PROPOSAL FOR A CAVITY POLARIMETER AT MIT-BATES

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Abstract

The possibility of successfully implementing a cavity polarimeter[1] has been greatly improved by the discovery[2] of a cavity mode for which the magnitude of the Stern-Gerlach force experienced by a magnetic moment traversing the cavity varies as the square of the relativistic factor gamma, so that the signal power varies as the fourth power of gamma. In addition, the interaction of this cavity mode with the beam charge varies as the inverse of the interaction with the magnetic moment, so that the background due to the beam charge varies as the inverse fourth power of gamma. If these gamma dependencies of moment and charge interaction with the pickup cavity do in fact exist, the possibility is opened for very fast, accurate, and inexpensive polarimetery at accelerators like MIT-Bates and RHIC. In addition, it might become possible to seriously consider Stern-Gerlach polarization of beams at LHC. We present details of a quick polarimeter test at the electron storage ring at MIT-Bates, and of an extension of this test to a working polarimeter in the RHIC rings.

1 POLARIMETER AT MIT-BATES

The cavity polarimeter consists of a passive resonant cavity, which accumulates spin-dependent energy from the beam via the Stern-Gerlach interaction. In the case of a resonant cavity the stored energy increases until equilibrium is attained, where the power supplied by the beam equals the power dissipated as a result of the finite cavity Q. We can write this as

$$Q \equiv \frac{stored\ energy}{loss} = \frac{\omega U}{\frac{dU}{dt}}$$
 (1)

where ω =2 π f, f=2.856GHz is the bunching frequency, and U is the stored energy. The rate at which polarization leaves energy in the cavity, or the signal power, is

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$$\frac{dU}{dt} = \mu B \tau b n f \gamma^2 P \tag{2}$$

where μ is the magnetic dipole moment of the electron, B is the cavity equilibrium field, τ ~0.7 is the transit time factor, b~0.7 is a bunch length factor, n=10⁸ is the number of electrons per bunch for all buckets filled and 50ma beam current, γ =2000 is the relativistic gamma factor at

1GeV, and P~0.4 is the beam polarization. Rearranging equation (1), we can write

$$U = \frac{Q}{\omega} \frac{dU}{dt} = \frac{\varepsilon_0}{4} B^2 c^2 V \tag{3}$$

Substituting (2) into (3) and solving for the equilibrium field B,

$$B = \frac{4Q\mu\tau bn\gamma^2 P}{\pi\varepsilon_0 Vc^2}$$

Where Q=5000, ε_0 is the permittivity of free space, c is the speed of light, and V~1000cm³ is the cavity volume. Substituting the value found for the equilibrium field into (2), the signal power is

$$\frac{\overline{dU}}{dt} = \frac{4Qf}{\pi \epsilon_0 V} \left(\frac{\mu \pi b n \gamma^2 P}{c^2}\right)^2 \approx 10^{-7} watt \tag{4}$$

or about -40dBm. In the world of small signal beam instrumentation, this is a whopping big signal. In comparison, Schottky signals are typically at -150dBm or less. It is interesting to note that the power varies as f^4 , because of the $1/f^3$ dependence in the cavity volume, so that there might be some advantage to operating the cavity at higher frequencies, and accepting the penalty in bunch length factor.

2 POLARIMETER AT RHIC

Performing the same calculation for conditions to be found at RHIC (55 bunches, 10^{11} per bunch, 50% polarization, γ =266, Q=5000) yields the result that the signal power is about -130dBm. Again, this is a good strong signal relative to Schottky signals.

3 CONCLUSION

While a number of issues remain to be addressed, some fairly straightforward and others thorny and difficult, the bvious course at this time is to build a simple cavity and perfom the measurement. We are proceeding with this as rapidly as funding permits.

4 REFERENCES

[1] Ya. S. Derbenev, "RF-resonance Beam Polarimeter: Part 1. Fundamental Concepts" NIM A 336, p.12-15 (1993).

[2] M. Conte et al, "The Stern-Gerlach Interaction between a travelling particle and a time varying magnetic field" http://xxx.lanl.gov/list/physics/0003, preprint 0003069

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