

ISOQUANT

SFB1225





based on work together with: J. Berges, M. Mace, S. Schlichting, S. Sharma, N. Tanji, R. Venugopalan PRL 117 (2016) 142301, PRD 93 (2016) 074507, arXiv:1612.02477, arXiv:1701.03331 7th Workshop of the APS Topical Group on Hadronic Physics Washington DC - **2017 / 02 / 02**



The Chiral Magnetic Effect

(Kharzeev, McLerran, Warringa 2007; Kharzeev, Fukushima, Warringa 2008)



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Topological Fluctuations (sphalerons)

- analogous to proposed 'electroweak baryogenesis'





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The axial anomaly

- axial non-conservation vs. topological density

$$\partial_{\mu} j_a^{\mu}(x) = 2m\eta_a(x) - \frac{g^2}{8\pi^2} \operatorname{Tr} F_{\mu\nu}(x) \tilde{F}^{\mu\nu}(x)$$



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Review: Kharzeev et al, Prog. Nucl. Part. Phys 88 (2016) 1

External Magnetic Fields

- 10¹⁸ Gauss, strongest ever produced in Lab!





- 1. Anomalous phenomena in heavy ion collisions
- 2. Why are early times so important?
- 3. Classical-statistical simulations
- 4. Towards kinetic theory and the bigger picture
- 5. Conclusions









from coherent fields (Tanji et al. 2016) and sphaleron transitions (Mace et al. 2016)



















- condensed matter systems (Li, Kharzeev et al, Nature 2015)

- heavy ions (charged particle correlations: RHIC 2009, LHC)





1

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Challenges:

Signal – How does the signal even look like?



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Signal – How does the signal even look like?

Background – Apples to apples?



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Challenges:

Signal – How does the signal even look like?



Earliest times! – How long does the magnetic field live?



?





45

PbPb centrality(%)

35

CMS

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Challenges:

Signal – How does the signal even look like?



Earliest times! – How long does the magnetic field live?



?

<u>×10⁻³</u>

65

55

Challenge for theorists:

Quantitative Predictions needed at various stages of the collision: Production at early time --- chiral transport --- hydrodynamics and observables

Clarification soon: Isobar run at RHIC!



Most important direct contribution to CME at early times!

- (Abelian) magnetic field 'still alive'





1

1.5

 τ (fm)

2

3

b = 4 fmb = 8 fm -----

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- Initial conditions via the **Color** -Glass Condensate (CGC), from coherent longitudinal color fields to over-occupied plasma



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occupation numbers of gluons ~ $1/g \rightarrow$ quark production O(1)





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b = 4 fmb = 8 fm ----- (Abelian) magnetic field 'still alive' $= 12 \, \text{fm}$ 10^{4} $B (MeV^2)$ Initial conditions via the **Color** 10^{3} -Glass Condensate (CGC), 10^{2} from coherent longitudinal color fields to over-occupied plasma 10¹ 100 0 2.50.51.53 - Quark production large τ (fm)

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→ correspondence principle --- mapping onto classical-statistical ensemble possible



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Classical-Statistical Simulations

(see Kasper 2015 for an pedagogical introduction)

Can be systematically derived within the Color Glass Framework (Mclerran, Venugopalan 1994)



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- → replace gauge field operators by ensemble of classical Yang-Mills fields
- → Fermions: solve (operator-) Dirac equation



<u>Remarkable progress at early times – classical statistical simulations</u>



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Fermions

- EW & Higgs on a lattice (Saffin & Tranberg 2011)
- Anomalous Effects in laser physics

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Many important contributions left out here --- my personal bias --- sorry!



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(NM, Hebenstreit, Berges PRL 117 (2016) 061601)

- Real-time simulations of the Chiral Magnetic Effect and Anomalous Dynamics (QCD+QED)

(Tanji, NM, Berges PRD93 (2016) 074507; Mace, NM, Schlichting, Sharma PRL 117 (2016) 061601 & arXiv:1612.02477)





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Will show you today:

Far-from-equilibrium anomalous and topological dynamics in HI collisions can be understood <u>quantitative</u> and <u>from first principles</u>, using real-time lattice simulations in the classical statistical regime



Anomalous fermion dynamics during a topologicial transition

 \rightarrow Clean test case for the development of our tools.



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simplified situation: setting up an isolated sphaleron transition





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Fermions: Challenging!

Solving Dirac operator equation

in mode-function expansion

$$i\gamma^{0}\partial_{t}\hat{\psi} = (-iD\!\!\!\!/_{W}^{s} + m)\hat{\psi}$$
$$\hat{\psi}_{\mathbf{x}}(t) = \frac{1}{\sqrt{V}}\sum_{\lambda} \left(\hat{b}_{\lambda}(0)\phi_{\lambda}^{u}(t,\mathbf{x}) + \hat{d}_{\lambda}^{\dagger}(0)\phi_{\lambda}^{v}(t,\mathbf{x})\right)$$


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Solving Dirac operator equation

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- \rightarrow extremely costly (~N⁶)
- → big obstacle so far and many attempts at reducing price (e.g. 'low-cost' techniques)

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Chiral Magnetic and Chiral Separation Effect – Chiral Magnetic Waves

NM, Schlichting, Sharma, PRL 117 (2016) 142301; Mace, NM, Schlichting, Sharma, arXiv:1612.02477



Chiral Magnetic and Chiral Separation Effect – **Chiral Magnetic Waves**

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lattice size:24x24x64







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abelian magnetic field along z





z [r_{sph}]

 $\times [r_{sob}]$

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6



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Initially: Vacuum (no fermions, no axial charge)

Chiral Magnetic Effect:

Chiral Separation Effect:

Electric current generated due to axial charge produced Axial current generated due to electric charge











z/r_{sph}

7





z/r_{sph}















Simulating <u>chiral</u> fermions in real-time: Overlap fermions



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Lattice simulations: exactly chiral fermions challenging (at least in strong coupling regime)



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Lattice simulations: exactly chiral fermions challenging (at least in strong coupling regime)

- in the classical-statistical limit **mass** renormalized perturbatively
 - → Wilson fermions arbitrary close to massless limit



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Benchmark:

Wilson-fermions vs. Overlap fermions

- mass matters! chiral instabilities (Yamamoto, Akamatsu, Kaplan, Reddy, Sen, Dvornikov...)
- real-time evolutio beyond early time



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How this became possible – <u>Technical breakthroughs</u> The challenge with fermions



Evolution of the fermion operators **extremely costly due to mode-function expansion**

$$\hat{\psi}_{\mathbf{x}}(t) = \frac{1}{\sqrt{V}} \sum_{\lambda} \left(\hat{b}_{\lambda}(0) \phi_{\lambda}^{u}(t, \mathbf{x}) + \hat{d}_{\lambda}^{\dagger}(0) \phi_{\lambda}^{v}(t, \mathbf{x}) \right)$$

(cost with lattice size: N⁶ for fermions vs. N³ for 'YM only')



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Our approach: - tree-level operator improvements (Eguchi and N. Kawamoto 1984)

- Wilson-averaging

→ works extremely well, convergence already on small lattices 16x16x16 for smooth gauge fields



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Computational Ressources:



Check *arXiv:1612.02477* for the current state-of-art!



The limits of classical statistical simulations





The limits of classical statistical simulations





The limits of classical statistical simulations



classical-statistical simulations become inaccurate as occupation numbers drop – **subsequent evolution of chiral and vector currents influenced by medium:**



The limits of classical statistical simulations



classical-statistical simulations become inaccurate as occupation numbers drop

- subsequent evolution of chiral and vector currents influenced by medium:
 - \rightarrow scattering of topological domains
 - \rightarrow axial transport
 - → E&M conductivity



The limits of classical statistical simulations



classical-statistical simulations become inaccurate as occupation numbers drop

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 - \rightarrow scattering of topological domains
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 - → E&M conductivity

Need to connect early time and hydrodynamic evolution!



Chiral Kinetic Theory

Connection between Berry phase (adiabatic limit) and quantum anomalies

 \rightarrow talk by Jinfeng Liao



Chiral Kinetic Theory

Connection between Berry phase (adiabatic limit) and quantum anomalies

$$egin{aligned} \dot{\mathbf{k}} &= \mathbf{E} + \mathbf{v} imes \mathbf{B} \ \dot{\mathbf{x}} &= rac{\partial \mathbf{H}}{\partial \mathbf{k}} - \dot{\mathbf{k}} imes \mathbf{\Omega} \end{aligned}$$

massless fermions, large chemical potential

→ talk by Jinfeng Liao
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UNIVERSITÄT HEIDELBERG ZUKUNFT SEIT 1386

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- Absence of level crossings (adiabaticity) vs. anomaly?





\rightarrow talk by Jinfeng Liao

4. Relativistic Chiral Kinetic Theory

Chiral Kinetic Theory

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massless fermions, large chemical potential

- Absence of level crossings (adiabaticity) vs. anomaly?
- topology of Berry's phase vs. topology of the anomaly? (Fujikawa 2005)

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massless fermions, large chemical potential

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- topology of Berry's phase vs. topology of the anomaly? (Fujikawa 2005)
- adiabatic approximation in heavy ion collision?
- Lorentz covariance (side jumps?)

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4. Relativistic Chiral Kinetic Theory



Chiral Kinetic Theory – just a teaser



We understand the anomaly very well in second quantization (Fujikawa 1979)

How can this understanding be extended to first quantization (x,p,...) – from first principles?



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→ World line quantization (Schwinger, Feynman, Bern, Kosower, Strassler ...)

$$Z_{S^{1}} = \int \frac{d^{D}p}{(2\pi)^{D}} \left(-\int_{0}^{\infty} \frac{dT}{T} e^{-(p^{2}+m^{2})T} \right) = \int \frac{d^{D}p}{(2\pi)^{D}} \ln(p^{2}+m^{2}) = \Gamma_{eff}^{QFT}$$
$$\int_{P} \mathcal{D}x e^{-\int_{0}^{1} d\tau \frac{1}{4T}\dot{x}^{2}} = \operatorname{Tr} e^{-\hat{p}^{2}T} = \int \frac{d^{D}p}{(2\pi)^{D}} e^{-p^{2}T}$$



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$$\int_{P} \mathcal{D}x e^{-\int_{0}^{1} d\tau \frac{1}{4T}\dot{x}^{2}} = \operatorname{Tr} e^{-\hat{p}^{2}T} = \int \frac{d^{D}p}{(2\pi)^{D}} e^{-p^{2}T}$$

- fermionic determinant and anomalies Witten & Alvarez-Gaume 1983, D'Hoker & Gagne 1996



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see NM, R. Venugopalan, arXiv:1701.03331 (+ longer paper very very soon); and NM, R. Venugopalan, Y.Yin, in prep.



- I have shown you real-time classical statistical simulations of fermion production during sphaleron transitions in background magnetic fields
- Axial anomaly realized in lattice simulations using Wilson fermions
- Chiral Magnetic and Chiral Separation Effect emerge dynamically
- Observation of the Chiral Magnetic Wave
- Have investigated finite mass and magnetic field dependence. Finite quark mass plays an important role: dissipation of anomalous currents
- Simulated chiral lattice fermions in real-time overlap!



The techniques are established for the quantitative first principle simulation of farfrom-equilibrium dynamics of fermions and gauge fields!

- → Realistic heavy ion collision CGC initial conditions & expansion (work in progress with J. Berges, M. Mace, S.Schlichting, S. Sharma, N. Tanji)
- → chiral plasma instabilities (with M. Mace, S.Schlichting, S. Sharma)
- → provide initial conditions for hydrodynamics
- → electromagnetic properties of the Glasma, how long does the magnetic field live?
- → connection to chiral kinetic theory & anomalous hydrodynamics (with R. Venugopalan, Y. Yin)
- → non-thermal photons and the direct photon puzzle (with J. Berges, O. Garcia-Montero, N. Tanji)







Thanks for your attention and thanks to





Studienstiftung des deutschen Volkes



Backup Classical Statistical Simulations

(see for example Kasper et al. Phys.Rev. D90 (2014) 2, 025016)

typical situation: large **coherent or highly occupied** gauge fields \rightarrow correspondence principle

- initial stages of a heavy ion collision: $A \sim 1/g$
- colliding laser beams \rightarrow large and coherent fields

The classical-statistical approximation is a systematic expansion of the 'quantum' fields around the 'classical' fields

Fermions are never "classical"

$$Z_{\mathcal{C}} = \int [dA] \rho_{G}(A) \exp\left(\operatorname{Tr}\log\Delta_{\mathcal{C}}[A]^{-1} + iS_{G}\right)$$
$$i\gamma^{0}\partial_{t}\hat{\psi} = (-iD_{W}^{s} + m)\hat{\psi}$$
$$\operatorname{Tr}\log\Delta_{c}^{-1}[A] = g^{2} \begin{bmatrix} \underline{\tilde{A}}' - \bigcirc \underline{\tilde{A}}' + \underbrace{\tilde{A}}' \\ \underline{\tilde{A}}' - \bigcirc \underline{\tilde{A}}' + \underbrace{\tilde{A}}' + \underbrace{\tilde{A}}' - \underbrace{\tilde{A}}' + \underbrace{\tilde{A}}' + \underbrace{\tilde{A}' - \underbrace{\tilde{A}}' + \underbrace{\tilde{A}}' + \underbrace{\tilde{A}}' + \underbrace{\tilde{A}}' + \underbrace{\tilde{A}}' + \underbrace{\tilde{A}' - \underbrace{\tilde{A}}' + \underbrace{\tilde{A}}' + \underbrace{\tilde{A}}' + \underbrace{\tilde{A}' - \underbrace{\tilde{A}' - \underbrace{\tilde{A}$$

Exact description via modefunctions up to 24x24x64 lattices





'classical' 'quantum'

 $A^{-}_{\mu,n} = \bar{A}_{\mu,\bar{n}} - \frac{1}{2}\tilde{A}_{\mu,\bar{n}} \,,$

 t_0

Backup B. Algorithmic Improvements



Fermions:

Exact description via **modefunctions** up to 24x24x64 lattices

Improved operators (NM, S.Schlichting, S.Sharma, arXiv:1606.00342, arXiv:1612.02477)

We use a tree-level improved version of the lattice Hamiltonian, which takes the form

$$H = \sum_{x} \psi_{x}^{\dagger} m \gamma^{0} \psi + \frac{1}{2} \sum_{n,x,i} C_{n} \psi_{x}^{\dagger} \gamma^{0} \left[\left(-i\gamma^{i} - nr_{w} \right) U_{x,+ni} \psi_{x+ni} + 2nr_{w} \psi_{x} - \left(-i\gamma^{i} + nr_{w} \right) U_{x,-ni} \psi_{x-ni} \right]$$

where r_w denotes the Wilson coefficient, the coefficients C_n are chosen to optimize the convergence, and we introduce the following short hand notation for the connecting gauge links

(1)
$$U_{x,+ni} = \prod_{k=0}^{n-1} U_{x+ki,i} , \qquad U_{x,-ni} = \prod_{k=1}^{n} U_{x-ki,i}^{\dagger}$$

Wilson-averaging (M.Mace, NM, S.Schlichting, S.Sharma, arXiv:1612.02477)

- → improvement of chiral properties
- \rightarrow extremely important for larger fermion masses
- \rightarrow average fermionic observables over Wilson parameters with opposite sign
- \rightarrow leading order errors in the anomaly equation cancel

Backup C. Magnetic Fields on the lattice



Magnetic fields break translation invariance \rightarrow magnetic translation group

- Magnetic fields on a torus very non-trivial (see Al-Hashimi & Wiese "Accidental Symmetries", also Bali et al.)

$$U_{y,n} = e^{ia^2qBn_x} ; U_{x,N_x-1,n_y,n_z} = e^{-ia^2qBN_xn_y}$$
$$U_{x,n} = \mathbf{1}, n_x \neq N_x - 1; U_{z,n} = \mathbf{1}$$

- Intriguing lattice artefacts!
- \rightarrow spoil the low-cost method
 - --- while there probably are field configurations where low-cost works, this is certainly not the case in magnetic fields

BackupD. Anomaly Realization on the Lattice



Chiral Symmetry + Fermion doubling + Chiral Anomaly = "one of the prettiest connections I have ever seen"

Lattice Fermions: Species Doubling, Chiral Invariance, and the Triangle Anomaly

Luuk H. Karsten (Stanford U., ITP) , Jan Smit (Amsterdam U.)

Sep 1980 - 38 pages

Nucl.Phys. B183 (1981) 103

In *Rebbi, C. (Ed.): Lattice Gauge Theories and Monte Carlo Simulations*, 495-532. (Nucl. Phys. B183 (1981) 103-140) and Stanford Univ. - ITP-677 (80, REC. NOV.) 71p

(1981) DOI: <u>10.1016/0550-3213(81)90549-6</u> Conference: <u>C81-06-01</u>, p.495-532 <u>Proceedings</u> ITP-677-STANFORD <u>Contributions</u>

- The axial anomaly and the fermion doubling problem are intimately related

- Lattice theory regularized on the basis of the action already

- Anomaly comes from the non-trivial continuum limit of any regulator you put in to remove doublers

