

#### EUROPEAN CENTRE FOR THEORETICAL STUDIES IN NUCLEAR PHYSICS AND RELATED AREAS TRENTO, ITALY

Institutional Member of the European Expert Committee NUPECC





How to test the mass decomposition?

#### **Jianwei Qiu**

Castello di Trento ("Tris Theory Center, palefferson Labay back from Venice (1495). British Museum, London

The Proton Mass: At the Heart of Most Visible Matter

Trento, April 3 - 7, 2017

### **Proton Mass**

□ Nucleon mass – dominates the mass of visible world:



Higgs mechanism is not enough!!!

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□ How does QCD generate the nucleon mass?

"... The vast majority of the nucleon's mass is due to quantum fluctuations of quark-antiquark pairs, the gluons, and the energy associated with quarks moving around at close to the speed of light. ..."

REACHING FOR THE HORIZON The 2015 Long Range Plan for Nuclear Science

### **Proton Mass**

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How to quantify and verify this, theoretically and experimentally?

### Mass vs. Spin

□ Mass – intrinsic to a particle:

= Energy of the particle when it is at the rest

 $\diamond\,$  QCD energy-momentum tensor in terms of quarks and gluons

$$T^{\mu\nu} = \frac{1}{2} \,\overline{\psi} i \vec{D}^{(\mu} \gamma^{\nu)} \psi + \frac{1}{4} \,g^{\mu\nu} F^2 - F^{\mu\alpha} F^{\nu}{}_{\alpha}$$

♦ Proton mass:

$$m = \frac{\langle p | \int d^3 x \, T^{00} | p \rangle}{\langle p | p \rangle} \sim \text{GeV}$$

X. Ji, PRL (1995)

when proton is at rest!

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❑ Spin – intrinsic to a particle:

= Angular momentum of the particle when it is at the rest

QCD angular momentum density in terms of energy-momentum tensor

$$M^{\alpha\mu\nu} = T^{\alpha\nu}x^{\mu} - T^{\alpha\mu}x^{\nu} \qquad \qquad J^{i} = \frac{1}{2}\epsilon^{ijk}\int d^{3}x M^{0jk}$$

♦ Proton spin:

$$S(\mu) = \sum_{z} \langle P, S | \hat{J}_{f}^{z}(\mu) | P, S \rangle = \frac{1}{2}$$

**Proton's spin:** 





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**Proton's spin:** 



#### **Current understanding:**



If we do not understand proton spin, we do not understand QCD

### □ Proton's mass: $\diamond$ QCD Lagrangian does not have mass dimension parameters, other than current quark masses $\diamond$ Asymptotic freedom $\leftarrow$ confinement: $\rightarrow$ A dynamical scale, $\Lambda_{QCD}$ , consistent with $\frac{1}{R} \sim 200 \text{ MeV}$







□ Nucleon mass from lattice QCD:

Martin Savage @ Temple meeting



 $M_N = 800 \text{ MeV} + m_{\pi}$ 

Unexpected behavior !!

See Richards' talk on excite states □ From Lattice QCD calculation: • • •  $H_{s}^{*}$  $B_c^*$  $H_{s}$  $B_c$  $H^{*}$ H2400 2200 2000 1800 1600 ..... 1400 (MeV) 1200 1000 800 600 <u>\_\_\_\_</u> 400 Input 200 -000 0 N Σ π K  $K^*$ η ω Λ Ξ Ω ρ η'  $\Sigma^*$  $\Xi^*$ φ Δ

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A major success of QCD – is the right theory for the Strong Interaction! How does QCD generate this? The role of quarks vs that of gluons?

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A major success of QCD – is the right theory for the Strong Interaction! How does QCD generate this? The role of quarks vs that of gluons? If we do not understand proton mass, we do not understand QCD

□ Three-pronged approach to explore the origin of hadron mass

- ♦ Lattice QCD
- ♦ Mass decomposition roles of the constituents
- ♦ Model calculation approximated analytical approach

# The Proton Mass

At the heart of most visible matter. Temple University, March 28-29, 2016

https://phys.cst.temple.edu/meziani/proton-mass-workshop-2016/

Philadelphia, Pennsylvania

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#### □ Roles of quarks and gluons?

♦ QCD energy-momentum tensor:

$$\begin{split} T^{\mu\nu} &= \overline{T^{\mu\nu}} + \widehat{T^{\mu\nu}} \\ \text{Traceless term:} & \overline{T^{\mu\nu}} \equiv T^{\mu\nu} - \frac{1}{4} g^{\mu\nu} T^{\alpha}_{\ \alpha} \\ \text{Trace term:} & \widehat{T^{\mu\nu}} \equiv \frac{1}{4} g^{\mu\nu} T^{\alpha}_{\ \alpha} \\ \text{with} & T^{\alpha}_{\ \alpha} = \frac{\beta(g)}{2g} F^{\mu\nu,a} F^{a}_{\mu\nu} + \sum_{\substack{q=u,d,s}} m_q (1+\gamma_m) \overline{\psi}_q \psi_q \\ \text{QCD trace anomaly} & \beta(g) = -(11-2n_f/3) g^3/(4\pi)^2 + \dots \end{split}$$

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♦ Invariant hadron mass (in any frame):

$$\langle p | T^{\mu\nu} | p \rangle \propto p^{\mu} p^{\nu} \qquad \longrightarrow \qquad \langle p | T^{\mu\nu} | p \rangle (g_{\mu\nu}) \propto p^{\mu} p^{\nu} (g_{\mu\nu}) = m^2$$

$$m^2 \propto \langle p | T^{\alpha}_{\ \alpha} | p \rangle \qquad \longrightarrow \qquad \frac{\beta(g)}{2g} \langle p | F^2 | p \rangle$$



At the chiral limit, the entire mass is from gluons!

See Dima's talk on Wednesday

#### □ Sum rules for Proton Mass:

X. Ji, PRL (1995) His talk this afternoon

Sum rules are only useful if individual terms can be measured independently



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Hadron state:

 $|P\rangle$  With the normalization:  $\langle P|P\rangle = (E/M)(2\pi)^3\delta^3(0)$ 

♦ Hamiltonian:

 $H_{\rm QCD} = \int d^3 \vec{x} \, T^{00}(0, \vec{x}) \qquad \langle P | H_{\rm QCD} | P \rangle = (E^2 / M_p) (2\pi)^3 \delta^3(0)$ 

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♦ QCD energy-momentum tensor:

$$T^{\mu\nu} = \overline{T^{\mu\nu}} + \widehat{T^{\mu\nu}} \qquad \langle P|T^{\mu\nu}|P\rangle = P^{\mu}P^{\nu}/M_{p} \qquad \text{No } g^{\mu\nu} \text{ term!}$$

$$\Rightarrow \langle P|\overline{T}^{\mu\nu}|P\rangle = (P^{\mu}P^{\nu} - \frac{1}{4}M_{p}^{2}g^{\mu\nu})/M_{p} \qquad \langle P|\widehat{T}^{\mu\nu}|P\rangle = \frac{1}{4}M_{p}g^{\mu\nu}$$

$$\Rightarrow \frac{\langle P|\int d^{3}x\,\overline{T}^{00}|P\rangle}{\langle P|P\rangle}\Big|_{\text{at rest}} = \frac{3}{4}M_{p} \qquad \frac{\langle P|\int d^{3}x\,\widehat{T}^{00}|P\rangle}{\langle P|P\rangle}\Big|_{\text{at rest}} = \frac{1}{4}M_{p}$$

$$\text{``Traceless'' term} \qquad \text{``Trace'' term}$$

#### □ Identities:

$$\frac{\langle P|\int d^3x \, T^{\alpha}_{\ \alpha}|P\rangle}{\langle P|P\rangle} = 4 \, \frac{\langle P|\int d^3x \, \widehat{T}^{00}|P\rangle}{\langle P|P\rangle} \bigg|_{\text{at rest}} = \frac{4}{3} \, \frac{\langle P|\int d^3x \, \overline{T}^{00}|P\rangle}{\langle P|P\rangle} \bigg|_{\text{at rest}}$$

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#### □ Traceless terms:

$$\begin{split} \overline{T}^{\mu\nu} &= \overline{T}_{q}^{\mu\nu} + \overline{T}_{g}^{\mu\nu} \qquad \overline{T}_{q}^{\mu\nu} = \frac{1}{2} \bar{\psi} i \overleftrightarrow{D}^{(\mu} \gamma^{\nu)} \psi - \frac{1}{4} g^{\mu\nu} \bar{\psi} m \psi, \qquad \overline{T}_{g}^{\mu\nu} = \frac{1}{4} g^{\mu\nu} F^{2} - F^{\mu\alpha} F^{\nu}_{\ \alpha}, \\ & \longrightarrow \quad \langle P | \overline{T}_{q}^{\mu\nu} | P \rangle \equiv a(\mu^{2}) (P^{\mu} P^{\nu} - \frac{1}{4} M_{p}^{2} g^{\mu\nu}) / M_{p} \\ & \langle P | \overline{T}_{g}^{\mu\nu} | P \rangle \equiv [1 - a(\mu^{2})] (P^{\mu} P^{\nu} - \frac{1}{4} M_{p}^{2} g^{\mu\nu}) / M_{P} \end{split}$$

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$$\Rightarrow \quad \langle P | \overline{T}_q^{\mu\nu} | P \rangle \equiv a(\mu^2) (P^{\mu} P^{\nu} - \frac{1}{4} M_p^2 g^{\mu\nu}) / M_p \\ \langle P | \overline{T}_g^{\mu\nu} | P \rangle \equiv [1 - a(\mu^2)] (P^{\mu} P^{\nu} - \frac{1}{4} M_p^2 g^{\mu\nu}) / M_P \end{aligned}$$

$$\text{Let} \quad \mu \to + \quad \nu \to + \\ \Rightarrow \quad a(\mu^2) = \sum_f \int_0^{-1} x [q_f(x, \mu^2) + \overline{q}_f(x, \mu^2)] dx \qquad \text{Total momentum fraction carried} \\ by the quarks - reasonably known! \\ \frac{\langle P | \int d^3 x \, \overline{T}_q^{00} | P \rangle}{\langle P | P \rangle} \bigg|_{\text{at rest}} = a(\mu^2) \, \frac{3}{4} \, M_p \qquad \frac{\langle P | \int d^3 x \, \overline{T}_g^{00} | P \rangle}{\langle P | P \rangle} \bigg|_{\text{at rest}} = [1 - a(\mu^2)] \, \frac{3}{4} \, M_p \end{aligned}$$

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$$\Rightarrow \langle P | \overline{T}_{q}^{\mu\nu} | P \rangle \equiv a(\mu^{2}) (P^{\mu} P^{\nu} - \frac{1}{4} M_{p}^{2} g^{\mu\nu}) / M_{p}$$

$$\langle P | \overline{T}_{g}^{\mu\nu} | P \rangle \equiv [1 - a(\mu^{2})] (P^{\mu} P^{\nu} - \frac{1}{4} M_{p}^{2} g^{\mu\nu}) / M_{P}$$
Let  $\mu \rightarrow + \nu \rightarrow +$ 

$$\Rightarrow a(\mu^{2}) = \sum_{f} \int_{0}^{1} x [q_{f}(x, \mu^{2}) + \overline{q}_{f}(x, \mu^{2})] dx \qquad \text{Total momentum fraction carried by the quarks - reasonably known!}$$

$$\frac{\langle P | \int d^{3}x \overline{T}_{q}^{00} | P \rangle}{\langle P | P \rangle} \Big|_{\text{at rest}} = a(\mu^{2}) \frac{3}{4} M_{p} \qquad \frac{\langle P | \int d^{3}x \overline{T}_{g}^{00} | P \rangle}{\langle P | P \rangle} \Big|_{\text{at rest}} = [1 - a(\mu^{2})] \frac{3}{4} M_{p}$$

$$\frac{\langle P| \int d^3x \, T_m^{00} |P\rangle}{\langle P|P\rangle} \bigg|_{\text{at rest}} = b \, \frac{1}{4} \, M_p \qquad \qquad \frac{\langle P| \int d^3x \, T_a^{00} |P\rangle}{\langle P|P\rangle} \bigg|_{\text{at rest}} = [1-b] \, \frac{1}{4} \, M_p$$

#### □ Roles of quarks and gluons?

♦ Quark energy contribution:

$$H_q = \int d^3 \vec{x} \ \bar{\psi}(-i\mathbf{D} \cdot \alpha)\psi,$$

♦ Gluon energy contribution:

$$H_g = \int d^3 \vec{x} \; \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2)$$

♦ Quark mass contribution:

$$H_m = \int d^3 \vec{x} \; \bar{\psi} m \psi$$

♦ Trace anomaly contribution:

$$H_a = \int d^3 \vec{x} \, \frac{9\alpha_s}{16\pi} \left( \mathbf{E}^2 - \mathbf{B}^2 \right)$$

$$M_q = \left. \frac{\langle P|H_q|P \rangle}{\langle P|P \rangle} \right|_{\text{at rest}} = (a-b) \frac{3}{4} M_p$$

$$M_g = \left. \frac{\langle P | H_g | P \rangle}{\langle P | P \rangle} \right|_{\text{at rest}} = (1-a) \frac{3}{4} M_p$$

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Need to find the "b" independently!

#### **Quark mass contribution – the** *b***-term:**

 $b M_p = \langle P | m_u \overline{u}u + m_d \overline{d}d | P \rangle + \langle P | m_s \overline{s}s | P \rangle + \dots$  (heavy flavors)

**D** The first term –  $\pi N \sigma$ -term:

 $\sigma_{\pi N} = \hat{m} \langle N | \bar{u}u + \bar{d}d | N \rangle$  with  $\hat{m} = (m_u + m_d)/2$ 

Both lattice QCD and phenomenological analyses give  $\pi = 45 - 50 \text{ M}_{\odot}\text{V}$ 

 $\sigma_{\pi N} \sim 45 - 50 \text{ MeV}$ 

See talks by Cloet, Liu, Roberts, Yang, ...

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Both lattice QCD and phenomenological analyses give  $\sigma_{\pi N} \sim 45-50~{\rm MeV}$ 

□ The second term – strange scalar charge:

Light strange quark:

$$\sigma_{KN} = (\hat{m} + m_s) \langle N | \bar{u}u + \bar{d}d + 2\bar{s}s | N \rangle / 4$$

Both lattice QCD and phenomenological analyses give

$$\sigma_{KN} \sim 360 - 400 \text{ MeV}$$

See talks by Cloet, Liu, Roberts, Yang, ...

□ Key to proton mass sum rule:

- ♦ Chiral symmetry breaking
- $\diamond$  Trace anomaly

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#### □ From Ji's original paper:

Mass type	$H_i$	$M_i$	$m_s \rightarrow 0 \; ({\rm MeV})$	$m_s \rightarrow \infty ({\rm MeV})$
Quark energy	$\psi^{\dagger}(-i\mathbf{D}\cdot\boldsymbol{\alpha})\psi$	3(a - b)/4	270	300
Quark mass	$\overline{\psi}m\psi$	b	160	110
Gluon energy	$\frac{1}{2}(\mathbf{E}^2 + \mathbf{B}^2)$	3(1 - a)/4	320	320
Trace anomaly	$\frac{9\tilde{\alpha}_s}{16\pi}$ ( $\mathbf{E}^2 - \mathbf{B}^2$ )	(1 - b)/4	190	210

$$a(\mu^2) = \sum_{f} \int_0^1 x[q_f(x,\mu^2) + \overline{q}_f(x,\mu^2)] dx$$
  
$$bM = \langle P|m_u \overline{u}u + m_d \overline{d}d|P \rangle + \langle P|m_s \overline{s}s|P \rangle$$

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17%

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		-	□ Trace	
$a(\mu^2) = \sum_{n=1}^{\infty} \int_{0}^{\infty} x[q_f(x,\mu^2) + \overline{q}_f(x,\mu^2)] dx$			Anomaly	Quark
$\frac{1}{f} J_0$			20%	Energy
$bM = \langle P   m_u u u \rangle$	$a + m_d a d   P \rangle + \langle P \rangle$		29%	
Lattice QC	<b>)</b> calculation:			
Total quark fr	actional moment			
$a(u^2) \approx 0$	) 55			
$u(\mu^{-}) \sim 0$				
Test the scale	e dependence on	lattice?	🗆 Gluon 🔪	Quark
			Energy	Mass

34%

### **Measurement of Trace Anomaly?**

□ Trace anomaly:

□ Theory background:

 $\diamond$  Recall:

$$m^2 \propto \langle p | T^{lpha}_{\ \ lpha} | p 
angle$$

**Quarkonium-proton interaction** 

$$\frac{\beta(g)}{2g} \langle p|F^2|p\rangle$$

See Kharzeev's talk, ...

Brodsky, Schmit de Teramond '90 Luke, Monohar, Savage '92

### **Measurement of Trace Anomaly?**

### □ Trace anomaly:

 $\Rightarrow \text{ Recall:} \qquad m^2 \propto \langle p | T^{\alpha}_{\ \alpha} | p \rangle \qquad \Longrightarrow \quad \frac{\beta(g)}{2g} \ \langle p | F^2 | p \rangle$ 

### Theory background:

**Quarkonium-proton interaction** 

Brodsky, Schmit de Teramond '90 Luke, Monohar, Savage '92

#### Quarkonium production near threshold:



$$\gamma N \to J/\psi N$$



#### Kharzeev, Satz, Syamtomov Zinovjev, EPJ '99

See Kharzeev's talk, ...

### **Measurement of Trace Anomaly?**

□ Trace anomaly:

See JP Chen's talk, ...

 $\Rightarrow \text{ Recall:} \qquad m^2 \propto \langle p | T^{\alpha}_{\ \alpha} | p \rangle \qquad \Longrightarrow \quad \frac{\beta(g)}{2g} \ \langle p | F^2 | p \rangle$ 

□ Charmonium production at Jlab near the threshold:



Wang, Liu, Zhang arxiv:1508.00339

Kubarovsky, Voloshin arxiv:1508.00888

But, the cross section may be much smaller – O(pb) Gobbi, Boffi, DK '94

# "Summary"

### The proton mass closely connected to quantum anomalies

Non-perturbative QCD generates a new scale:  $\langle 0|F^2|0\rangle \neq 0$ 

#### □ Three-pronged approach to explore the origin hadron mass

lattice QCD mass decomposition – roles of the constituents approximated analytical approach

#### **Questions:**

- What can lattice QCD do to explore the role of "individual" constituent in making up the proton mass?
- $\diamond\,$  What can the mass decomposition teach us?
- How well can we control the approximation of the analytical or model approaches



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∻ ...

### Let's work together at this workshop