

Iterative Monte Carlo analysis of spin-dependent parton distributions

(arXiv:1601.07782)

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In Collaboration with:

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Trento Workshop 2016

Motivations

Spin structure of nucleons

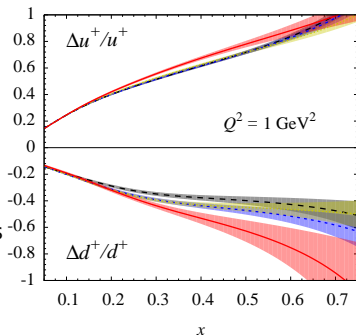
- spin carried by a quark of type $q \rightarrow \frac{1}{2}\Delta q^{(1)} = \frac{1}{2} \int_0^1 dx \Delta q(x, Q^2)$
- spin sum rule $\rightarrow \frac{1}{2} = \frac{1}{2}\Delta\Sigma^{(1)} + \Delta g^{(1)} + \mathcal{L} \rightarrow$ How big is \mathcal{L} ?
- From existing global analysis $\rightarrow \Delta\Sigma_{[10^{-3},1]}^{(1)} \sim 0.3, \Delta g_{[0.05,0.2]}^{(1)} \sim 0.1$

High x

- $SU(6)$ spin-flavor symmetry:
 $\rightarrow \Delta u/u \rightarrow 2/3, \Delta d/d \rightarrow -1/3$
- pQCD $\rightarrow \Delta q/q \rightarrow 1$

Higher twists

- d_2 matrix element
 $\rightarrow d_2 = 2g_1^{(3)}(Q^2) + 3g_2^{(3)}(Q^2)$
- Color forces experienced by struck quarks
 $\rightarrow \tilde{F}_E = 2d_2 + f_2, \tilde{F}_B = 4d_2 - f_2$



Data

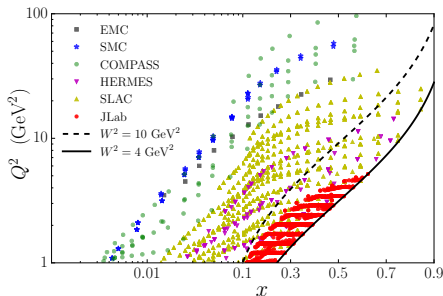
- ✓ Polarized DIS $\rightarrow \Delta u^+, \Delta d^+$
- Polarized SIDIS: $\rightarrow \Delta \bar{d}, \Delta \bar{u}, \Delta s$
- Inclusive Jets/ π^0 : $\rightarrow \Delta g$
- W production $\rightarrow \Delta \bar{d}, \Delta \bar{u}$

Theory

- ✓ Target mass corrections
- ✓ Twist-3 and twist-4 contributions in polarized structure functions
- ✓ Nuclear corrections for ^3He and deuteron targets
- Threshold resummation $\rightarrow (\alpha_S^m \log(1-x))^n$

Tools

- ✓ Numerical codes developed within python framework
- ✓ Development of DGLAP evolution equations in Mellin space
- ✓ Fast calculation of observables \rightarrow Mellin space techniques



Asymmetries

$$A_{\parallel} = \frac{\sigma^{\uparrow\downarrow} - \sigma^{\downarrow\downarrow}}{\sigma^{\uparrow\downarrow} + \sigma^{\downarrow\downarrow}} = D(A_1 + \eta A_2)$$

$$A_{\perp} = \frac{\sigma^{\uparrow\Rightarrow} - \sigma^{\downarrow\Rightarrow}}{\sigma^{\uparrow\Rightarrow} + \sigma^{\downarrow\Rightarrow}} = d(A_2 - \xi A_1)$$

$$A_1 = \frac{(g_1 - \gamma^2 g_2)}{F_1} \quad A_2 = \gamma \frac{(g_1 + g_2)}{F_1} \quad \gamma^2 = \frac{4M^2 x^2}{Q^2}$$

Polarized structure functions

$$g_1(x, Q^2) = g_1^{\text{LT+TMC}}(\Delta u^+, \Delta d^+, \Delta g, \dots) + g_1^{\text{T3+TMC}}(D_u, D_d) + g_1^{\text{T4}}(H_{p,n})$$

$$g_2(x, Q^2) = g_2^{\text{LT+TMC}}(\Delta u^+, \Delta d^+, \Delta g, \dots) + g_2^{\text{T3+TMC}}(D_u, D_d)$$

Parametrization

- $xf(x) = Nx^a(1-x)^b(1+c\sqrt{x}+dx)$
- LT quark distributions $\rightarrow \Delta u^+, \Delta d^+, \Delta s^+, \Delta g$
- T3 quark distributions $\rightarrow D_u, D_d$
- T4 structure functions $\rightarrow H_p, H_n$

Chi-squared minimization \rightarrow with correlated systematic uncertainties

$$\chi^2 = \sum_i \left(\frac{D_i - T_i(1 - \sum_k r^k \beta_i^k / D_i)^{-1}}{\alpha_i} \right)^2 + \sum_k (r^k)^2$$

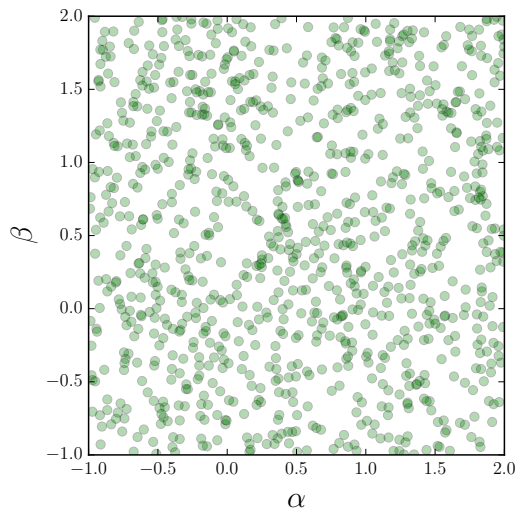
Issues

- Stability in the moments (e.g. $\Delta\Sigma^{(1)}$)
- Is the solution given by a single fit unique? \rightarrow False minima
- Is over-fitting present in our fits?
- Which parameters should be fixed and at which value?
- Determination of uncertainty bands.

Solution \rightarrow MC approach

Iterative Monte Carlo Analysis (IMC)

Toy example \rightarrow fitting 2 model parameters α, β

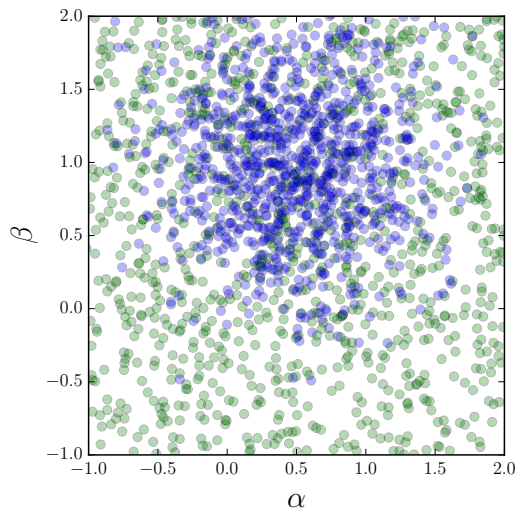


I. Flat sampling

Initial priors $\{(\alpha, \beta)\}$

Iterative Monte Carlo Analysis (IMC)

Toy example \rightarrow fitting 2 model parameters α, β



I. Flat sampling

Initial priors $\{(\alpha, \beta)\}$

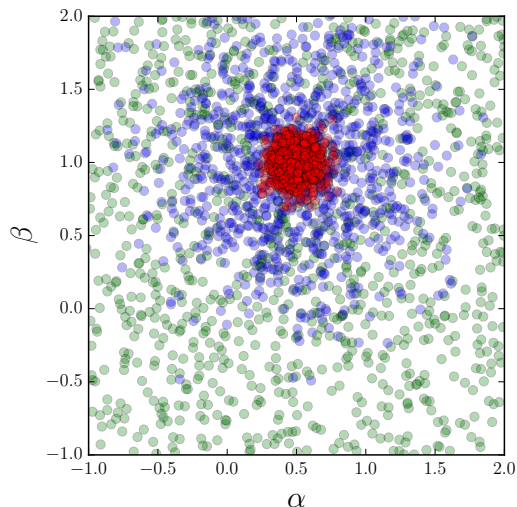
II. First iteration

priors $\{(\alpha, \beta)\}$

posteriors $\{(\alpha, \beta)\}$

Iterative Monte Carlo Analysis (IMC)

Toy example \rightarrow fitting 2 model parameters α, β



I. Flat sampling

Initial priors $\{(\alpha, \beta)\}$

II. First iteration

priors $\{(\alpha, \beta)\}$

posteriors $\{(\alpha, \beta)\}$

III. Second iteration

priors $\{(\alpha, \beta)\}$

posteriors $\{(\alpha, \beta)\}$

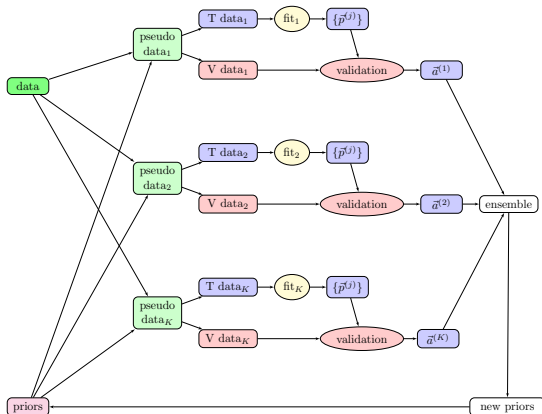
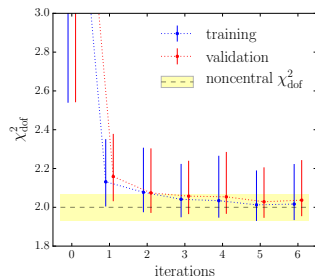
... until convergence

Iterative Monte Carlo Analysis (IMC)

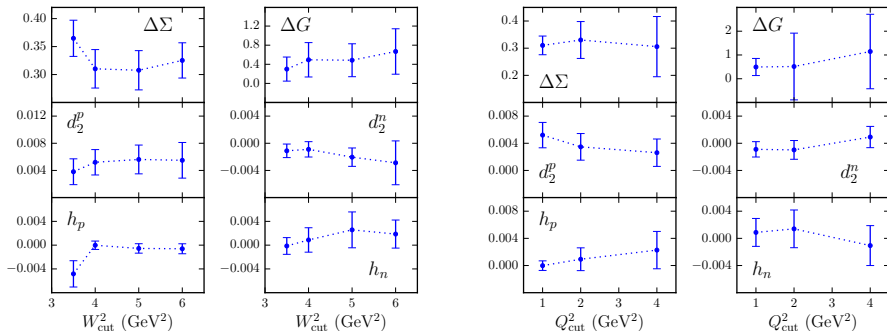
Each iteration

- Generate pseudo data sets via data resampling
- Random data partition \rightarrow Training & Validation
- Fit the training set
- Validation

of fits: 10000



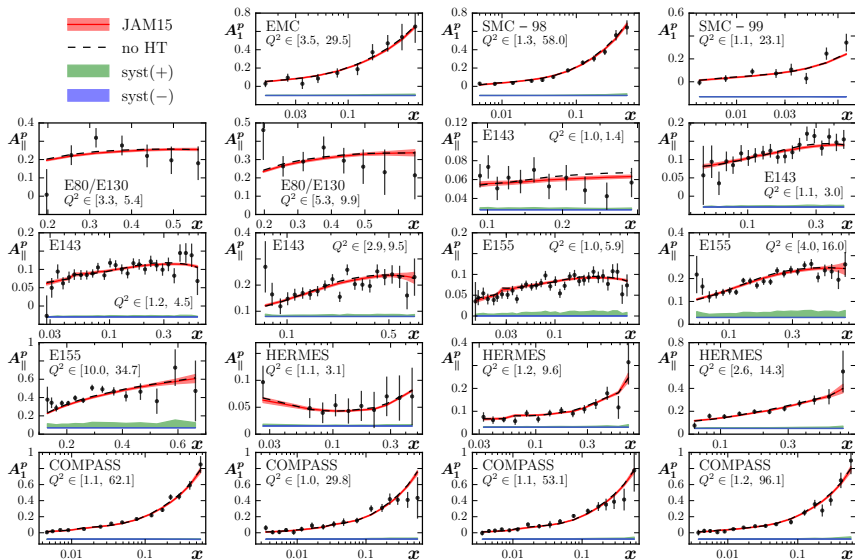
W_{cut}^2 and Q_{cut}^2



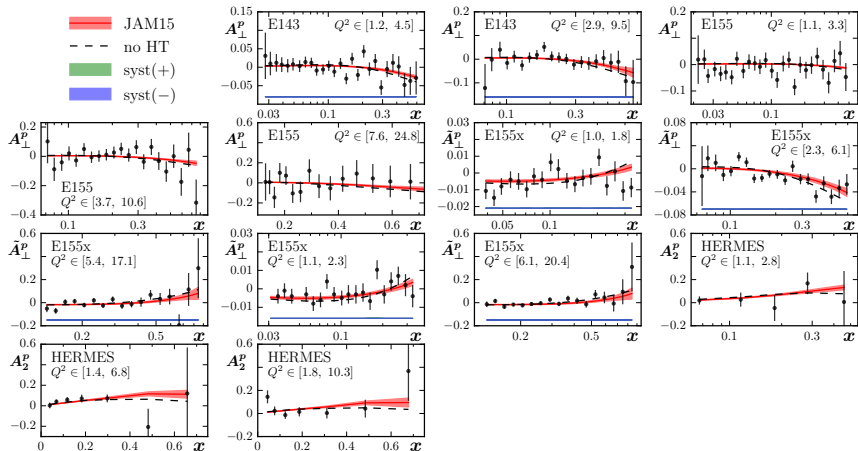
W_{cut}^2 (GeV ²)	3.5	4	5	6	8	10
# points	2868	2515	1880	1427	943	854
χ_{dof}^2	1.20	1.07	1.03	1.02	0.99	0.97

Q_{cut}^2 (GeV ²)	1.0	2.0	4.0
# points	2515	1421	611
χ_{dof}^2	1.07	1.08	0.95

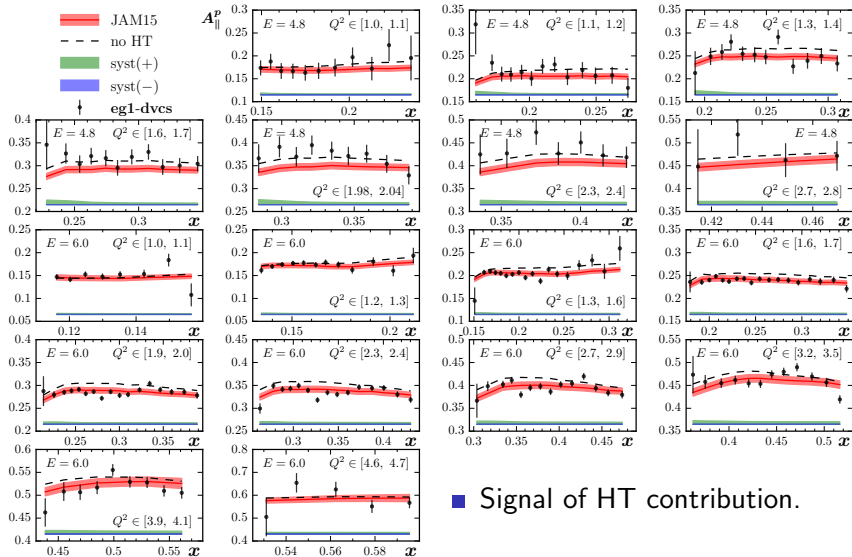
Data vs theory: proton



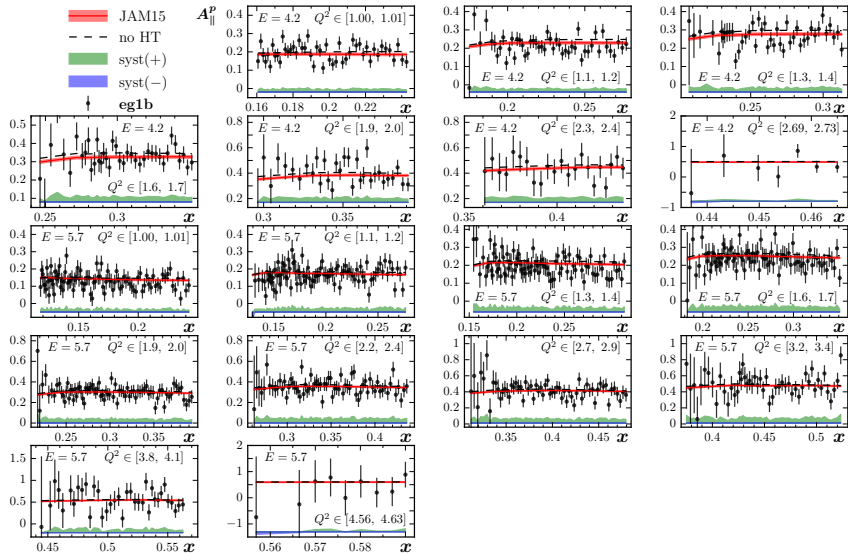
Data vs theory: proton



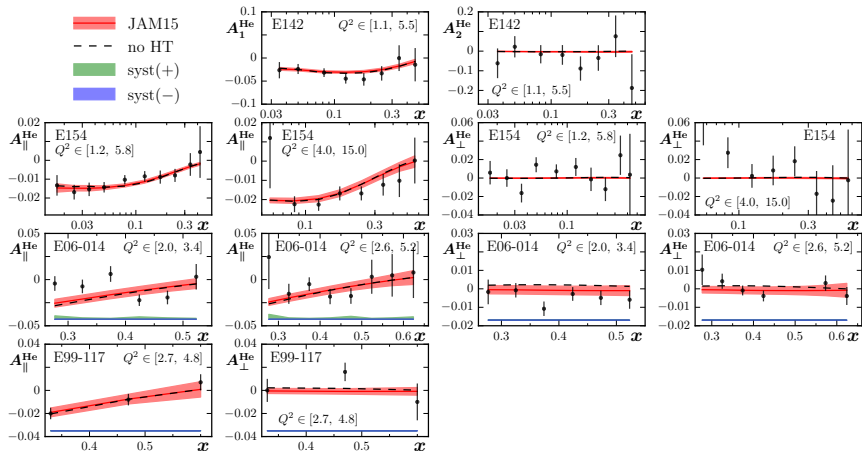
Data vs theory: proton JLab eg1b-dvcs



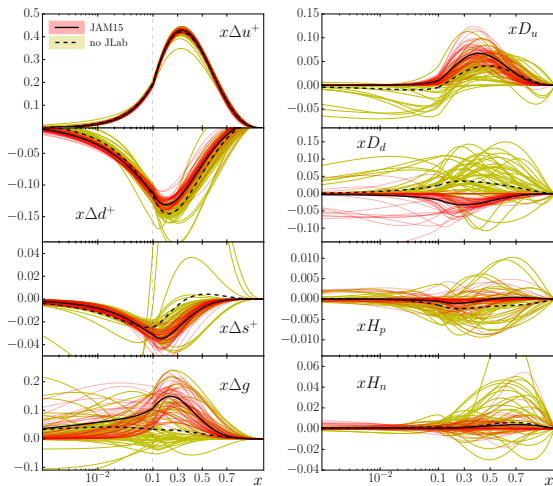
Data vs theory: proton JLab eg1b



Data vs theory: ^3He

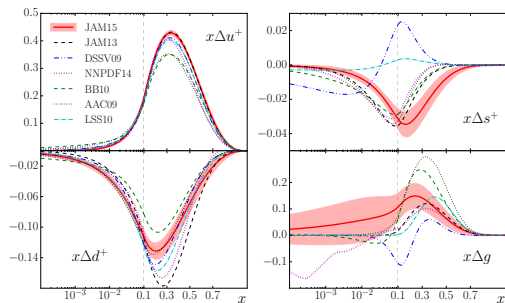
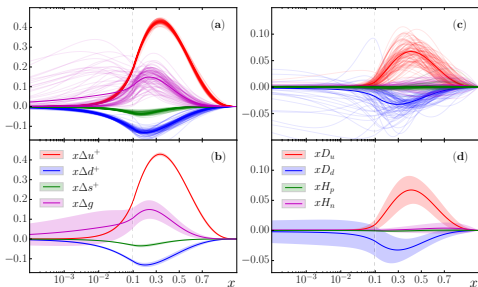


Impact of JLab data



- JLab data $\rightarrow 0.1 < x < 0.7$
- Constraints on small x from large $x \rightarrow$ weak baryon decay constraints
- Large uncertainties in Δs^+ , Δg removed by JLab data
- Non vanishing T3 quark distributions
- T4 distributions consistent with zero

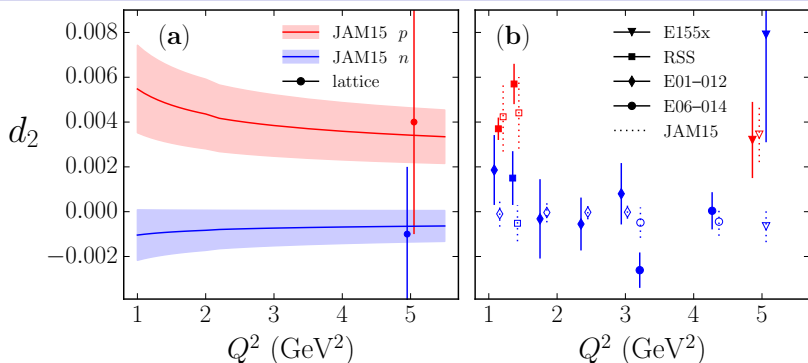
Results



moment	truncated	full
Δu^+	0.82 ± 0.01	0.83 ± 0.01
Δd^+	-0.42 ± 0.01	-0.44 ± 0.01
Δs^+	-0.10 ± 0.01	-0.10 ± 0.01
$\Delta \Sigma$	0.31 ± 0.03	0.28 ± 0.04
ΔG	0.5 ± 0.4	1 ± 15
d_2^p	0.005 ± 0.002	0.005 ± 0.002
d_2^n	-0.001 ± 0.001	-0.001 ± 0.001
h_p	-0.000 ± 0.001	0.000 ± 0.001
h_n	0.001 ± 0.002	0.001 ± 0.003

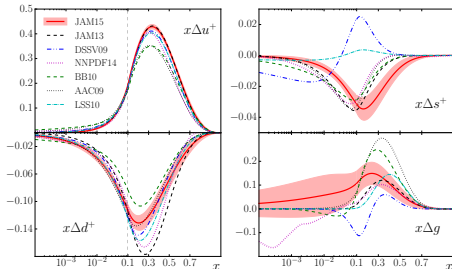
- Significant constraints on Δs^+ and Δg
- Non zero T3 quark distributions
- T4 contribution to g_1 consistent with zero
- **Negative Δs^+**
- JAM15 Δg compatible with recent DSSV fits.

d_2 matrix element



- $d_2(Q^2) \equiv \int_0^1 dx x^2 [2g_1^{T^3}(x, Q^2) + 3g_2^{T^3}(x, Q^2)]$
- d_2 is related to “color polarizability” or the “transverse color force” acting on quarks.
- Existing measurements of d_2 are in the resonance region (contains TMC T4 and beyond.)
- Agreement with data indicates quark-hadron duality

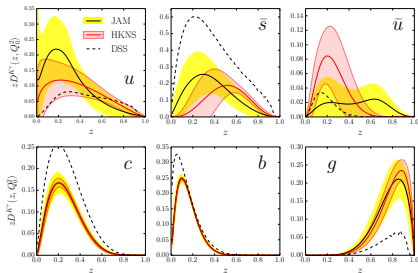
The strange puzzle



- The sign of ΔS^+ from combined DIS and SIDIS depends on Kaon fragmentation: positive for DSS and negative for HKNS. (Leader, etal)

**NEW IMC FF analysis
(to be published soon)**

- Only SIA data is used : npts=245, $\chi^2 = 305.2$ (identical to HKNS)
- Size of Δs^+ similar to HKNS but smaller than DSS
- Combined DIS and SIDIS analysis unlikely to change Δs^+

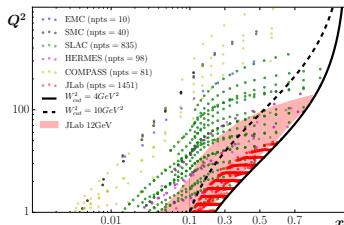


JAM

- ✓ New JAM15 analysis to study impact of all JLab 6 GeV inclusive DIS data at low W and high x
- ✓ New extraction of LT & HT distributions
 - **Upcoming JAM16** analysis to study polarization of sea quarks & gluons.
 - SIDIS for flavor separation.
 - polarized pp cross sections (inclusive jet & π production) for Δg
 - W boson asymmetries
 - Threshold resummation impacts on large x
 - Combined analysis of all inclusive (un)polarized DIS data
 - Fits to helicity distributions

JLab 12

- Measurements at high- $x \rightarrow \Delta q/q$
- Wider coverage in $Q^2 \rightarrow \Delta g$
- Determination of pure twist-3 d_2 in DIS



TMD Cross sections in SIDIS

In Collaboration with:

T. Rogers, B. Wang,

A. Prokudin, L. Gamberg, D.B Clarki, O. Gonzalez

Trento Workshop 2016

Formal QCD description of TMD cross sections

$$\begin{aligned}\Gamma &= \mathbf{T}_{\text{TMD}}\Gamma + [\Gamma - \mathbf{T}_{\text{TMD}}\Gamma] \\ &\approx \underbrace{\mathbf{T}_{\text{TMD}}\Gamma}_{\text{W}} + \underbrace{\mathbf{T}_{\text{coll}}[\Gamma - \mathbf{T}_{\text{TMD}}\Gamma]}_{\text{Y}}\end{aligned}$$

- Γ = A generic differential Cross section (i.e $d\sigma/dq_T$)
- \mathbf{T}_{TMD} = Small q_T approximant
- \mathbf{T}_{coll} = Large q_T approximant

Region of $q_T \ll Q$

- TMD approx. dominates $\rightarrow \Gamma \approx \mathbf{T}_{\text{TMD}}\Gamma$
- Y term small



Region of $q_T \gtrsim Q$

- Collinear approx. dominates $\rightarrow \Gamma \approx \mathbf{T}_{\text{coll}}\Gamma$
- At large Q $\mathbf{T}_{\text{TMD}}\Gamma$ is mostly perturbative



Formal QCD description of TMD cross sections

$$\begin{aligned}\Gamma &= \mathbf{T}_{\text{TMD}}\Gamma + [\Gamma - \mathbf{T}_{\text{TMD}}\Gamma] \\ &\approx \underbrace{\mathbf{T}_{\text{TMD}}\Gamma}_{\mathbf{W}} + \underbrace{\mathbf{T}_{\text{coll}}[\Gamma - \mathbf{T}_{\text{TMD}}\Gamma]}_{\mathbf{Y}}\end{aligned}$$

Standard recipe:

- Use CSS to calculate \mathbf{W} (TMD evolution)
- Use standard collinear factorization for $\mathbf{FO} = \mathbf{T}_{\text{coll}}\Gamma$
- Use Asymptotic expansion of \mathbf{FO} to get $\mathbf{ASY} = \mathbf{T}_{\text{coll}}\mathbf{T}_{\text{TMD}}\Gamma$

The state of the art for SIDIS \rightarrow Gaussian approach

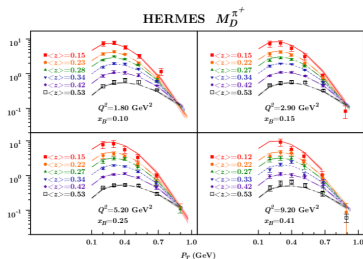
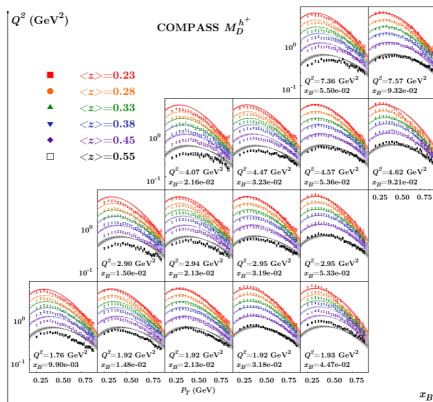
- Don't use CSS to calculate \mathbf{W} (TMD evolution)
 \rightarrow use Gaussian model for \mathbf{W} (DGLAP evolution)
- Not include $\mathbf{FO} = \mathbf{T}_{\text{coll}}\Gamma$
- Not include $\mathbf{ASY} = \mathbf{T}_{\text{coll}}\mathbf{T}_{\text{TMD}}\Gamma$

Success of the gaussian approach: \rightarrow Anselmino *et al.*

HERMES \rightarrow

Points: ~ 497

$$\chi^2_{dof} = 1.69$$



\leftarrow **COMPASS**

Points: ~ 5385

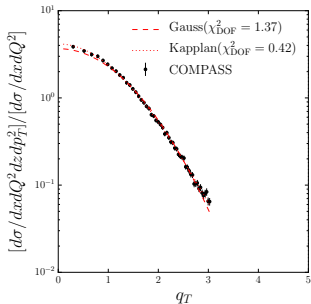
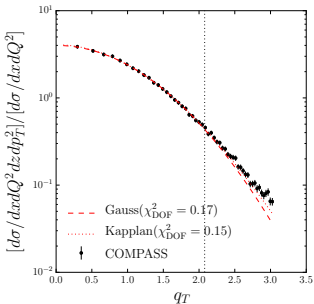
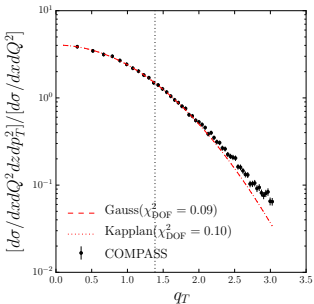
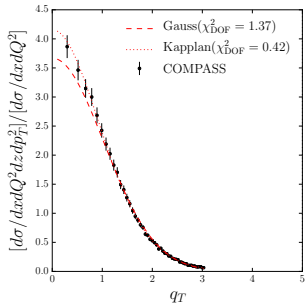
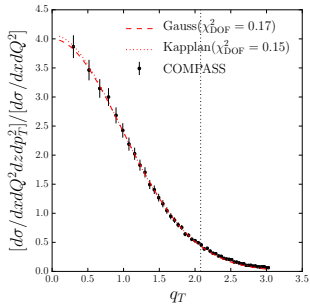
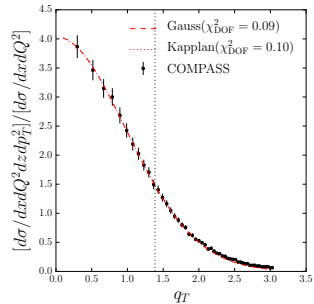
$$\chi^2_{dof} = 8.54$$

$$M \propto \sum_q f_{q/p}(x_B) D_{h/q}(z_h) \frac{e^{-P_T^2 / \langle P_T^2 \rangle}}{\pi \langle P_T^2 \rangle}$$

Despite the success some questions remains

- How consistent is to use DGLAP evolution in the Gaussian approach?
- Can we ignore the information of the high q_T tail? Are the TMD widths independent from the tails?
- Is the gaussian model the best behavior for $W(b_T)$?
 - Gaussian $\sim \exp[-mb_T^2]$
 - Kaplan $\sim \exp[-mb_T] \rightarrow$ More consistent with field theory
- Can data really discriminate between gaussian and other form? (i.e. Kaplan form)

Open questions



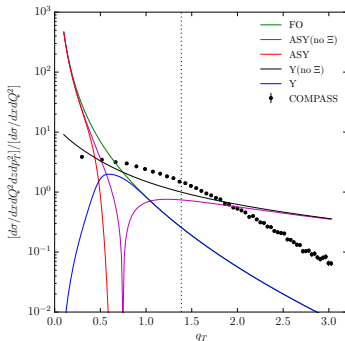
Formal QCD description of TMD cross sections

$$\Gamma = \mathbf{T}_{\text{TMD}}\Gamma + [\Gamma - \mathbf{T}_{\text{TMD}}\Gamma]$$

$$\approx \underbrace{\mathbf{T}_{\text{TMD}}\Gamma}_{\text{W}} + \underbrace{\mathbf{T}_{\text{coll}}[\Gamma - \mathbf{T}_{\text{TMD}}\Gamma]}_{\text{Y}}$$

Issues with the standard recipe:

- FO is too small. The NLO calculation needed. (A. Daleo et al.)
- $Y = \text{FO} - \text{ASY}$ is too big.
- Incomplete cancellation between W and ASY at large $q_T \rightarrow$ new definition of $\mathbf{T}_{\text{TMD}}^{\text{New}}$ (T. Rogers talk)



$$Q^2 = 1.92 \text{ GeV}^2, x = 0.0318, z = 0.375$$

Backup

Curse of dimensionality → Mellin trick (Stratmann, Vogelsang)

$$\begin{aligned}
 I(x) &= \int_x^1 \frac{dy}{y} f(y) \int_y^1 \frac{dz}{z} g\left(\frac{x}{yz}\right) \quad \leftarrow \quad g(\xi) = \frac{1}{2\pi i} \int dN \xi^{-N} g_N \\
 &= \frac{1}{2\pi i} \int dN g_N \left[\int_x^1 \frac{dy}{y} f(y) \int_y^1 \frac{dz}{z} \left(\frac{x}{yz}\right)^{-N} \right] \\
 &= \frac{1}{2\pi i} \int dN g_N \mathcal{M}_N \\
 &= \sum_{i,k} w_i^k j^k \operatorname{Im} \left(e^{i\phi} g_{N_j^k} \mathcal{M}_{N_j^k} \right) \quad \leftarrow \quad \text{Gaussian quadrature}
 \end{aligned}$$

- Time consuming part can be precalculated prior to the fit
- Extensible to higher dimensional integrals.
- A single fit that takes about 4 days → ≈ **20 mins**

$$\begin{aligned} \rightarrow \xi &= \frac{2x}{1+(1+4\mu^2 x^2)^{1/2}} \\ \rightarrow \mu^2 &= M^2/Q^2 \end{aligned}$$

Leading twist structure functions:

$$\begin{aligned} g_1^{\text{LT+TMC}}(x, Q^2) &= \frac{x}{\xi} \frac{g_1^{\text{LT}}(\xi)}{(1+4\mu^2 x^2)^{3/2}} + 4\mu^2 x^2 \frac{x+\xi}{\xi(1+4\mu^2 x^2)^2} \int_{\xi}^1 \frac{dz}{z} g_1^{\text{LT}}(z) \\ &\quad - 4\mu^2 x^2 \frac{2-4\mu^2 x^2}{2(1+4\mu^2 x^2)^{5/2}} \int_{\xi}^1 \frac{dz}{z} \int_{z'}^1 \frac{dz'}{z'} g_1^{\text{LT}}(z') \\ g_2^{\text{LT+TMC}}(x, Q^2) &= -\frac{x}{\xi} \frac{g_1^{\text{LT}}(\xi)}{(1+4\mu^2 x^2)^{3/2}} + \frac{x}{\xi} \frac{(1-4\mu^2 x\xi)}{(1+4\mu^2 x^2)^2} \int_{\xi}^1 \frac{dz}{z} g_1^{\text{LT}}(z) \\ &\quad + \frac{3}{2} \frac{4\mu^2 x^2}{(1+4\mu^2 x^2)^{5/2}} \int_{\xi}^1 \frac{dz}{z} \int_{z'}^1 \frac{dz'}{z'} g_1^{\text{LT}}(z') \end{aligned}$$

In the Bjorken limit ($Q^2 \rightarrow \infty$):

$$g_1^{\text{LT+TMC}}(x, Q^2) \simeq g_1^{\text{LT}}(x), \quad g_2^{\text{LT+TMC}}(x, Q^2) \simeq -g_1^{\text{LT}}(x) + \int_{\xi}^1 \frac{dz}{z} g_1^{\text{LT}}(z)$$

Leading twist quark distributions:

$$g_1^{\text{LT}}(x) = \frac{1}{2} \sum_q e_q^2 [\Delta C_{qq} \otimes \Delta q(x) + \Delta C_{qg} \otimes \Delta g(x)]$$

Twist-3 structure functions:

$$\begin{aligned}
 g_1^{\text{T3+TMC}}(x, Q^2) &= 4\mu^2 x^2 \frac{D(\xi)}{(1 + 4\mu^2 x^2)^{3/2}} - 4\mu^2 x^2 \frac{3}{(1 + 4\mu^2 x^2)^2} \int_{\xi}^1 \frac{dz}{z} D(z) \\
 &\quad + 4\mu^2 x^2 \frac{2 - 4\mu^2 x^2}{(1 + 4\mu^2 x^2)^{5/2}} \int_{\xi}^1 \frac{dz}{z} \int_{z'}^1 \frac{dz'}{z'} D(z') \\
 g_2^{\text{T3+TMC}}(x, Q^2) &= \frac{D(\xi)}{(1 + 4\mu^2 x^2)^{3/2}} - \frac{1 - 8\mu^2 x^2}{(1 + 4\mu^2 x^2)^2} \int_{\xi}^1 \frac{dz}{z} D(z) \\
 &\quad - \frac{12\mu^2 x^2}{(1 + 4\mu^2 x^2)^{5/2}} \int_{\xi}^1 \frac{dz}{z} \int_{z'}^1 \frac{dz'}{z'} D(z')
 \end{aligned}$$

Bjorken limit ($Q^2 \rightarrow \infty$):

$$g_1^{\text{T3+TMC}}(x, Q^2) \simeq 0 \quad g_2^{\text{T3+TMC}}(x, Q^2) \simeq D(x) - \int_{\xi}^1 \frac{dz}{z} D(z)$$

Twist-3 quark distributions:

$$D(x, Q^2) = \frac{4}{9} D_u(x, Q^2) + \frac{1}{9} D_d(x, Q^2)$$

Twist-4 structure function (Nucleon d.o.f.):

$$g_1^{\text{T4}(p,n)}(x, Q^2) = H^{(p,n)}(x)/Q^2$$

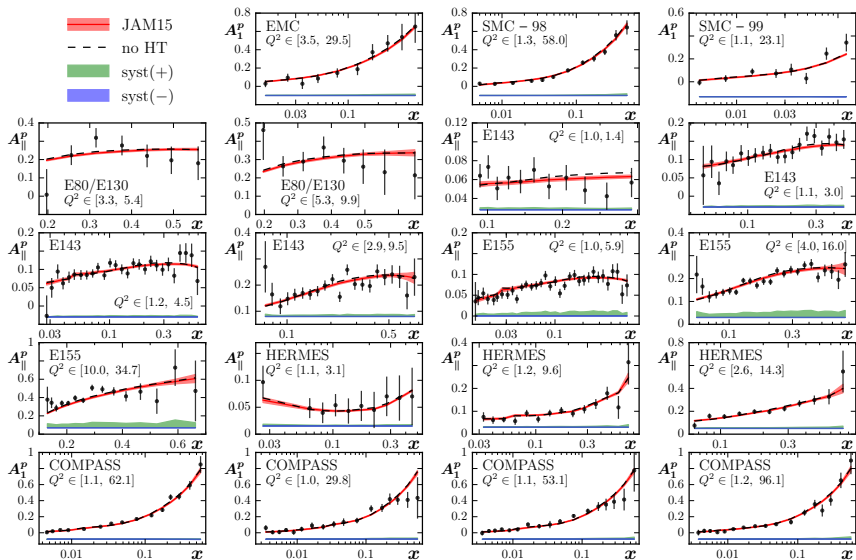
Nuclear corrections: → nuclear smearing functions

$$g_i^A(x, Q^2) = \sum_N \int \frac{dy}{y} f_{ij}^N(y, \gamma) g_j^N(x/y, Q^2)$$

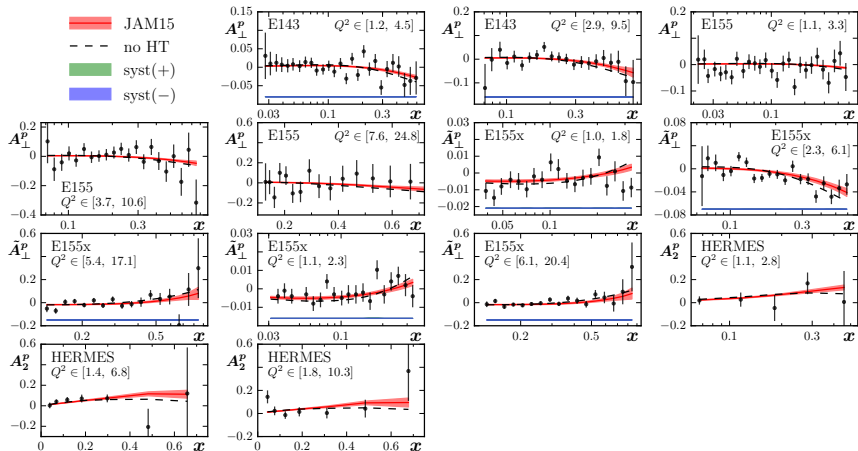
Addition constraints: → weak neutron and hyperon decay constants

- $\Delta u^{+(1)} - \Delta d^{+(1)} = F + D = 1.269(3)$
- $\Delta u^{+(1)} + \Delta d^{+(1)} - 2\Delta s^{+(1)} = 3F - D = 0.586(31)$

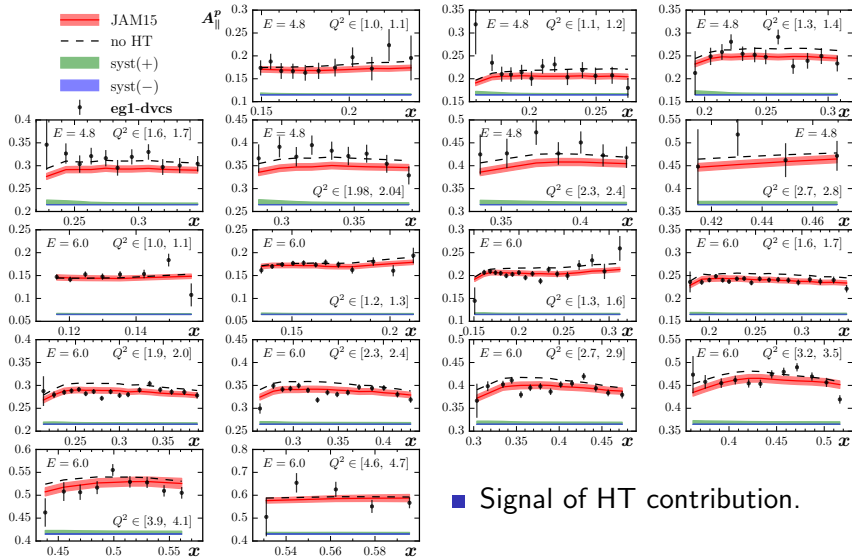
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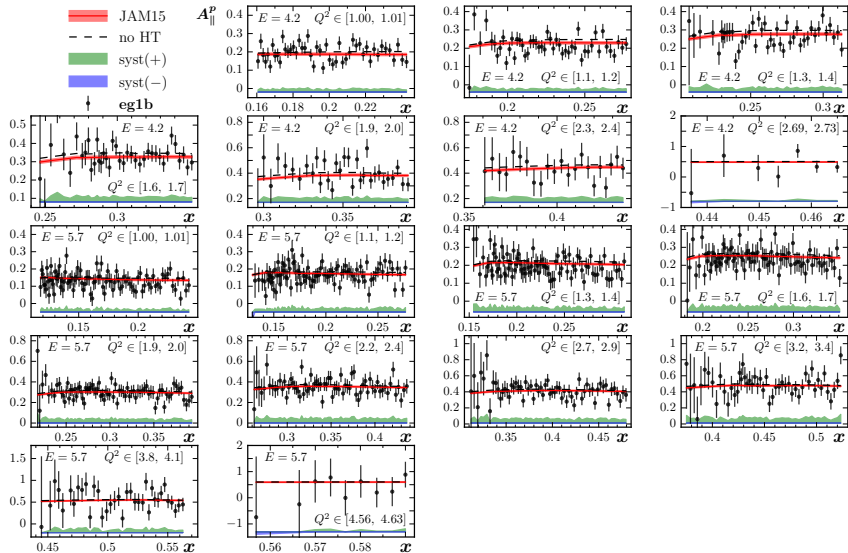
Data vs theory: proton



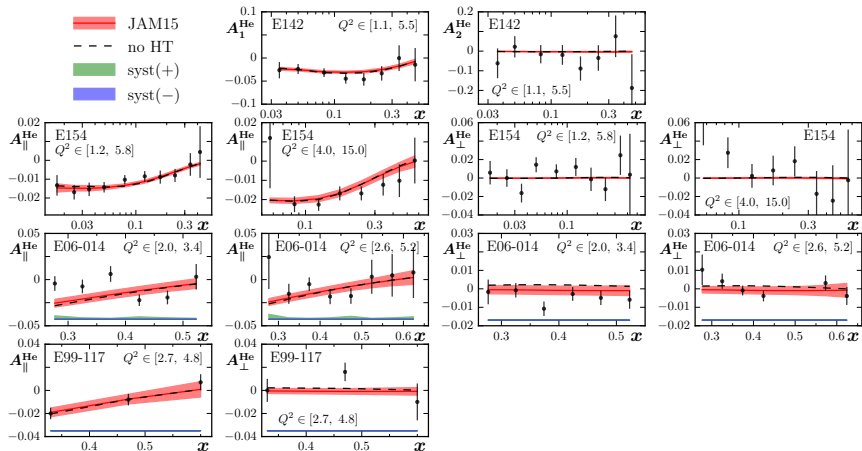
Data vs theory: proton JLab eg1b-dvcs



Data vs theory: proton JLab eg1b



Data vs theory: ^3He



JAM: moments

