

Pion electromagnetic form factor at high Q^2 from lattice QCD

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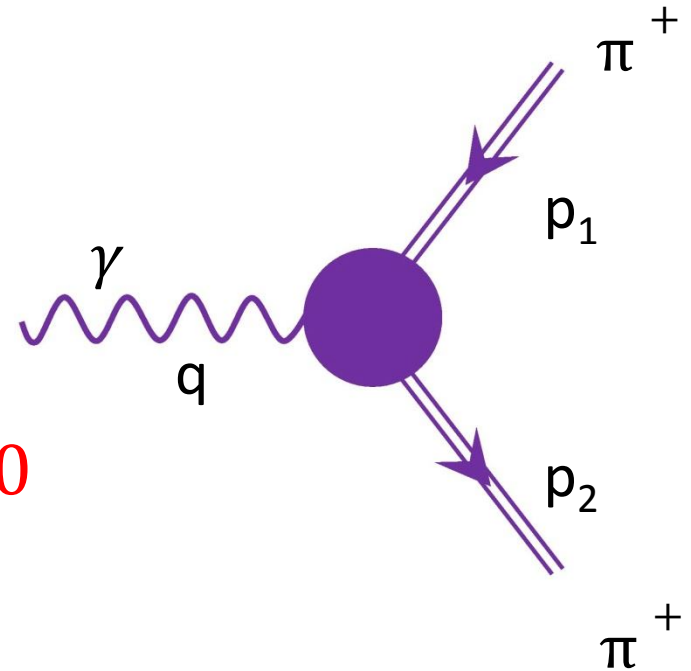
Definition

Simplest hadron

Space like “ q ”:

$$q^2 = (p_2 - p_1)^2 \leq 0$$

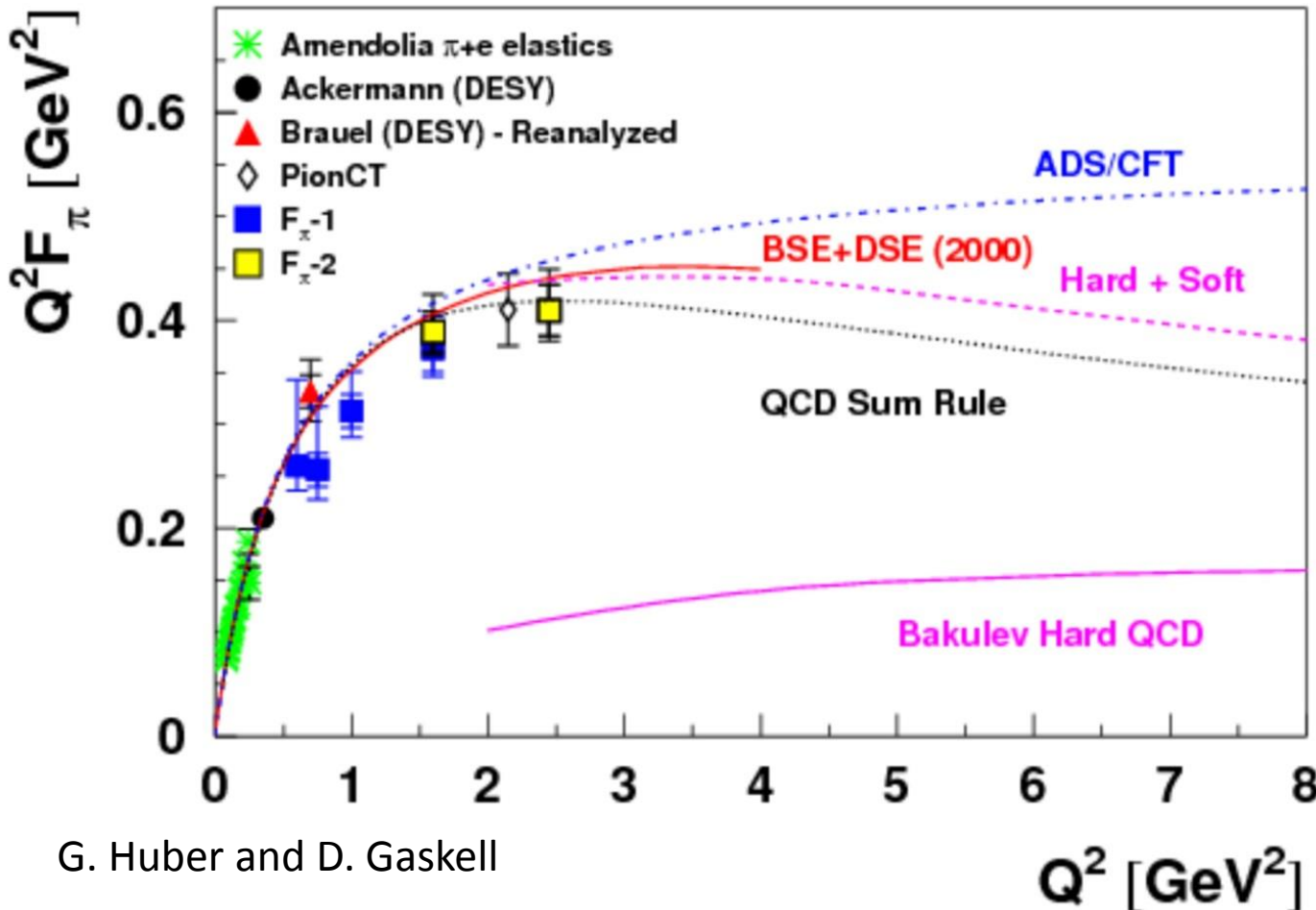
$$Q^2 = -q^2$$



$$\langle \pi^+(\vec{p}') | j^\mu | \pi^+(\vec{p}) \rangle = (p + p')^\mu F_\pi(Q^2)$$

(in units of ‘ e ’)

Interplay between hard and soft scales



G. Huber and D. Gaskell

Hard tail ($Q^2 \rightarrow \infty$)
from pQCD:

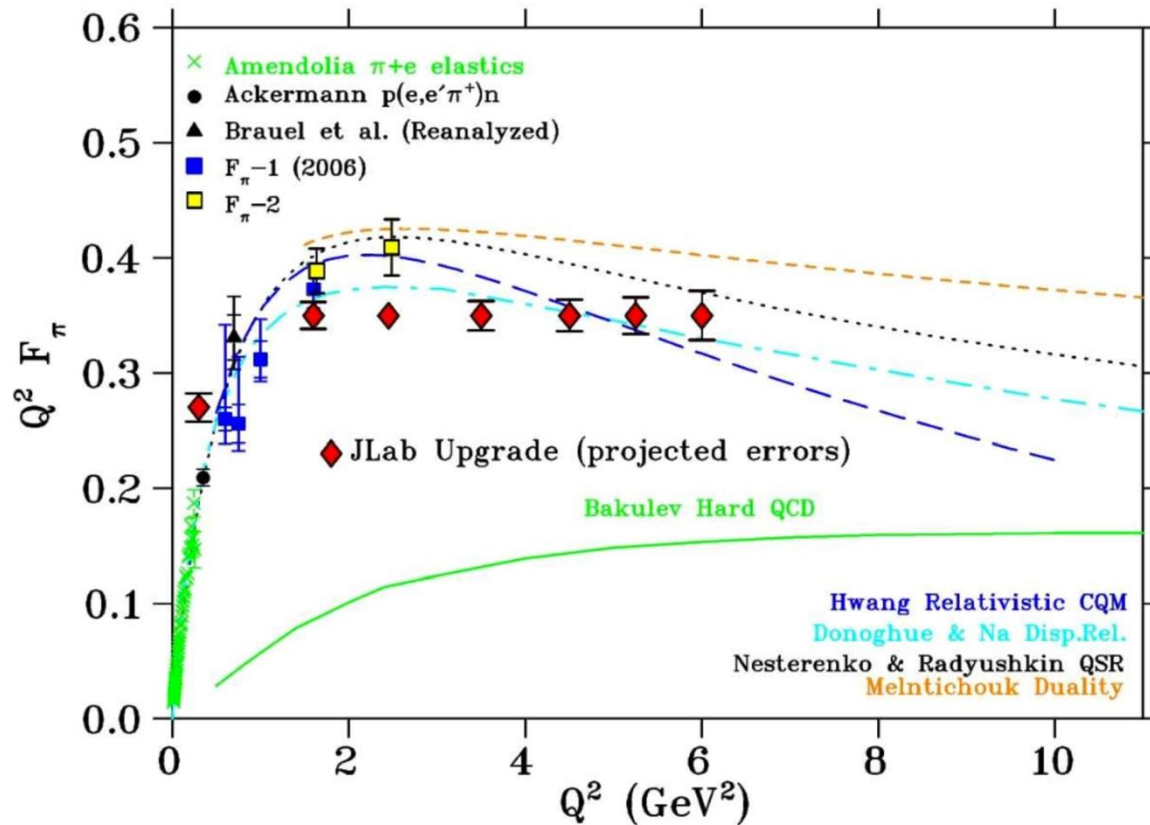
$$F_\pi(Q^2) \rightarrow \frac{16\pi\alpha_s(Q^2)f_\pi^2}{Q^2}$$

G. P. Lepage, S.J. Brodsky,
Phys. Lett. 87B(1979)359

Soft part
($Q^2 < 1 \text{ GeV}^2$):
vector meson
dominance with
 $F_\pi(0) = 1$,
data fits well

Need better understanding of the **transition to the asymptotic region**

JLAB 12 GeV upgrade



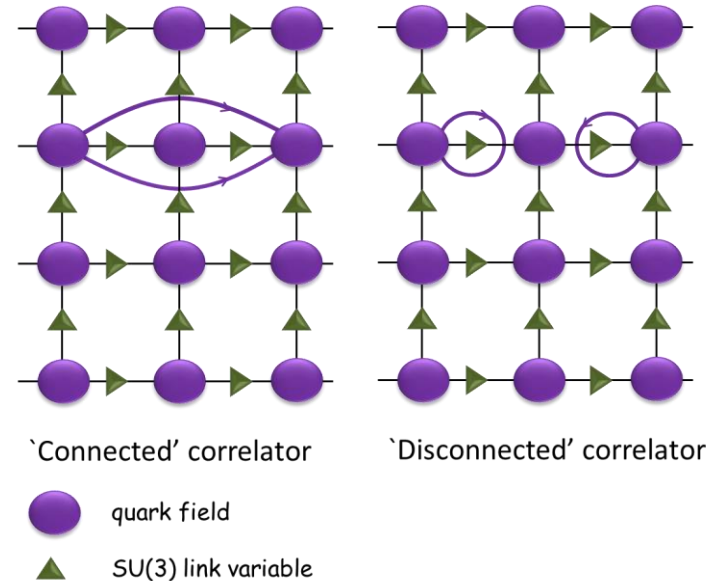
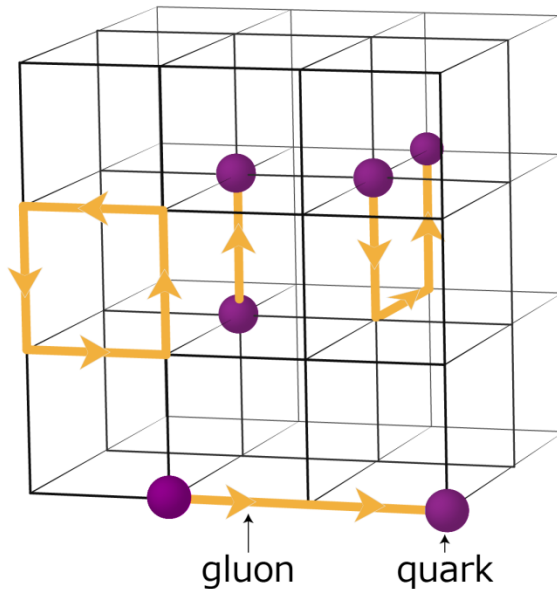
G. Huber and E. Gaskell

F_π measurements
at $Q^2 \sim 6 \text{ GeV}^2$:
E12-06-101 at
JLAB Hall C

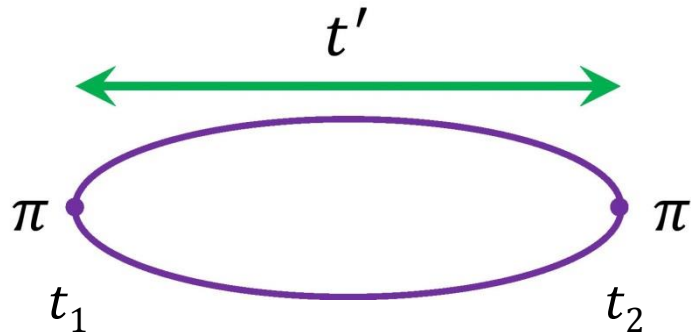
Can we get some insight from first principles lattice QCD calculations to the question - where does the transition to pQCD happen?

Lattice recipe for meson correlators

- Expectation values of observables : $\int DUD\psi D\bar{\psi} \exp(-\int L_{QCD} d^4x)$
- 4-D space-time lattice
- Gauge configurations : gluons + sea quarks



- Discretise : $L_q \equiv \bar{\psi}(\gamma_\mu D^\mu + m)\psi \rightarrow \bar{\psi}(\gamma.\Delta + ma)\psi$
- Inversion of Dirac matrix : propagator
- 2-point, 3-point correlation functions : extract meson properties
- Corrections for lattice artifacts



$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^\dagger(0) | 0 \rangle$$

- Basis of operators

$$\mathcal{O} \sim \bar{\psi} \Gamma \overleftrightarrow{D} \dots \overleftrightarrow{D} \psi$$

- Optimized operator for state $|n\rangle$

$$\Omega_n^\dagger = \sum_i w_i^{(n)} \mathcal{O}_i^\dagger$$

in a variational sense by solving generalized eigenvalue problem-

$$C(t) v^{(n)} = \lambda_n(t) C(t_0) v^{(n)}$$

- Diagonalize the correlation matrix – eigenvalues

$$\lambda_n(t) = \exp[-E_n(t - t_0)]$$

Correlator Construction: smearing of quark fields - 'distillation' with

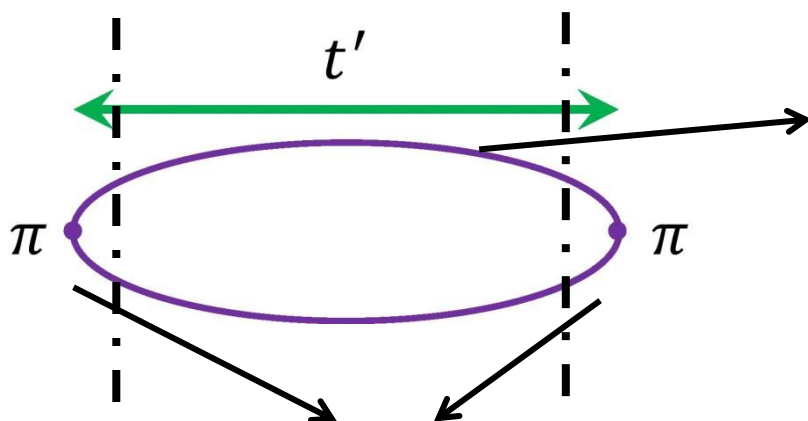
$$\square_{\vec{x}\vec{y}}(t) = \sum_{n=1}^{N_D} \xi_{\vec{x}}^{(n)}(t) \xi_{\vec{y}}^{(n)\dagger}(t)$$



Low lying hadron states

Meson creation operator :

$$\mathcal{O}^\dagger(\vec{p}) = \bar{\psi}_{\vec{x}} \square_{\vec{x}\vec{y}} e^{-i\vec{p}\cdot\vec{y}} \Gamma_{\vec{y}\vec{z}} \square_{\vec{z}\vec{w}} \psi_{\vec{w}}$$

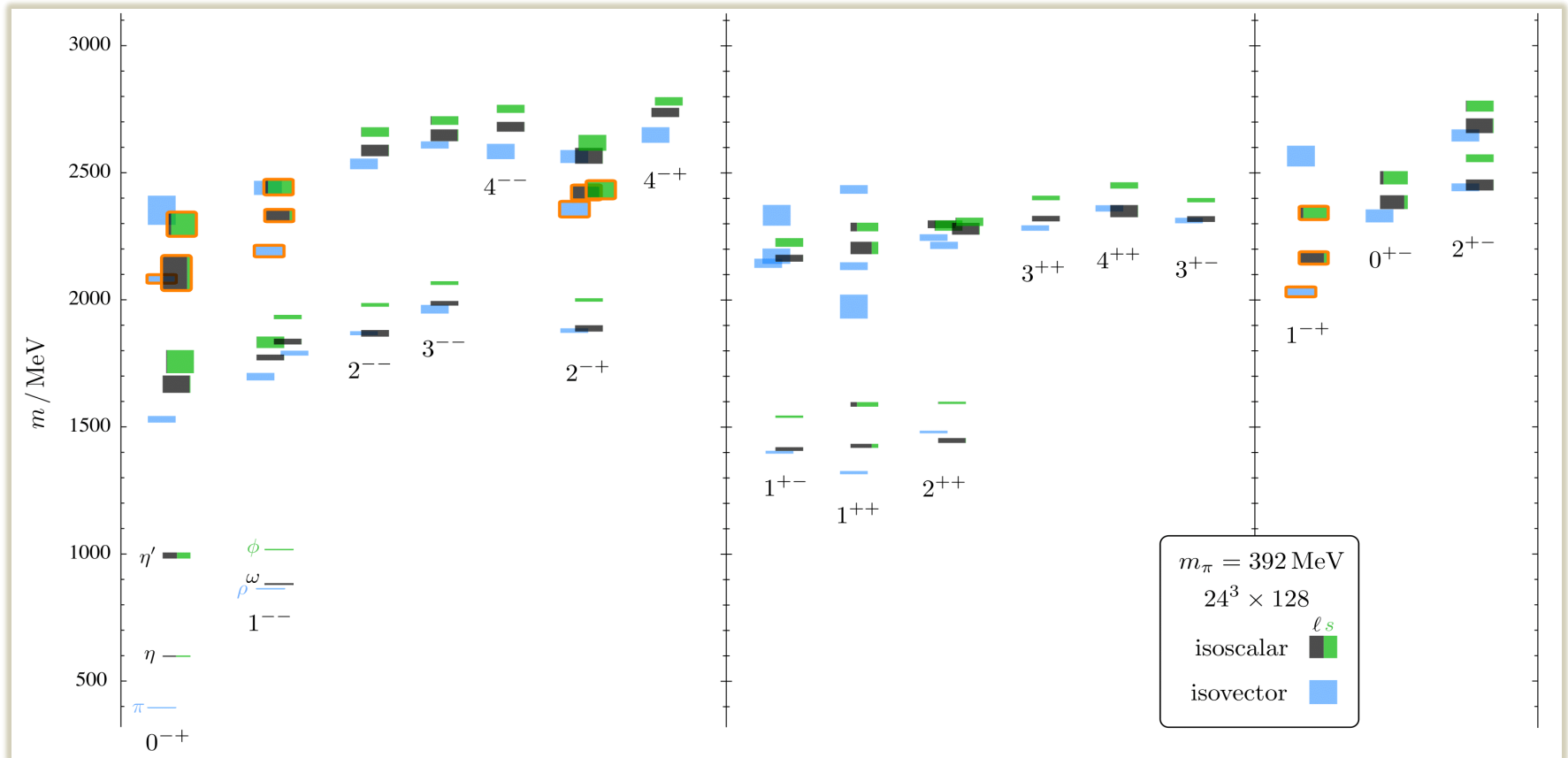


Parambulators by inverting the Dirac matrix

+

Operator construction with momentum projection

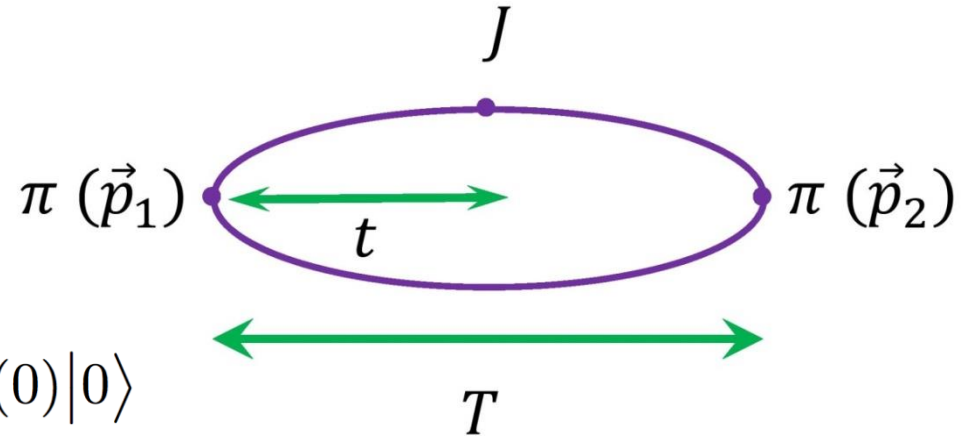
Tools well established for spectroscopy



Hadron Spectrum Collaboration

Jozef J. Dudek *et. al.* Phys.Rev. D88 (2013)

Need three-point correlator



$$C_{f\mu i}(\Delta t, t) = \langle 0 | \mathcal{O}_f(\Delta t) j_\mu(t) \mathcal{O}_i^\dagger(0) | 0 \rangle$$

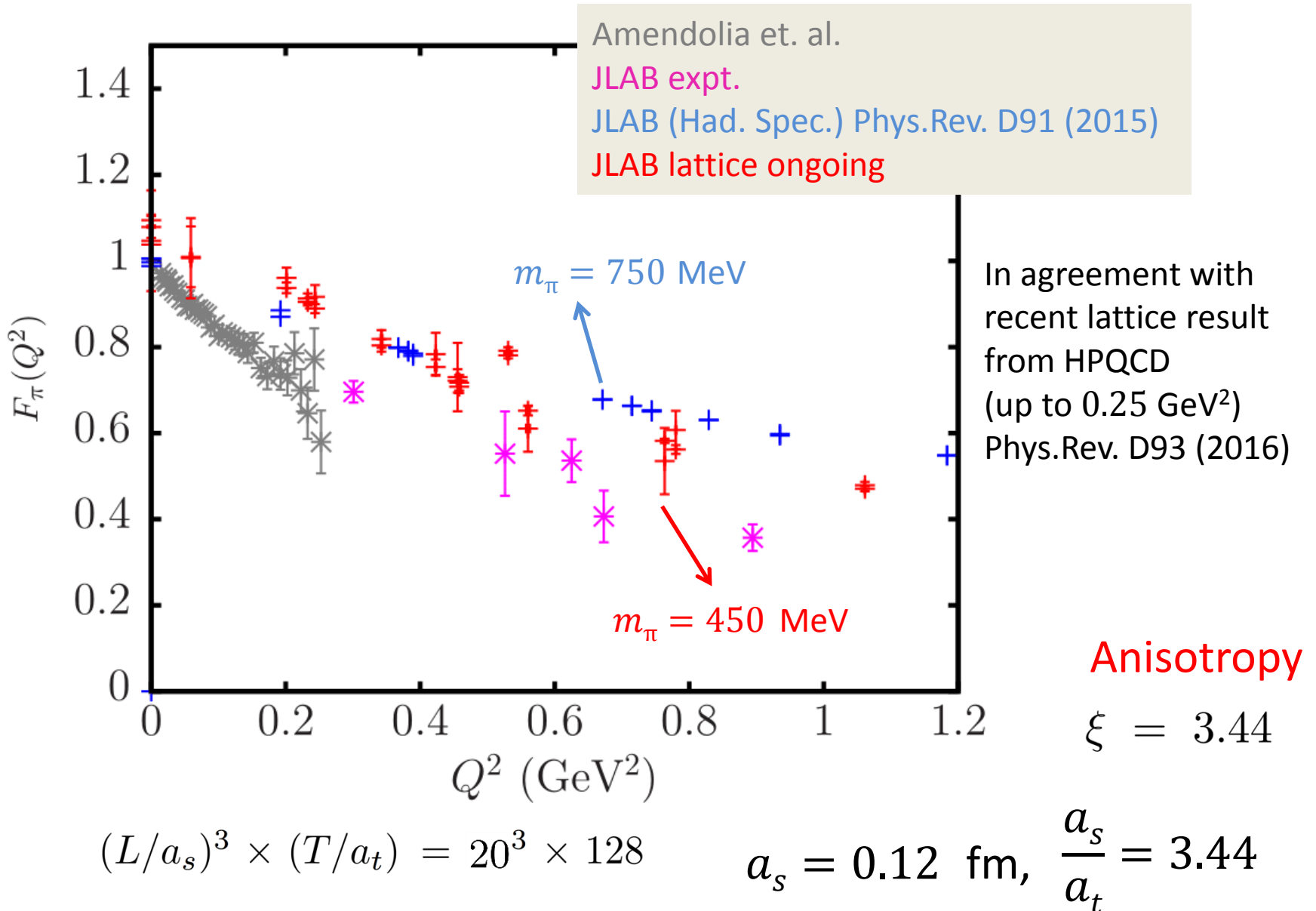


$$Z_V \langle \pi^+(p_2) | J_\mu^\pi(0) | \pi^+(p_1) \rangle = e(p_1 + p_2)^\mu F_\pi(q^2)$$

Clover discretised
fermion action

Z_V calculated using $F_\pi(q^2 = 0) = 1$

Pion electromagnetic form factor: up to $Q^2 = 1 \text{ GeV}^2$



Towards higher Q^2

More difficult on lattice for higher momenta

Signal-to-noise ratio:

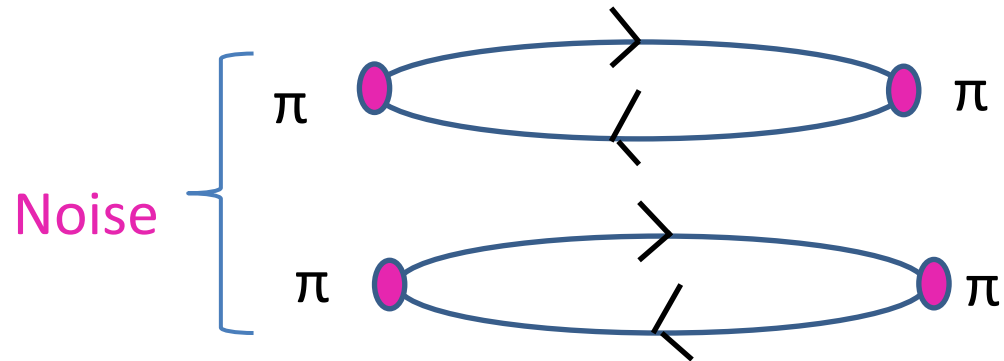
2-point correlators :

$$\exp[-(E_\pi(p) - 2m_\pi)t]$$

3-point correlators :

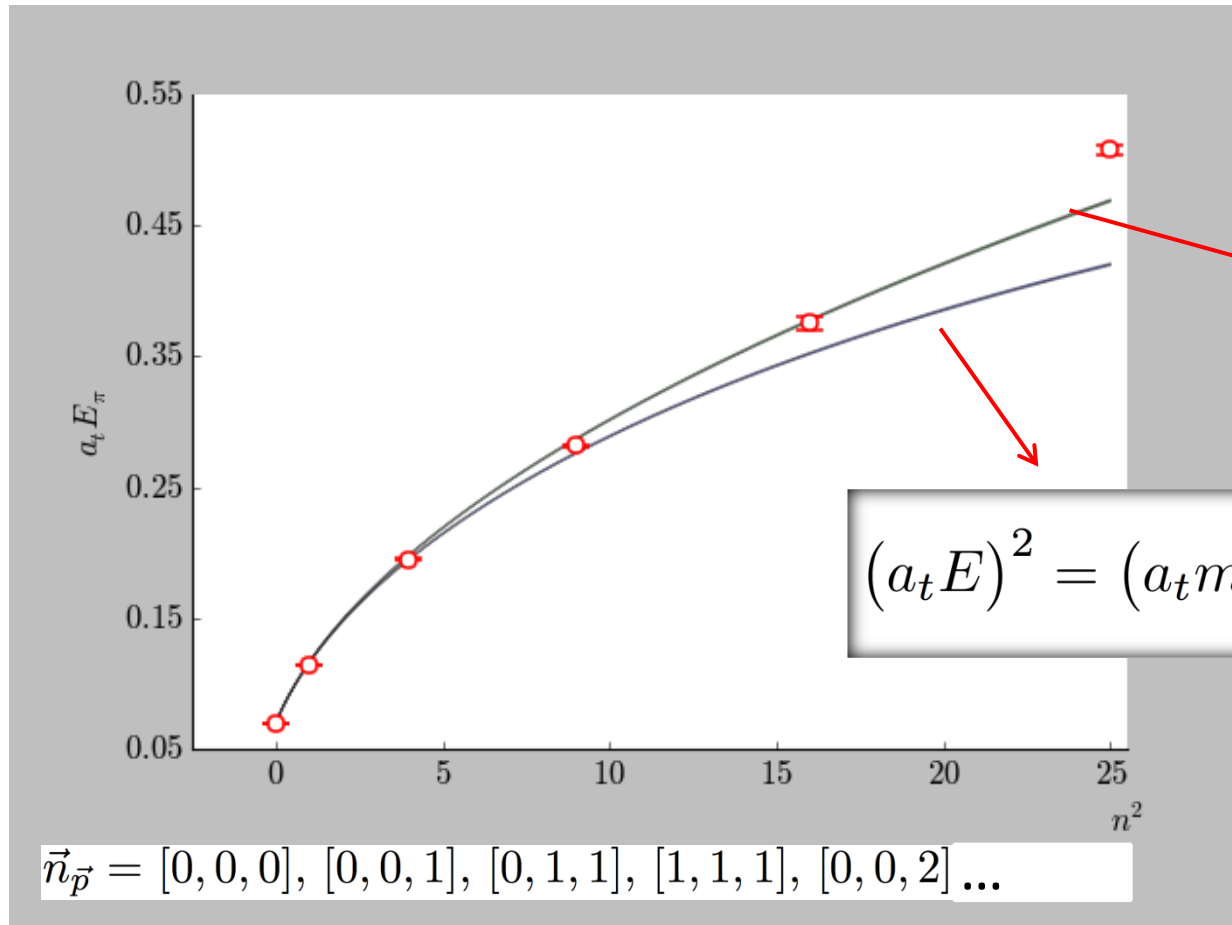
$$\exp[-(E_\pi(pi) + E_\pi(pf) - 2m_\pi)t/2]$$

in the middle of the plateau



Minimize energies
for a given Q^2
to get better signal

Towards higher Q^2



Dispersion relation:

$$E^2 = m^2 + p^2$$



$$(a_t E)^2 = (a_t m)^2 + \left(\frac{2\pi}{\xi(L/a_s)} \right)^2 |\vec{n}_{\vec{p}}|^2$$

Achieve maximum Q^2 by using Breit frame : $\vec{P}_f = -\vec{P}_i$

Outlook

Immediate goals:

- Pion form factor at $Q^2 \geq 6 \text{ GeV}^2$
- Extend to more ensembles with lighter pion masses , multiple volumes, multiple lattice spacings
- Take care of lattice artifacts

Long term goals:

- Hadron structure program – distribution amplitude, PFDs, Quasi PDFs
- Extend to nucleons & more – charges, moments, TMDs, GPDs