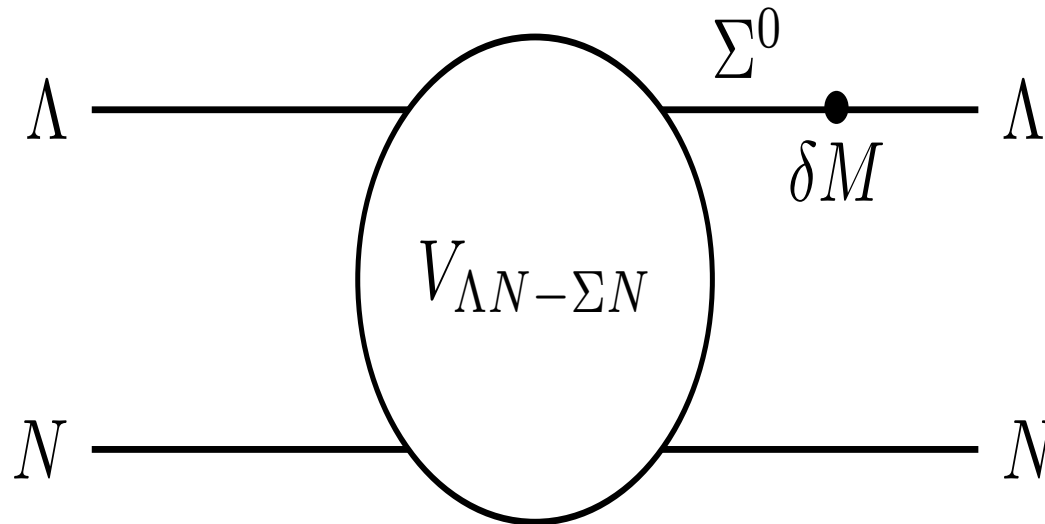
J-PARC's E13,  ${}^4\text{He}(K^-, \pi^- \gamma)$ , PRL 115 (2015) 222501



CSB due to  $\Lambda$ - $\Sigma^0$  mixing, strongly spin dependent, dominantly in  $0^+_{\text{g.s.}}$ , large w.r.t.  $\approx -70$  keV in  ${}^3\text{H}-{}^3\text{He}$ .

Re-measure  ${}^4_{\Lambda}\text{He}_{\text{g.s.}}$  (**E13**  $\rightarrow$  **E63**).

## Relating $\Lambda$ - $\Sigma^0$ CSB mixing to $\Lambda\Sigma$ SI coupling



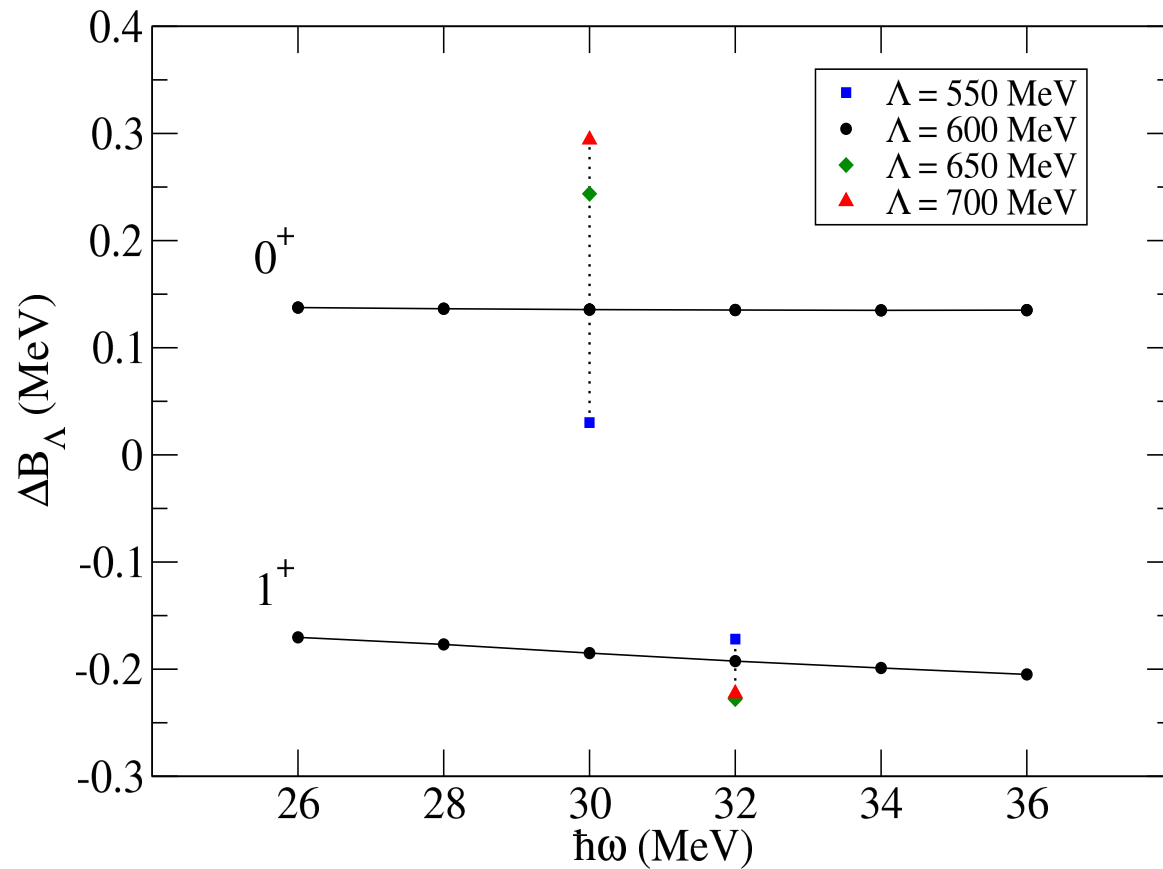
**Dalitz-von Hippel (1964):** “applies to any isovector meson exchange,  $\pi, \rho\dots$ ” & also to  $\chi$ EFT contact interactions.

$$\langle N\Lambda | V_{\Lambda N}^{\text{CSB}} | N\Lambda \rangle = -0.0297 \tau_{Nz} \frac{1}{\sqrt{3}} \langle N\Sigma | V^{\text{SI}} | N\Lambda \rangle.$$

Applied systematically by A. Gal, PLB 744 (2015) 352

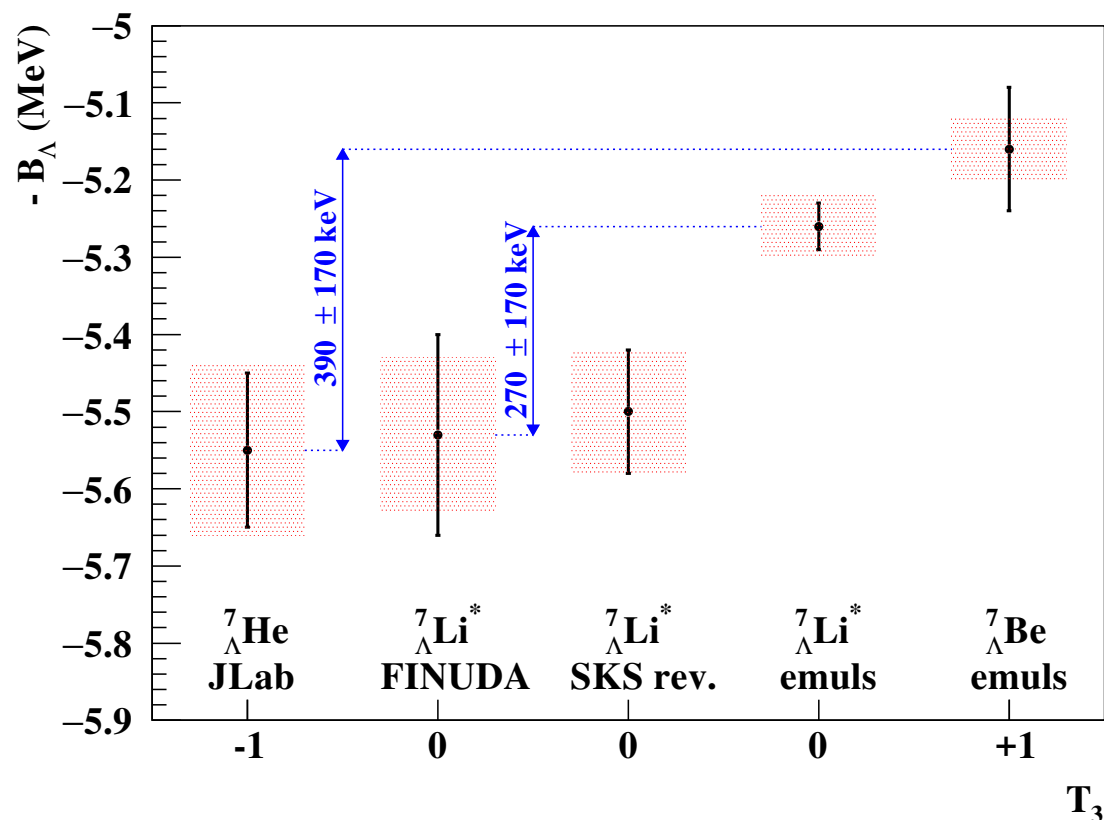
A=4: D.Gazda A.Gal (2016), PRL 116 122501; NPA 954 161.

Latest summary: J. Phys. Conf. Series 966 (2018) 012006.



NCSM HO  $\hbar\omega$  dependence of  $\Delta B_\Lambda({}^4_\Lambda\text{He}-{}^4_\Lambda\text{H})$  for  $0^+$  &  $1^+$ .  
 Note  $\pm$  sign pattern resulting from  ${}^1\text{S}_0$   $\Lambda$ - $\Sigma$  contact term dominance at LO [see OPE discussion NPA 954 (2016) 161].  
 $\Lambda=600$  MeV:  $\Delta E_\gamma = \Delta(\Delta B_\Lambda) = 0.33 \pm 0.03$  MeV compared to a measured  $\Delta E_\gamma = 0.32 \pm 0.02$  MeV.

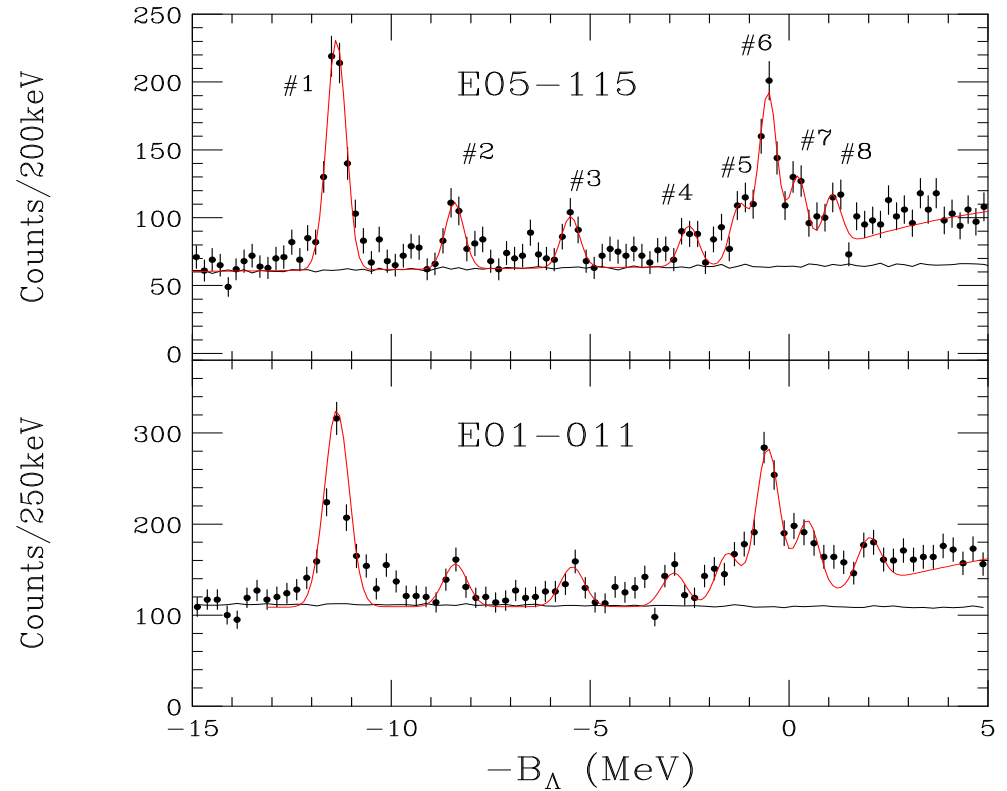
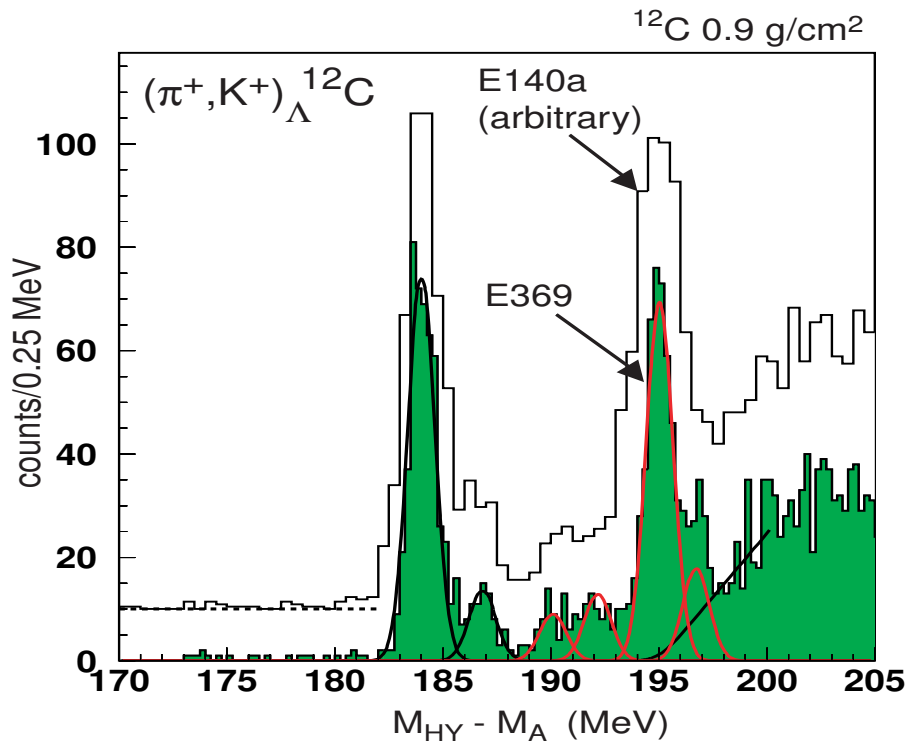
# CSB in p-shell hypernuclei



E. Botta, T. Bressani, A. Feliciello, NPA 960 (2017) 165-179

CSB appears to be much weaker in the  $A=7$  isotriplet than in the  $A=4$  isodoublet **provided** counter experiments are not compared directly with old emulsion results.

# Room for hypernuclear spectroscopy



H. Hotchi et al., PRC 64 (2001) 044302

L. Tang et al., PRC 90 (2014) 034320

$1s_{\Lambda}$ - $1p_{\Lambda}$  intermediate structure

Jlab:  $^{12}C(e, e'K^+)_{\Lambda}^{12}B$  (HKS)

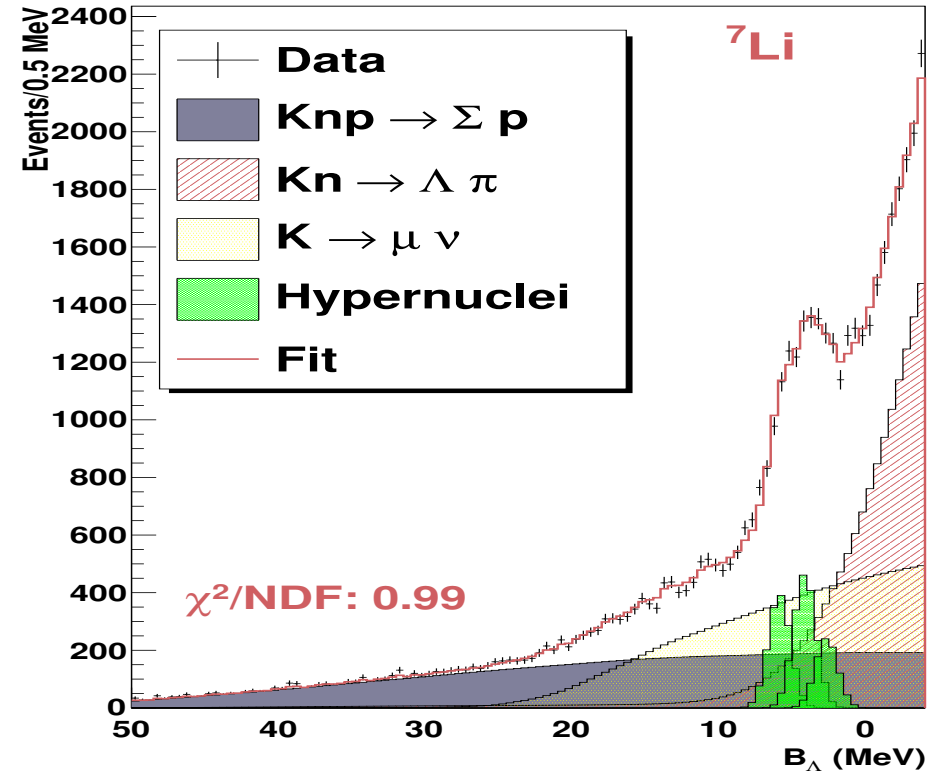
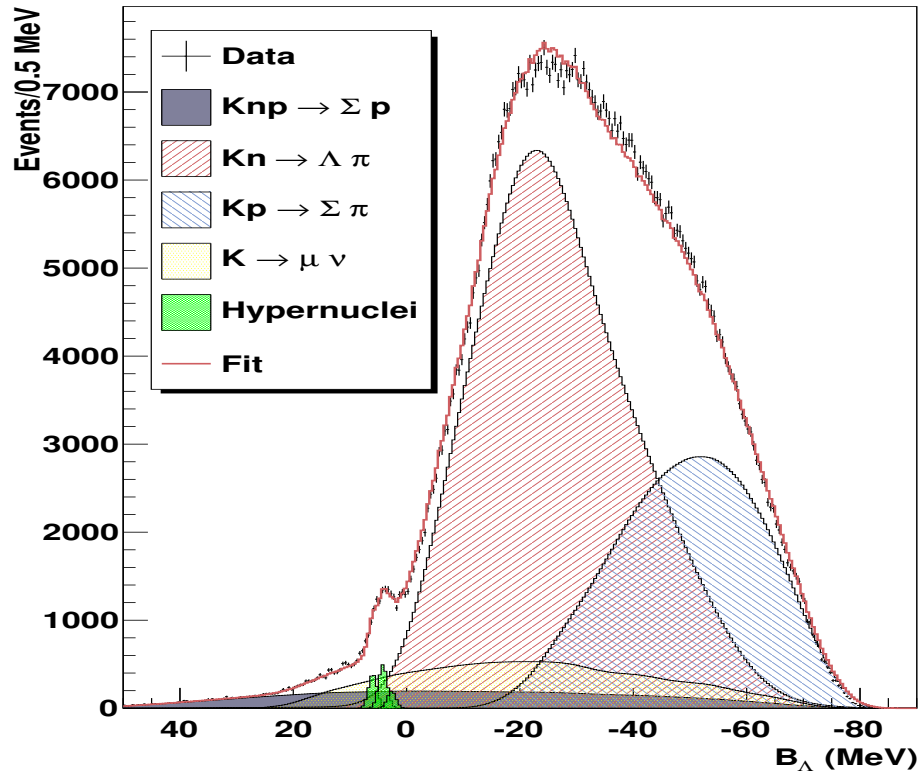
Spin-nonflip exc.,  $\Delta S=0$

Spin-flip excitations,  $\Delta S=1$

energy resolution 1.6 MeV  $\rightarrow$  0.6 MeV in Hall-C E05-115



# Hypernuclear production in $(K_{\text{stop}}^-, \pi^-)$ , PLB 698 (2011) 219 & 226



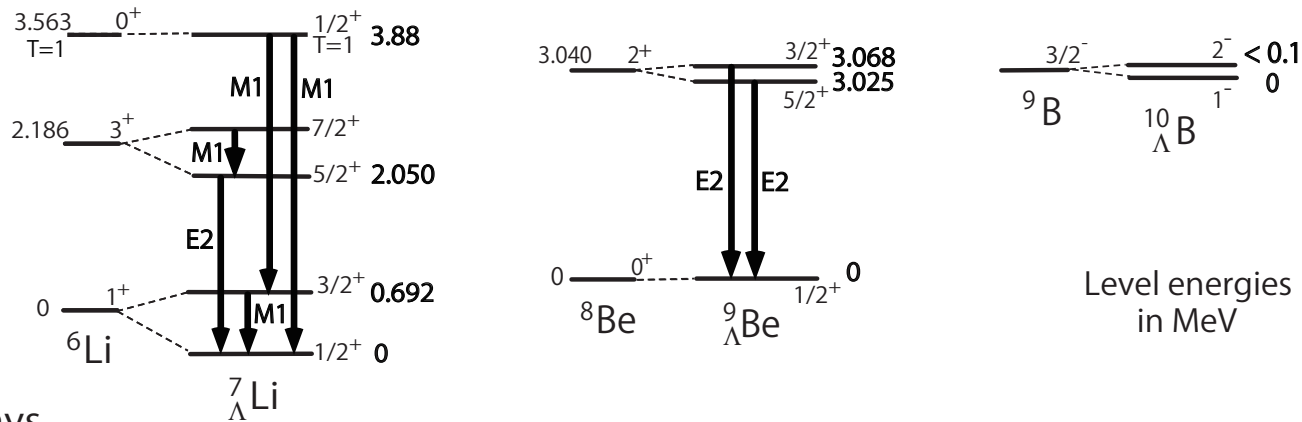
Production spectrum on  ${}^7\text{Li}$

FINUDA, DAΦNE, Frascati

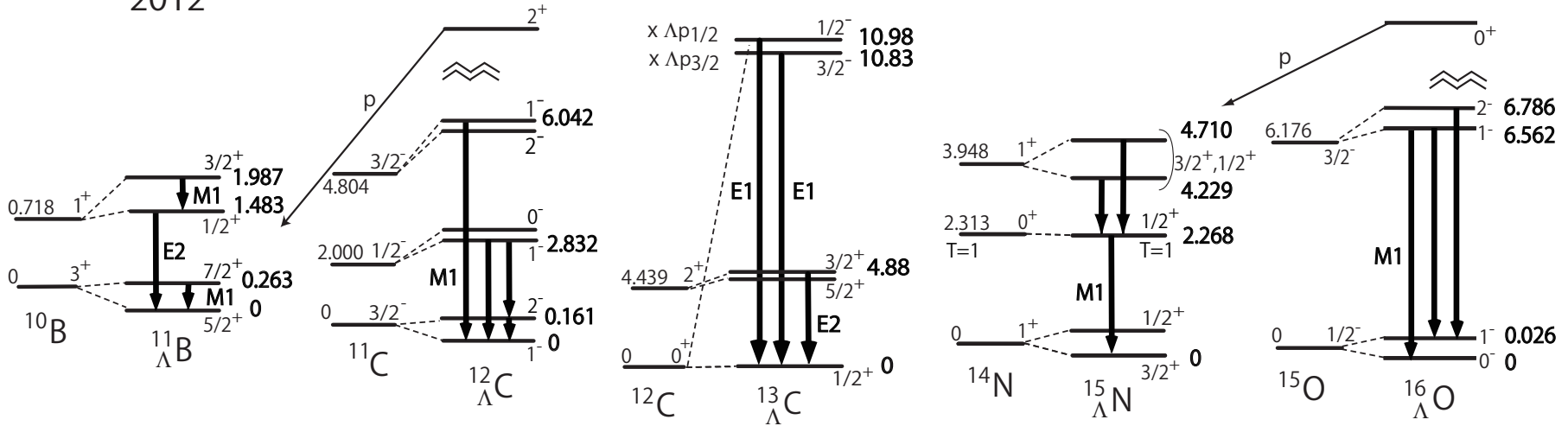
$A=7-16$  data also indicate DEEP  $K^-$  nuclear potential.

Three  ${}^7_\Lambda\text{Li}$  levels,  $\delta B_\Lambda=0.4$  MeV

Formation rate  $1 \cdot 10^{-3}/K_{\text{stop}}^-$



Hypernuclear  $\gamma$  rays  
2012



## p-shell spectra from $(K^-, \pi^- \gamma)$ ( $+^{19}_{\Lambda}\text{F}$ @HYP2018)

H. Tamura et al., Nucl. Phys. A 835 (2010) 3, updated at HYP12

$\Lambda$  spin-orbit splitting (keV): 150 in  $^{13}_{\Lambda}\text{C}$  & related 43 in  $^9_{\Lambda}\text{Be}$

# p-shell $\Lambda$ Hypernuclei

$$V_{\Lambda N} = V_0(r) + V_\sigma(r) s_N \cdot s_\Lambda + V_{LS}(r) l_{N\Lambda} \cdot (s_\Lambda + s_N) + V_{ALS}(r) l_{N\Lambda} \cdot (s_\Lambda - s_N) + V_T(r) S_{12}$$

For  $p_N s_Y$  :  $V_{\Lambda N} = \bar{V} + \Delta s_N \cdot s_\Lambda + S_\Lambda l_N \cdot s_\Lambda + S_N l_N \cdot s_N + T S_{12}$

R.H Dalitz, A. Gal, Ann. Phys. 116 (1978) 167

D.J. Millener, A. Gal, C.B. Dover, R.H. Dalitz, PRC 31 (1985) 499

$N\Lambda-N\Lambda$	$\bar{V}$	$\Delta$	$S_\Lambda$	$S_N$	$T$	(MeV)
$A = 7 - 9$	(-1.32)	0.430	-0.015	-0.390	0.030	fit
$A = 11 - 16$	(-1.32)	0.330	-0.015	-0.350	0.024	fit
$N\Lambda-N\Sigma$	1.45	3.04	-0.085	-0.085	0.157	input

D.J. Millener, Nucl. Phys. A 804 (2008) 84

# Doublet spacings in $p$ -shell hypernuclei (in keV)

D.J. Millener, NPA 881 (2012) 298

	$J_u^\pi$	$J_l^\pi$	$\Lambda\Sigma$	$\Delta$	$S_\Lambda$	$S_N$	$T$	$\Delta E^{\text{th}}$	$\Delta E^{\text{exp}}$
${}^7_\Lambda\text{Li}$	$3/2^+$	$1/2^+$	72	628	-1	-4	-9	693	692
${}^7_\Lambda\text{Li}$	$7/2^+$	$5/2^+$	74	557	-32	-8	-71	494	471
${}^8_\Lambda\text{Li}$	$2^-$	$1^-$	151	396	-14	-16	-24	450	(442)
${}^9_\Lambda\text{Be}$	$3/2^+$	$5/2^+$	-8	-14	37	0	28	44	43
${}^{11}_\Lambda\text{B}$	$7/2^+$	$5/2^+$	56	339	-37	-10	-80	267	264
${}^{11}_\Lambda\text{B}$	$3/2^+$	$1/2^+$	61	424	-3	-44	-10	475	505
${}^{12}_\Lambda\text{C}$	$2^-$	$1^-$	61	175	-22	-13	-42	153	161
${}^{15}_\Lambda\text{N}$	$3/2_2^+$	$1/2_2^+$	65	451	-2	-16	-10	507	481
${}^{16}_\Lambda\text{O}$	$1^-$	$0^-$	-33	-123	-20	1	188	23	26
${}^{16}_\Lambda\text{O}$	$2^-$	$1_2^-$	92	207	-21	1	-41	248	224

$\Lambda\Sigma$  coupling contributions normally are below 100 keV

# $\Lambda N$ interaction matrix elements in Nijmegen models

- G-Matrix elements from  $N\Lambda$ - $N\Sigma$  calculation fitted with sums of Gaussians, Yukawas, OBEP forms, ...
- $p_N s_\Lambda$  matrix elements (MeV) calculated with WS wave functions.

		<i>p</i> -shell					<i>s</i> -shell	
		$\bar{V}$	$\Delta$	$S_\Lambda$	$S_N$	$T$	$\bar{V}_s$	$\Delta_s$
fit-DJM	${}^7_\Lambda\text{Li}$	-1.142	0.438	-0.008	-0.414	0.031		
	${}^{16}_\Lambda\text{O}$	-1.161	0.441	-0.007	-0.401	0.030		
NSC97f	${}^7_\Lambda\text{Li}$	-1.086	0.421	-0.149	-0.238	0.055	-1.725	0.775
ESC04a	${}^7_\Lambda\text{Li}$	-1.287	0.381	-0.108	-0.236	0.013	-1.577	0.850
ESC08a	${}^7_\Lambda\text{Li}$	-1.221	0.146	-0.074	-0.241	0.055	-1.796	0.650

- Fitted matrix elements are roughly constant with A - same YNG interaction, WS wells have  $R=r_0A^{1/3}$ , but rms radii of *p*-shell nuclei are roughly constant.

# YN interaction contributions to g.s. binding energies

	${}^7_{\Lambda}\text{Li}$	${}^8_{\Lambda}\text{Li}$	${}^9_{\Lambda}\text{Be}$	${}^9_{\Lambda}\text{Li}$	${}^{10}_{\Lambda}\text{B}$	${}^{11}_{\Lambda}\text{B}$	${}^{12}_{\Lambda}\text{B}$	${}^{13}_{\Lambda}\text{C}$	${}^{15}_{\Lambda}\text{N}$	${}^{16}_{\Lambda}\text{N}$
keV	$1/2^+$	$1^-$	$1/2^+$	$3/2^+$	$1^-$	$5/2^+$	$1^-$	$1/2^+$	$3/2^+$	$1^-$
$\Lambda\Sigma$	78	160	4	183	35	66	103	28	59	62
$\Delta$	419	288	0	350	125	203	108	-4	40	94
$S_{\Lambda}$	0	-6	0	-10	-13	-20	-14	0	12	6
$S_N$	94	192	207	434	386	652	704	841	630	349
$T$	-2	-9	0	-6	-15	-43	-29	-1	-69	-45
sum	589	625	211	952	518	858	869	864	726	412
Exp	5.58	6.80	6.71	8.50	8.89	10.24	11.37	11.69		13.76
$\bar{V}$	-0.94	-1.02	-0.84	-1.06	-1.05	-1.04	-1.05	-0.96		-0.93

$$B_{\Lambda}^{\text{exp}}(\text{g.s.}) = [B_{\Lambda}^{\text{exp}}({}^5_{\Lambda}\text{He}) = 3.12 \pm 0.02 \text{ MeV}] - (A - 5)\bar{V} + \text{'sum'}$$

**Note  ${}^9_{\Lambda}\text{Be}$  anomaly.** Improve fit by adding a  $\Lambda NN$  term  
see Millener-Gal-Dover-Dalitz, PRC 31 (1985) 499

$\Lambda\Sigma$  matrix elements & contributions to  $B_{\Lambda}^{\text{g.s.}}$   
(in MeV) across the periodic table

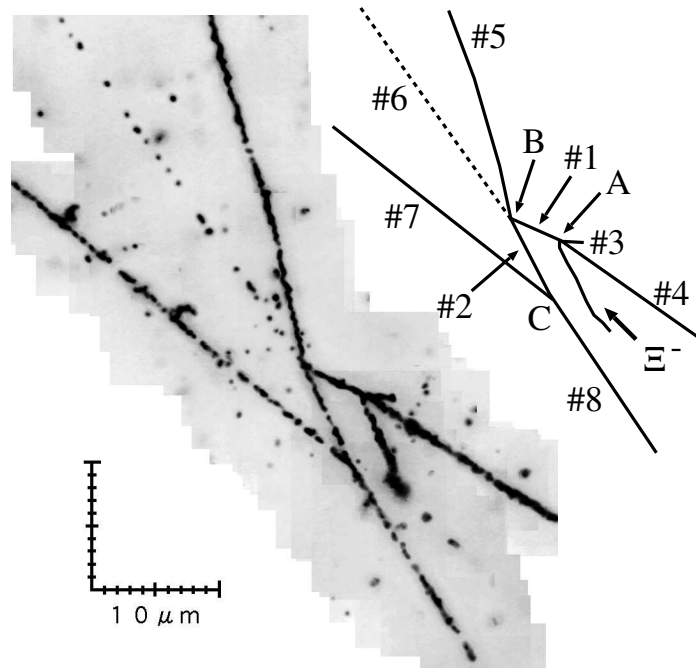
A. Gal & D.J. Millener PLB 725 (2013) 445

$N-Z$	${}^A_{\Lambda}Z$	$V_{\Lambda\Sigma}$	$\Lambda\Sigma(V)$	$\Delta_{\Lambda\Sigma}$	$\Lambda\Sigma(\Delta)$	$\Delta B_{\Lambda}^{\text{g.s.}}(\Lambda\Sigma)$
4	${}^9_{\Lambda}\text{He}$	1.194	0.143	4.070	0.104	0.246
8	${}^{49}_{\Lambda}\text{Ca}$	0.175	0.010	0.946	0.014	0.024
22	${}^{209}_{\Lambda}\text{Pb}$	0.0788	0.052	0.132	0.001	0.053

- $\Lambda\Sigma$  from Halderson, following NSC97f.
- $V_{\Lambda\Sigma}$  &  $\Delta_{\Lambda\Sigma}$  decrease drastically as overlap between 0s hyperon and high- $\ell$  excess neutrons becomes poorer with  $A$  ( $0f_{7/2}$  in  ${}^{49}_{\Lambda}\text{Ca}$ ,  $0h_{9/2}$  &  $0i_{13/2}$  in  ${}^{209}_{\Lambda}\text{Pb}$ ).
- **Conclusion:**  $\Lambda\Sigma$  contributes less than 100 keV to binding of medium & heavy n-rich hypernuclei.

# $\Lambda\Lambda$ hypernuclei





Nagara event,  $_{\Lambda\Lambda}^6\text{He}$ , (KEK-E373) PRL 87 (2001) 212502

$B_{\Lambda\Lambda}({}_{\Lambda\Lambda}^6\text{He}_{\text{g.s.}}) = 6.91 \pm 0.16 \text{ MeV}$ , unambiguously determined.

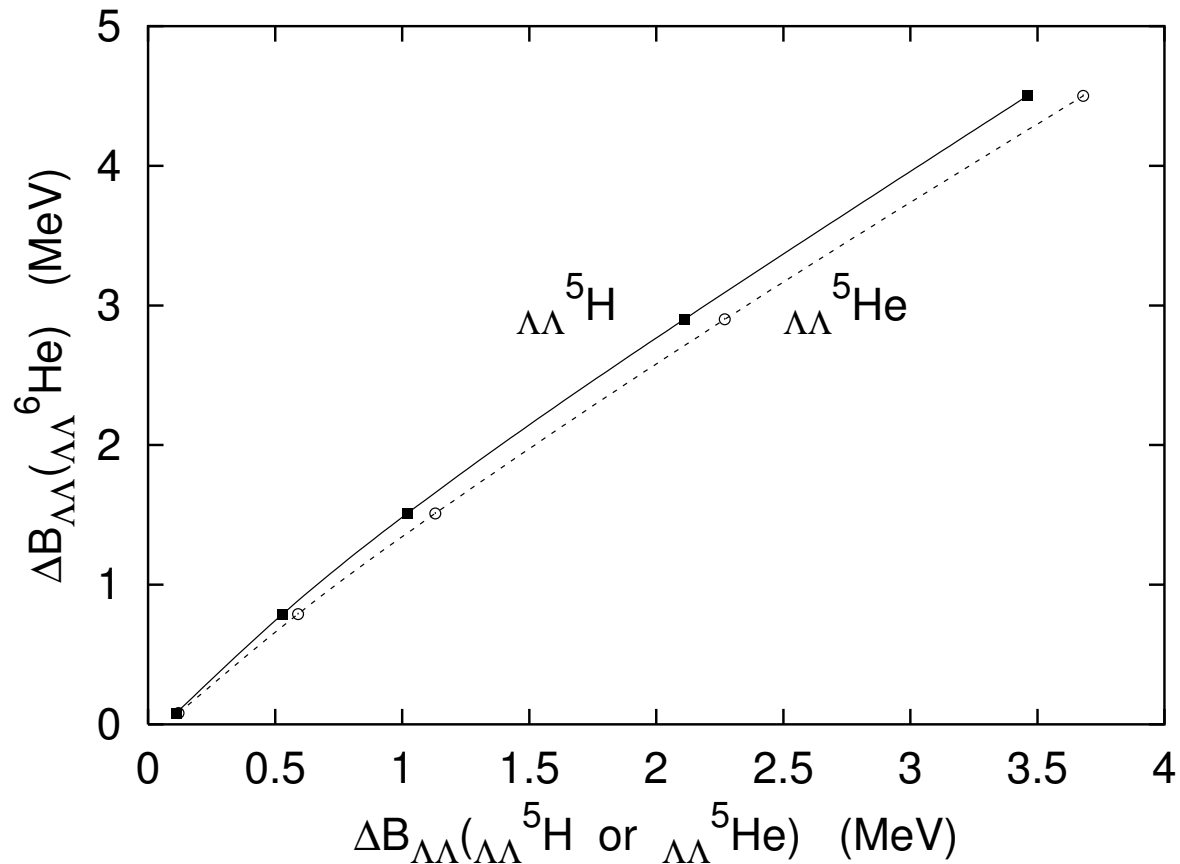
- A:  $\Xi^-$  capture  $\Xi^- + {}^{12}\text{C} \rightarrow {}_{\Lambda\Lambda}^6\text{He} + t + \alpha$
- B: weak decay  $_{\Lambda\Lambda}^6\text{He} \rightarrow {}^5_{\Lambda}\text{He} + p + \pi^-$  (no  $_{\Lambda\Lambda}^6\text{He} \rightarrow {}^4\text{He} + H$ )
- C:  ${}^5_{\Lambda}\text{He}$  nonmesic weak decay to 2  $Z=1$  recoils + n.

# The elusive H dibaryon

Jaffe's  $\mathbf{H}(uuddss)$  [PRL 38 (1977) 195] predicted stable

$$\mathbf{H} \sim \mathcal{A}[\sqrt{1/8} \Lambda\Lambda + \sqrt{1/2} N\Xi - \sqrt{3/8} \Sigma\Sigma, ]_{I=S=0}$$

- To forbid  ${}_{\Lambda\Lambda}^6\text{He} \rightarrow \mathbf{H} + {}^4\text{He}$ , impose  $B(\mathbf{H}) \leq 7$  MeV.  
A bound H most likely overbinds  ${}_{\Lambda\Lambda}^6\text{He}$   
[Gal, PRL 110 (2013) 179201].
- Weakly bound H in Lattice QCD calculations.  
SU(3)<sub>f</sub> breaking pushes it to  $\approx N\Xi$  threshold,  
 $\approx 26$  MeV in  $\Lambda\Lambda$  continuum [HALQCD, NPA 881  
(2012) 28; Haidenbauer & Meißner, ibid. 44].
- $\Lambda\Lambda$  correlation femtoscopy in pp at  $\sqrt{s}=7$  TeV  
(ALICE exp. at the LHC, arXiv:1805.12455)  
rules out a bound H (Laura Fabbietti @HYP18).



Faddeev calc. by I.N. Filikhin, A. Gal, NPA 707 (2002) 491

$$\Delta B_{\Lambda\Lambda}(\Lambda\Lambda^6\text{He}) \equiv B_{\Lambda\Lambda}(\Lambda\Lambda^6\text{He}) - 2B_{\Lambda}(\Lambda^5\text{He}) \approx 0.7 \text{ MeV}$$

implying that  $\Lambda\Lambda^5\text{H}$  &  $\Lambda\Lambda^5\text{He}$  are also bound.

With  $\Lambda\Lambda^4\text{H}$  likely unbound,  $\Lambda\Lambda$  binding onset is  $\Lambda\Lambda^5\text{H}$  &  $\Lambda\Lambda^5\text{He}$ .

# Binding energy consistency of $\Lambda\Lambda$ hypernuclei

event	${}_{\Lambda\Lambda}^AZ$	$B_{\Lambda\Lambda}^{\text{exp}}$	$B_{\Lambda\Lambda}^{\text{CM}} \dagger$	$B_{\Lambda\Lambda}^{\text{SM}} \dagger\dagger$
E373-Nagara	${}_{\Lambda\Lambda}^6\text{He}$	$6.91 \pm 0.16$	$6.91 \pm 0.16$	$6.91 \pm 0.16$
E373-DemYan	${}_{\Lambda\Lambda}^{10}\text{Be}$	$14.94 \pm 0.13 \ddagger$	$14.74 \pm 0.16$	$14.97 \pm 0.22$
E373-Hida	${}_{\Lambda\Lambda}^{11}\text{Be}$	$20.83 \pm 1.27$	$18.23 \pm 0.16$	$18.40 \pm 0.28$
E373-Hida	${}_{\Lambda\Lambda}^{12}\text{Be}$	$22.48 \pm 1.21$	–	$20.72 \pm 0.20$
E176	${}_{\Lambda\Lambda}^{13}\text{B}$	$23.4 \pm 0.7^*$	–	$23.21 \pm 0.21$

$\dagger$  E. Hiyama et al., PRL 104 (2010) 212502, & refs. therein

$\dagger\dagger$  A. Gal, D.J. Millener, PLB 701 (2011) 342, assuming that

$$\langle V_{\Lambda\Lambda} \rangle \approx \Delta B_{\Lambda\Lambda}({}_{\Lambda\Lambda}^6\text{He}) = 0.67 \pm 0.16 \text{ MeV}$$

$\ddagger$  Assuming production in  ${}_{\Lambda\Lambda}^{10}\text{Be}$  non g.s.  $2^+(3.04 \text{ MeV})$

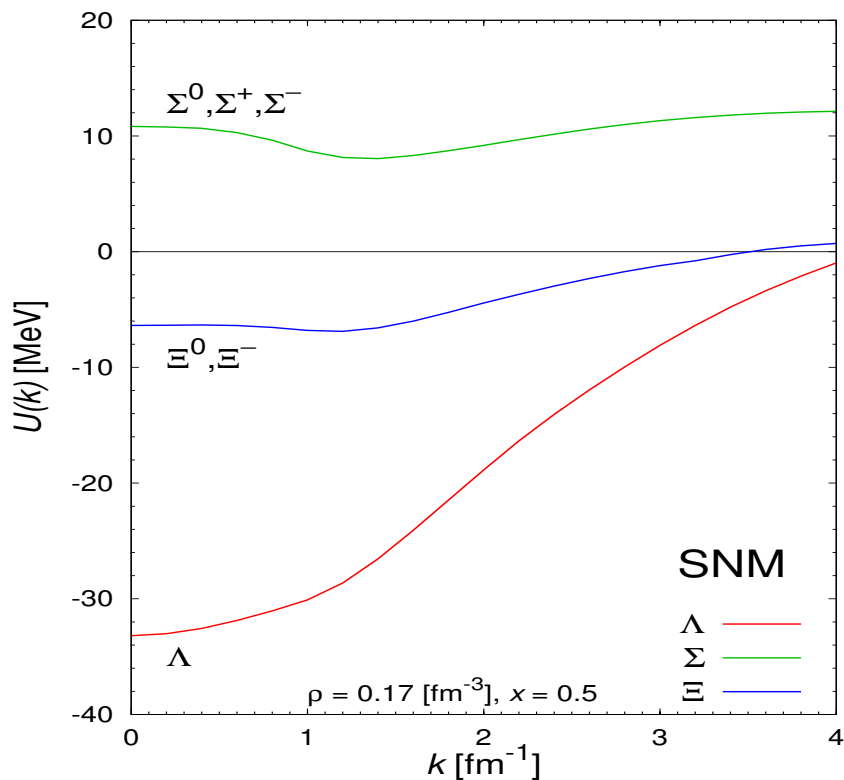
\* Assuming  ${}_{\Lambda\Lambda}^{13}\text{B}_{\text{g.s.}}$  decay to  ${}_{\Lambda}^{13}\text{C}^*(5/2^+, 3/2^+; 4.8 \text{ MeV}) + \pi^-$

- Unassigned Hida event [PTPS 185 (2010) 335]

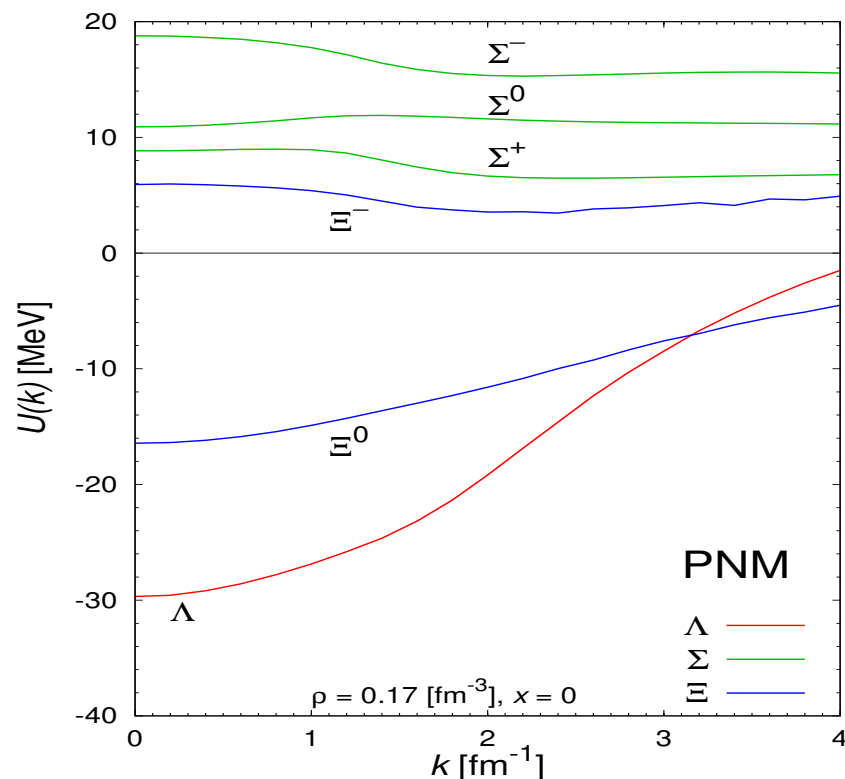
- $B_{\Lambda\Lambda}^{\text{SM}} \approx B_{\Lambda\Lambda}^{\text{CM}}$ , but SM spans a wider  $A$  range

# Other Strange Hadrons in Matter

# Hyperon-Nucleus potentials from LQCD



Symmetric nuclear matter



Pure neutron matter

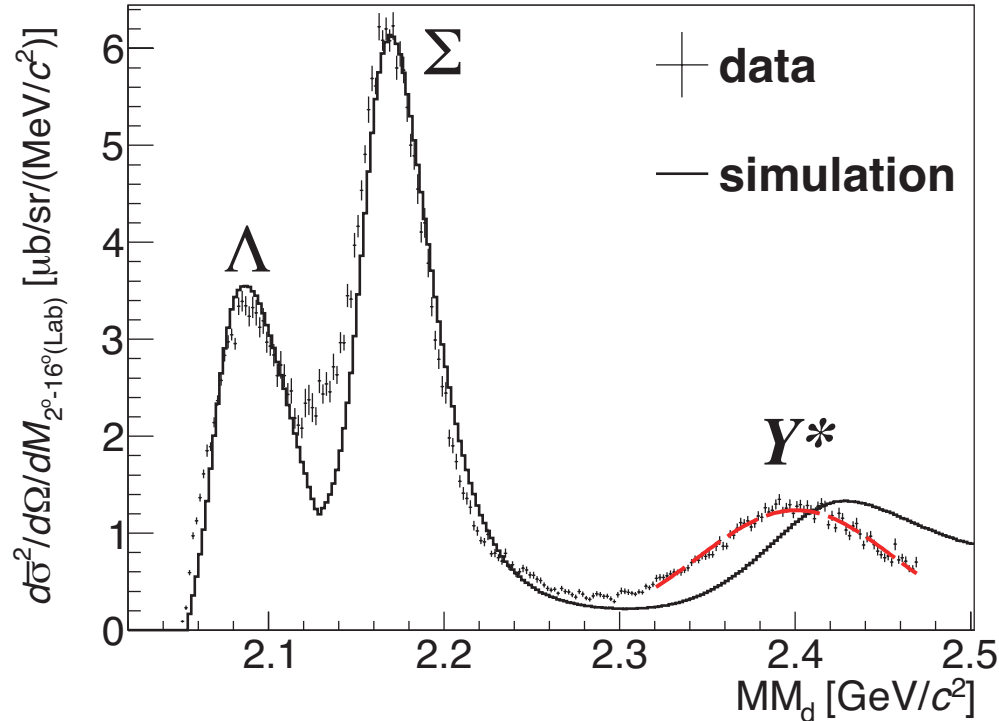
T. Inoue, for HAL QCD Collab., arXiv:1612.08399

**BHF applied to Lattice YN potentials**

$\Sigma$  – repulsion,  $\Xi$  – weak attraction

How about  $YN^*$  potentials, e.g.  $\Lambda^*(1405)N$ , related to  $K^-pp$ ?

## J-PARC E27 $d(\pi^+, K^+)$ missing-mass spectrum



PTEP 2014, 101D03

$Y^*$  quasi-free peak shifted by  $\approx -22$  MeV,  
indicating  $Y^*N$  attraction [ $Y^* = \Sigma^*(1385)$  &  $\Lambda^*(1405)$ ].

Two dibaryons likely below  $K^-pp$  threshold:

(i)  $\Lambda^*N$  observed in E15 (ii)  $\Sigma^*N$  formed in E27 ?

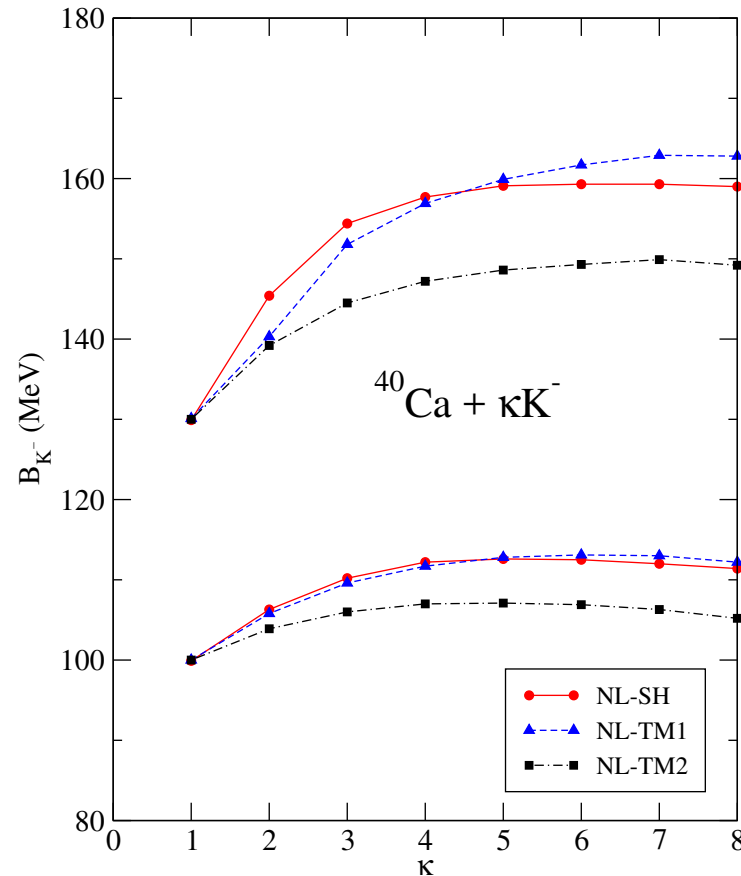
# $\Lambda^*(1405)N$ & $\Sigma^*(1385)N$ dibaryons?

- $\Lambda(1405)N$  is a doorway to a  $I=1/2, J^P=0^-$   $\bar{K}NN$ , found quasibound in all calculations & in E15. It is coupled to  $\pi\Lambda N$  and  $\pi\Sigma N$ , but  $\pi\Lambda N$  cannot support any strongly attractive meson-baryon s-wave interaction.
- The  $\pi\Lambda N$  system can benefit from strong meson-baryon  $p$ -wave interactions fitted to  $\Delta(1232) \rightarrow \pi N$  and  $\Sigma(1385) \rightarrow \pi\Lambda$  form factors. Maximize isospin and angular momentum couplings by full alignment:  $I=3/2, J^P=2^+$ , a good example of a Pion Assisted Dibaryon.

Gal-Garcilazo, NPA 897 (2013) 167



# $\Lambda^*$ : do antikaons condense on earth?



D. Gazda, E. Friedman, A. Gal, J. Mareš, PRC 77 (2008) 045206

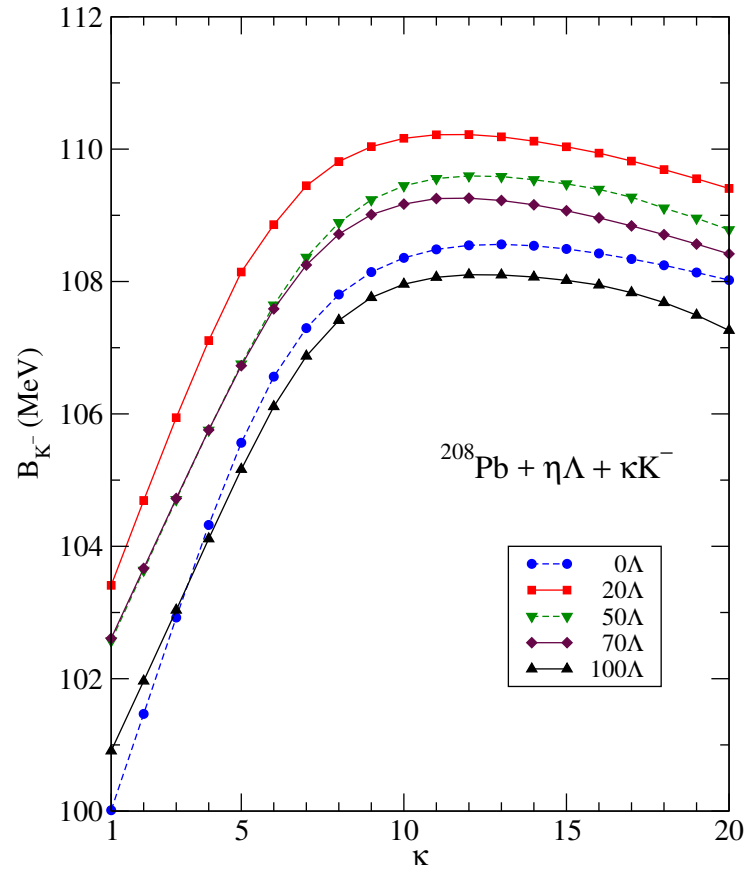
**RMF:  $B_{\bar{K}}$  saturates in multi- $K^-$   $^{40}\text{Ca}$  nuclei**

Large binding of 100, 130 MeV assumed for  $\kappa=1$

Vector-meson repulsion among  $\bar{K}$  mesons

$$B_{\bar{K}}(\kappa \rightarrow \infty) \ll (m_K + M_N - M_\Lambda) \approx 320 \text{ MeV}$$

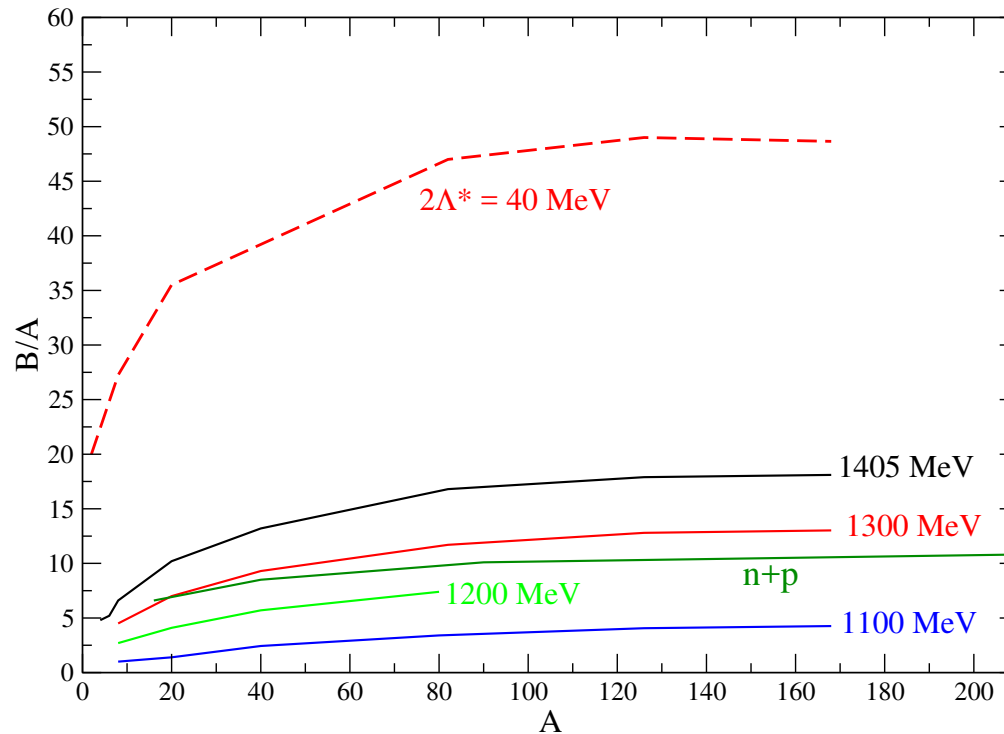
# ...and adding $\Lambda$ hyperons



Gazda-Friedman-Gal-Mareš, Phys. Rev. C 80 (2009) 035205

Saturation of  $B_{\bar{K}}$  in RMF for  $^{208}\text{Pb} + \eta\Lambda + \kappa\bar{K}^-$

Hyperons dominate stable self-bound strange matter  
No kaon condensation on earth...



## $\Lambda^*(1405)$ matter: stable or unstable?

**Stable:** Y. Akaishi T.Yamazaki, PLB 774 (2017) 522

**Unstable:** J. Hrtánková et al., arXiv:1805.11368, and in HYP2018

- RMF calculations demonstrate that  $B/A$  saturates, at values allowing  $\Lambda^*\Lambda^* \rightarrow \Lambda\Lambda$  &  $\Sigma\Sigma$  strong decays.

Central densities saturate at about  $2\rho_0$ .

# Summary & Outlook

- $\Lambda N$  hypernuclear spin dependence deciphered.
- How small is  $\Lambda$  spin-orbit splitting and why?
- Role of 3-body  $\Lambda NN$  interactions in hypernuclei & neutron stars?
- Resolve the  ${}^3_{\Lambda}\text{H}$  lifetime puzzle from HIC.
- Re-measure the  ${}^4_{\Lambda}\text{H}-{}^4_{\Lambda}\text{He}$  complex (E13→E63).
- Search for n-rich  ${}^A_{\Lambda}\text{Z}$ ; ( ${}^6_{\Lambda}\text{H}$  not seen in E10).
- Repulsive  $\Sigma$ -nuclear interaction; how repulsive? (relevant to neutron star matter scenarios).
- Search for  $H$  dibaryon in  $(K^-, K^+)$  (E42).
- Onset of  $\Lambda\Lambda$  binding:  ${}_{\Lambda\Lambda}{}^4\text{H}$  or  ${}_{\Lambda\Lambda}{}^5\text{Z}$ ? (E07).

- Shell model works well **for g.s.** beyond  ${}_{\Lambda\Lambda}{}^6\text{He}$ .
- Study **excited states** by slowing down  $\Xi^-$  from  $\bar{p}p \rightarrow \Xi^- \bar{\Xi}^+$  in **FAIR (PANDA)**.
- Do  $\Xi$  hyperons quasi-bind in nuclei ( $\Xi N \rightarrow \Lambda\Lambda$ )? No quasibound  $\Xi$  established yet (**E05**).
- Onset of  $\Xi$  stability:  ${}_{\Lambda\Xi}{}^6\text{He}$  or  ${}_{\Lambda\Lambda\Xi}{}^7\text{He}$ ?
- No  $\bar{K}$  condensation in self-bound matter.  $\{N, \Lambda, \Xi\}$  provides Strange-Hadronic-Matter g.s.
- Search for a  $\Sigma(1385)\text{N}$  ( $I=3/2, J^P=2^+$ ) dibaryon  $\mathcal{Y}$  at J-PARC or GSI in  $\pi^- + d \rightarrow \mathcal{Y}^- + K^+, \mathcal{Y}^- \rightarrow \Sigma^- + n$ .

**Thanks for your attention!**