

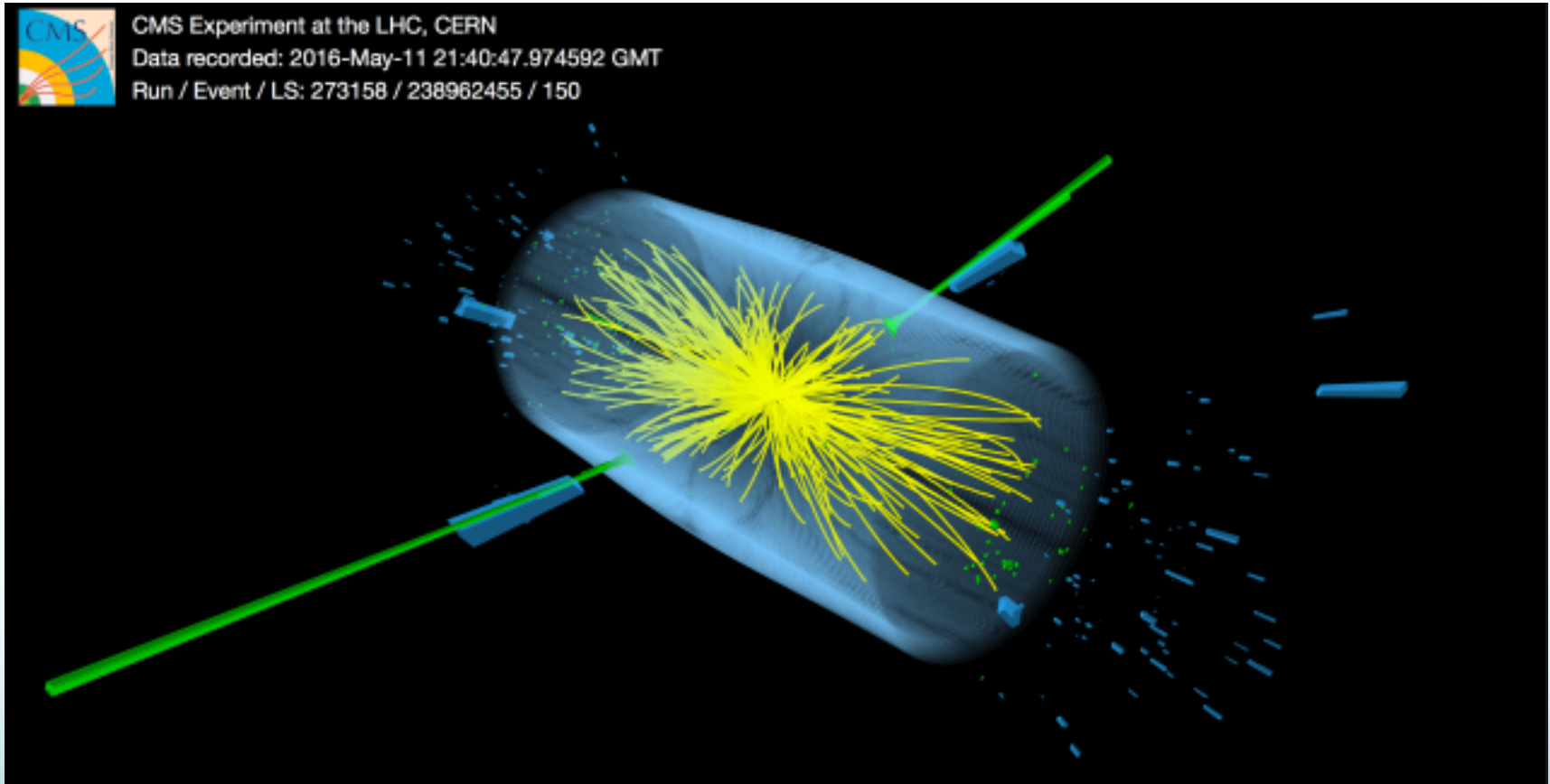
Probing collinear and TMD fragmentation functions through  
hadron distribution inside the jet

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The 7<sup>th</sup> Workshop of the APS Topical Group on  
Hadronic Physics  
February 1 - 3, 2017

# Jets are abundantly produced at the LHC

- They are most common at the LHC



# Jets and its internal substructure as new tools

- Jesse Thaler, 2015

## Jets as a Tool for (B)SM Physics

Importance/relevance of jet radius variation, multiple jet algorithms  
Making jet substructure part of everyday analyses (e.g. pileup mitigation, jet shapes)  
Improved VBF tagging, jet vetoes for Higgs physics

...

## Jets as a Precision Probe of QCD

Wishlist of jet shape measurements (e.g. angularities)  
Interplay between fixed order and resummation for jet observables (esp. PS/ME matching)  
IRC Unsafe but Sudakov Safe observables where resummation is essential  
Analytic handles on soft QCD (e.g. underlying event, hadronization)

...

*Many points of contact with other working groups*

# Hadron distribution inside the jet

- Study a hadron distribution inside a fully reconstructed jet

$$p + p \rightarrow \text{jet} (h) + X$$

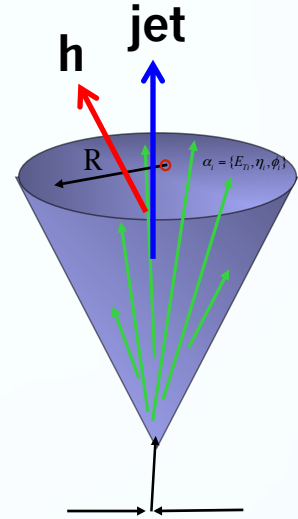
$$F(z_h; p_T) = \frac{d\sigma^h}{dp_T d\eta dz_h} \bigg/ \frac{d\sigma}{dp_T d\eta}$$

$$F(z_h, j_T; p_T) = \frac{d\sigma^h}{dp_T d\eta dz_h d^2 j_T} \bigg/ \frac{d\sigma}{dp_T d\eta}$$

$$z_h = p_T^h / p_T^{\text{jet}}$$

$j_T$  : hadron transverse momentum with respect to the jet direction

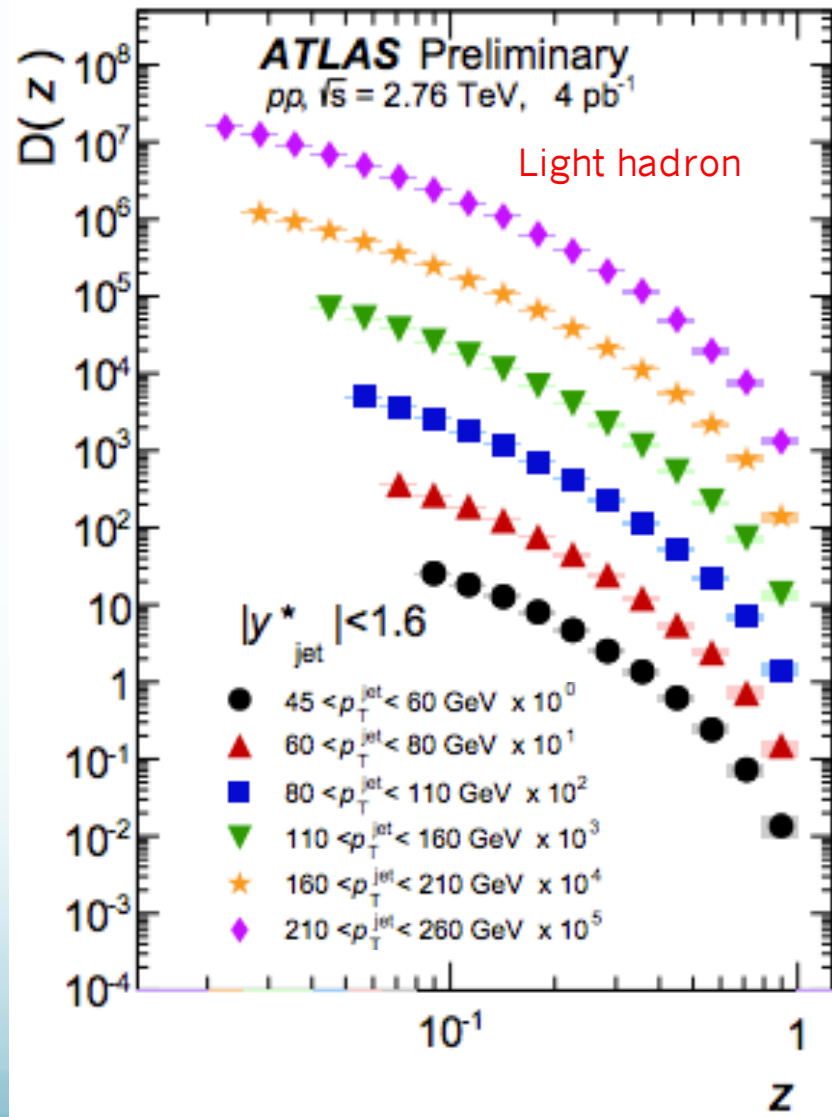
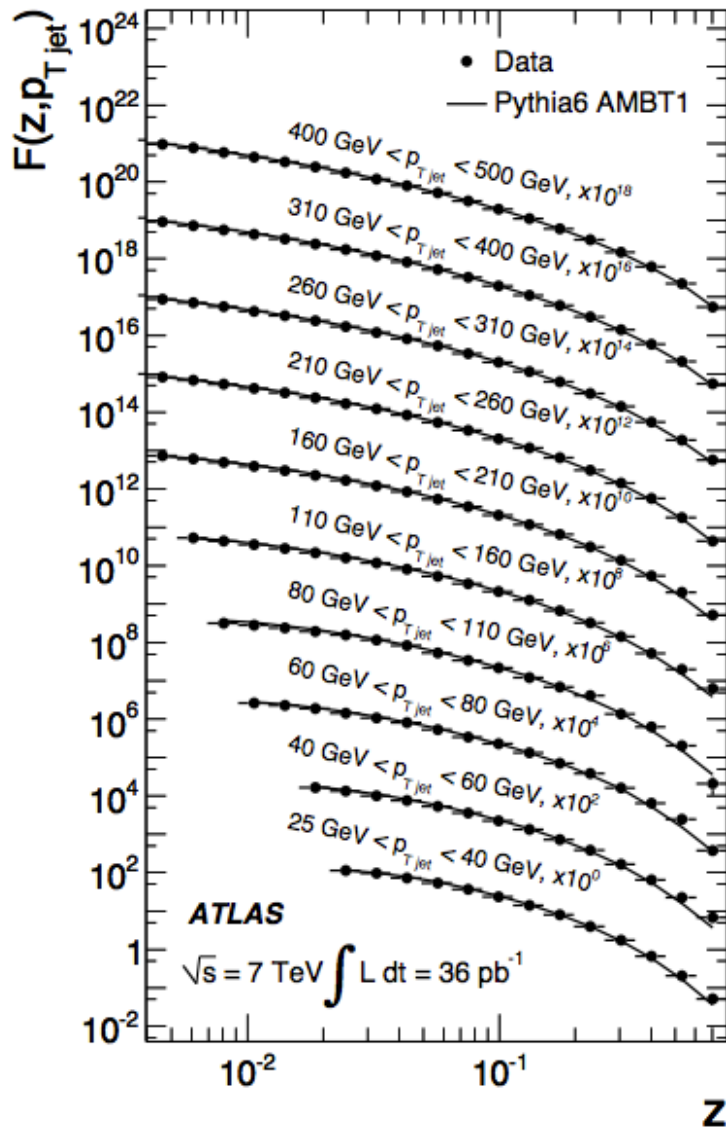
- The 1<sup>st</sup> observable is like collinear fragmentation function, while the 2<sup>nd</sup> observable is more like a TMD fragmentation function
- LHC did a great deal of all kinds of measurements, and compared with Pythia simulation



# Collinear z-dependence: light hadron

- ATLAS measurements at 7 TeV and 2.76 TeV

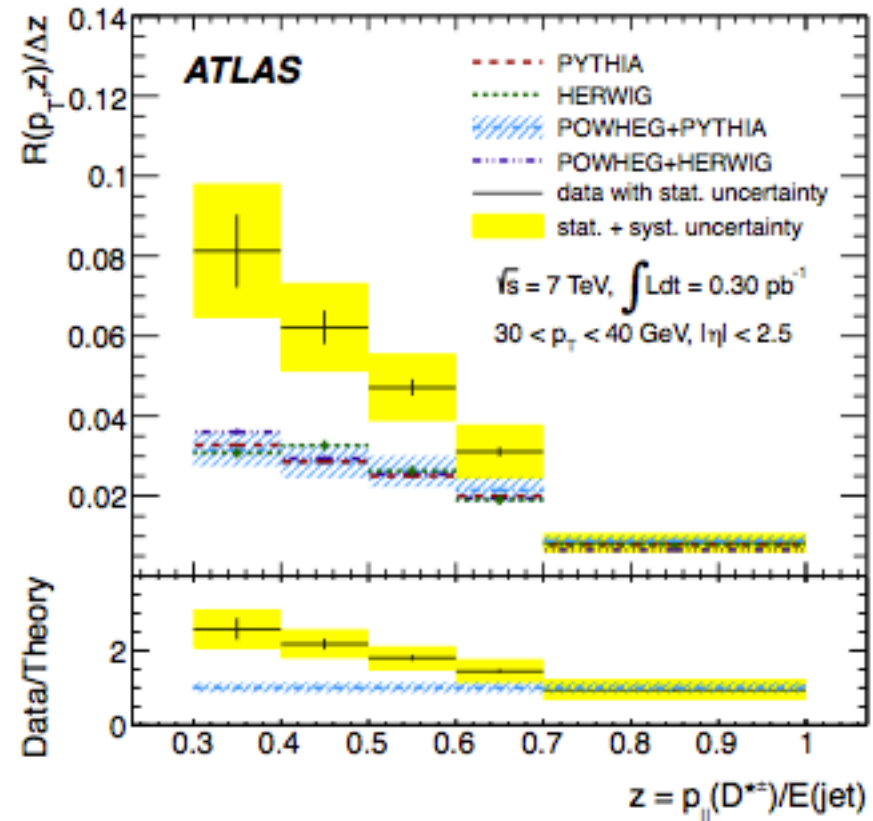
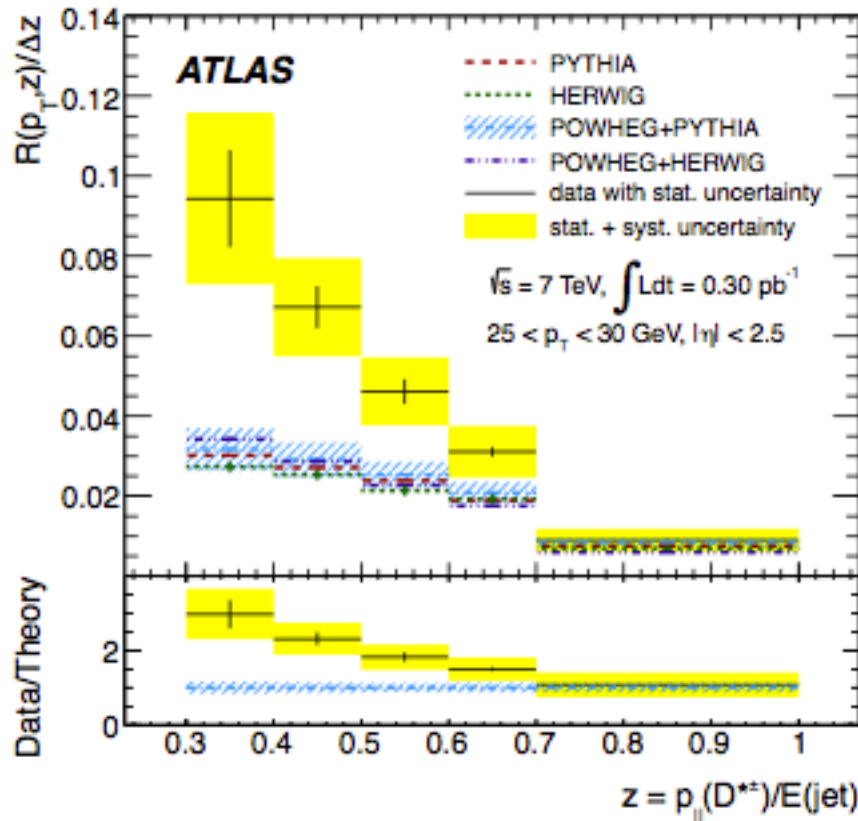
1109.5816, ATLAS-CONF-2015-022



# Collinear z-dependence: heavy meson

- D meson production inside a jet

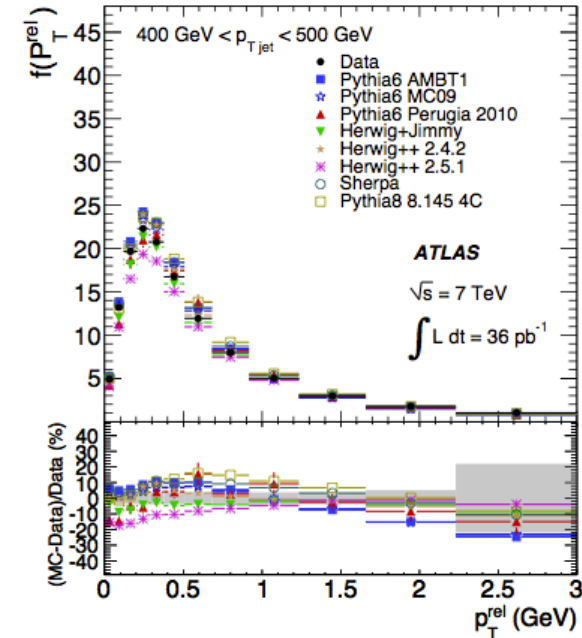
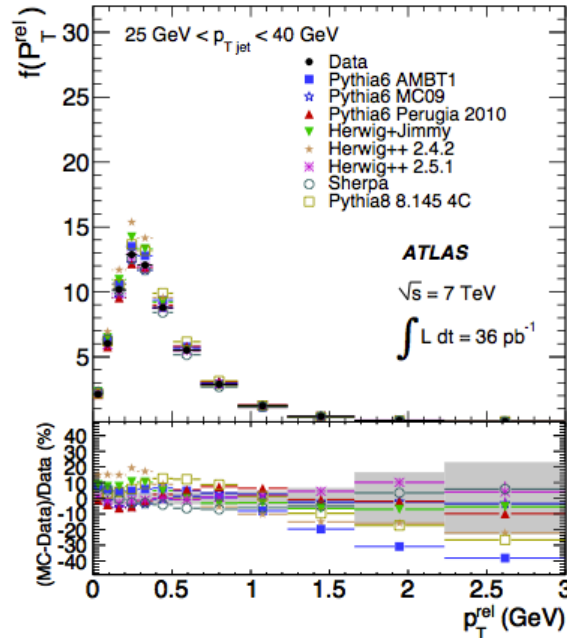
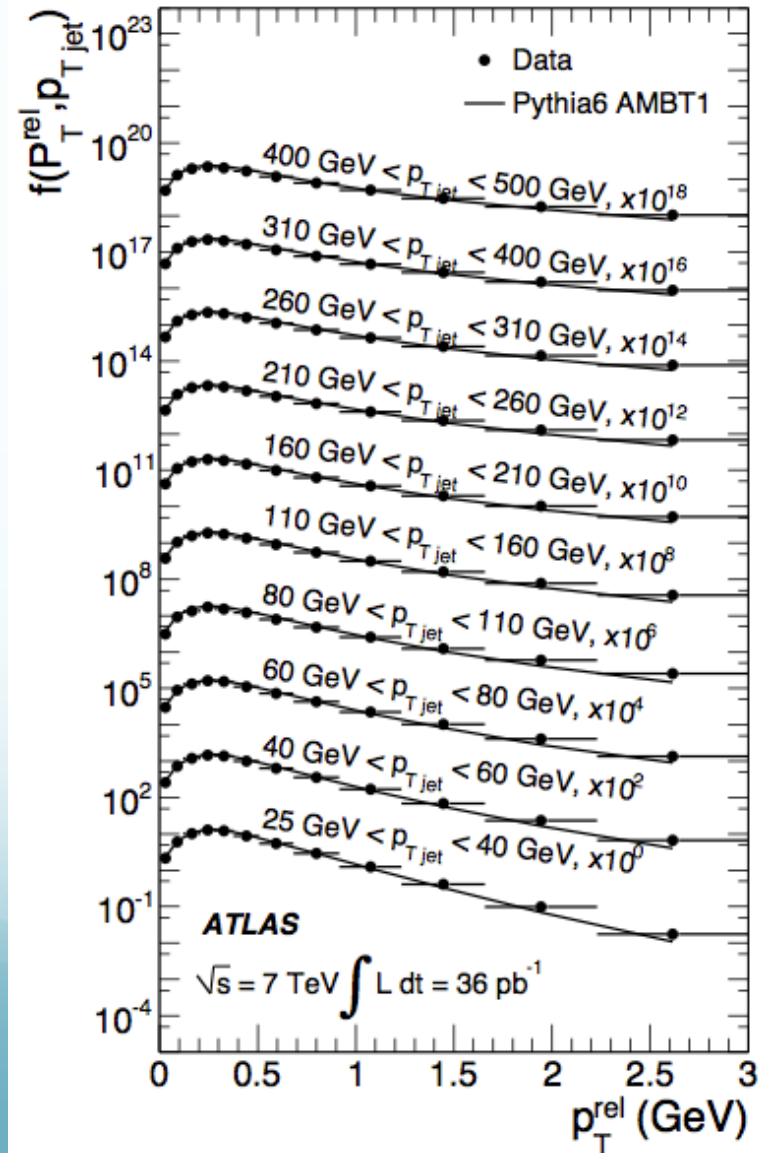
ATLAS, arXiv:1112.4432



# Relative momentum $j_T$ dependence

- $j_T$  shape does not change much: how to link to TMD evolution

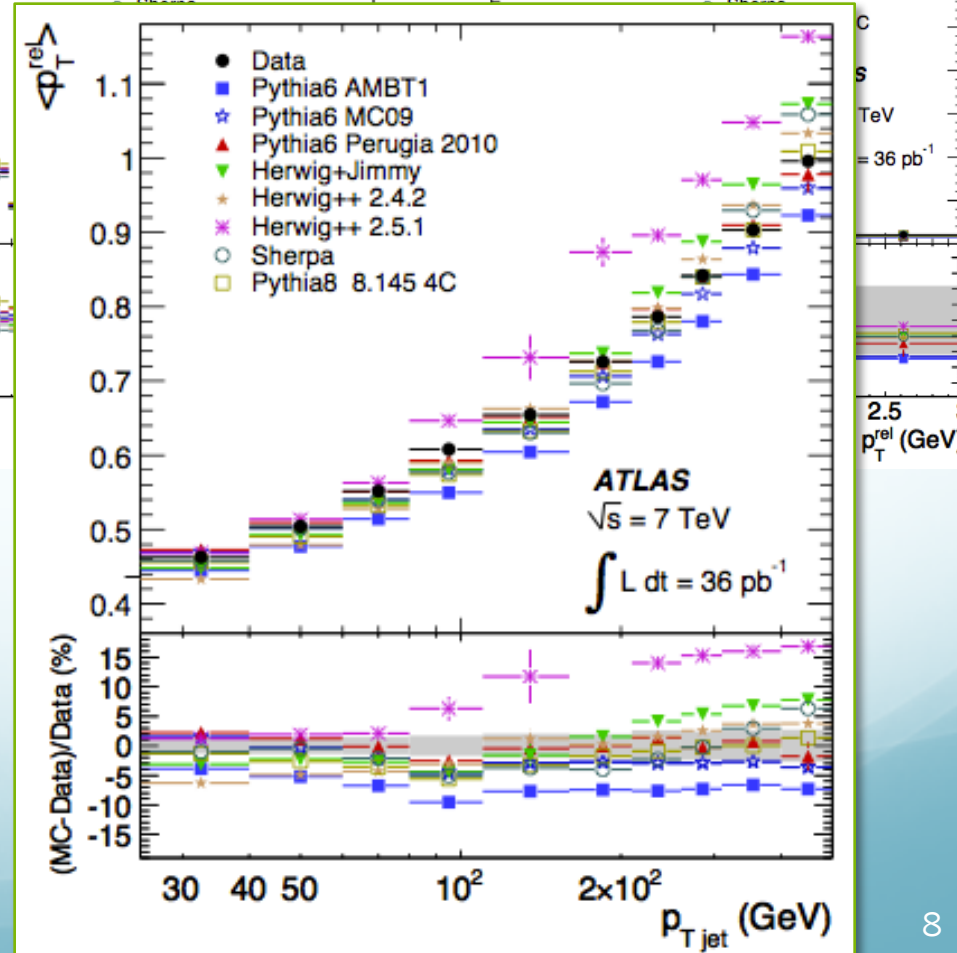
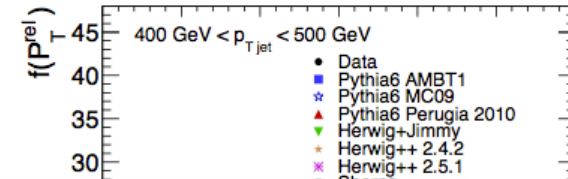
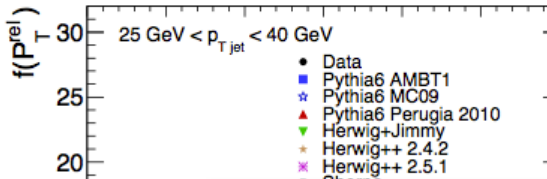
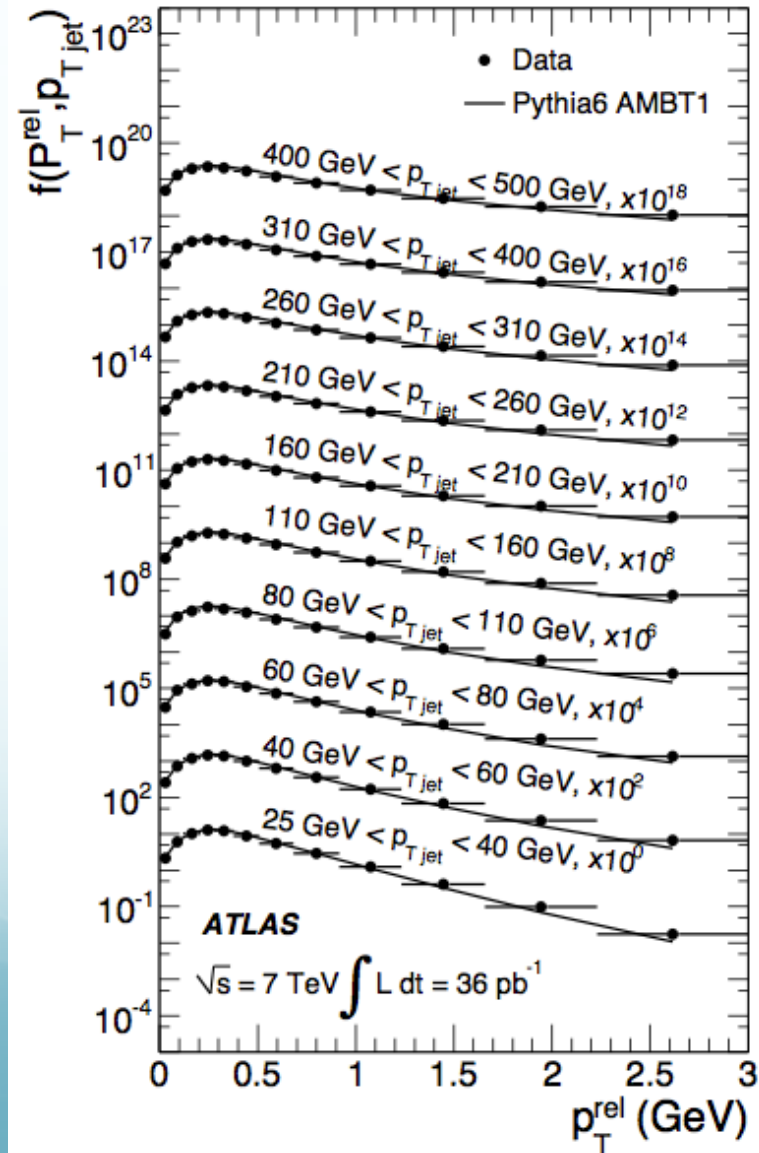
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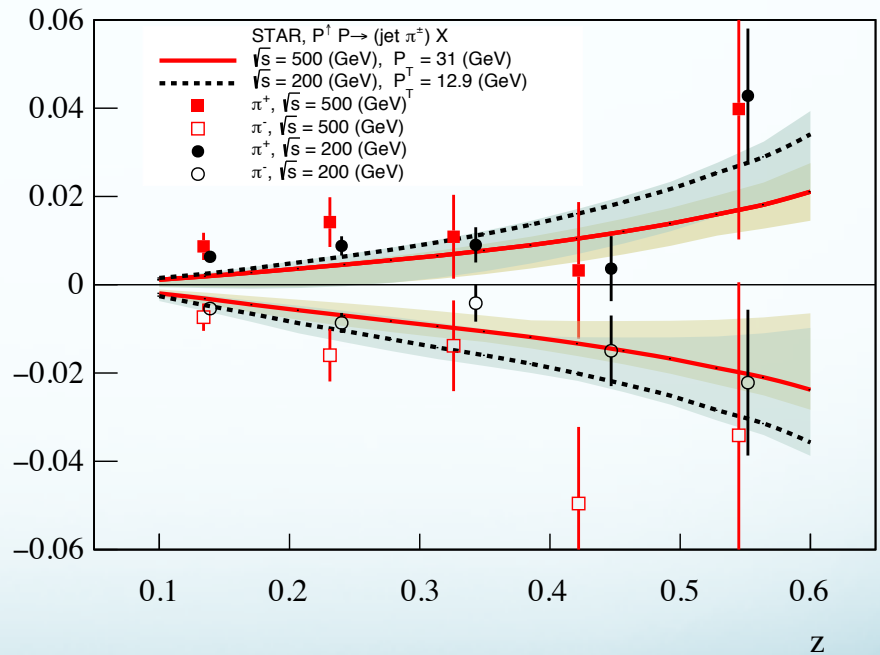
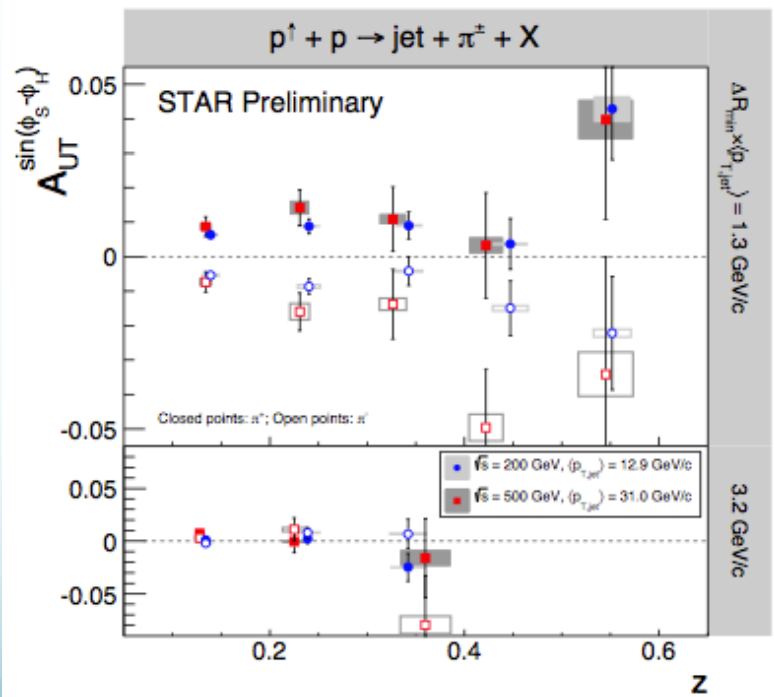


# RHIC measurements

- Hadron azimuthal distribution inside the jet in transversely polarized p+p collisions: spin dynamics

$$p^\uparrow \left[ \vec{S}_\perp(\phi_S) \right] + p \rightarrow [\text{jet } h(\phi_H)] + X$$

STAR, in arXiv:1501.01220



See Prokudin's talk on Wednesday  
Kang, Prokudin, Ringer, Yuan, to appear

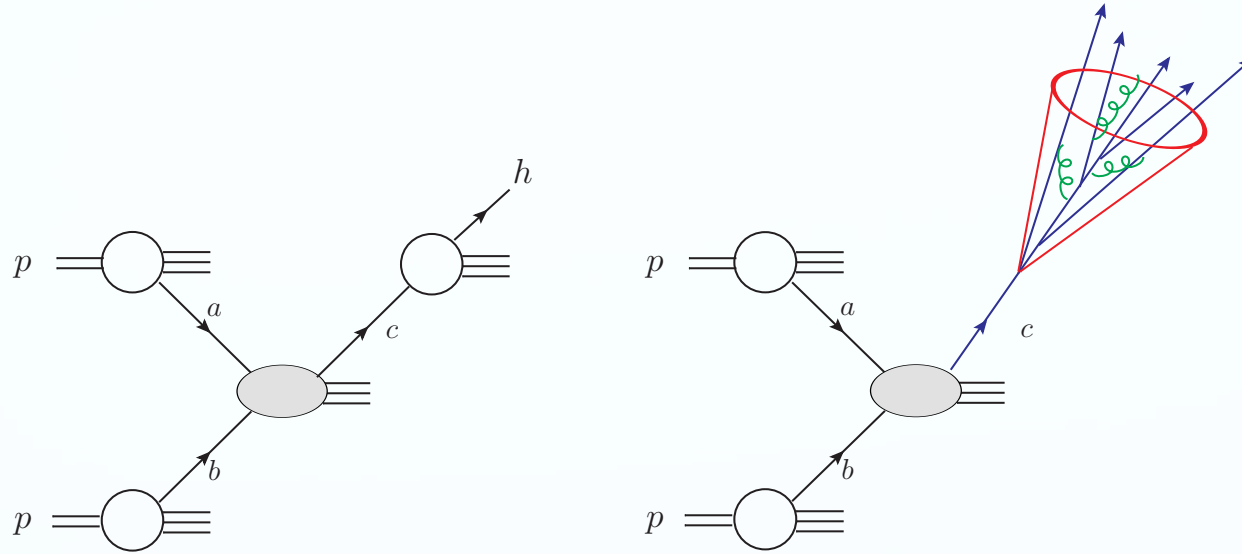
# Questions

- How does the factorization formalism look like?
- How is the collinear  $z$ -distribution of hadrons in the jet related to the standard collinear fragmentation function?
- How is the transverse momentum dependent  $j_T$ -distribution of hadrons in the jet related to the usual TMD fragmentation function as measured in SIDIS and  $e^+e^-$ ?

Lots of work have been performed along these directions recently, and very active developments e.g., Kaufmann, Mukherjee, Vogelsang; Bain, Makris, Mehen, Leibovich; Kang, Ringer, Vitev; Neill, Scimenmi, Waalewijn; ...

# A further re-factorization for jet and jet substructure

- For cross section or substructure of single **inclusive** jet production



$$\frac{d\sigma^{pp \rightarrow hX}}{dp_T d\eta} = \sum_{a,b,c} f_a \otimes f_b \otimes H_{ab \rightarrow c} \otimes D_c^h$$

$$\frac{d\sigma^{pp \rightarrow \text{jet}(v)X}}{dp_T d\eta dv} = \sum_{a,b,c} f_a \otimes f_b \otimes H_{ab \rightarrow c} \otimes \mathcal{G}_c(\mu \sim p_T R, v)$$

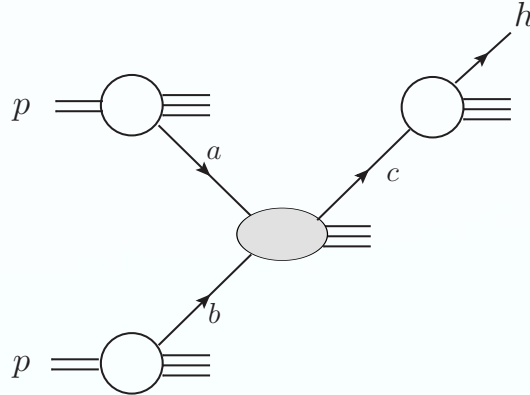
Fragmentation function

Semi-inclusive jet function

$$D_c^h \Rightarrow \mathcal{G}_c(\mu \sim p_T R, v)$$

# Recall single hadron production

- Illustration of single hadron production:  $p + p \rightarrow h + X$

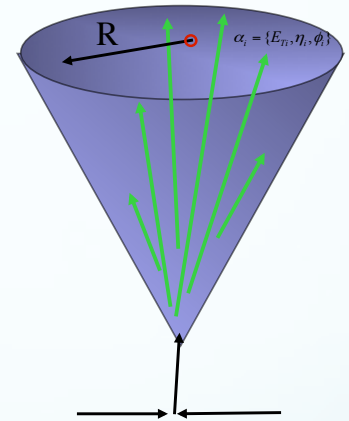
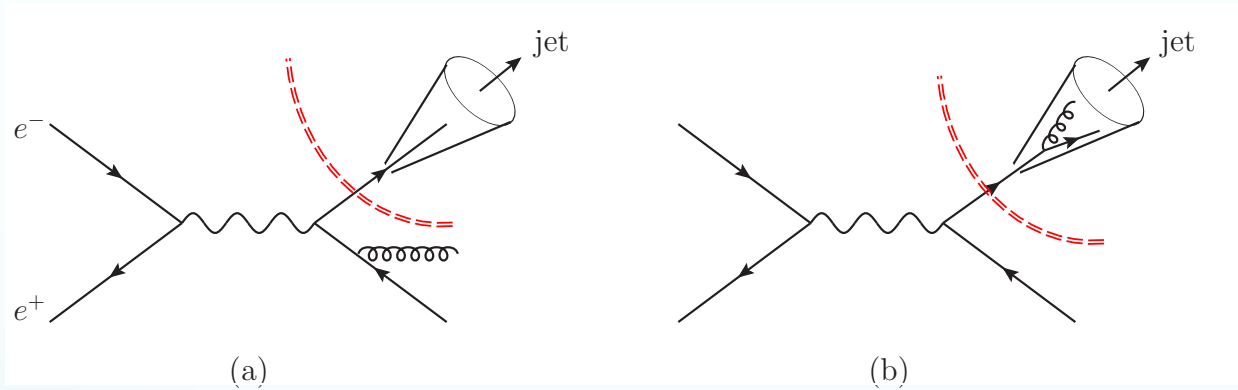


$$\frac{d\sigma^{pp \rightarrow hX}}{dp_T d\eta} = \sum_{a,b,c} f_a \otimes f_b \otimes H_{ab \rightarrow c} \otimes D_c^h$$

- QCD factorization can be reviewed from the spirit of the effective field theory: physics at very different scales do not affect each other
  - Hard collision happens at scale  $\sim p_T$
  - Hadronization/fragmentation happens at a much lower scale  $\sim m_h$
  - The interference between these two scales should be suppressed by  $m_h/p_T$

# QCD factorization makes things simple

- Think of QCD factorization using the spirit of effective field theory
  - What are the relevant scales for single jet production?
  - Two momenta: (1) hard collision:  $p_T$  (2) jet radius can build one:  $p_T R$
  - In the small- $R$  limit, one can actually factorize the jet cross section into two steps, just like single hadron production



$$\frac{d\sigma^{pp \rightarrow \text{jet} X}}{dp_T d\eta} = \sum_{a,b,c} f_a \otimes f_b \otimes H_{ab \rightarrow c} \otimes J_c(\mu \sim p_T R)$$

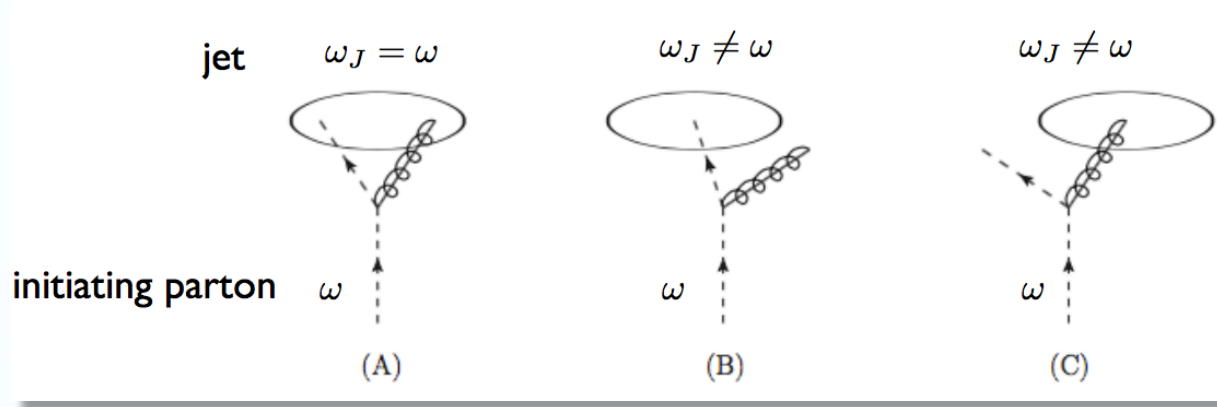
- Good thing: semi-inclusive jet function  $J_{q,g}(z, R, w)$  are purely perturbative

Kang, Ringer, Vitev, arXiv:1606.06732, Dai, Kim, Leibovich, 1606.07411, see also, Kaufmann, Mukherjee, Vogelsang, 1506.01415

# Semi-inclusive jet function

- Describe how a parton (q or g) is transformed into a jet (with a jet radius R) and energy fraction z

$$J_q(z, \omega_J, \mu) = \frac{z}{2N_c} \text{Tr} \left[ \frac{\vec{n}}{2} \langle 0 | \delta(\omega - \vec{n} \cdot \mathcal{P}) \chi_n(0) | JX \rangle \langle JX | \bar{\chi}_n(0) | 0 \rangle \right] \quad z = \omega_J / \omega$$



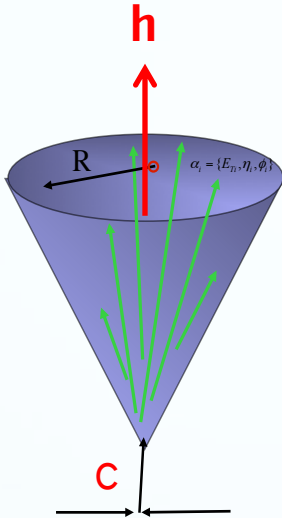
**Semi-inclusive quark/gluon jets follow DGLAP evolution equation, just like hadron fragmentation functions**

$$\mu \frac{d}{d\mu} J_i(z, \omega_J, \mu) = \frac{\alpha_s(\mu)}{\pi} \sum_j \int_z^1 \frac{dz'}{z'} P_{ji} \left( \frac{z}{z'}, \mu \right) J_j(z', \omega_J, \mu)$$

# Collinear hadron distribution inside the jet

- First produce a jet, and then further look for a hadron inside the jet

Kang, Ringer, Vitev, arXiv:1606.07063, JHEP



$$F(z_h, p_T) = \frac{d\sigma^h}{dp_T d\eta dz_h} \bigg/ \frac{d\sigma}{dp_T d\eta}$$

$$z_h = p_T^h / p_T$$

- Just like the single inclusive jet production, we have
  - Semi-inclusive fragmenting jet function

$$z = p_T / p_T^c$$

$$\frac{d\sigma}{dp_T d\eta dz_h} \propto \sum_{a,b,c} f_a \otimes f_b \otimes H_{ab \rightarrow c} \otimes \mathcal{G}_c^h(z, z_h, \mu)$$

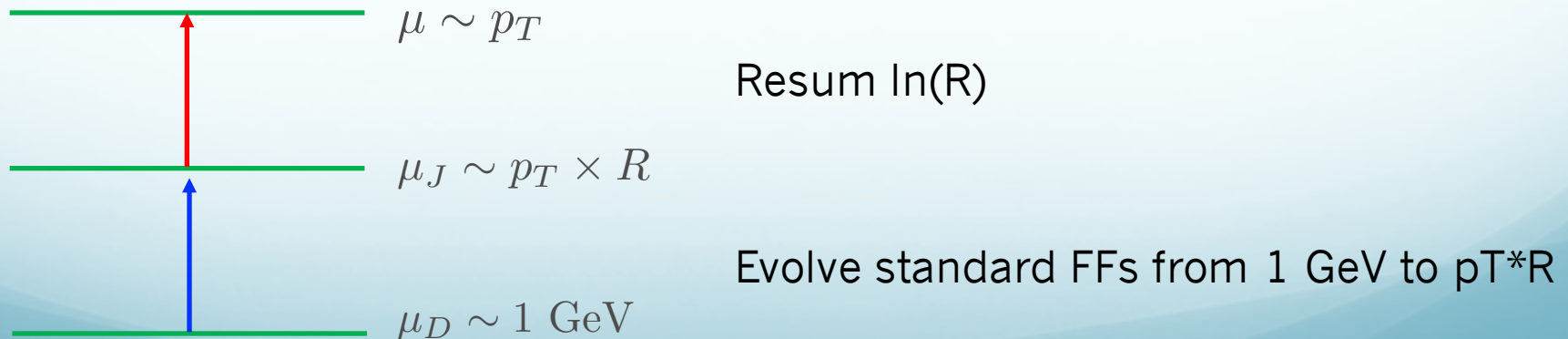
# Two DGLAPs

- Parton-to-jet part: evolution is for variable  $z$

$$\mu \frac{d}{d\mu} \mathcal{G}_i^h(z, z_h, \mu) = \frac{\alpha_s(\mu)}{\pi} \sum_j \int_z^1 \frac{dz'}{z'} P_{ji} \left( \frac{z}{z'} \right) \mathcal{G}_j^h(z', z_h, \mu)$$

- Substructure of the jet: collinear hadron distribution in the jet, relevant to variable  $z_h$

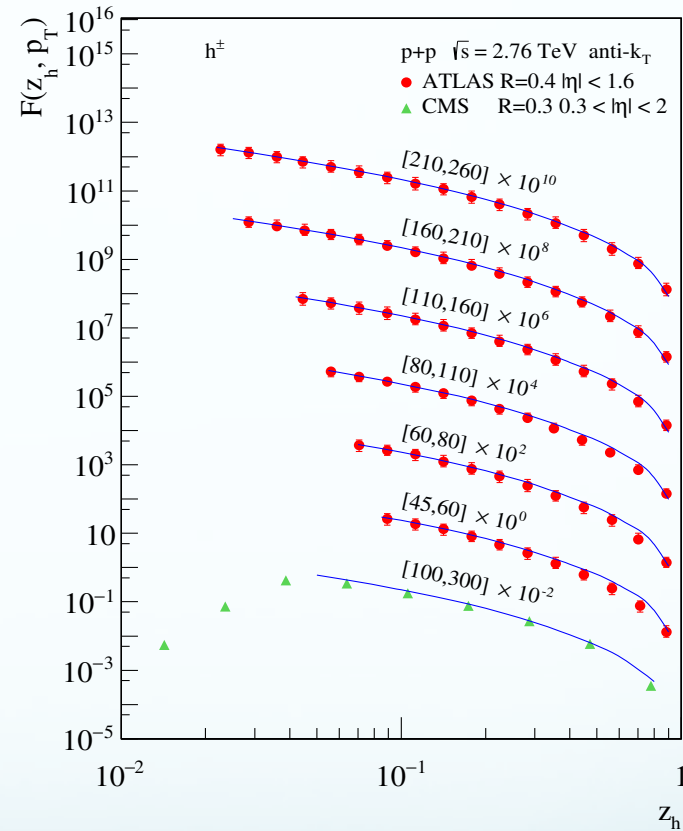
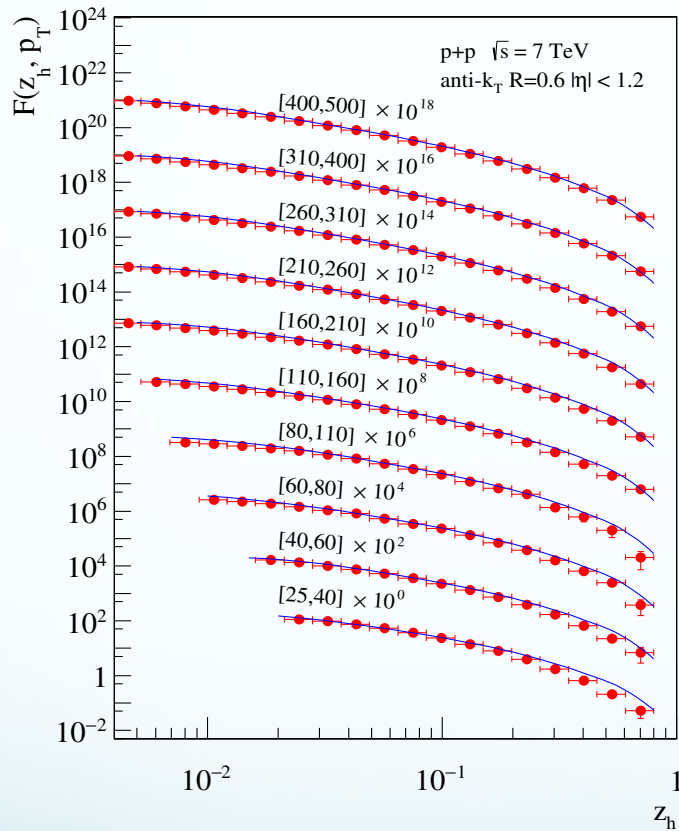
$$\mathcal{G}_i^h(z, z_h, \mu) = \sum_j \int_{z_h}^1 \frac{dz'_h}{z'_h} \mathcal{J}_{ij}(z, z'_h, \mu) D_j^h \left( \frac{z_h}{z'_h}, \mu \right)$$





# Great probe for collinear FFs

- Works pretty well in comparison with experimental data



- Could be used for better constraining gluon-to-hadron FFs, large- $z$  region and etc

Kang, Ringer, Vitev, arXiv:1606.07063

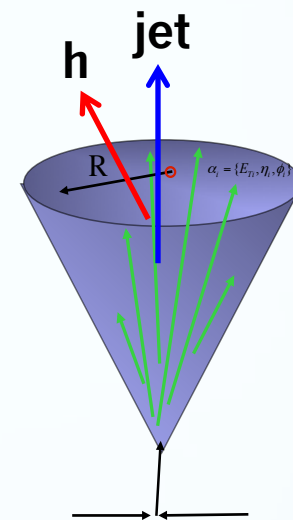
# What about TMD FFs?

- TMD hadron distribution inside the jet

$$F(z_h, j_T; p_T) = \frac{d\sigma^h}{dp_T d\eta dz_h d^2 j_T} \bigg/ \frac{d\sigma}{dp_T d\eta}$$

$$z_h = p_T^h / p_T^{\text{jet}}$$

$j_T$  : hadron transverse momentum with respect to the jet direction



- Factorization formalism

Kang, Liu, Ringer, Xing, in preparation

$$\frac{d\sigma}{dp_T d\eta dz_h d^2 j_T} \propto \sum_{a,b,c} f_a \otimes f_b \otimes H_{ab \rightarrow c} \otimes \mathcal{G}_c^h(z, z_h, j_T, \mu)$$

- Re-factorization of semi-inclusive fragmenting jet function

$$\mathcal{G}_c^h(z, z_h, j_T, \mu) = \mathcal{C}_{c \rightarrow d}(z, R) \int d^2 \lambda_T d^2 k_T \delta^2(j_T - \lambda_T - k_T) S(\lambda_T, R) D_d^h(z_h, k_T)$$

# A couple of main points

- One soft function + one TMD FFs
  - How do the rapidity divergences cancel between them?
  - Recall: standard TMD factorization for SIDIS, DY,  $e^+e^-$ , which always involve one soft function + TWO TMDs
- What sets the scale for the TMD evolution of TMD FFs?

# TMD factorization for DY: $p + p \rightarrow [\gamma^* \rightarrow \ell^+ \ell^-] + X$

- Factorized form and mimic “parton model”

$$\frac{d\sigma}{dQ^2 dy d^2q_\perp} \propto \int d^2k_{1\perp} d^2k_{2\perp} d^2\lambda_\perp H(Q) f(x_1, k_{1\perp}) f(x_2, k_{2\perp}) S(\lambda_\perp) \delta^2(k_{1\perp} + k_{2\perp} + \lambda_\perp - q_\perp)$$

$$= \int \frac{d^2b}{(2\pi)^2} e^{iq_\perp \cdot b} H(Q) f(x_1, b) f(x_2, b) S(b)$$

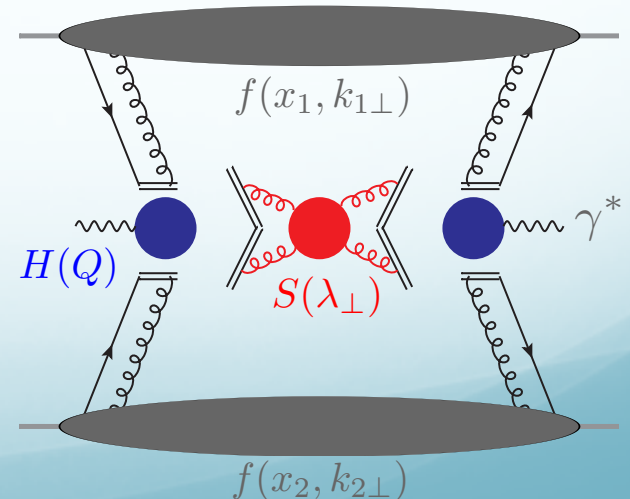


$$F(x, b) = f(x, b) \sqrt{S(b)}$$

$$= \int \frac{d^2b}{(2\pi)^2} e^{iq_\perp \cdot b} H(Q) F(x_1, b) F(x_2, b)$$

- Rapidity divergences cancel between

$$F(x, b) = f(x, b) \sqrt{S(b)}$$



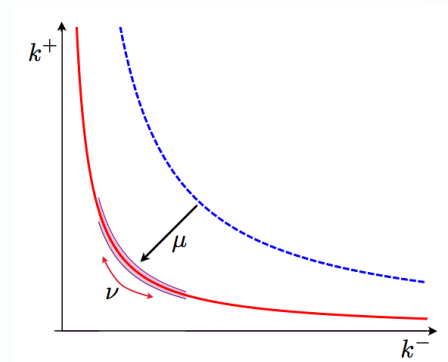
# TMDs in b-space at NLO

- Quark TMD at one loop

$$f_{q/q}(x, b) = \frac{\alpha_s}{2\pi} C_F \left\{ \left[ \frac{2}{\eta} \left( \frac{1}{\epsilon} + \ln \frac{\mu^2}{\mu_b^2} \right) + \frac{2}{\epsilon} \ln \frac{\nu}{p^+} + \frac{3}{2} \frac{1}{\epsilon} \right] \delta(1-x) \right. \\ \left. + \left( -\frac{1}{\epsilon} - \ln \frac{\mu^2}{\mu_b^2} \right) P_{qq}(x) \right. \\ \left. + \left[ 2 \ln \frac{\mu^2}{\mu_b^2} \ln \frac{\nu}{p^+} + \frac{3}{2} \ln \frac{\mu^2}{\mu_b^2} \right] \delta(1-x) + (1-x) \right\}$$

- Soft factor

$$S(b) = \frac{\alpha_s}{2\pi} C_F \left\{ \frac{4}{\eta} \left( -\frac{1}{\epsilon} - \ln \frac{\mu^2}{\mu_b^2} \right) + \frac{2}{\epsilon^2} + \frac{2}{\epsilon} \left( \ln \frac{\mu^2}{\mu_b^2} - \ln \frac{\nu^2}{\mu_b^2} \right) \right. \\ \left. + \left[ -2 \ln \frac{\mu^2}{\mu_b^2} \ln \frac{\nu^2}{\mu_b^2} + \ln^2 \frac{\mu^2}{\mu_b^2} - \frac{\pi^2}{6} \right] \right\}$$



$$\mu_b = 2e^{-\gamma_E} / b$$

- Interesting features

- Rapidity divergence cancels in  $F_{q/q}^{\text{sub}}(x, b) = f_{q/q}(x, b) \sqrt{S(b)}$
- $f_{q/q}(x, b)$  and  $S(b)$  lives in the same  $\mu \sim \mu_b$ , but different rapidity scale  $\nu \sim p^+, \mu_b$

Kang, Spin 2016 conference

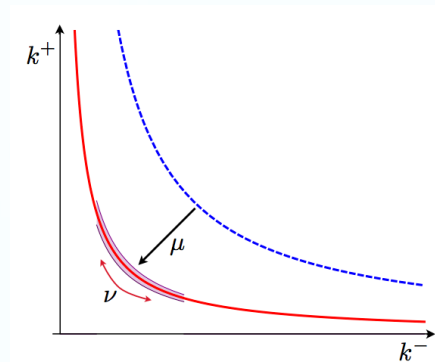
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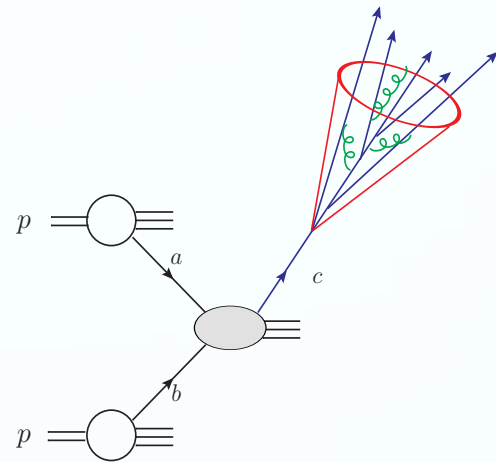
Kang, Spin 2016 conference

# What's different for hadron in the jet?

- Soft radiation has to happen inside the jet
  - For single inclusive jet production, first we produce a high-pt jet
  - This process only involves hard-collinear factorization, and such a process is not sensitive to any soft radiation
  - This is the usual standard “collinear factorization”

$$\int_0^\infty \frac{dy}{y} \Rightarrow \int_0^{\tan^2 \frac{R}{2}} \frac{dy}{y}$$

$$y \sim \frac{\ell^+}{\ell^-}$$

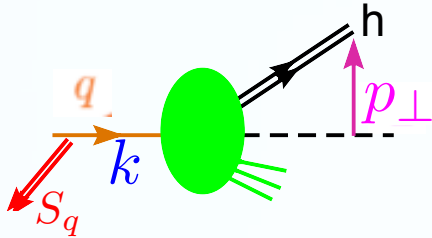


- Once such a high-pt jet is produced, we further observe a hadron inside the jet
  - At this step, we measure the relative transverse momentum of hadron w.r.t the jet. For such a step, soft radiation matters
  - However, only those soft radiation that happens inside the jet matters
  - Restricts soft radiation to be within the jet: cuts half of the rapidity divergence
- Rapidity divergence cancel between restricted “soft factor” and TMD FFs
  - At least up to this order, the combined evolution is the same as the usual TMD evolution in SIDIS, DY, e+e-; justify the use of same TMD evolution here

$$\sqrt{S(b)} D_c^h(z_h, b)_{e^+e^-} \Rightarrow S(b, R) D_c^h(z_h, b)_{pp}$$

# Collins function: universal

- Collins function: unpolarized hadron from a transversely polarized quark

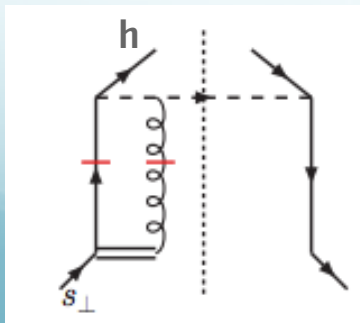


$$D_{h/q}(z, p_{\perp}) = D_1^q(z, p_{\perp}^2) + \frac{1}{zM_h} H_1^{\perp q}(z, p_{\perp}^2) \vec{S}_q \cdot (\hat{k} \times p_{\perp})$$

Spin-independent

Spin-dependent

- ✓ 2002: A. Metz studied the universality property of Collins function in a model-dependent way – very subtle – finally found it is universal between SIDIS and e+e-
- ✓ 2004: Collins and Metz have general arguments
- ✓ 2008: Yuan generalizes to pp
- ✓ Collins function is universal: concern on collinear gauge link (unsubtracted TMDs)
- ✓ Now soft function seems to be fine, too



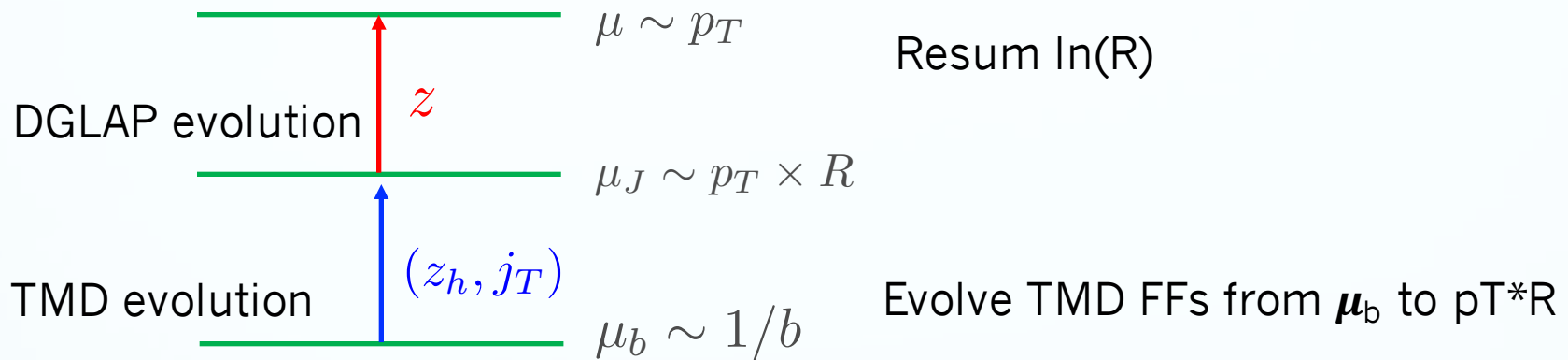
$$H_1^{\perp \text{SIDIS}}(z, p_{\perp}^2) = H_1^{\perp e^+e^-}(z, p_{\perp}^2) = H_1^{\perp \text{pp}}(z, p_{\perp}^2)$$

Metz 02, Collins, Metz 04, Yuan 08,  
Gamberg, Mulders, 10,  
Boer, Kang, Vogelsang, Yuan, 10, ...



# TMD + DGLAP evolution

- Evolution structure



- TMD FFs thus are related to the usual TMD FFs in SIDIS at scale  $p_T \times R$
- Thus hadron TMD distribution inside the jet could be used to test the universality of TMD FFs from SIDIS,  $e^+e^-$  processes

# Summary

- jet cross section and jet substructure for inclusive jet production follow a two-step factorization
  - First step: parton-to-jet production
  - Second step: jet internal substructure
- The hard function associated with the 1<sup>st</sup> step is the same as that for single inclusive hadron production
- For jet substructure, one could then concentrate on the 2<sup>nd</sup> step
- Collinear and TMD distribution of hadron in a jet are great processes to probe collinear and/or TMD FFs
  - Factorization seems to be okay

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Thank you!