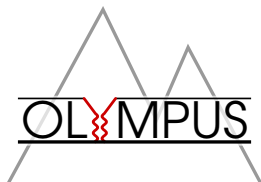


Two-Photon Exchange Experiments: Present and Future

D.K. Hasell



Hadronic Physics with Lepton and Hadron Beams
CEBAF Center
September 5–8, 2017



Introduction

~ 100 years ago Rutherford named the nucleus of hydrogen the proton

But there are still many puzzles remaining for this fundamental particle

- proton spin

C.A. Aidala *et al.*, *Reviews of Modern Physics*, 85 (2013) 655–691

- proton mass

S. Dürr *et al.*, *Science*, 322 (2008)1224–1227

- proton radius

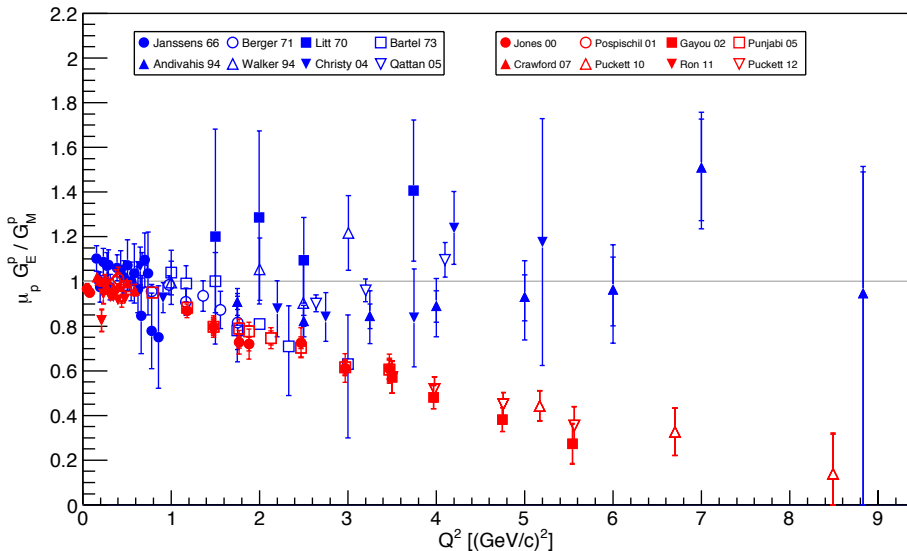
R. Pohl *et al.*, *Annual Review of Nuclear and Particle Science*, 63 (2013) 175–204

My talk addresses still another problem, namely the proton form factors

A. Afanasev *et al.*, *Progress in Particle and Nuclear Physics*, 95 (2017) 245–274



Problem with Form Factor Ratio $\mu_p G_E^p / G_M^p$ Measurements



Proposed Explanation - Two Photon Exchange (TPE)

Thought to be a small effect

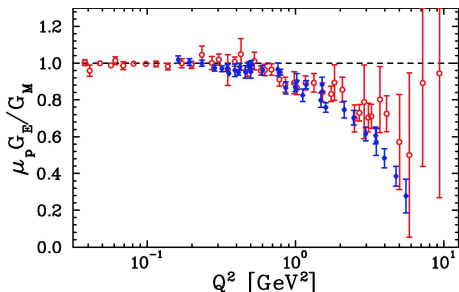
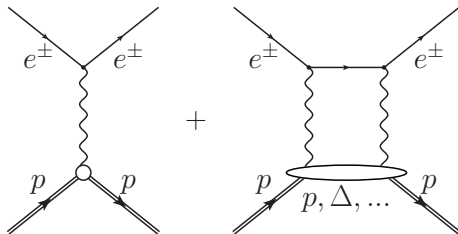
- suppressed by order α
- “soft” radiative corrections included

But “hard” TPE difficult

- model dependent intermediate p, Δ, \dots

Calculations suggest it can resolve the discrepancy

Need a definitive experiment



J. Arrington *et al.*, Phys. Rev. C 76 (2007) 035205



How to Measure “Hard” Two-Photon Contribution

$$\begin{aligned}
 \frac{d\sigma}{d\Omega} \propto |\mathcal{M}|^2 &= \left| \text{Diagram 1} + \text{Diagram 2} + \dots \right|^2 \\
 &= \left| \text{Diagram 1} \right|^2 + \left| \text{Diagram 2} \right|^2 + \\
 &\quad 2\text{Re} \left[\text{Diagram 1}^\dagger \text{Diagram 2} \right] + \dots
 \end{aligned}$$

Interference term has a factor z^3 , where z is the lepton charge

⇒ Interference term changes sign between e^+p and e^-p scattering

Definitive Measure of Two-Photon Contribution

Measure $\sigma_{e^+p}/\sigma_{e^-p}$

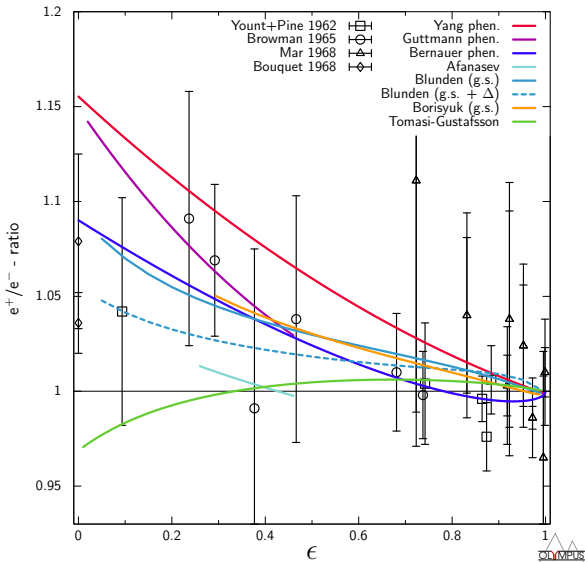
$$\frac{\sigma_{e^+p}}{\sigma_{e^-p}} \approx 1 + 4 \frac{\text{Re}(\mathcal{M}_{1\gamma}^\dagger \mathcal{M}_{2\gamma})}{\mathcal{M}_{1\gamma}^2}$$

Existing data

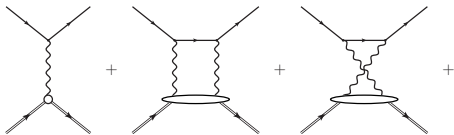
- low Q^2
- large uncertainties

Three recent experiments

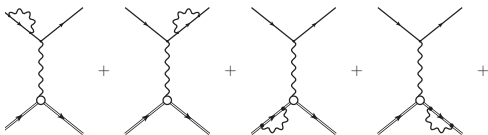
- VEPP-3 - Novosibirsk
- CLAS - JLab
- OLYMPUS - DESY



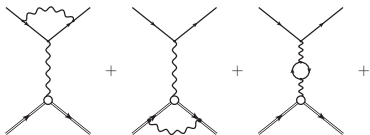
Radiative Corrections that Need to be Included



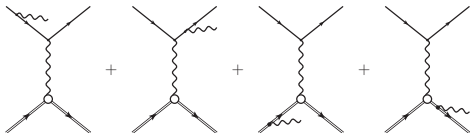
1 and 2 photon exchange



self-energy



vertex corrections and
vacuum polarization



bremstrahlung



Radiative Corrections in Elastic Cross Section

$$\begin{aligned}
 \frac{d\sigma}{d\Omega} = & \left| \begin{array}{c} \text{tree-level diagram} \end{array} \right|^2 \\
 & + 2 \operatorname{Re} \left[\begin{array}{c} \text{tree-level diagram}^\dagger \end{array} \left(\begin{array}{c} \text{tree-level diagram} \\ + \text{one-loop diagram} \\ + \text{tree-level diagram with photon} \\ + \text{tree-level diagram with photon} \end{array} \right) \right] \quad \left. \vphantom{\frac{d\sigma}{d\Omega}} \right\} z^2 \\
 & + 2 \operatorname{Re} \left[\begin{array}{c} \text{tree-level diagram}^\dagger \end{array} \left(\begin{array}{c} \text{two-photon exchange diagram} \\ + \text{two-photon exchange diagram} \end{array} \right) \right] \quad \left. \vphantom{\frac{d\sigma}{d\Omega}} \right\} z^3 \\
 & + 2 \operatorname{Re} \left[\begin{array}{c} \text{tree-level diagram}^\dagger \end{array} \left(\begin{array}{c} \text{tree-level diagram with photon} \\ + \text{tree-level diagram with photon} \\ + \text{tree-level diagram with photon} \end{array} \right) \right] \quad \left. \vphantom{\frac{d\sigma}{d\Omega}} \right\} z^4 \\
 & + \mathcal{O}(\alpha^4)
 \end{aligned}$$

Rebecca Russell, MIT

Even powers of z same for e^+p and e^-p scattering must be included



Radiative Corrections from Inelastic Processes

$$\begin{aligned}
 \frac{d\sigma_{\text{inel}}}{d\Omega} = & \left| \begin{array}{c} \text{Diagram 1} \\ + \\ \text{Diagram 2} \end{array} \right|^2 & \left. \vphantom{\frac{d\sigma_{\text{inel}}}{d\Omega}} \right\} z^2 \\
 & + 2 \operatorname{Re} \left[\left(\begin{array}{c} \text{Diagram 3} \\ + \\ \text{Diagram 4} \end{array} \right)^\dagger \left(\begin{array}{c} \text{Diagram 5} \\ + \\ \text{Diagram 6} \end{array} \right) \right] & \left. \vphantom{\frac{d\sigma_{\text{inel}}}{d\Omega}} \right\} z^3 \\
 & + \left| \begin{array}{c} \text{Diagram 7} \\ + \\ \text{Diagram 8} \end{array} \right|^2 & \left. \vphantom{\frac{d\sigma_{\text{inel}}}{d\Omega}} \right\} z^4 \\
 & + \mathcal{O}(\alpha^4)
 \end{aligned}$$

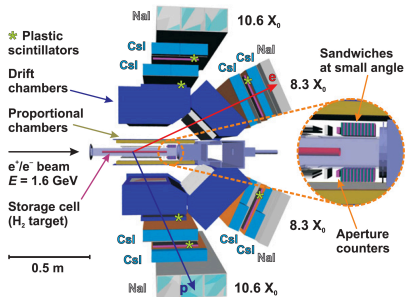
Rebecca Russell, MIT

Inelastic IR divergences cancel with elastic divergences

- must separate “hard” and “soft” parts in two-photon exchange
- “soft” part included in radiative corrections, “hard” part measured
- prescriptions defining “soft” - e.g. Mo - Tsai, Maximon - Tjon

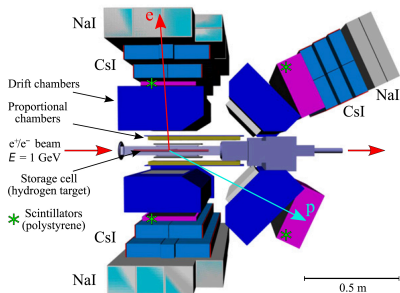


VEPP-3 Detector Configuration



Run 1 (2009)

$$E_{Beam} = 1.594 \text{ GeV}$$



Run 2 (2011–2012)

$$E_{Beam} = 0.998 \text{ GeV}$$

Large acceptance, non-magnetic detector configuration

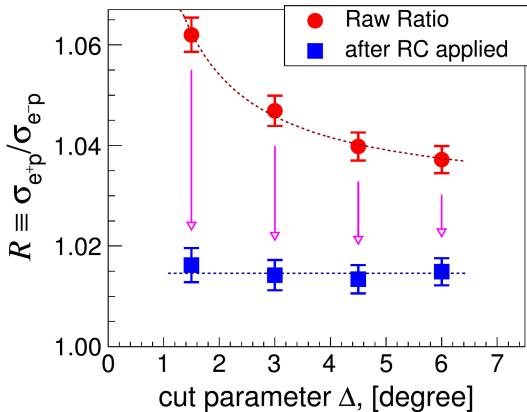
- same acceptance, efficiency for both electrons and positrons
- lepton and proton detected in coincidence
- forward angle measurement used for luminosity normalization

VEPP-3 Radiative Corrections

Dedicated event generator

- ESEPP
- full radiative corrections
- GEANT4 detector simulation

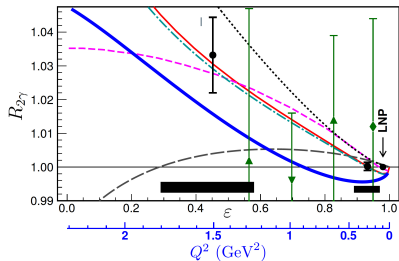
Sensitivity of ratio to radiative corrections



I.A. Rachek *et al.* Phys. Scr. **T166** 014017 (2015).

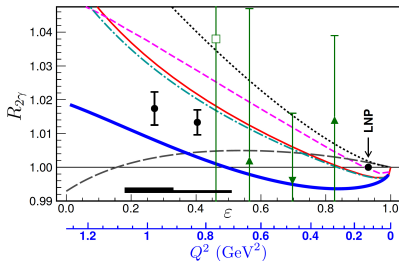


VEPP-3 Results



$$E_{Beam} = 1.594 \text{ GeV}$$

- *I. A. Qattan, et al.,*
- *P. G. Blunden, et al.,*
- · - · - · *D. Borisyuk and A. Kobushkin,*
- - - - *E. Tomasi-Gustafsson, et al.,*
- - - - *J. Arrington and I. Sick,*
- *J. C. Bernauer, et al.,*



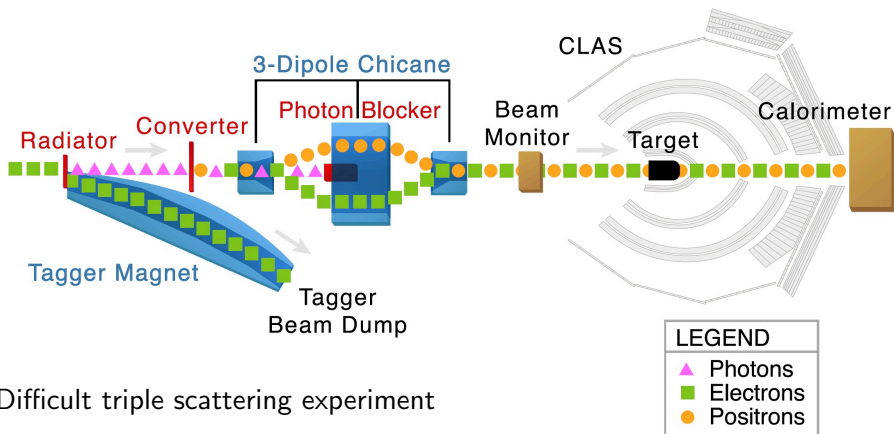
$$E_{Beam} = 0.998 \text{ GeV}$$

- Phys. Rev. C **84** (2011) 054317
- Phys. Rev. C **72** (2005) 034612
- Phys. Rev. C **78** (2008) 025208
- Phys. Atom. Nucl. **76** (2013) 937
- Phys. Rev. C **70** (2004) 028203
- Phys. Rev. C **90** (2014) 015206

I.A. Rachek *et al.* Phys. Rev. Lett. **114** 062005 (2015).



CLAS Detector Configuration

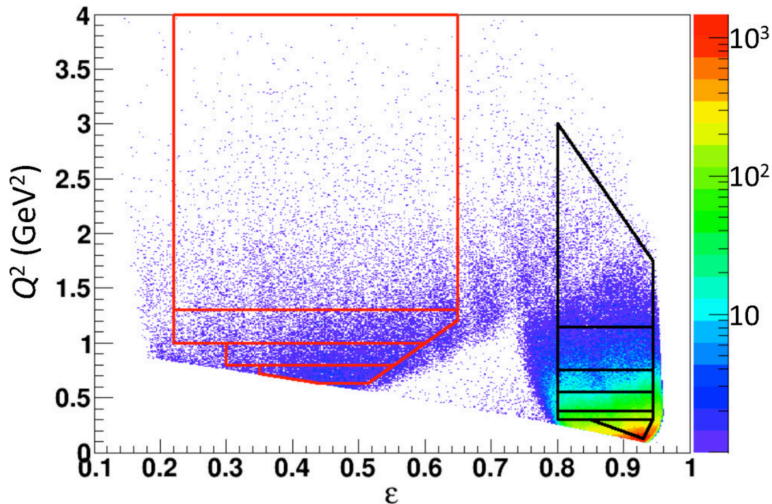


Difficult triple scattering experiment

Must reconstruct beam energy by measuring both lepton and proton

D. Rimal *et al.* Phys. Rev. **C95** 065291 (2017).

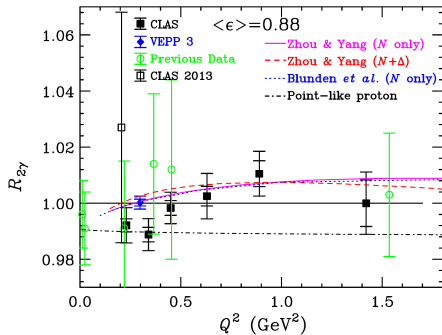
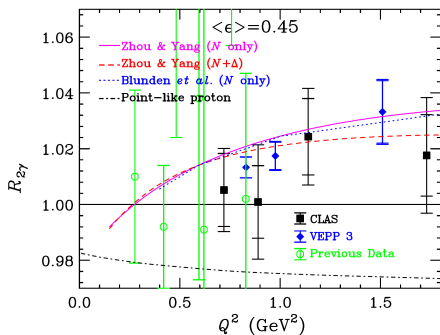


CLAS Bins for ϵ Dependence

D. Rimal *et al.* Phys. Rev. C95 065291 (2017).

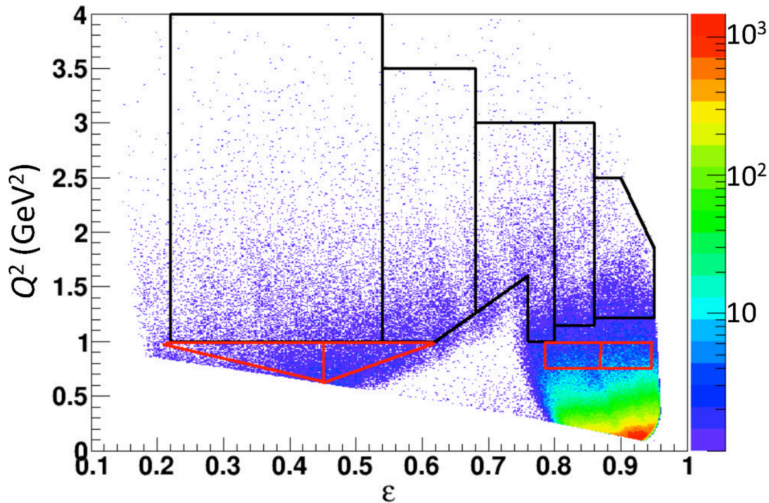


CLAS, VEPP-3, and Previous Results versus ϵ



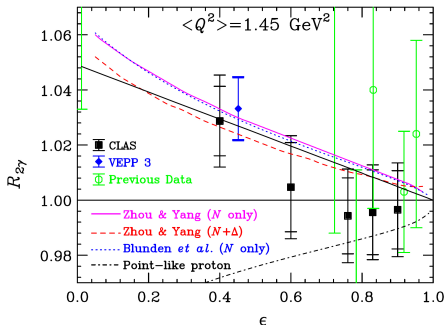
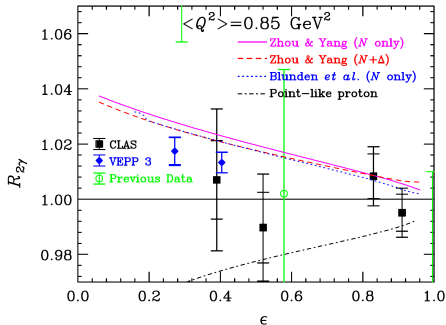
D. Rimal *et al.* Phys. Rev. **C95** 065291 (2017).



CLAS Bins for Q^2 Dependence

D. Rimal *et al.* Phys. Rev. C95 065291 (2017).

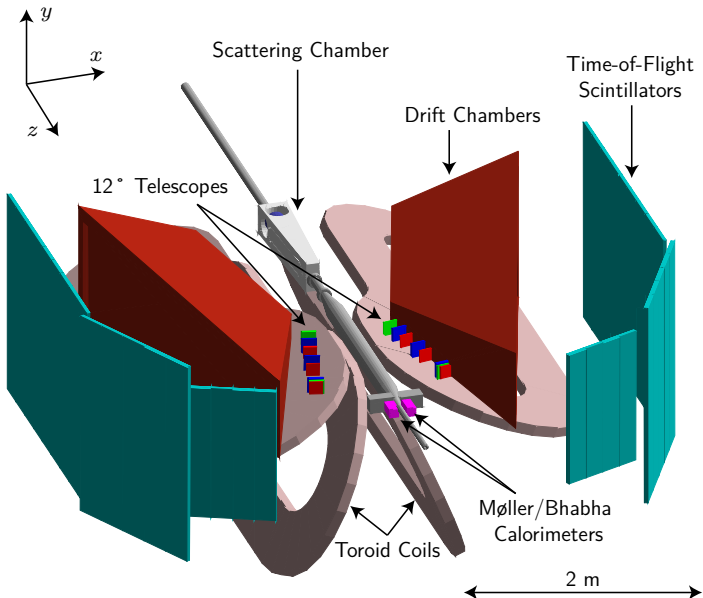


CLAS, VEPP-3, and Previous Results versus ϵ 

D. Rimal *et al.* Phys. Rev. **C95** 065291 (2017).



Detector Overview



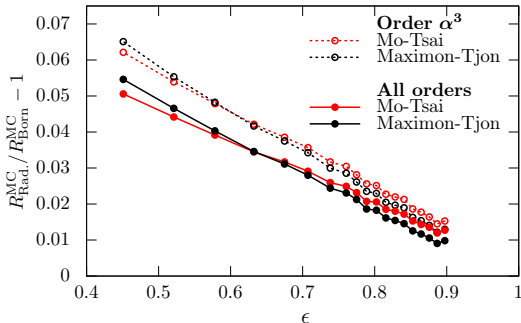
Radiative Corrections in OLYMPUS

Dedicated event generator

- checked against ESEPP - VEPP-3
- full radiative corrections
- GEANT4 detector simulation

Small differences between results with different “Soft” photon definitions

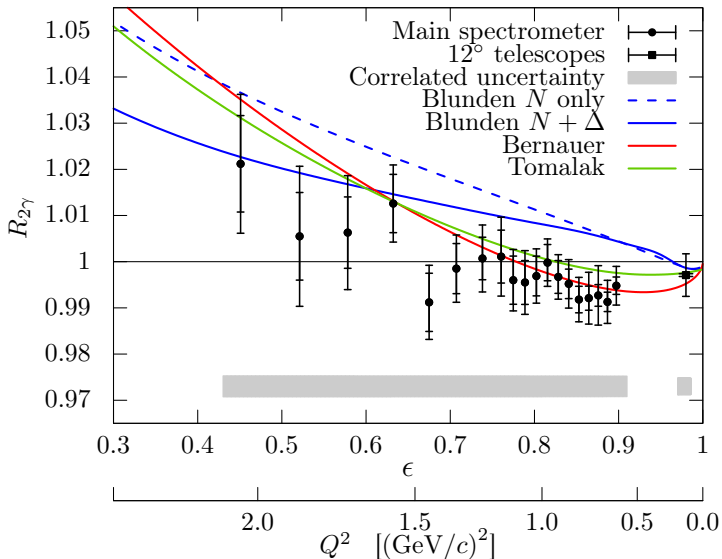
Some sensitivity to α^3 or exponentiation



B. Henderson *et al.* Phys. Rev. Lett. **118** 092501 (2017).



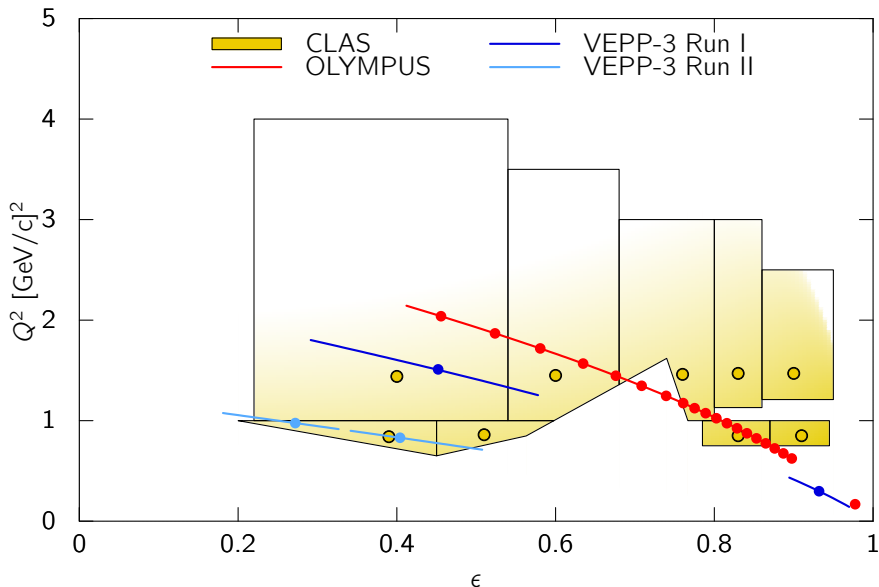
OLYMPUS Results

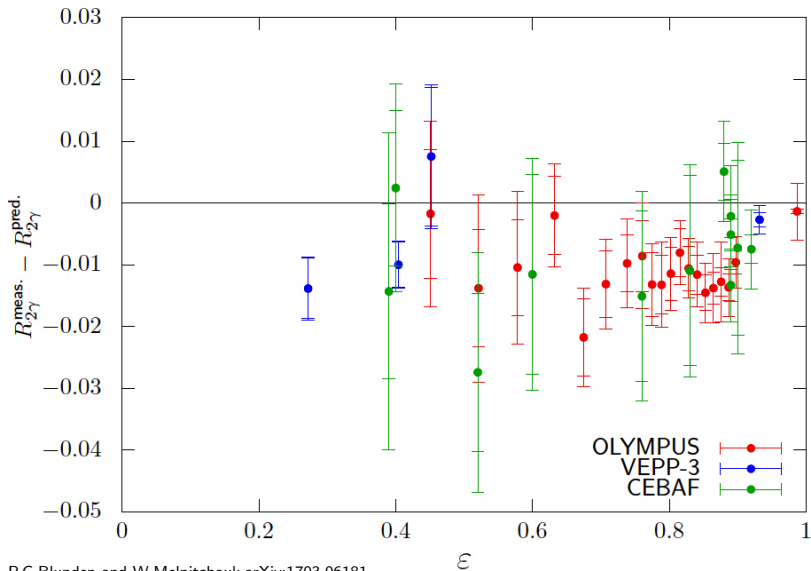


B. Henderson *et al.* Phys. Rev. Lett. **118** 092501 (2017).



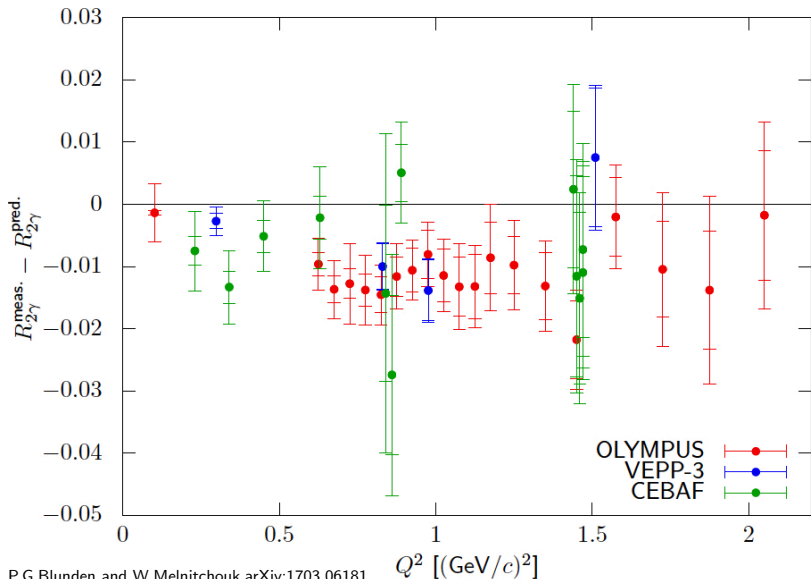
Comparing the Three Experiments - (ϵ, Q^2) Reach



Comparison with Blunden $N + \Delta$ 

P.G.Blunden and W.Melnitchouk, arXiv:1703.06181.

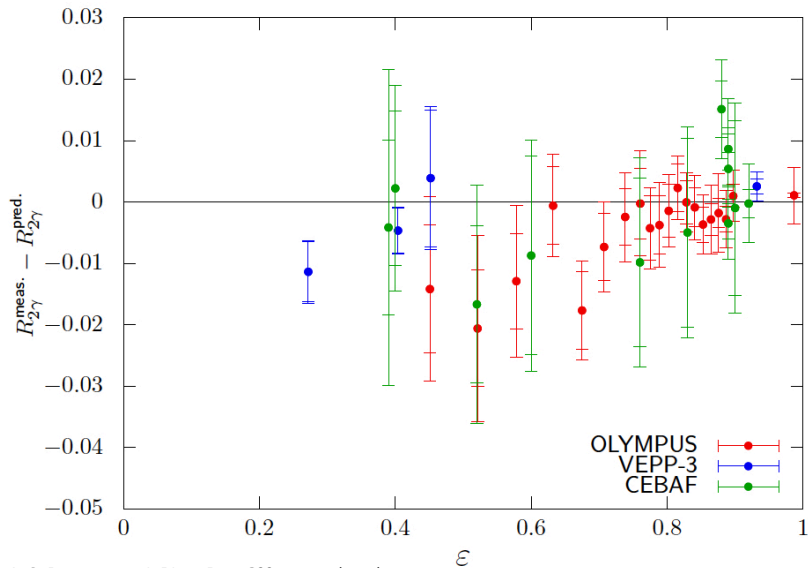


Comparison with Blunden $N + \Delta$ 

P.G.Blunden and W.Melnitchouk, arXiv:1703.06181.



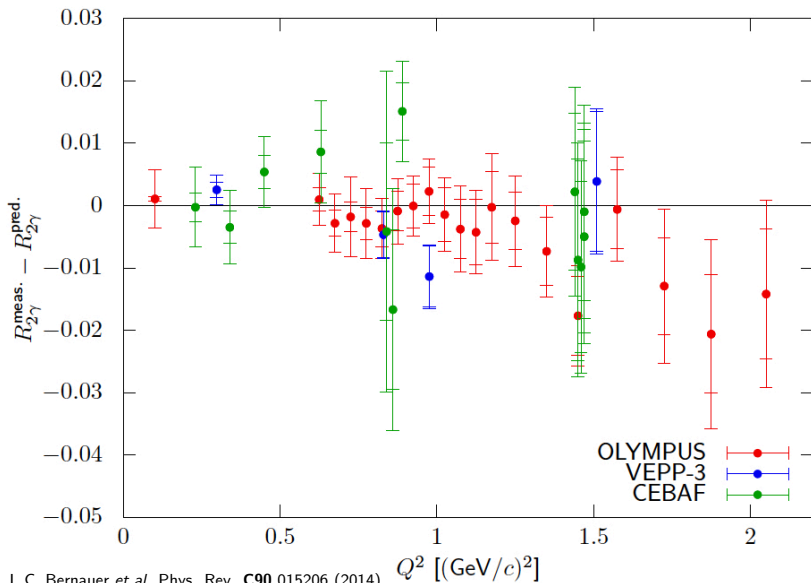
Comparison with Bernauer



J. C. Bernauer *et al.* Phys. Rev. **C90** 015206 (2014).

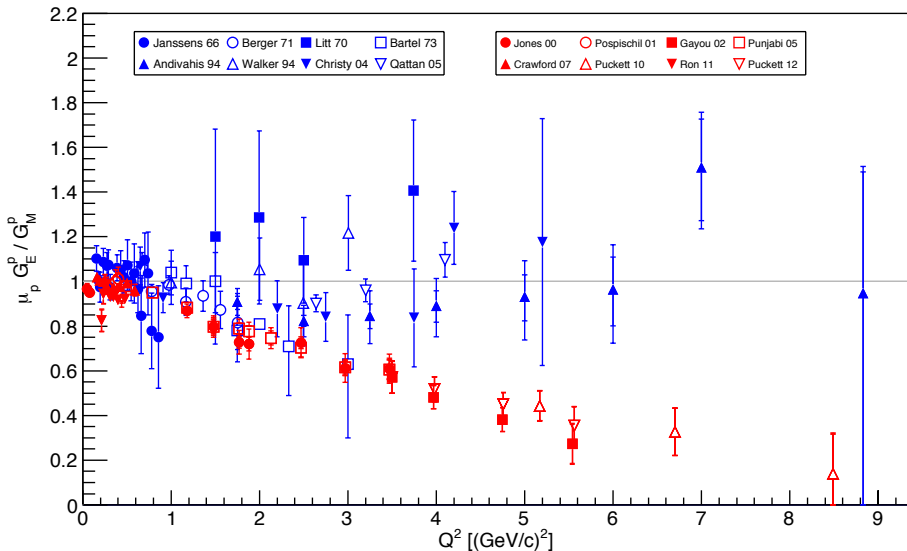


Comparison with Bernauer

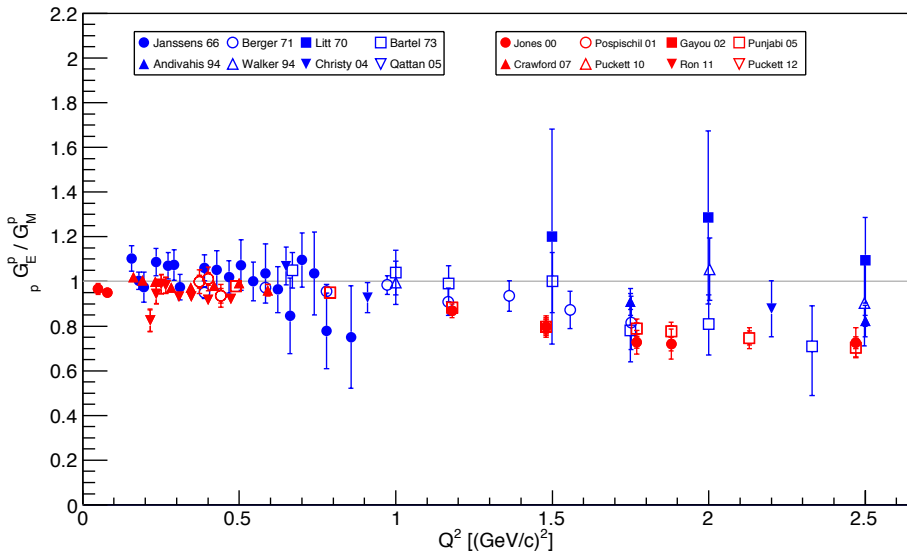


J. C. Bernauer *et al.* Phys. Rev. **C90** 015206 (2014).

Should we be surprised ?



Discrepancy in $\mu_p G_E^p / G_M^p$ at $Q^2 < 2.5$ (GeV/c)²



Summary of Experimental Results

$R_{2\gamma}$ measured for $Q^2 < 2.3 \text{ (GeV/c)}^2$

Radiative corrections and prescription for handling TPE important.

Small, $< 1\%$, hard two-photon exchange observed

Evidence for effect increasing with increasing Q^2 (decreasing ϵ)

Results less than expected from theoretical calculations

In better agreement with phenomenological predictions

Further theoretical effort on radiative corrections needed

Experiments at higher energy required to resolve discrepancy

- When ? Where ?



Possible Future Two-Photon Experiments at JLab

CLAS12 ?

Positrons ?

- J. Grames (NSTAR2017), polarized positrons, 100 nA
- higher unpolarized positron current ?
- JPos17 - September 12-15, JLab
 - good idea and not just for two-photon exchange
 - DVCS, full flavor decomposition of structure functions

Hall A

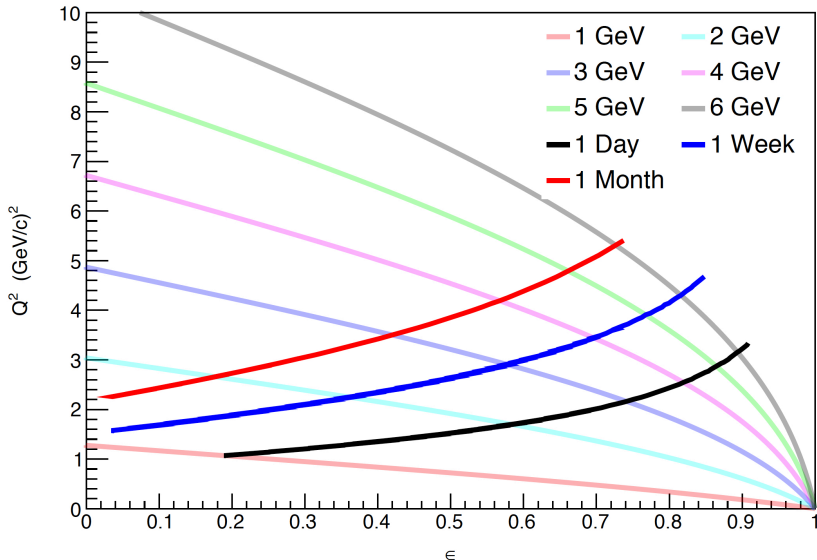
- 10 cm liquid hydrogen target
- HRS 6 msr

Event rate $\sim 1.6 \times 10^{-6}$ per second per femtobarn



Time to Measure $R_{2\gamma}$ at One Point in (ϵ, Q^2)

Time to collect 0.5 % statistics for both e^- and e^+ with 50% livetime



Possible Future Two-Photon Experiments at DESY

OLYMPUS 2 ?

Electrons and positrons available directly from DESY synchrotron

- 1–6 GeV electrons or positrons
- 30 nA
- higher current ?

Test hall

- 10 cm liquid hydrogen target
- high resolution, fine granularity calorimeter (PbWO_4) 10 msr
 - increase coverage with more calorimeter modules

Event rate $\sim 7.9 \times 10^{-7}$ per second per femtobarn

$\sim 2 \times$ slower than JLab Hall A with HRS



Recent and Future Meetings

Two parallel session on experiment and theory of two-photon exchange

- NSTAR Conference, Columbia, SC, August 20-23

Hadronic Physics with Lepton and Hadron Beams

- JLab, Newport News, VA September 5-8

International Workshop on Physics with Positrons at Jefferson Lab

- JLab, Newport News, VA September 12-15

Workshop on Two-Boson Exchange

- UMass, Amherst, MA September 28-30



Thank You

Work supported by the United States Department of Energy.

Nucleon Form Factors from Elastic Electron Scattering

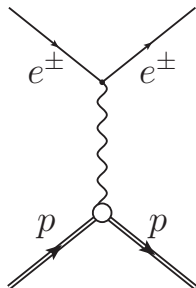
One photon exchange approximation

$$\gamma^\mu F_1^N(Q^2) + i\sigma^{\mu\nu} q_\nu \frac{\kappa}{2M} F_2^N(Q^2)$$

Electric and magnetic form factors

$$G_E^N(Q^2) = F_1^N(Q^2) - \tau\kappa F_2^N(Q^2)$$

$$G_M^N(Q^2) = F_1^N(Q^2) + \kappa F_2^N(Q^2)$$



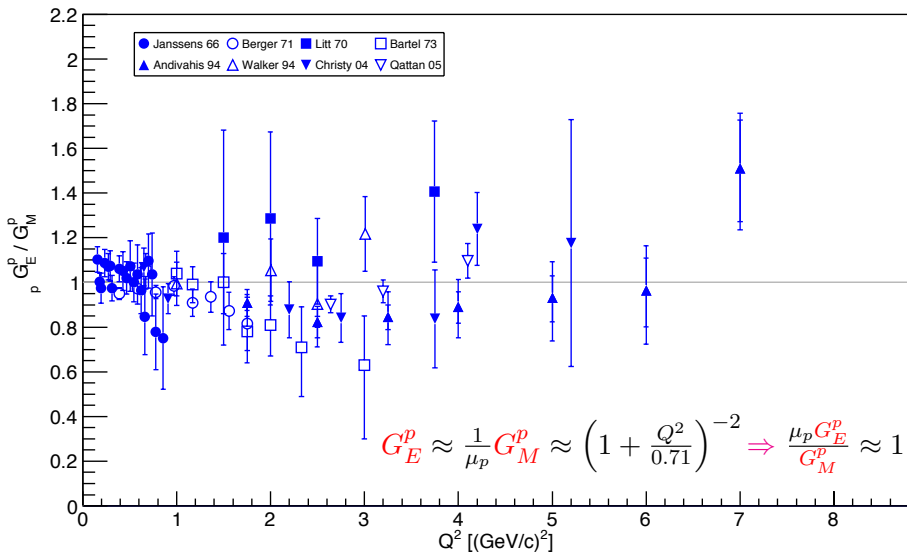
Rosenbluth cross section

$$\left(\frac{d\sigma}{d\Omega}\right)_{Mott} \left[\left(\frac{G_E^{N2} + \tau G_M^{N2}}{1 + \tau} \right) + 2\tau G_M^{N2} \tan^2 \frac{\theta}{2} \right] \quad \tau = \frac{Q^2}{4M_N^2}$$

$$\left(\frac{d\sigma}{d\Omega}\right)_{Mott} \frac{\tau G_M^{N2} + \epsilon G_E^{N2}}{\epsilon(1 + \tau)} \quad \epsilon = \left(1 + 2(1 + \tau) \tan^2 \frac{\theta}{2} \right)^{-1}$$



Form Factor Ratio $\mu_p G_E^p / G_M^p$ - Rosenbluth Technique



Measuring Form Factors - Polarized Techniques

Advent of polarized beams and targets provided another technique

In polarization transfer experiments $\vec{e}p \rightarrow e\vec{p}$

$$\mu_p \frac{G_E}{G_M} = -\mu_p \sqrt{\frac{\tau(1+\epsilon)}{2\epsilon}} \frac{P_T}{P_L} = -\mu_p \frac{E+E'}{2M_p} \tan \frac{\theta_e}{2} \frac{P_T}{P_L}$$

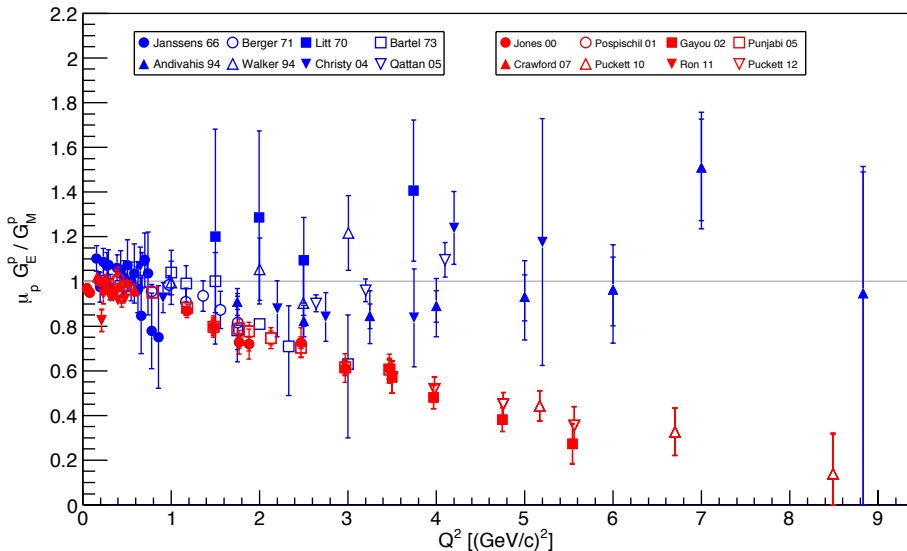
where P_T and P_L are the polarizations of the recoil proton.

This is a simpler and more accurate measurement for $\mu_p G_E/G_M$ particularly at higher Q^2

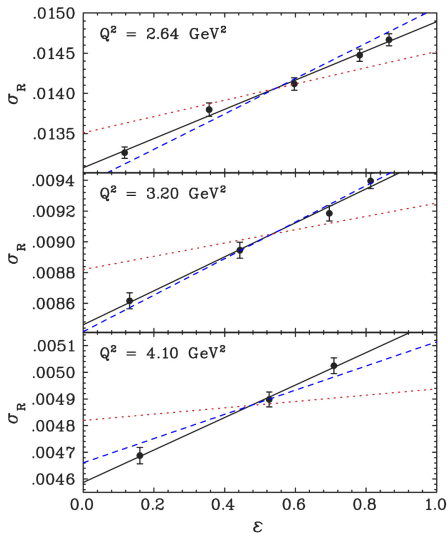
It is also possible to determine $\mu_p G_E/G_M$ from $\vec{e}\vec{p} \rightarrow ep$ by measuring the asymmetries (see Crawford 07).



Discrepancy in Form Factor Ratio $\mu_p G_E^p / G_M^p$?



Measuring Form Factors - Rosenbluth Technique



$$\begin{aligned}\sigma_R &= \epsilon(1 + \tau) \left(\frac{d\sigma}{d\Omega} \right) / \left(\frac{d\sigma}{d\Omega} \right)_{Mott} \\ &= \tau G_M^N{}^2 + \epsilon G_E^N{}^2\end{aligned}$$

Vary E and θ to measure σ_R at different ϵ but same Q^2 and plot:

- Slope $\rightarrow G_E^N{}^2$
- Intercept $\rightarrow G_M^N{}^2$
- G_M^N dominates at high Q^2
- σ_R decreases quickly with Q^2

Blue dashed \rightarrow FF ratio = 1

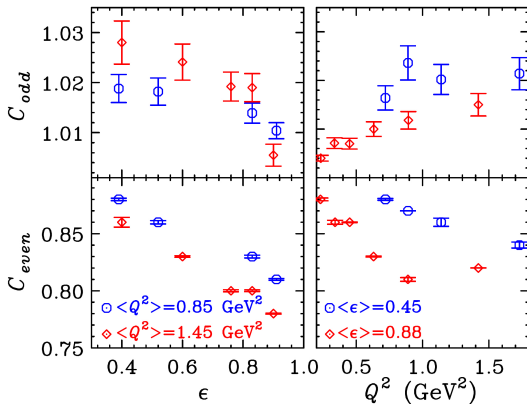
Red dotted \rightarrow polarized measure

CLAS Radiative Corrections

Calculated following
R. Ent *et al.*

$$\frac{1 + \delta_{\text{even}} - \delta_{2\gamma} - \delta_{\text{epbrem}}}{1 + \delta_{\text{even}} + \delta_{2\gamma} + \delta_{\text{epbrem}}}$$

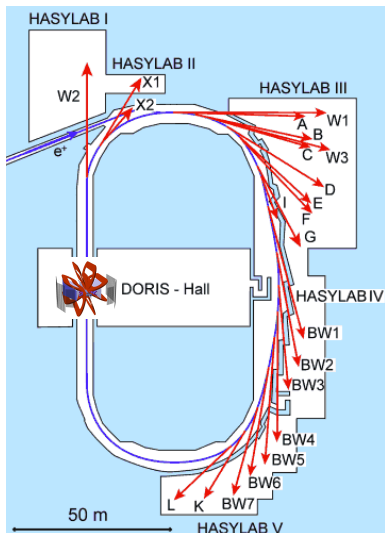
$$R_{2\gamma} = 1 - 2\delta_{2\gamma}$$



D. Rimal *et al.* Phys. Rev. **C95** 065291 (2017).



DORIS Storage Ring at DESY, Hamburg, Germany



Extensive modifications to DORIS

- move RF cavities, ARGUS
- provide cooling water, power
- open pit, move shielding walls
- optics, synchrotron radiation
- automated polarity switches

Great support from DESY !

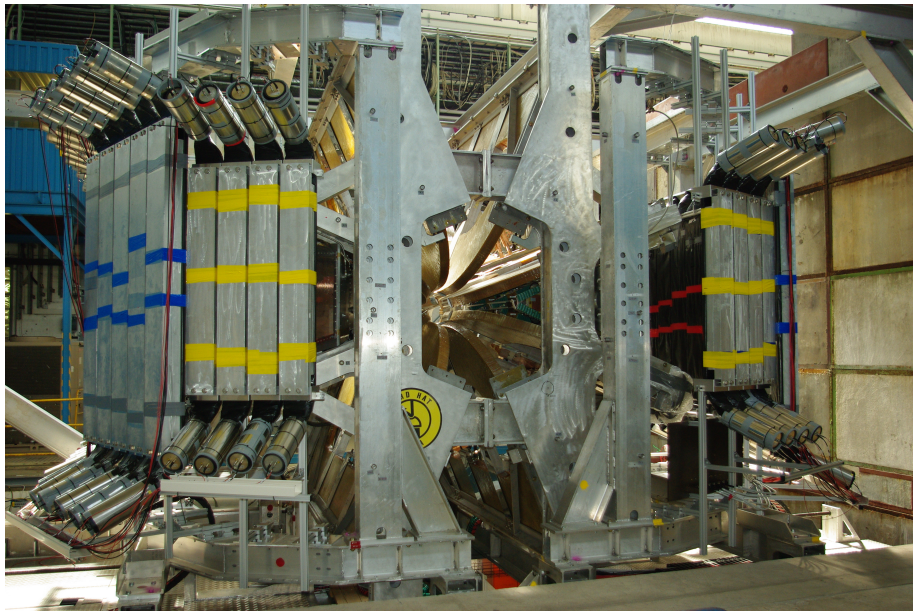
- MEA, MKK, DORIS operators
- Jan Hausschildt, Frank Brinker

Tight schedule shutdown end 2012

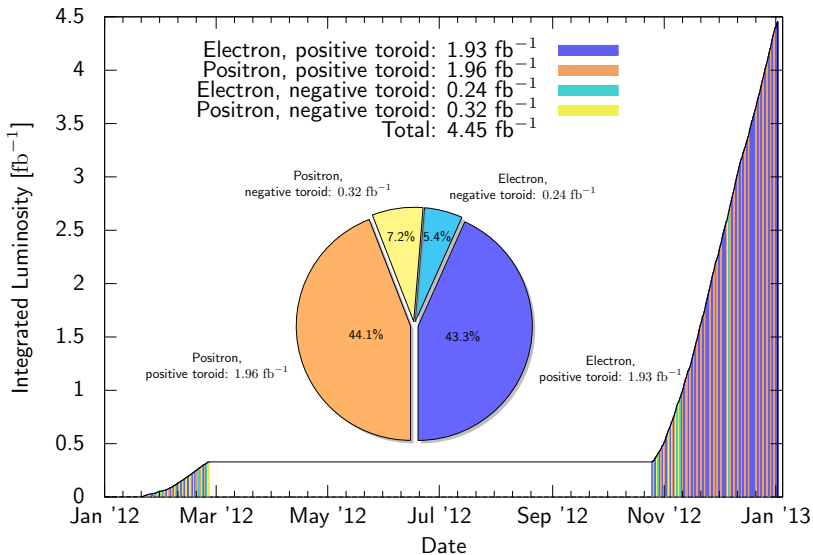
OLYMPUS funded end 2009 !



OLYMPUS Detector



Luminosity



Luminosity

Three independent and consistent measures of luminosity:

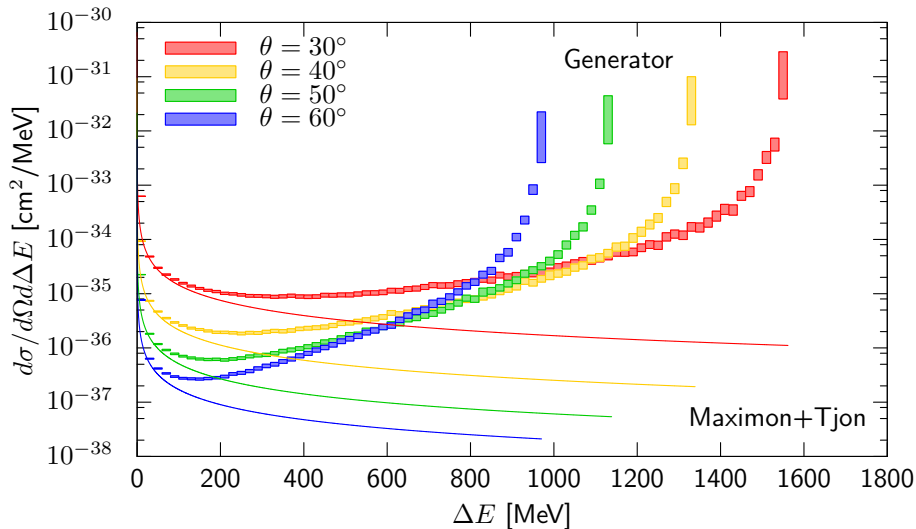
- slow control using molecular flow calculation
 - 2 % between beam species, 5 % absolute
- 12° MWPC with coincident proton in WC
 - 0.46 % between beam species, 2.4 % absolute
- multi-interaction events $(e^\pm e \rightarrow e^\pm e) + (e^\pm p \rightarrow e^\pm p)$ in SYMB
 - 0.1 % statistical, 0.36 % systematic

Chose to use multi-interaction events, MIE, as the most accurate:

- negligible TPE at 1.29°
 - $\langle Q^2 \rangle = 0.002 \text{ GeV}^2$, $\langle \epsilon \rangle = 0.99975$
- allows additional measurement of TPE at 12°
 - $R_{2\gamma} = 0.9975 \pm 0.010 \pm 0.0053$
 - $\langle Q^2 \rangle = 0.165 \text{ GeV}^2$, $\langle \epsilon \rangle = 0.98$



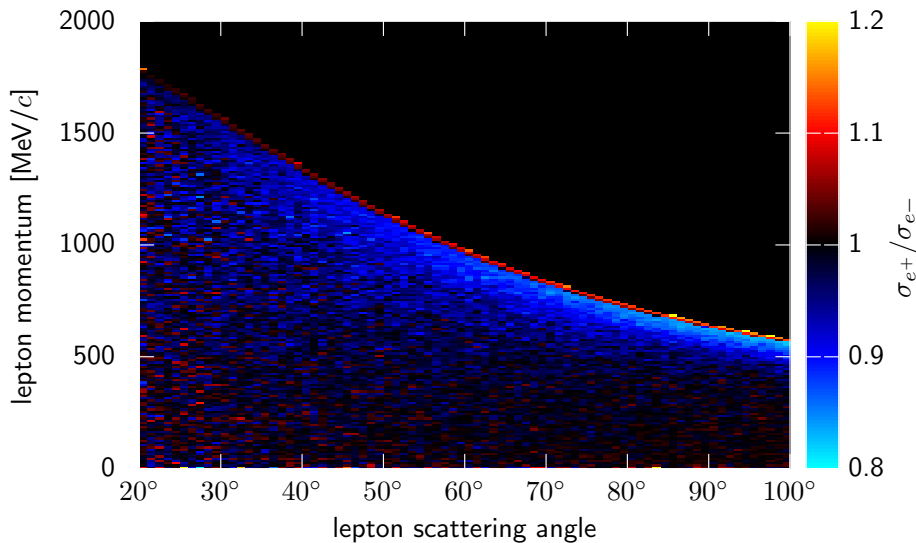
Radiative Corrections and Initial State Radiation



Jan Bernauer, Rebecca Russell, and Axel Schmidt, MIT



Radiative Corrections



Analysis Procedure

All analyses share the following:

- based on the same run list and same tracked data files
- use the same tracked, radiatively generated, MC files
 - based on the same detector calibration, simulation, and digitization
- results normalized with MIE and binned in the same Q^2 and ϵ bins

Analyses are independent in the following:

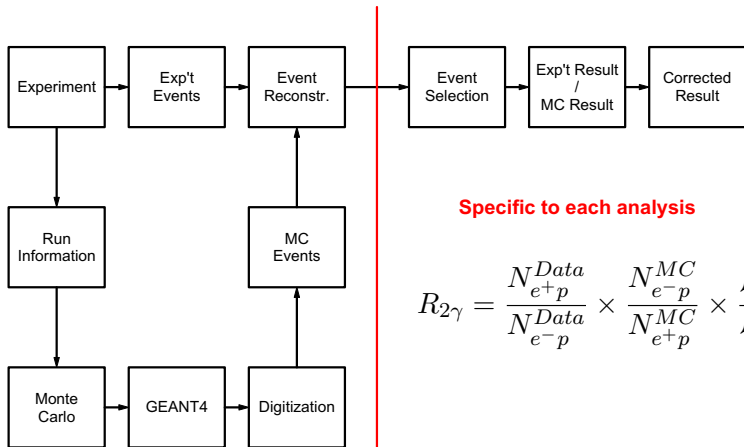
- philosophy in selecting elastic candidates vary
- selection and size of applying cuts are different

Four analyses combined for final result (Axel, Rebecca, Brian, and Jan)

- results simply averaged
- variance added to uncorrelated uncertainty in quadrature



Schematic of Analysis Procedure



Common to all analyses

Specific to each analysis

$$R_{2\gamma} = \frac{N_{e^+p}^{Data}}{N_{e^-p}^{Data}} \times \frac{N_{e^-p}^{MC}}{N_{e^+p}^{MC}} \times \frac{\mathcal{L}_{e^+p}}{\mathcal{L}_{e^-p}}$$

Systematic Uncertainties

OLYMPUS control of systematics

- left / right symmetric detector \rightarrow two independent measurements
- $R_{2\gamma}$ is a ratio so many efficiencies cancel
- four independent analyses that can be examined and combined

Correlated systematic uncertainties

- luminosity (MIE) - 0.36%
- beam energy - 0.04%–0.13%
- beam and detector geometry - 0.25%
- **total - 0.46%**

Uncorrelated systematic uncertainties

- track efficiency - 0.25%
- event selection and background subtraction - 0.25%–1.17%
- **total - 0.37%–1.20%**



Timeline

2005

- May - BLAST Experiment ends
- November - BLAST@ELSA, @DORIS

2007

- May - seminars DESY, Zeuthen, and PRC
- June - Letter of Intent

2008

- September - OLYMPUS proposal
- December - cond. approval DESY + PRC

2009

- August - Technical Design Report
- September - technical review

2010

- January - approval and funding
- February - disassemble BLAST and ship
- July - start modifications and assembly

2011

- January - install target and test
- February - ring run tests
- July - roll into DORIS ring
- August–December - service day test runs

2012

- February - first data run
- July - repair target, other improvements
- October - December - second data run

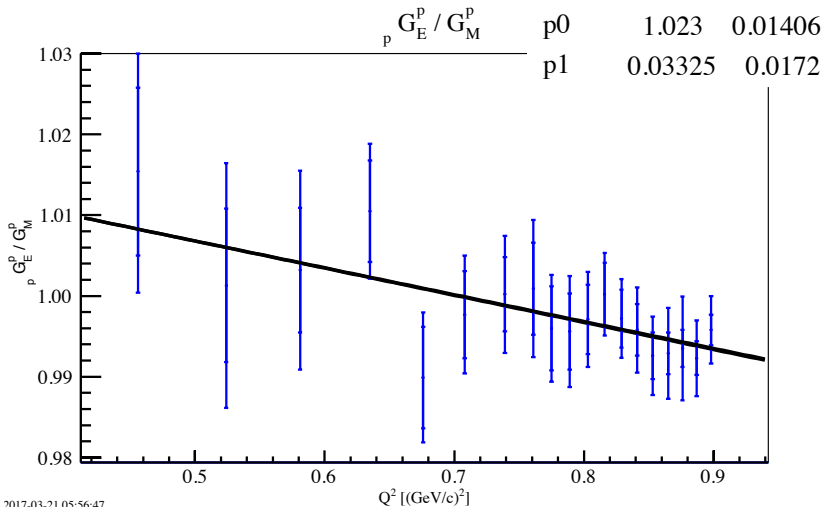
2013

- January - collected cosmic data
- February–May - optical survey, field map
- June–July - disassemble OLYMPUS

2016

- October - most of the analysis complete
- 7 PhD's



Fit to OLYMPUS $R_{2\gamma}$ 

Fit to Rebinned OLYMPUS $R_{2\gamma}$ 