

# IC at IC: IceCube can probe the intrinsic charm of the proton

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arXiv 1607.08240



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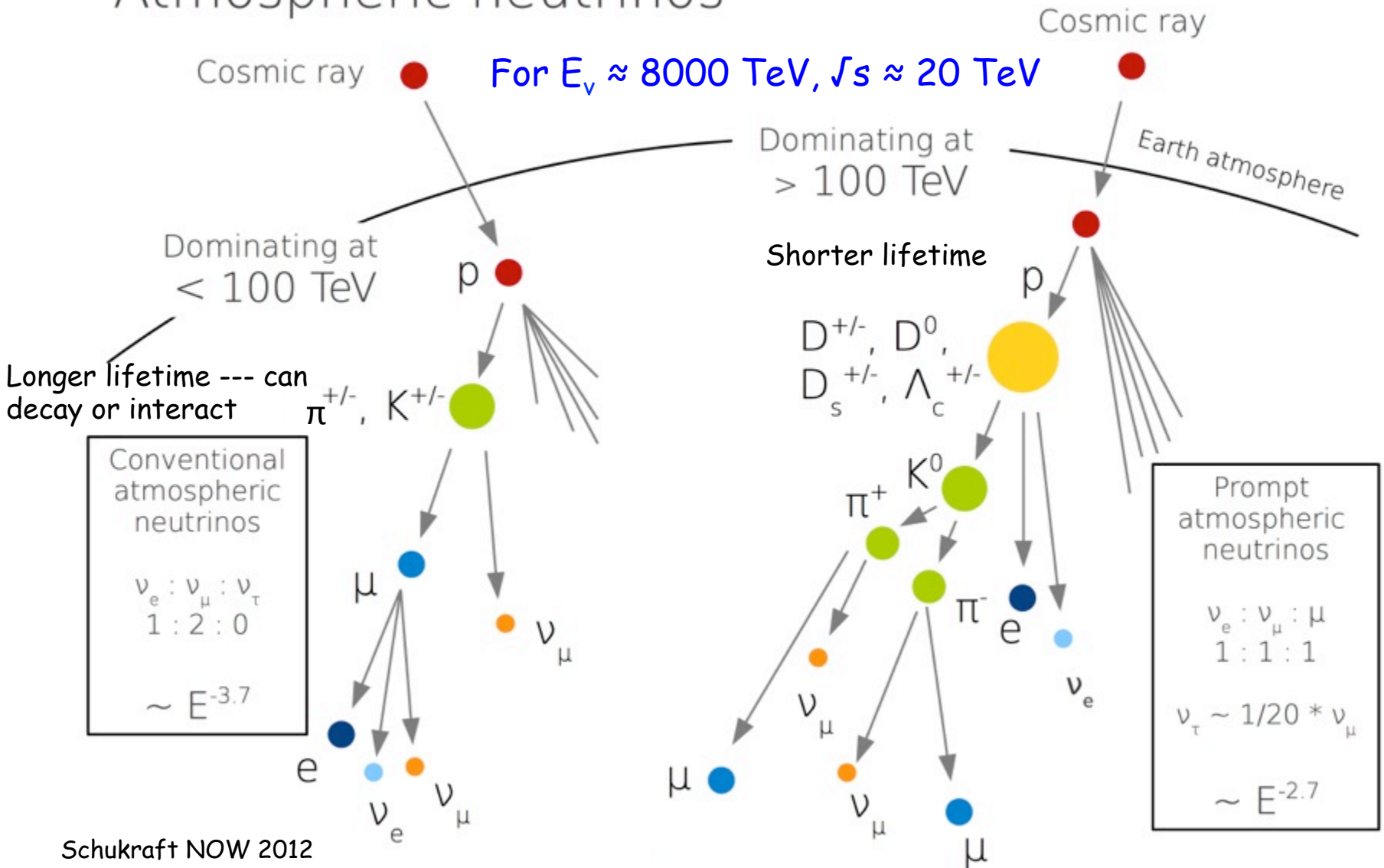
- Atmospheric neutrinos and IceCube
- Prompt atmospheric neutrinos
- Intrinsic charm contribution to prompt atmospheric neutrinos

# Atmospheric neutrinos and IceCube

Neutrinos  $\equiv$  neutrinos + antineutrinos

# Atmospheric neutrinos

## Atmospheric neutrinos



# IceCube neutrino telescope

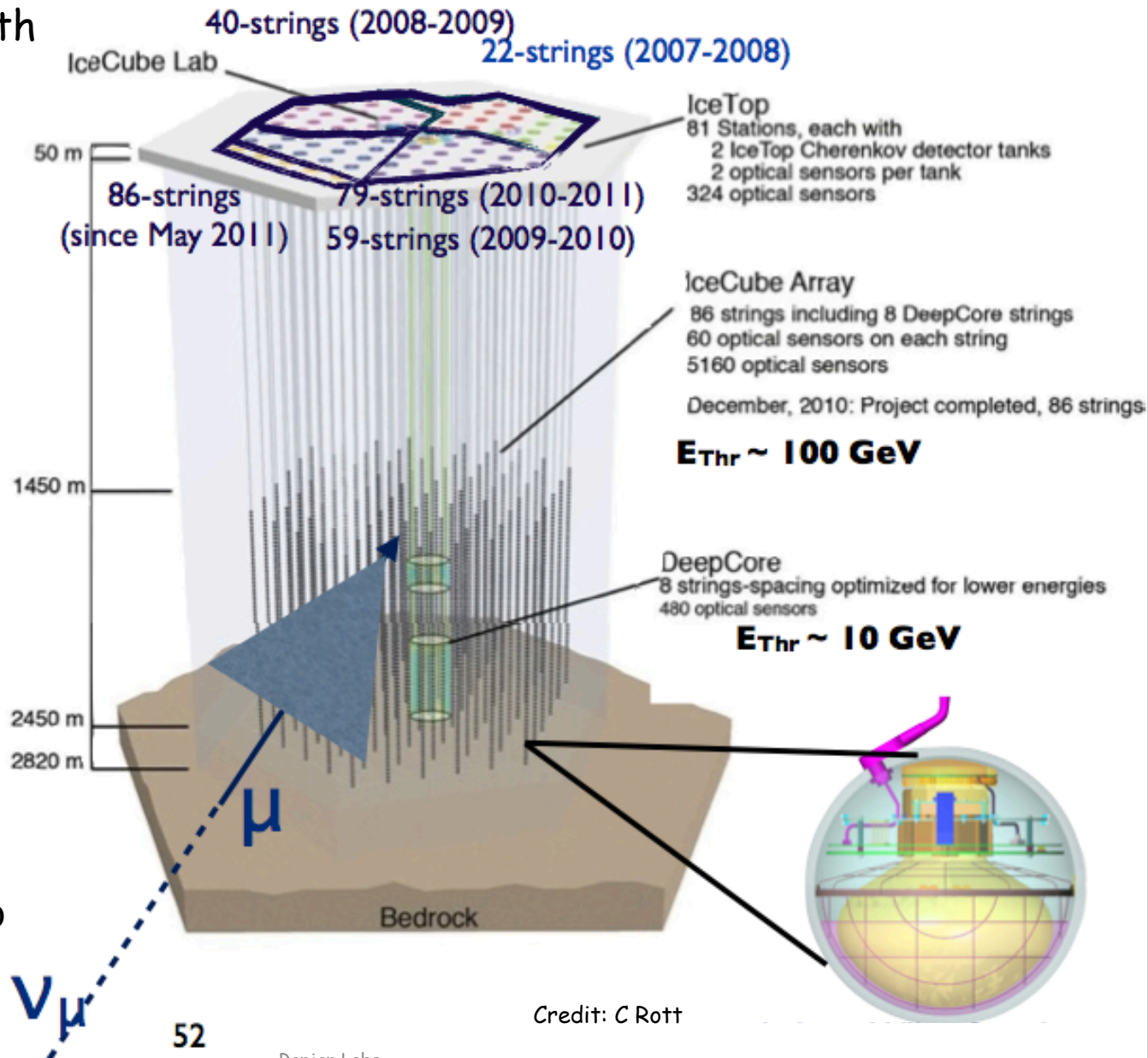
Gigaton effective volume  
neutrino detector at South  
Pole

5160 Digital Optical  
Modules distributed over  
86 strings

Completed in Dec 2010;  
data in full configuration  
from May 2011

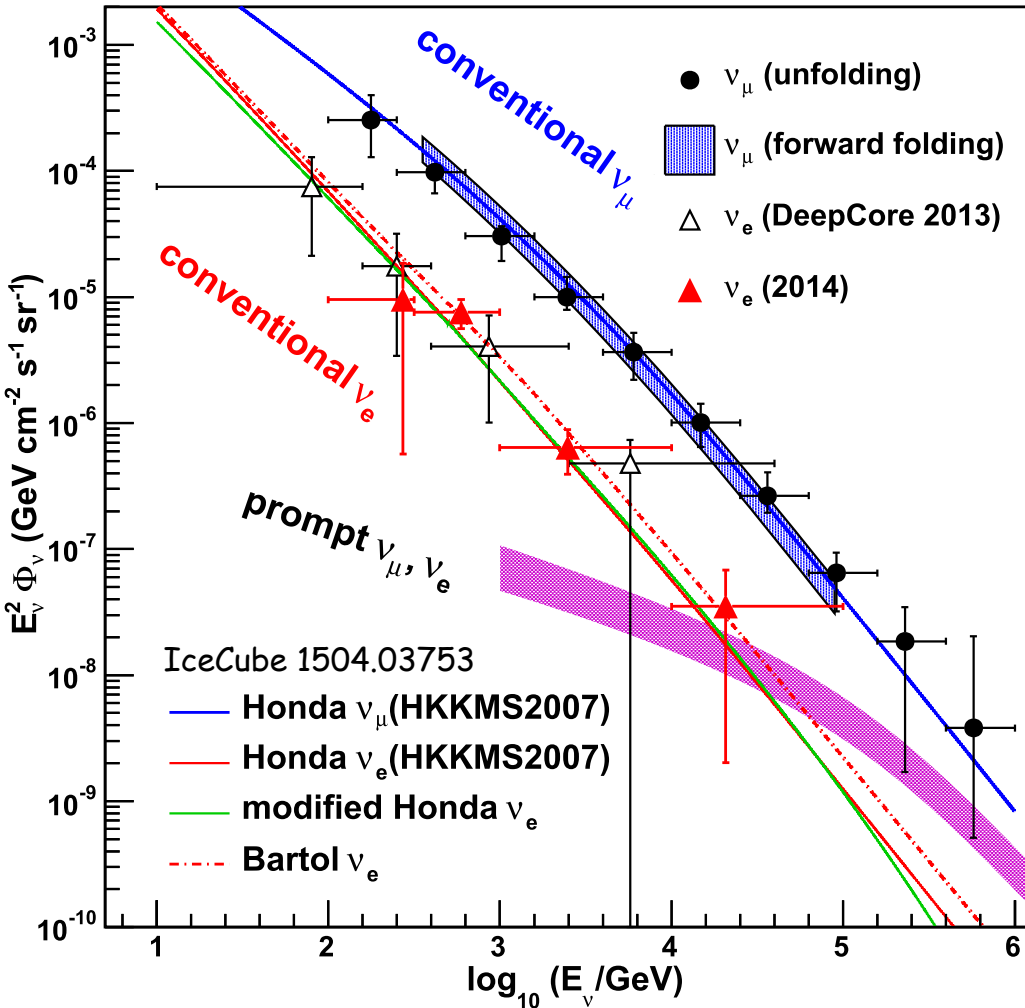
Data acquired during  
construction phase is  
analyzed

Neutrino detected  
through Cherenkov light  
emission from charged  
particles produced due to  
neutrino CC/NC  
interactions



$\nu_\mu$

# Conventional atmospheric neutrinos



CR + air  $\rightarrow \pi^\pm, K^\pm \rightarrow \nu_\mu, \bar{\nu}_\mu, \nu_e, \bar{\nu}_e$

Very important measurement ---  
neutrino oscillation discovery

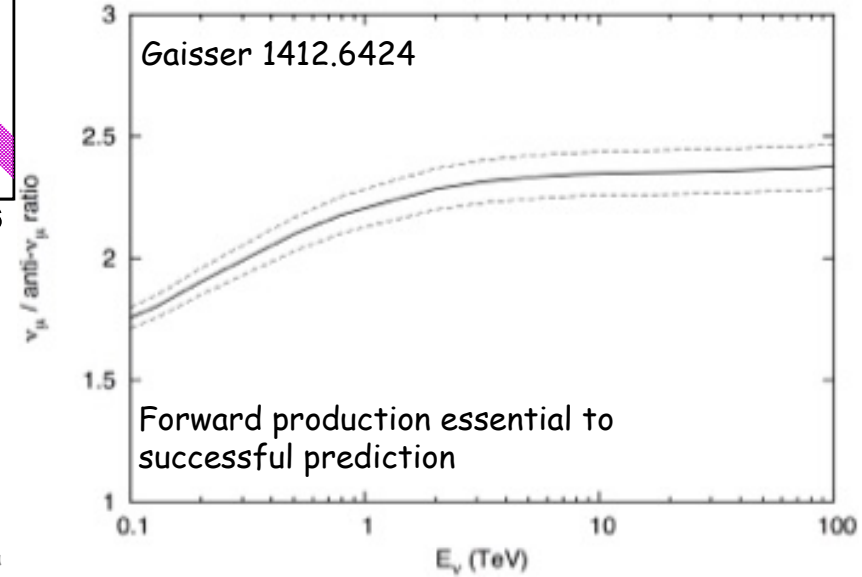


2015

The flux is lower in the vertical region, compared to the horizontal region

Negligible amount of  $\nu_\tau + \bar{\nu}_\tau$

$\nu_\mu : \nu_e$  (+ anti - particles)  $\approx 2 : 1$



Lower energy measurements also exist by  
Frejus, AMANDA, Super-Kamiokande.  
Measurements at similar energy by  
ANTARES

# Prompt atmospheric neutrinos

# Prompt atmospheric neutrinos

- Production of charmed hadrons in the atmosphere

$$\text{CR} + \text{air} \rightarrow D^{\pm}, D^0, \bar{D}^0, D_s^{\pm}, \Lambda_c^{\pm}, \dots$$

$$\text{charmed hadrons} \rightarrow \ell^{\pm}, \nu_{\ell}, \bar{\nu}_{\ell}, \dots$$

$$\nu_e : \nu_{\mu} : \nu_{\tau} \approx 1 : 1 : 0.1 \quad \text{and} \quad \nu : \bar{\nu} = 1 : 1$$

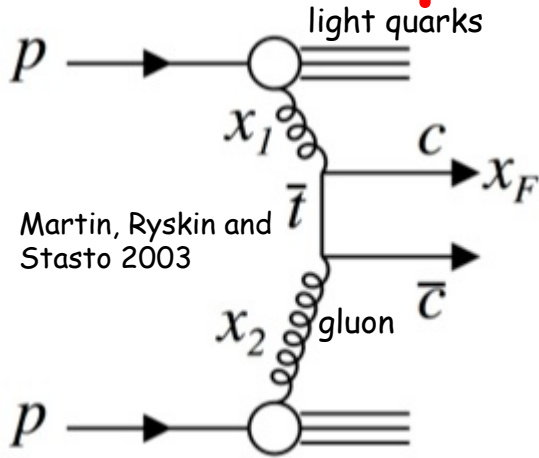
“prompt” = the charmed hadrons decay faster compared to  $\pi/K$

Isotropic flux till  $10^7$  GeV

Prompt atmospheric neutrino flux shape  $\approx$  cosmic ray flux shape



# Prompt atmospheric neutrinos



Most calculations are performed in **perturbative QCD**

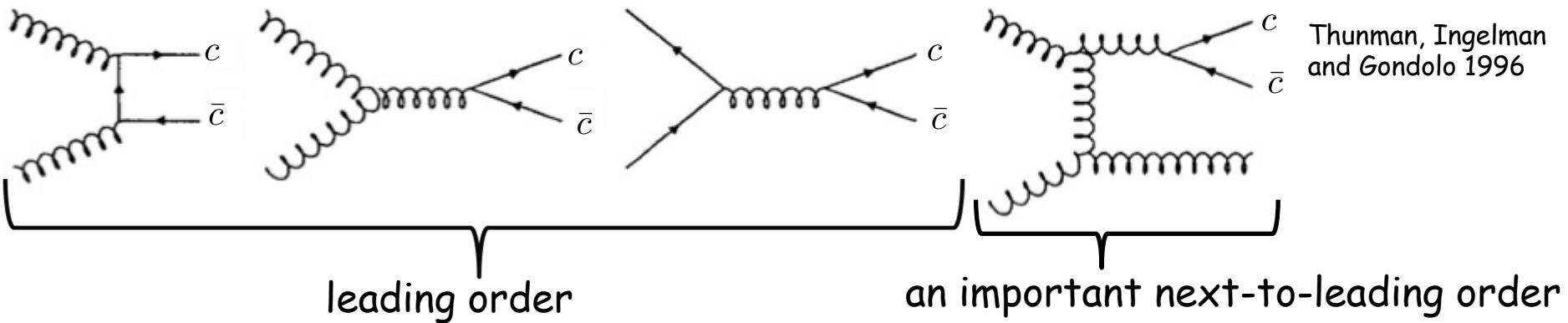
**Significant uncertainties** due to

- (i) Charm mass,
- (ii) Factorization and renormalization scale, and
- (iii) choice of the parton distribution function

Bhattacharya et al., 1502.01076

$$p + p \rightarrow c + \bar{c} + X$$

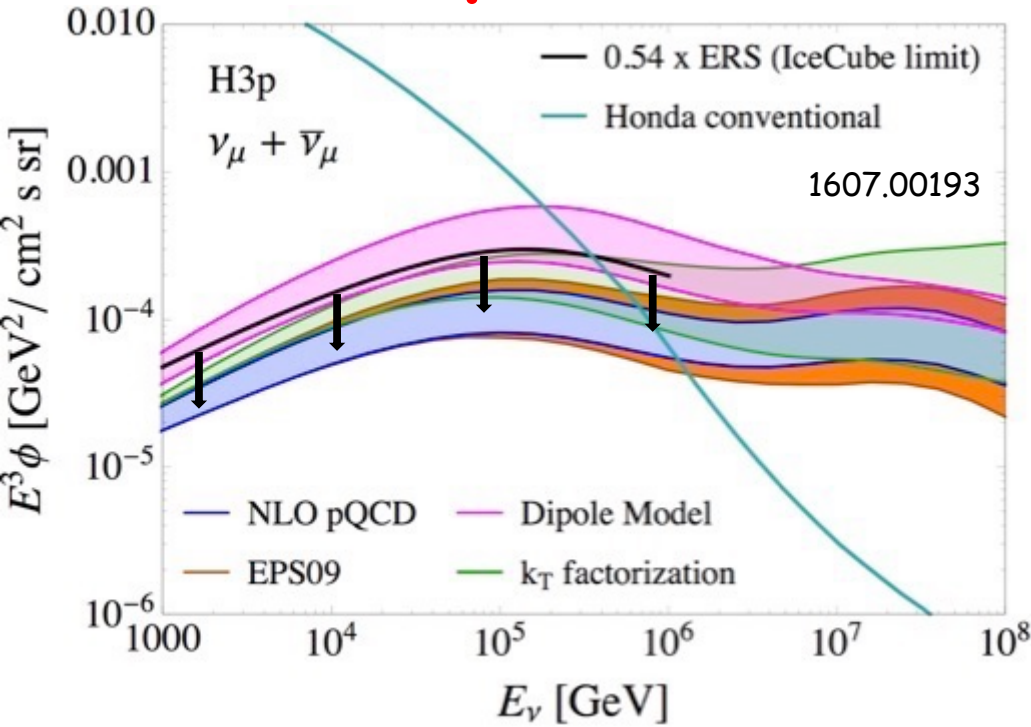
Additional **uncertainty** due to the cosmic ray input spectrum



Sensitive to QCD mechanisms in regions **beyond the reach of LHC**

At high  $s$ , the interaction is very sensitive to the gluon distribution ( $x \sim 10^{-8} - 10^{-4}$ )

# Flux predictions from pQCD



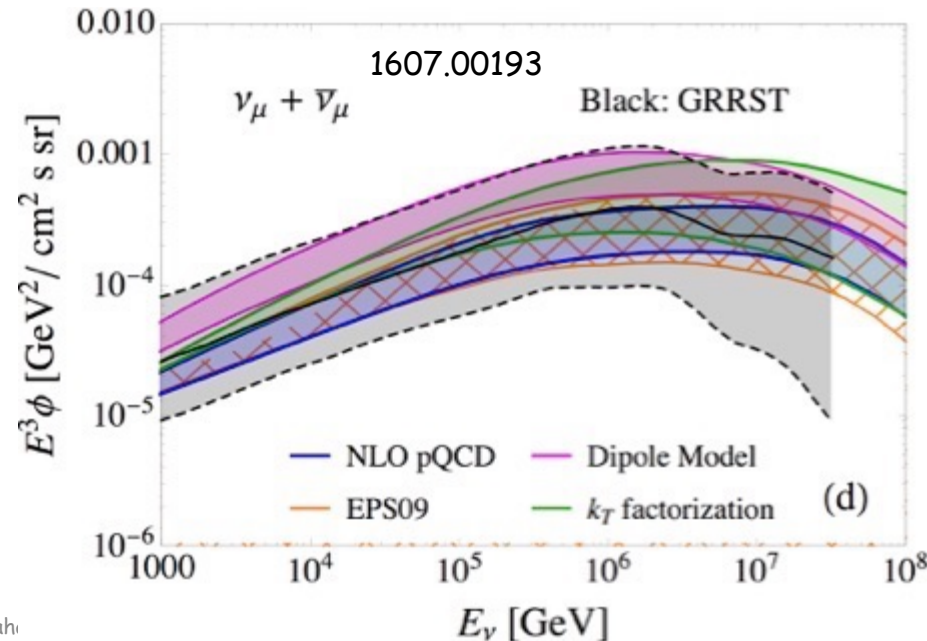
Various approaches in pQCD

Large uncertainty

Calculations by various groups in recent times: 1502.010776, 1507.01570, 1511.06346, 1607.00193, 1611.03815

The central values of various approaches largely agree

Present IceCube upper limits can constrain some of these approaches  
--- unique constraints on pQCD parameter space



# Intrinsic charm contribution to the prompt atmospheric neutrinos

arXiv 1607.08240

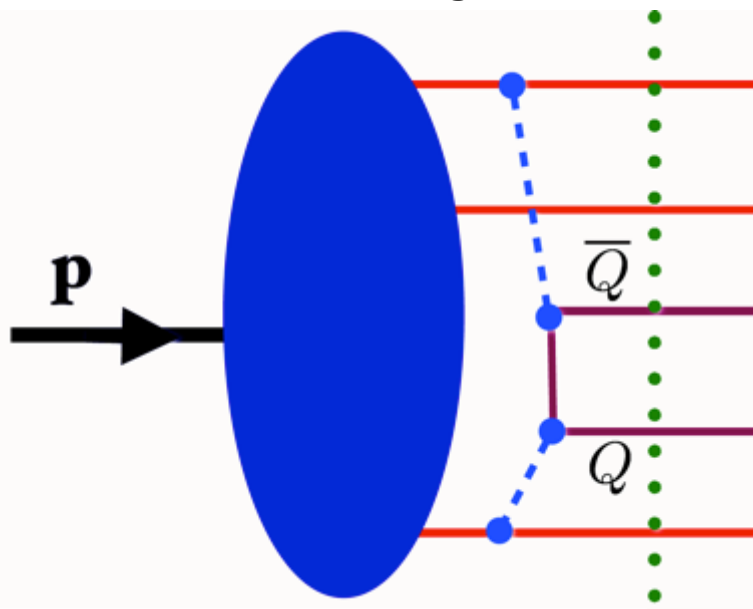
# Intrinsic charm

Brodsky, Hoyer, Peterson, and Sakai 1980

- A rigorous prediction of QCD
- Non-perturbative component
- Flattish  $d\sigma/dx_F$  observed at SELEX, ISR
- Dominates at high  $x_F$

See talks in this meeting by Gardner, Signori, Ilten, Dulat, Hobbs, and Sato

$$|p\rangle = A |uud\rangle + B |uudc\bar{c}\rangle + \dots$$



Quantum fluctuation of the proton

Probability for the proton to contain an intrinsic charm and anti-charm quark is related to  $|B|^2$

During an interaction, the u, d and c quark can combine to form a  $\Lambda_c$

Brodsky et al., 1504.06287

There are a number of fixed-target experiments like SMOG at LHCb and AFTER@LHC which aim to confirm or constrain the intrinsic charm of the proton

The normalization constant B has to be deduced from experiments

# Intrinsic charm contribution

- The measurement of prompt atmospheric neutrinos is a **forward measurement**  $\Rightarrow$  **intrinsic charm** can play an important role

$$x_F \approx E_c / E$$

$E_c$  = outgoing charm quark energy

$E$  = incident proton energy

- **Intrinsic charm** uses the incident proton energy more efficiently. It can play an important role since the cosmic rays have a steeply falling spectrum
- **Nuclear dependence** ( $\sim A^{0.7}$ ) is important since the atmospheric target is mostly nitrogen

# Cascade equations for protons, charm hadron and leptons

Atmospheric column depth

Equivalent all-proton flux

$$\frac{d\phi_p(E, X)}{dX} = -\frac{\phi_p(E, X)}{\lambda_p(E)} + Z_{pp}(E) \frac{\phi_p(E, X)}{\lambda_p(E)}$$

Charm hadron flux

$$\frac{d\phi_m(E, X)}{dX} = -\frac{\phi_m(E, X)}{\rho(X) \delta_m(E)} - \frac{\phi_m(E, X)}{\lambda_m(E)} + Z_{mm}(E) \frac{\phi_m(E, X)}{\lambda_m(E)} + Z_{pm}(E) \frac{\phi_p(E, X)}{\lambda_p(E)}$$

Lepton flux

$$\frac{d\phi_\ell(E, X)}{dX} = \sum_m Z_{m\ell}(E) \frac{\phi_m(E, X)}{\rho(X) \delta_m(E)}$$

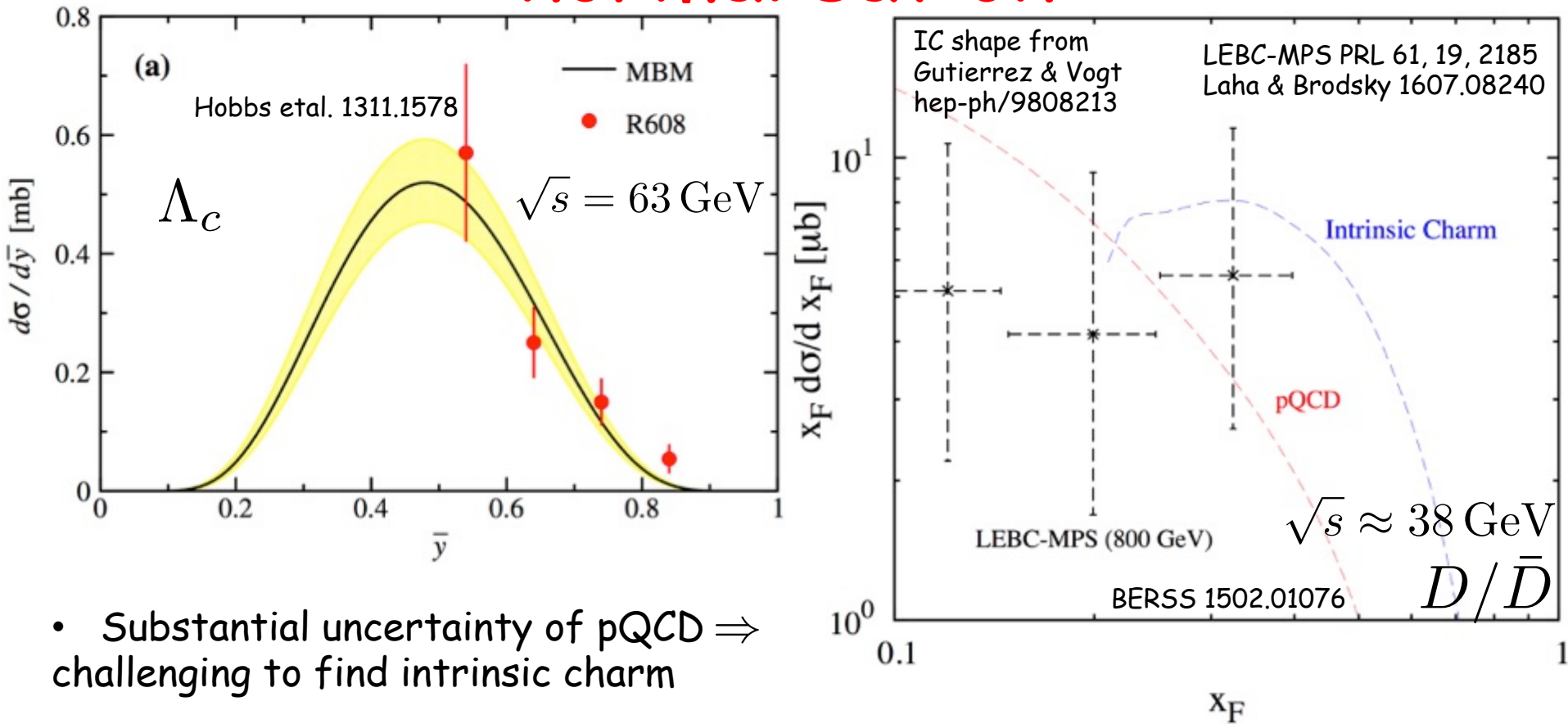
Cross-section weighted moment

$$Z_{kj}(E) = \int_0^1 \frac{dx_E}{x_E} \frac{\phi_k\left(\frac{E}{x_E}, 0\right)}{\phi_k(E, 0)} \frac{\lambda_k(E)}{\lambda_k\left(\frac{E}{x_E}\right)} \frac{dn_{kj}}{dx_E}(E/x_E)$$

$$\frac{dn}{dx_E} \equiv \frac{1}{\sigma} \frac{d\sigma}{dx_E}$$

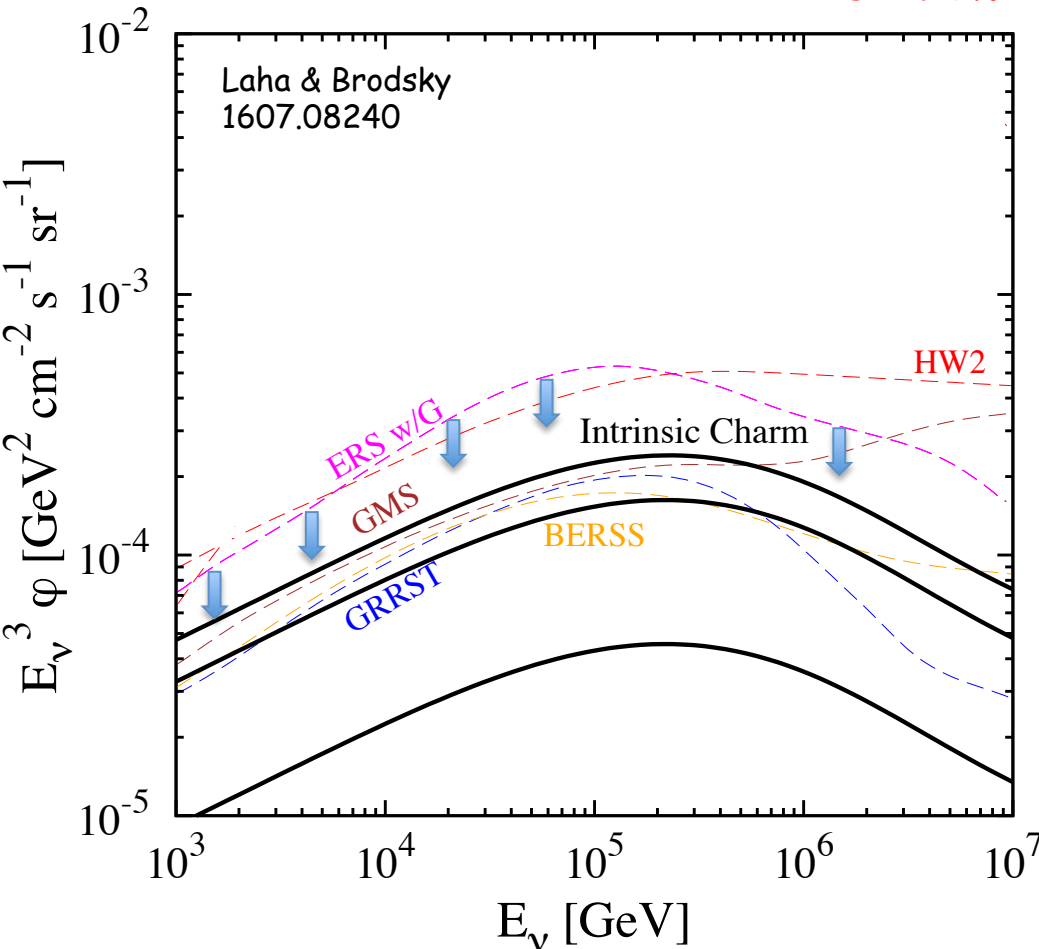
$$x_E = \frac{E_{\text{hadron}}}{E_{\text{beam}}}$$

# Intrinsic charm cross section normalisation



- Substantial uncertainty of pQCD  $\Rightarrow$  challenging to find intrinsic charm
- We assume the **best-fit pQCD** cross section and then use **LEBC-MPS** data to normalize the D cross sections
- ISR cross section by itself produces **too large** atmospheric prompt neutrino flux

# Prompt atmospheric flux due to intrinsic charm



Depending on the normalization, the **contribution due to intrinsic charm** can be as large as that due to perturbative QCD.

The **important charm hadrons** that contribute towards this flux are  $D^0, \bar{D}^0, D^\pm, D_s^\pm, \Lambda_c$

The **neutrino flavor ratio** is

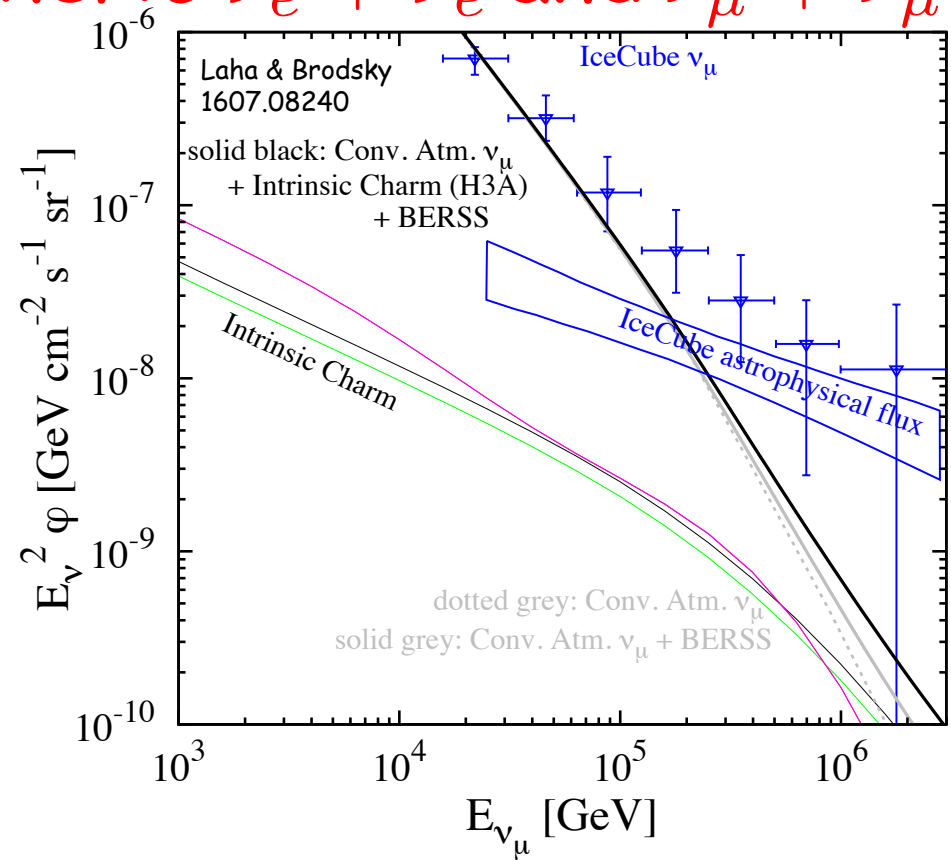
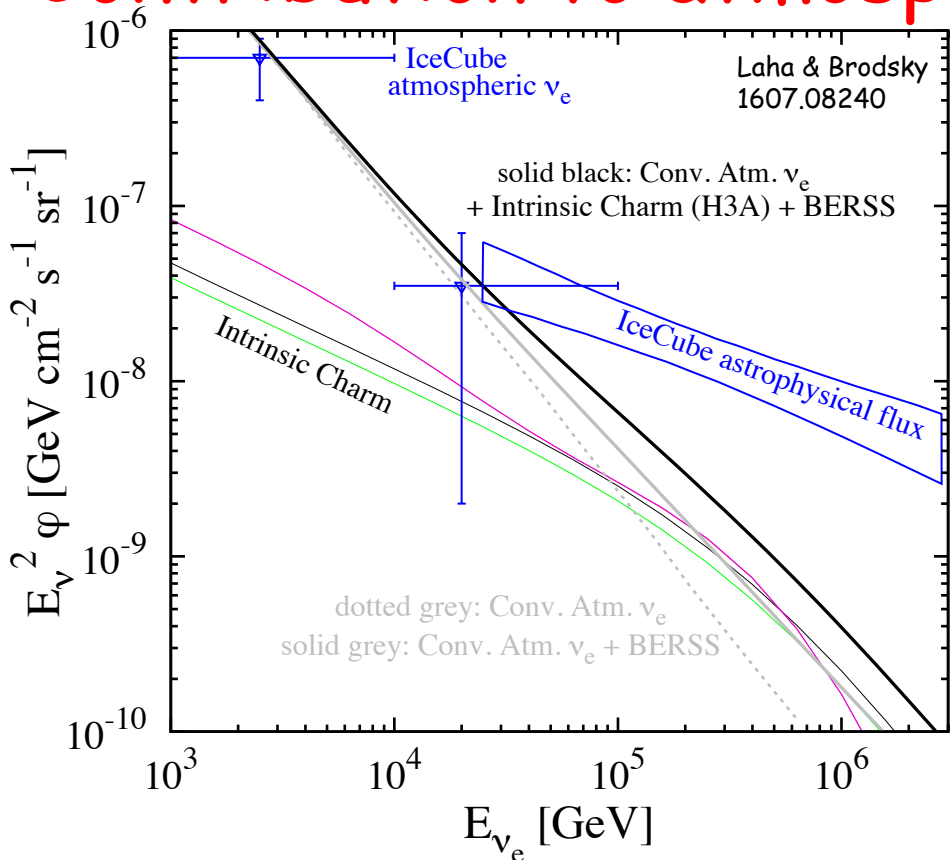
$$\nu_e : \nu_\mu : \nu_\tau \approx 1 : 1 : 0.1$$

This is an **additional contribution** to the prompt atmospheric flux

IceCube **upper limits** (near ERS w/G flux) are very close to the contribution due to intrinsic charm



# Contribution to atmospheric $\nu_e + \bar{\nu}_e$ and $\nu_\mu + \bar{\nu}_\mu$



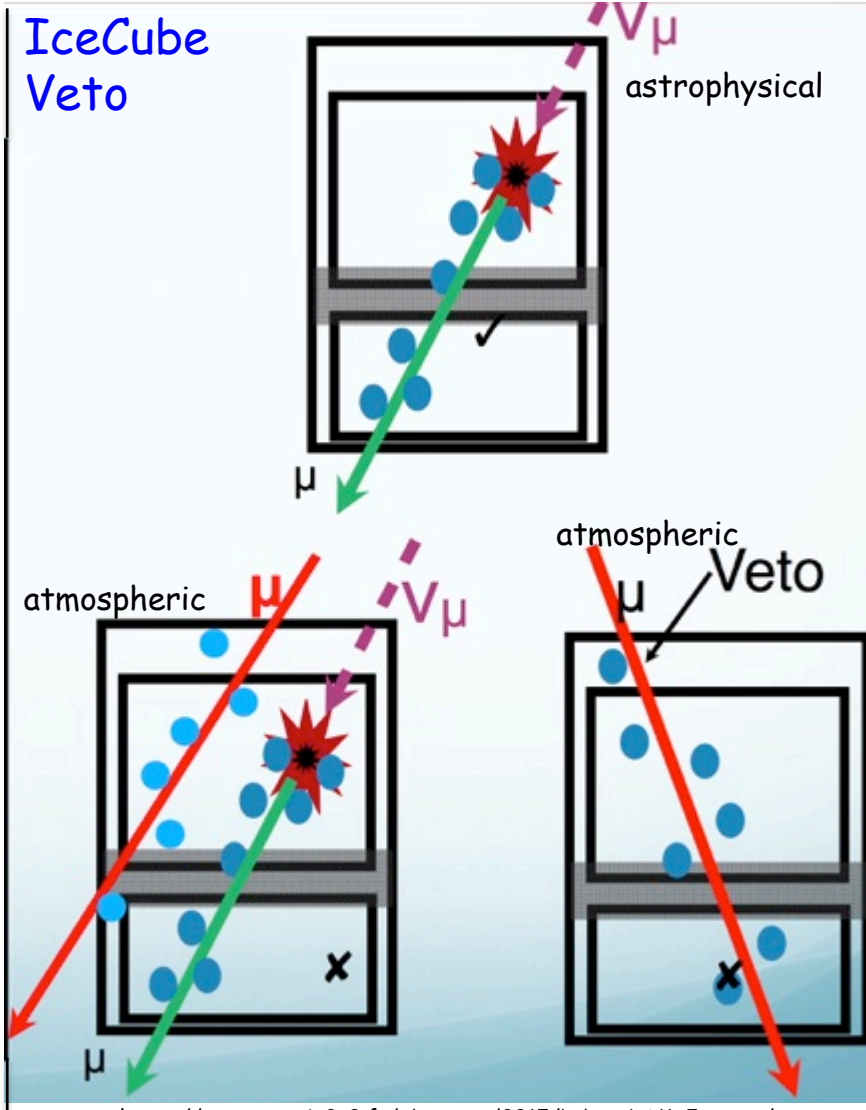
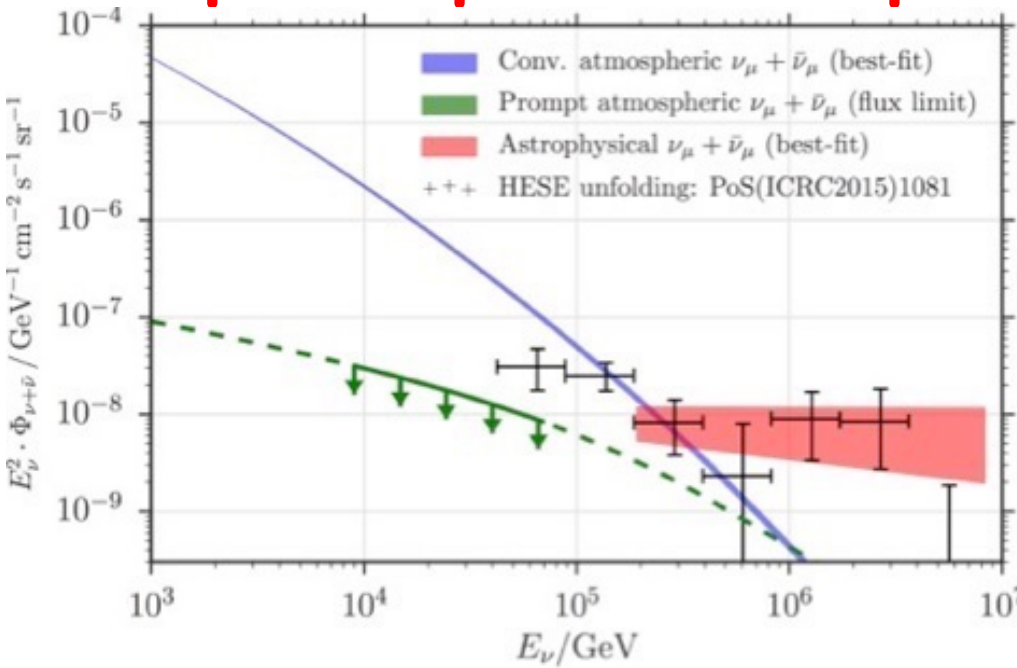
Conventional atmospheric  $\nu_e + \bar{\nu}_e$  flux is lower 😊 but the statistics are poor 😞

Conventional atmospheric  $\nu_\mu + \bar{\nu}_\mu$  flux is higher 😞 but the statistics are larger 😊

Present IceCube upper limits on atmospheric prompt neutrinos already constrain models of pQCD

Future limits on atmospheric prompt neutrinos will get stronger

# Astrophysical neutrinos $\nu/s$ prompt atmospheric neutrinos



## Difference in spectrum

Astrophysical neutrinos have a harder spectrum compared to prompt atmospheric neutrinos

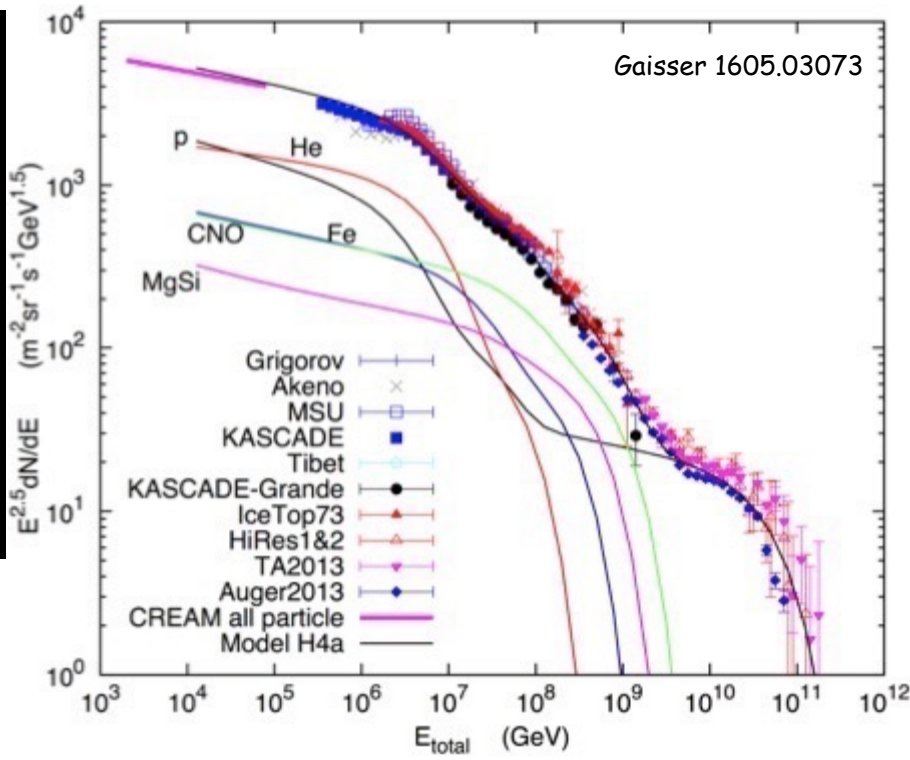
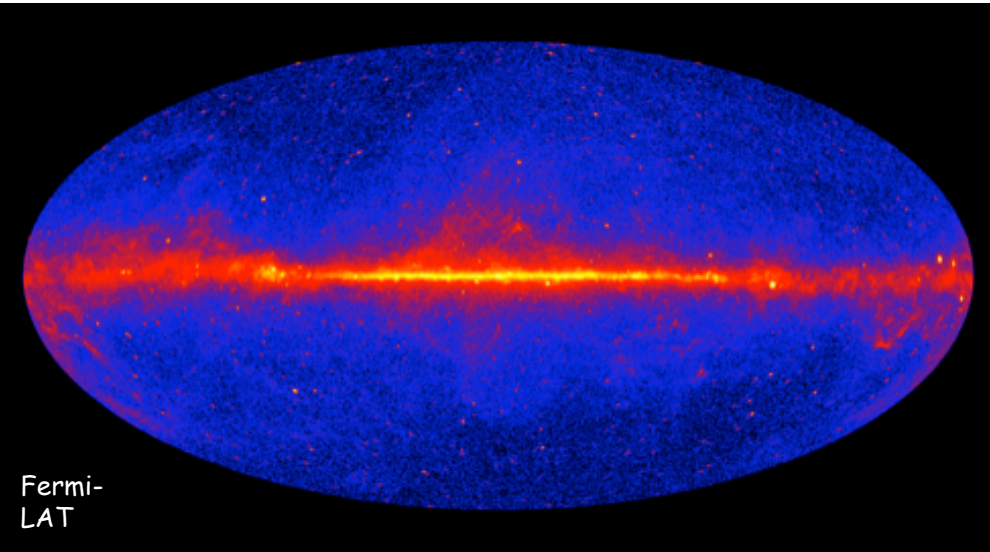
Conventional atmospheric  $\nu_e$  spectrum is lower  $\Rightarrow$  might be easier to search in that channel

[https://www.cppm.in2p3.fr/~brunner/2015/Labex-LAM-January/Talk\\_IceCube\\_Antares.pdf](https://www.cppm.in2p3.fr/~brunner/2015/Labex-LAM-January/Talk_IceCube_Antares.pdf)

# Conclusions

- High  $x_F$  physics important for astroparticle experiments (see for e.g., 1601.03044, 1605.01409, 1701.08451)
- Prompt atmospheric neutrinos are an important signal for IceCube
- Intrinsic charm can have important contribution to the prompt atmospheric neutrinos
- IceCube can constrain the intrinsic charm of the proton

# Puzzling questions about the high energy astrophysical universe



Gamma-ray sky: what process produces them?

Leptonic:  $e^- + \gamma \rightarrow e^- + \gamma$

Hadronic:  $p + p \rightarrow \pi^0 \rightarrow \gamma + \gamma$   
 $p + p \rightarrow \pi/K + \dots \rightarrow \nu/\bar{\nu}$

Cosmic rays observed over a huge energy range

Neutrinos are inevitably produced in cosmic ray interactions

The key difference are the neutrinos

# Neutrinos as cosmic messengers

- + No deflection from source
- + Can escape from very dense sources
- + No interaction on the way from source to detector
- + Complementary to gamma-rays

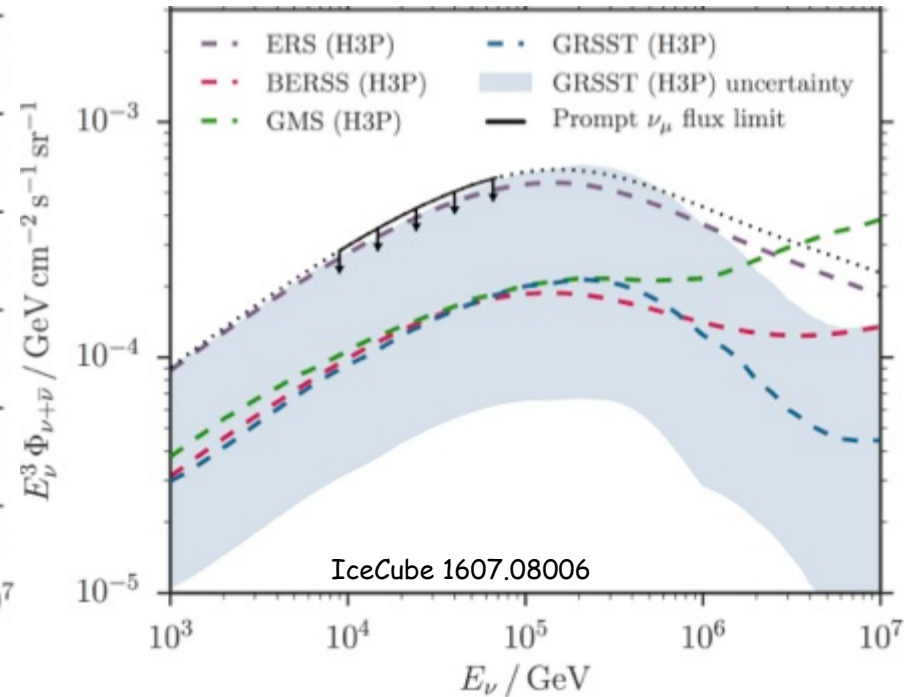
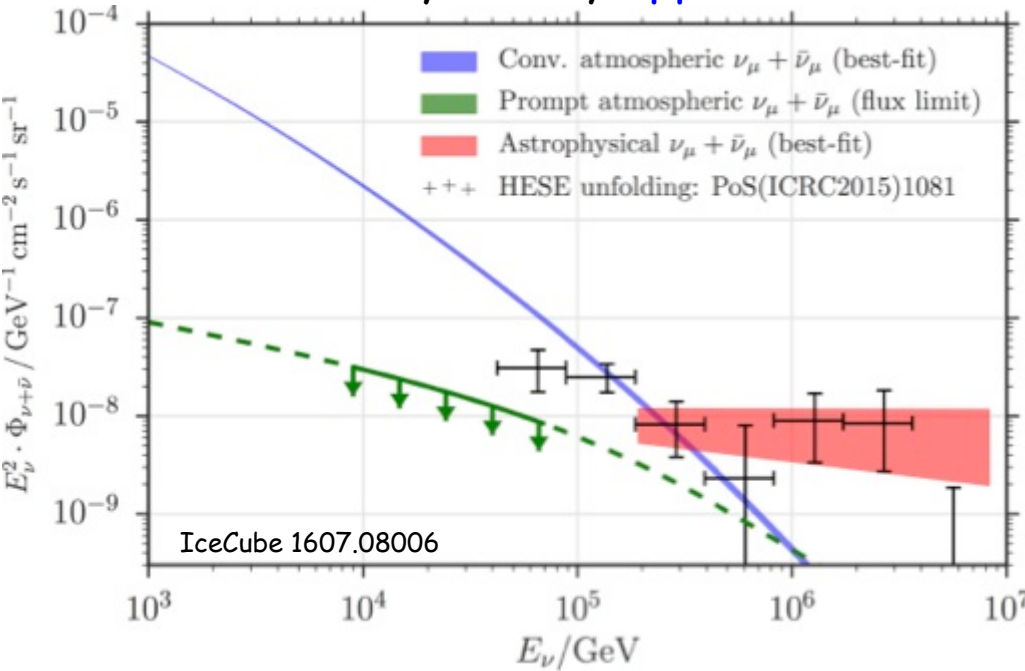
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- Large detectors required
- Very long time required to collect signal

.....

# Prompt atmospheric neutrinos

No detection yet, only **upper limits** exist on prompt atmospheric neutrinos



Upper limits are at the **upper end** of present perturbative QCD predictions

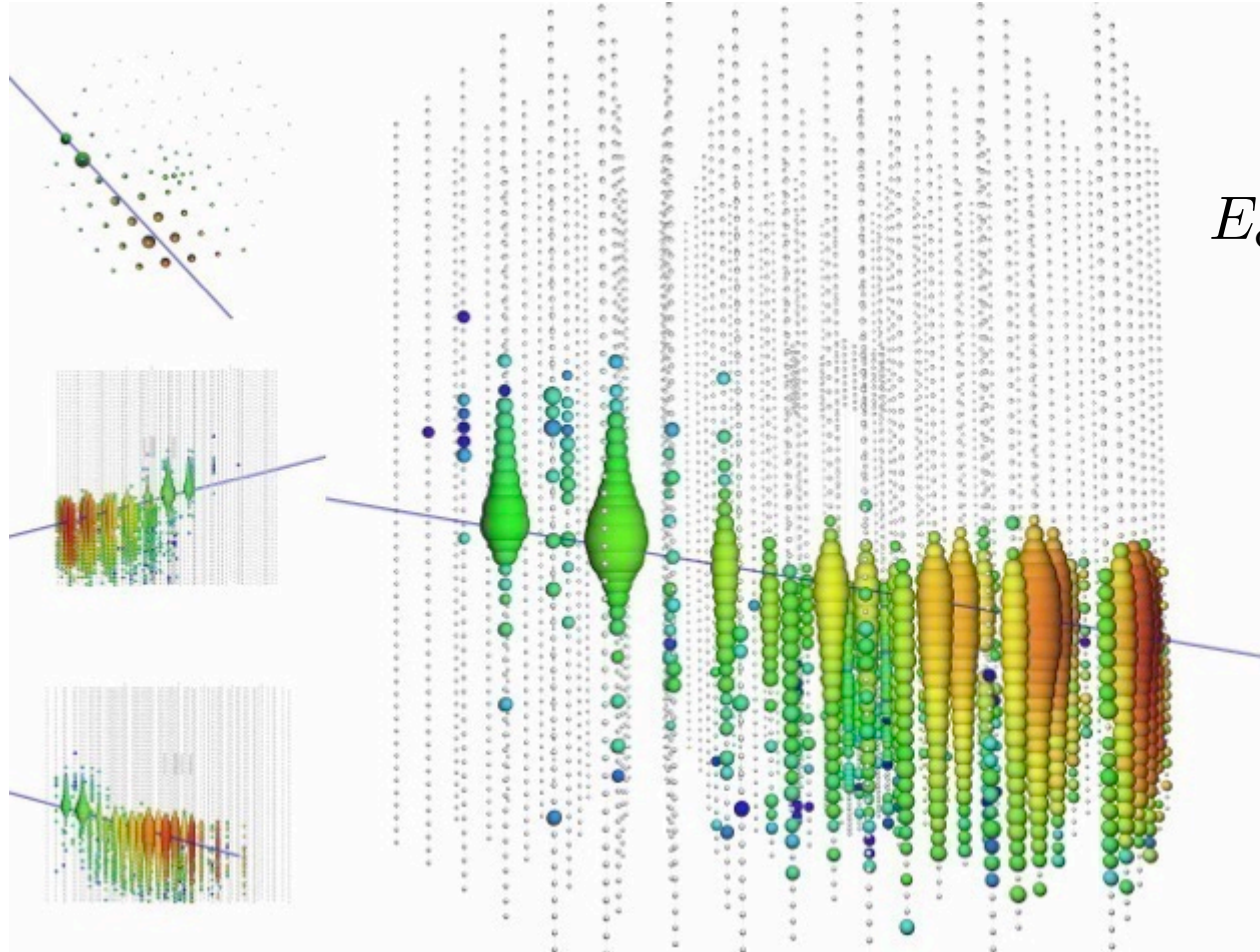
Already **constrains** mechanisms which predict a higher prompt atmospheric fluxes

The **lower** and **higher energy** neutrino spectrum dominated by conventional atmospheric neutrinos and astrophysical neutrinos respectively

Limits on the prompt atmospheric neutrinos depend on **astrophysical neutrino flux** modeling

# IceCube neutrino telescope

# Most energetic astrophysical neutrino event ever!!!



Deposited energy

$$E_{\text{dep}} = 2.6 \pm 0.3 \text{ PeV}$$

Where does it come from?

Does not correspond to any TeV source

Equatorial coordinates:  
dec  $11.42^\circ$  RA  $110.63^\circ$

IceCube 1607.08006

What is the energy of the parent neutrino?

If  $\nu_\mu$  produces this track, the neutrino energy is  $\sim 10 \text{ PeV}$

If  $\nu_\tau$  produces this track, the neutrino energy is  $\sim 100 \text{ PeV}$



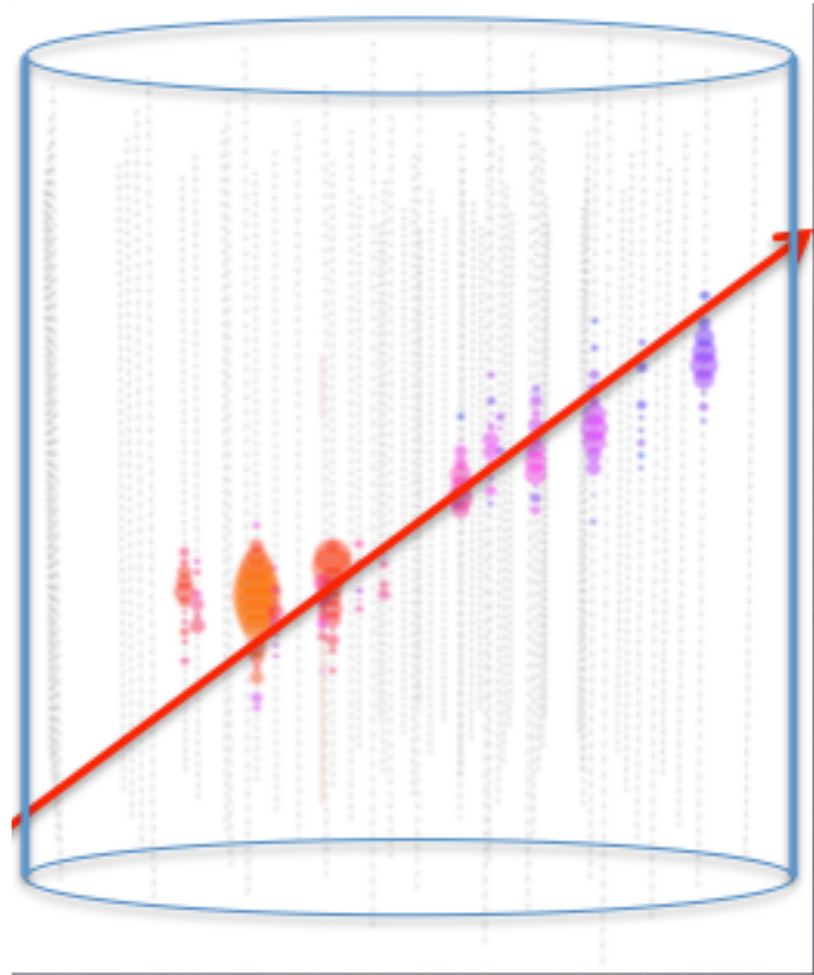
# Muon Tracks

Caused by **muons** produced in **CC** interaction of  $\nu_\mu/\bar{\nu}_\mu$

+ Long range implies **larger effective volume**

+ Better Angular resolution  $\sim 1^\circ$

- **Higher atmospheric neutrino backgrounds**



$$\left. \frac{dN_\mu}{dE_\mu} \right|_{\text{tracks}} = \frac{N_A \rho T A_{\text{det}}}{\rho(\alpha + \beta E_\mu)} \times \int_{E_\mu}^{\infty} dE_\nu \frac{d\Phi_\nu}{dE_\nu} \sigma_{\text{CC}}(E_\nu) e^{-\frac{L}{\lambda}}$$

# Cascades

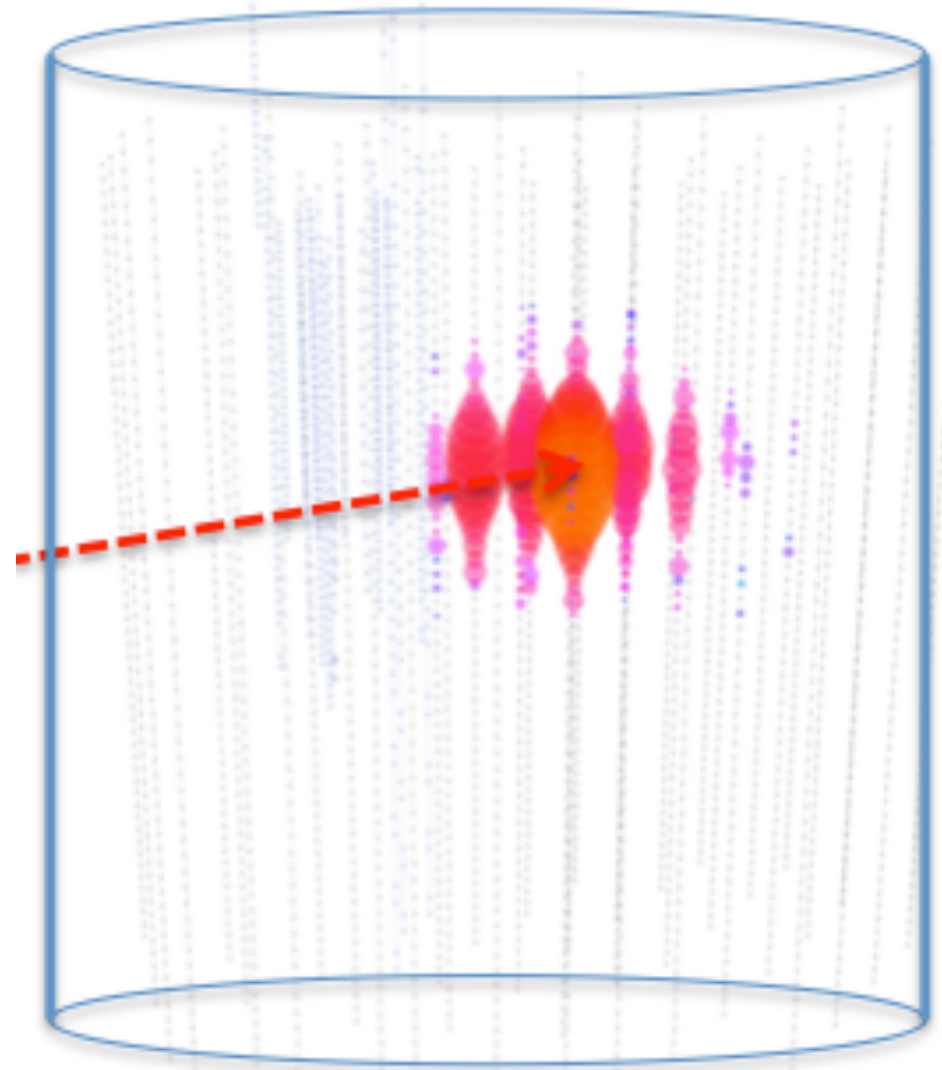
Caused by **CC** interactions of  $\nu_e/\nu_\tau$  and their antiparticles and **NC** interactions of neutrinos of all flavors

+ **Calorimetric**

+ **Lower** atmospheric neutrino background

- **Smaller** effective volume

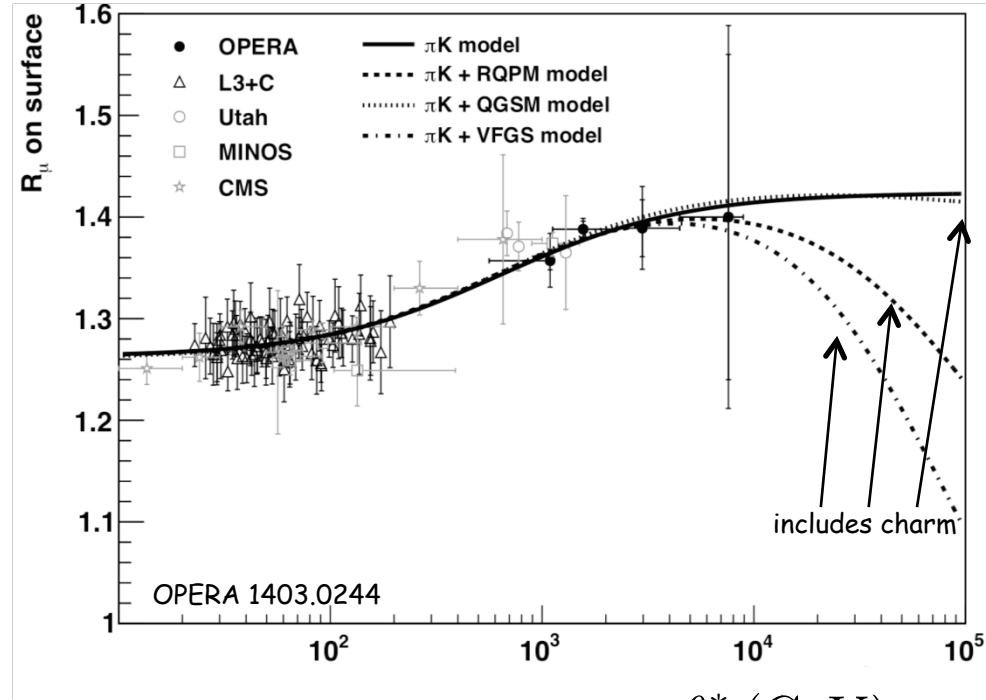
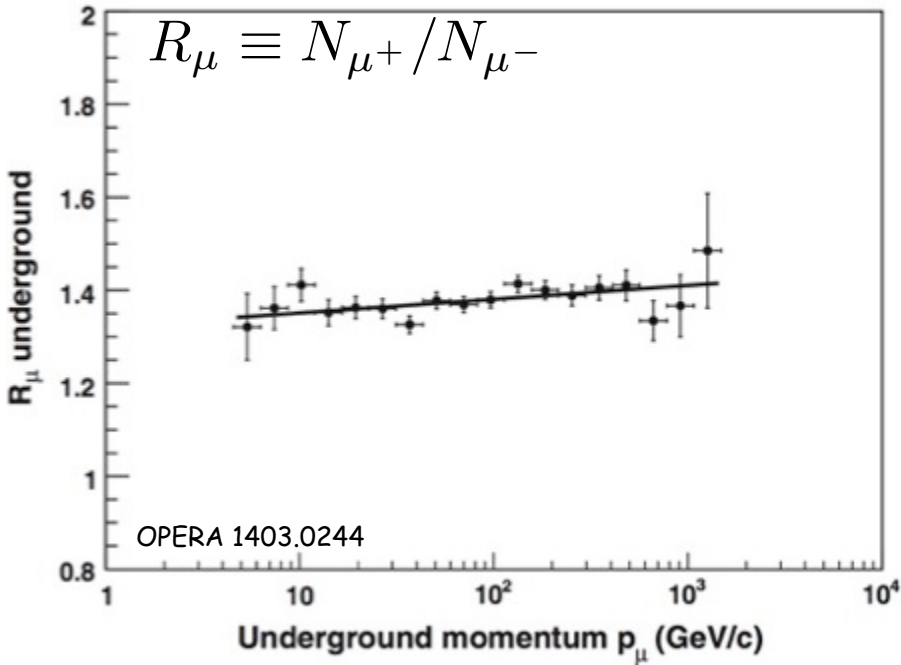
- **Poor** angular resolution  $\sim 50^\circ$



$$\left. \frac{dN_\nu}{dE_\nu} \right|_{\text{casc}} = N_A T V_{\text{casc}} \times \left( \sigma_{\text{CC}}(E_\nu) \frac{d\Phi_{\nu_{e,\tau}}}{dE_\nu} + \sigma_{\text{NC}}(E_\nu) \frac{d\Phi_{\nu_{e,\mu,\tau}}}{dE_\nu} \right)$$

# Atmospheric muon charge ratio and high $x_F$ physics

# Atmospheric muon charge ratio



CR + air  $\rightarrow \pi^\pm / K^\pm \dots \rightarrow \mu^\pm \nu_\mu / \bar{\nu}_\mu$

Muon energy at the surface  $\rightarrow \varepsilon_\mu \cos \theta^*$  (GeV)

Zenith angle at the muon production point

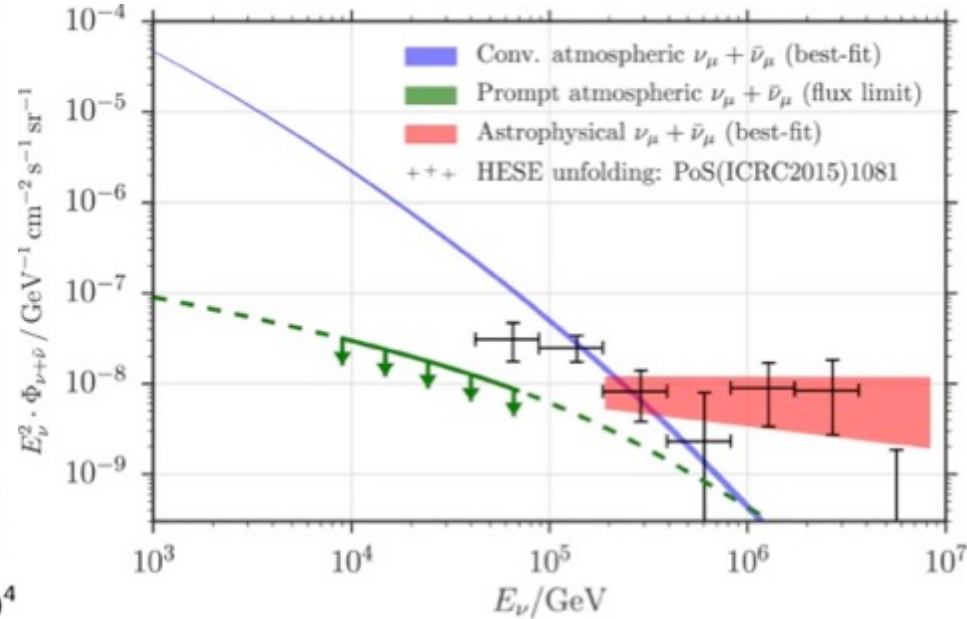
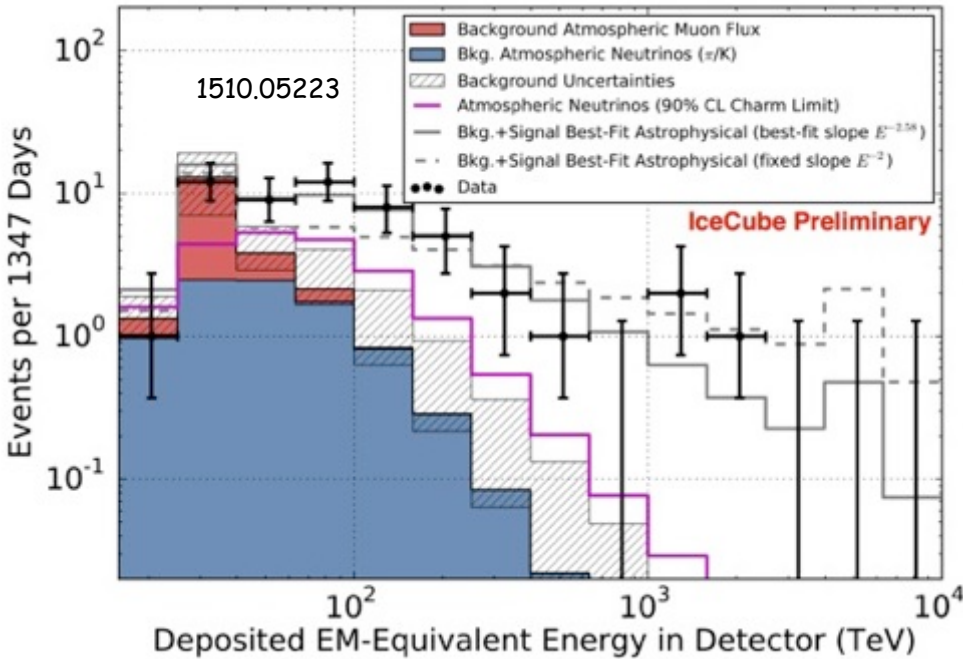
**Sensitive probe** of cosmic ray physics and hadronic interactions

Plateau implies asymptotic kaon contribution  $\Rightarrow$  **validity of Feynman scaling** up to primary energies/ nucleon around 200 TeV

Associated production of  $\Lambda K^+$  important; no analogue for  $K^-$ ; constrains the cross section in the **high  $x_F$  region** (difficult to probe via laboratory experiments) Gaisser 1111.6675

Interplay of **high  $x_F$  physics** and **astroparticle experiments**

# IceCube data



Diffuse spectrum of neutrinos

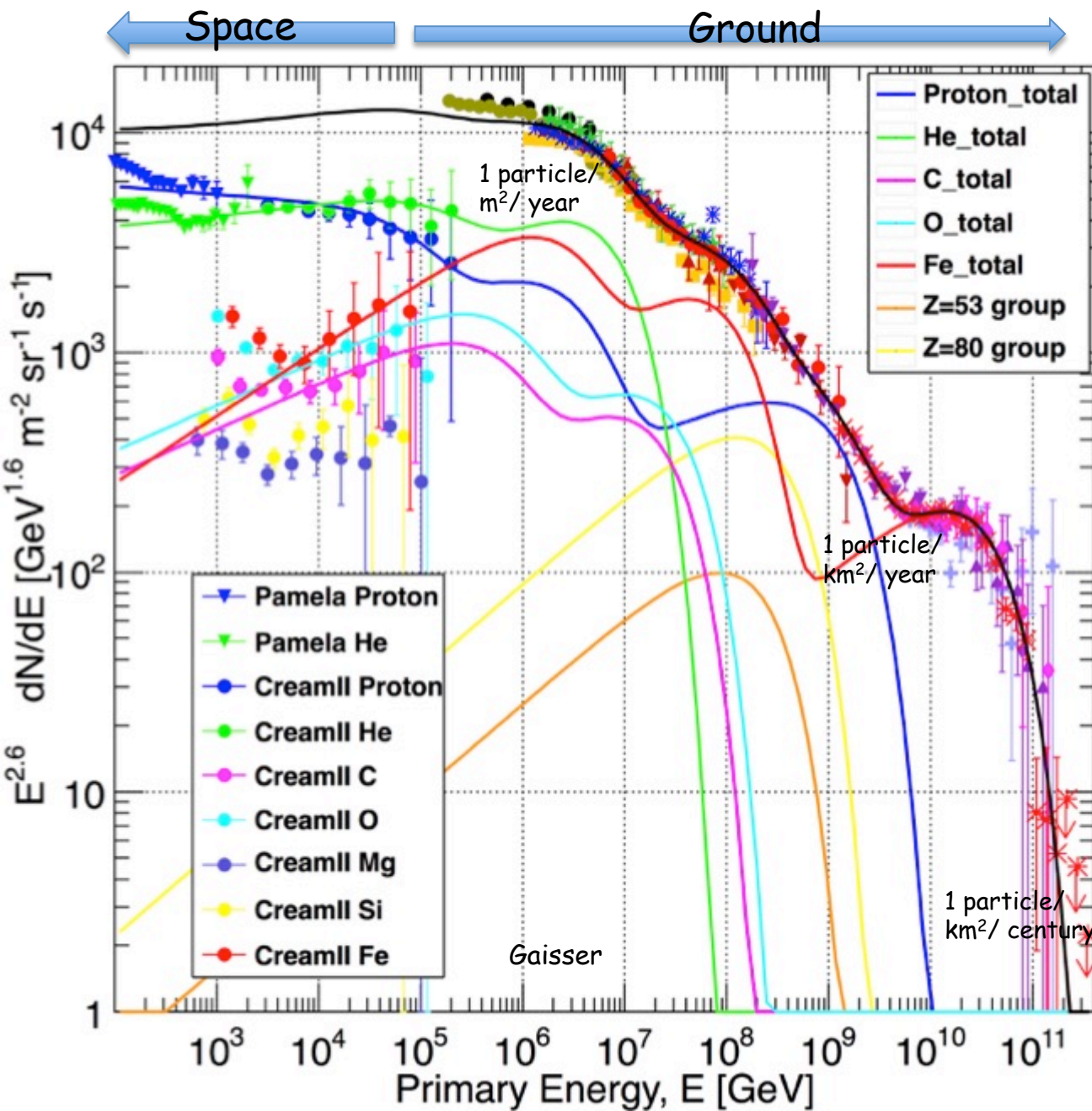
Time-independent

Clear evidence of the **astrophysical** nature of these neutrinos

None of them **point** to a specific source

# Cosmic rays

# Observed spectrum of cosmic rays



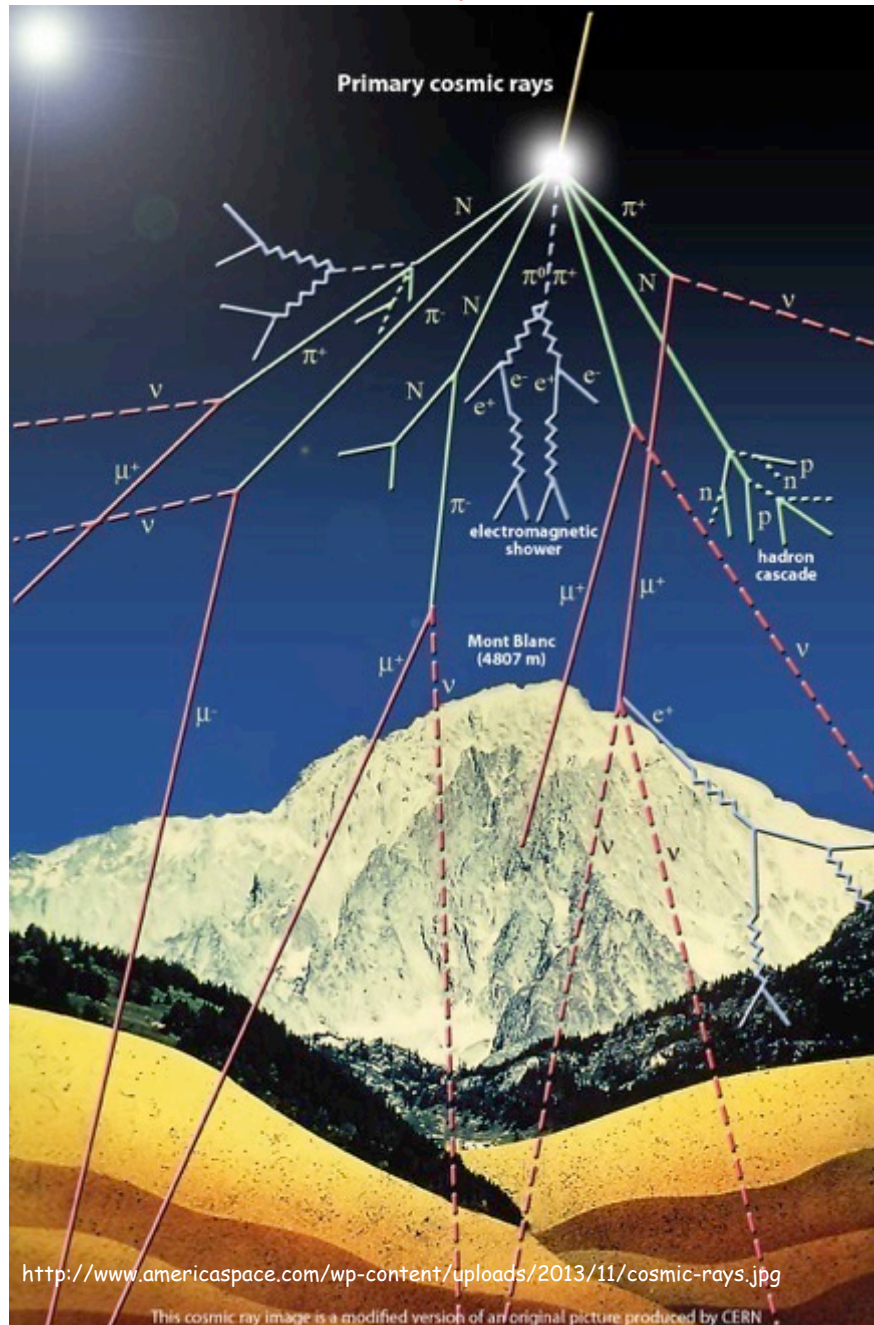
Spectrum of nuclei observed over many orders of magnitude in energy and flux

Observed via direct measurements (balloon/space observatory) and air shower arrays

Tremendous amount of data to understand the acceleration mechanism

Knowledge of QCD needed to understand the interaction of cosmic rays with the atmosphere

# Cosmic ray interactions in the atmosphere



Cosmic ray interaction in the atmosphere is a **fixed target experiment**

Detailed knowledge of **QCD** required to understand the rates of various interaction products

Sensitive to QCD parameter space which is difficult to access in the laboratory

For e.g., for  $E_\nu \approx 8000 \text{ TeV}$ ,  $\sqrt{s} \approx 20 \text{ TeV}$

High energy pp interactions sensitive to gluon density in the **small  $x$**  ( $\approx 10^{-8} - 10^{-4}$ ) region

Due to the **flattish  $d\sigma/dx_F$** , **intrinsic charm** uses the incoming proton energy quite efficiently  $\Rightarrow$  important since cosmic ray spectrum is falling