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# JPAC activities

Vladyslav Pauk

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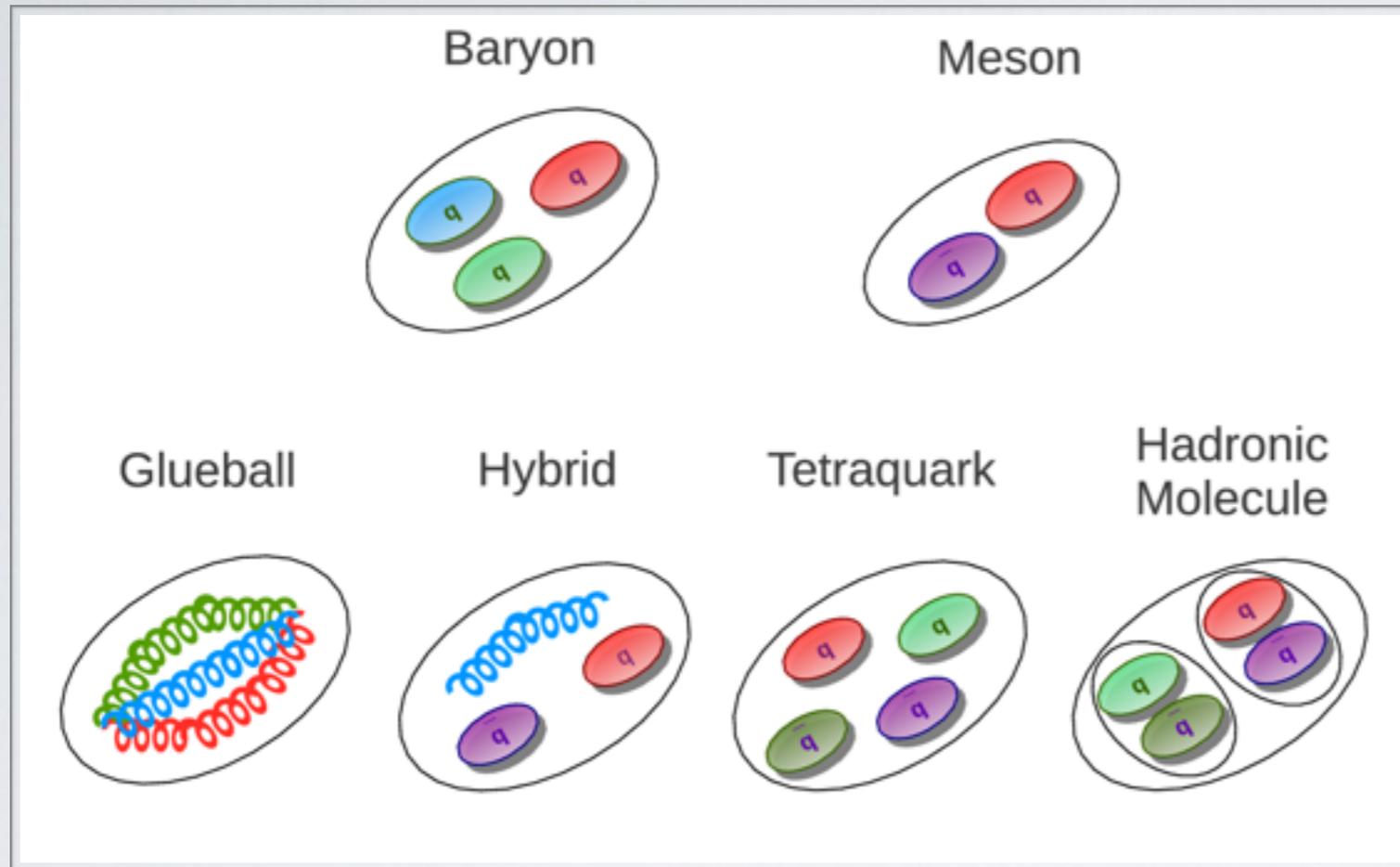
CLAS collaboration meeting

JLab

February 25, 2016

# Hadron spectroscopy

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hadron spectroscopy  
lattice QCD

amplitude  
analysis

spectrum of resonances

**JPAC**

- Extract values of fundamental parameters
- Extract properties of resonances
- Searches for new resonances/new states of matter
- Understand fundamental laws (matching QCD with exp.)

# Amplitude analysis

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**unitarity**

$$2\text{Im } T = iTT^\dagger$$

**crossing symmetry**

$$\begin{array}{l} s \quad 1 + 2 \rightarrow 3 + 4 \\ t \quad 1 + \bar{3} \rightarrow \bar{2} + 4 \\ u \quad 1 + \bar{4} \rightarrow 2 + \bar{3} \end{array}$$

**Lorentz symmetry,  
quantum numbers, etc.**

**analyticity**

$$A(s, t) = \frac{2}{\pi} \int_{s_{th}}^{\infty} \frac{\text{Im}(s', t)}{s' - s} ds'$$

consistent fitting functions  
constraints from existing data  
Finite Energy Sum Rules (FESR)

**model predictions  
data  
lattice simulations**

# Finite Energy Sum Rules

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$$A(\nu, t) = \frac{2}{\pi} \int_{\nu_{th}}^{\infty} \frac{\text{Im}A(\nu', t)}{\nu'^2 - \nu^2} \nu d\nu'$$

fixed-t dispersion relation

$$\nu = \frac{s - u}{2}$$

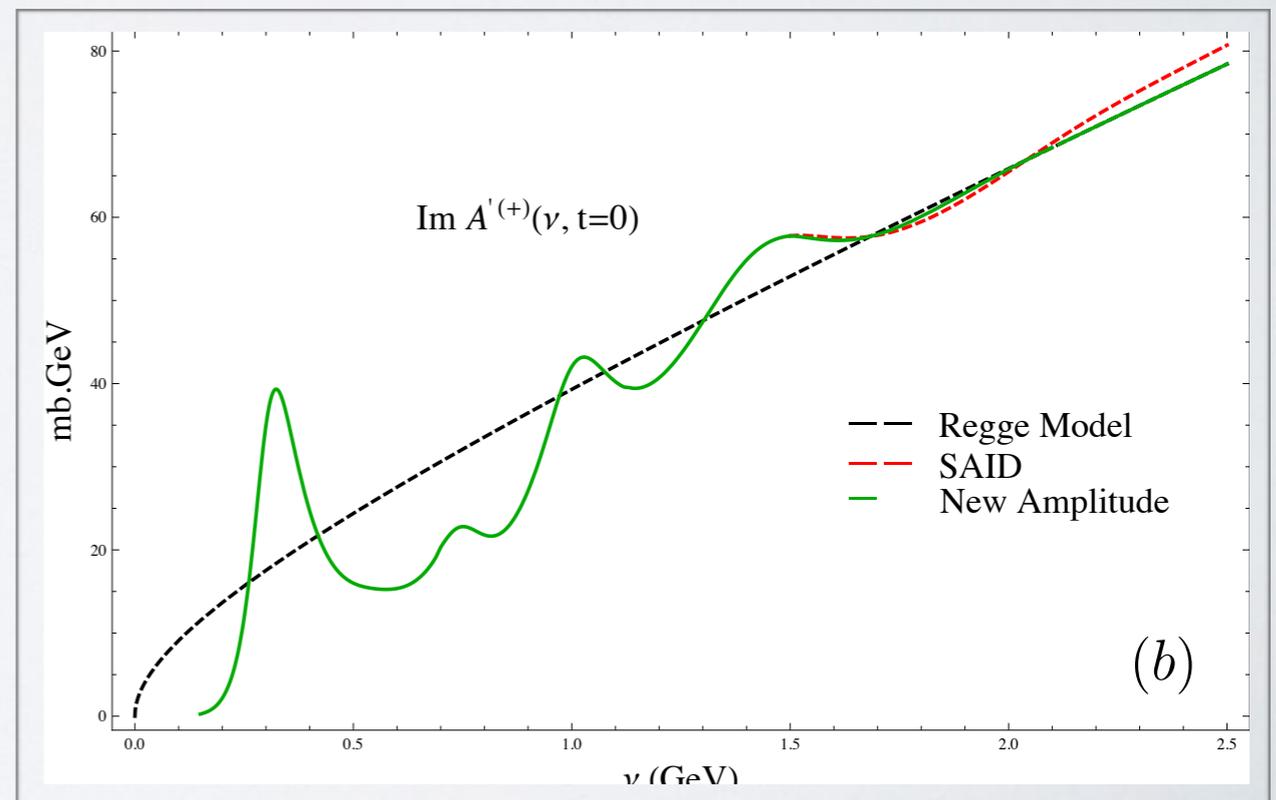
$$\frac{1}{\Lambda^n} \int_{\nu_{th}}^{\Lambda} \text{Im}A(\nu, t) \nu^n d\nu = \frac{\beta(t) \Lambda^{\alpha(t)+1}}{\alpha(t) + n + 1}$$

**Finite Energy Sum Rules**

$$A(\nu, t) \rightarrow A(\nu, t) - A(\nu \rightarrow \infty, t)$$

$$A(\nu, t) \xrightarrow{\nu \rightarrow \infty} \beta(t) \nu^{\alpha(t)}$$

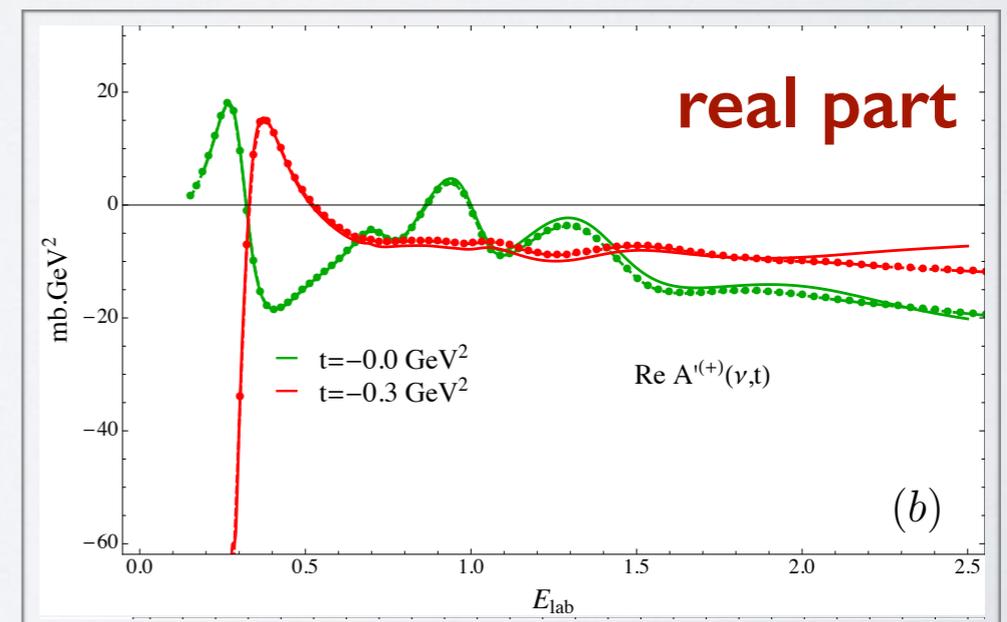
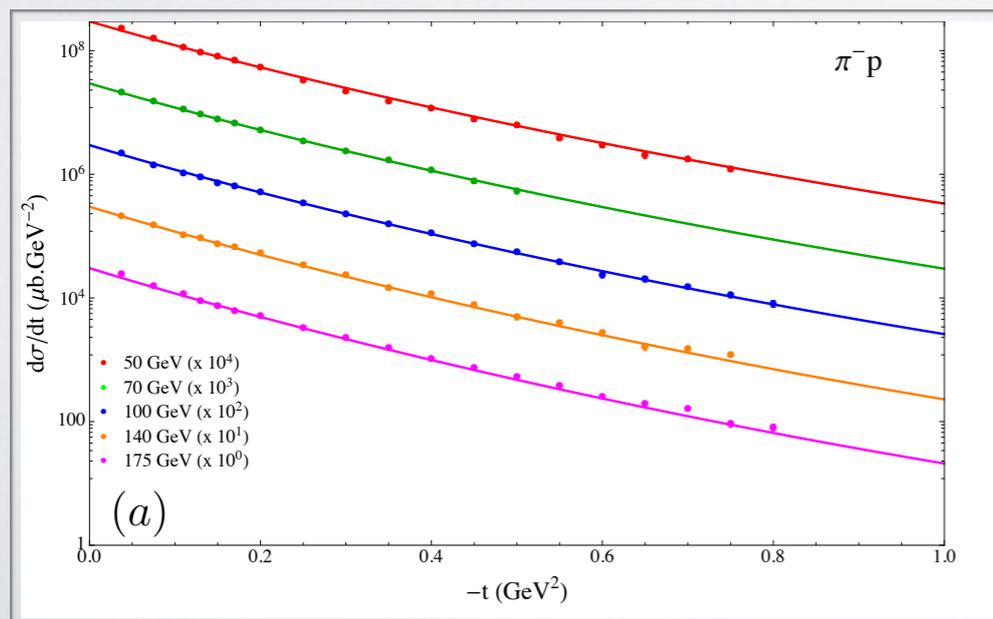
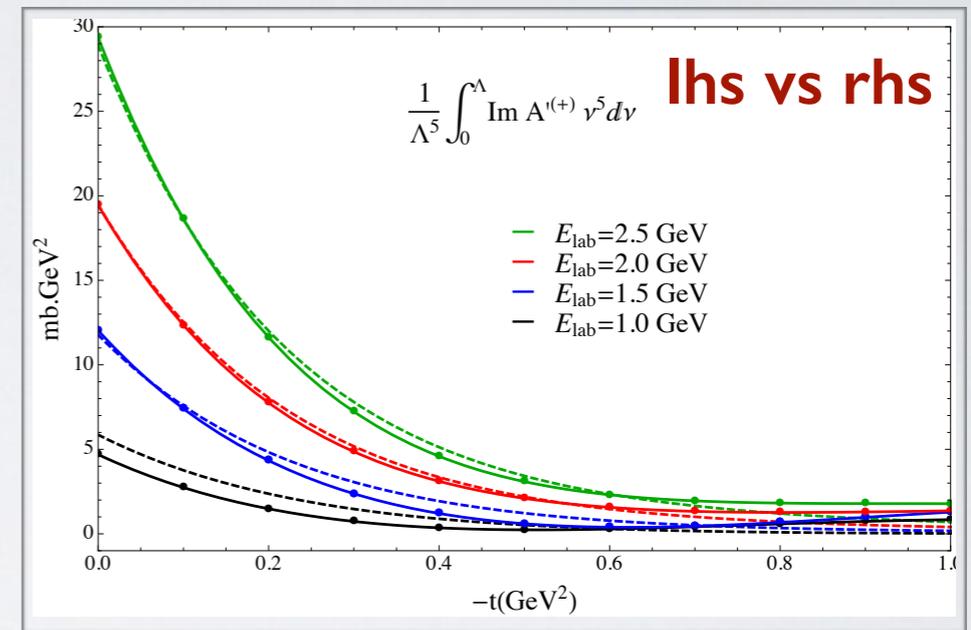
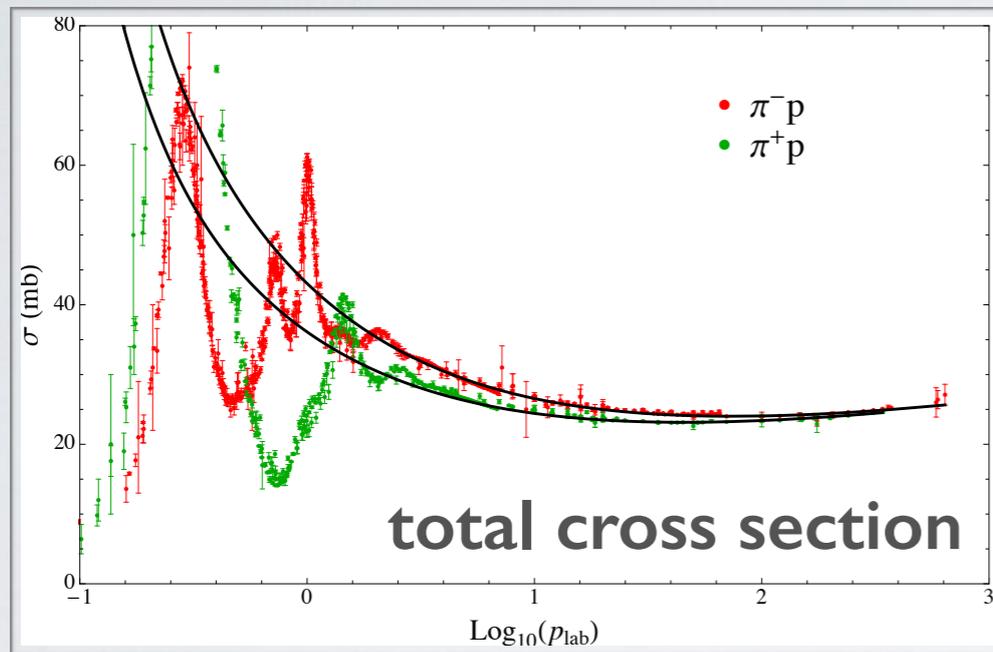
high-energy behavior



# Pion-Nucleon Amplitudes

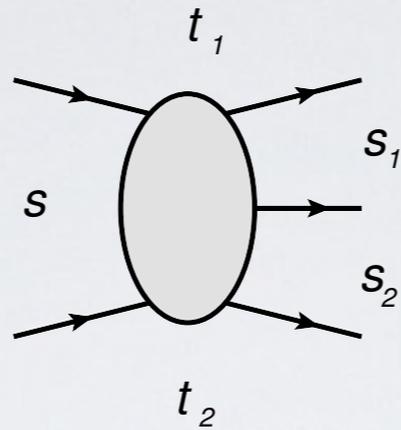
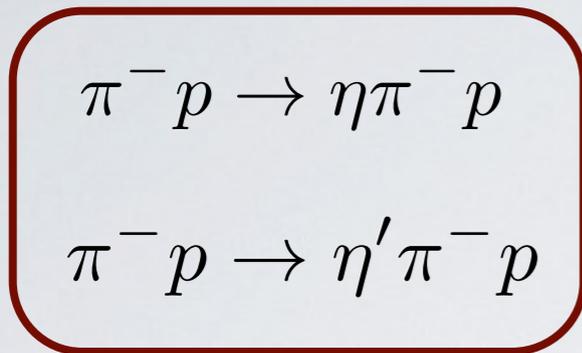
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high-energy fit to total  
and differential cross sections



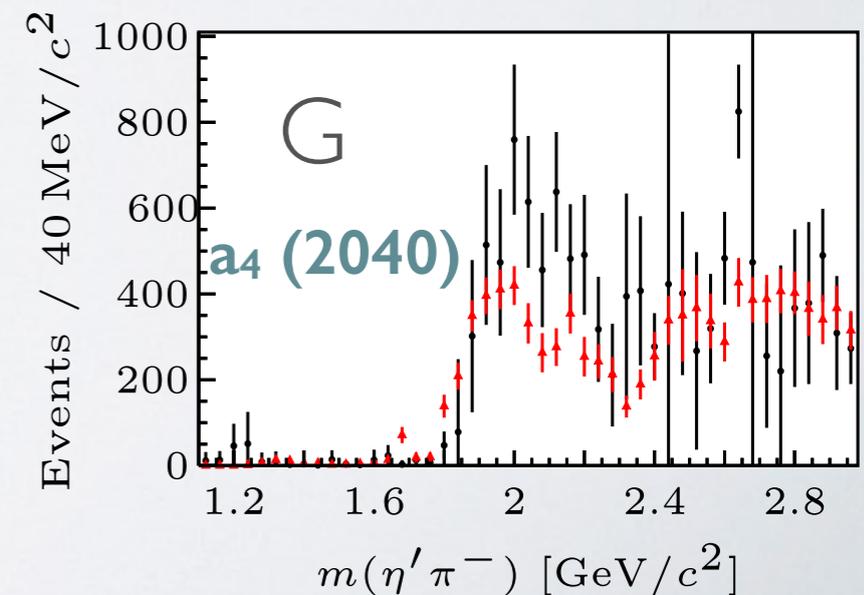
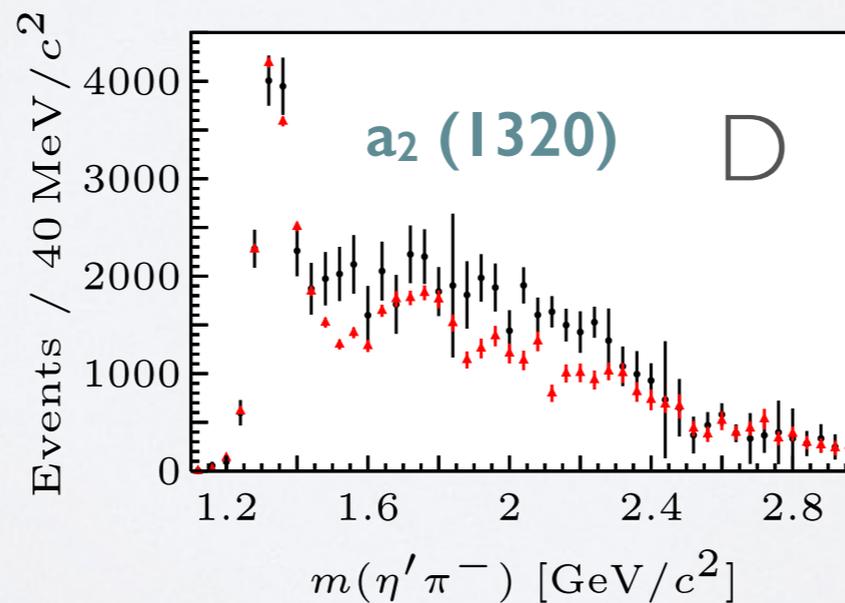
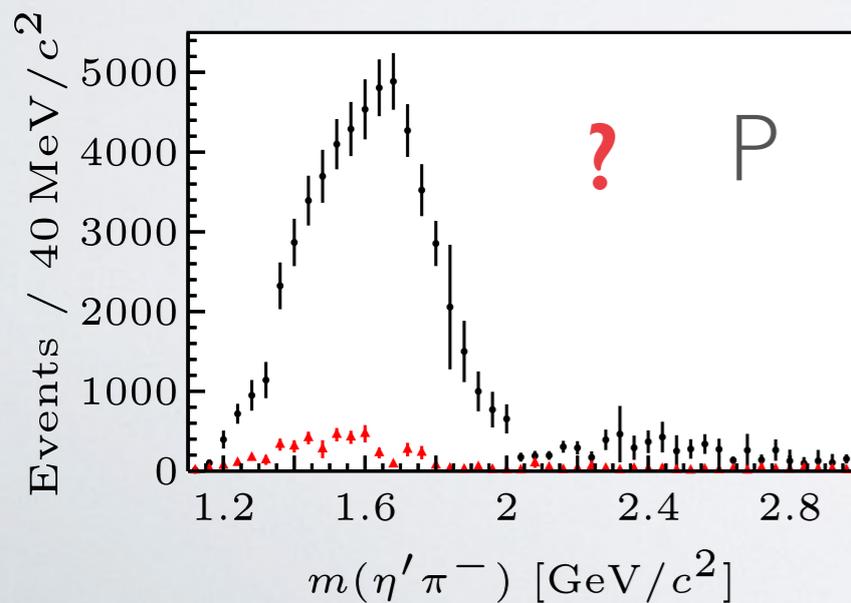
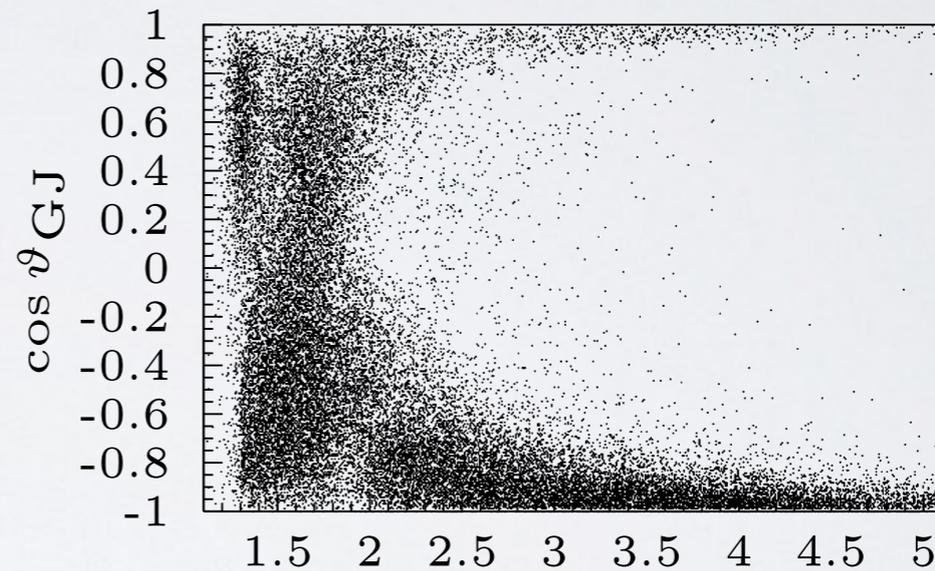
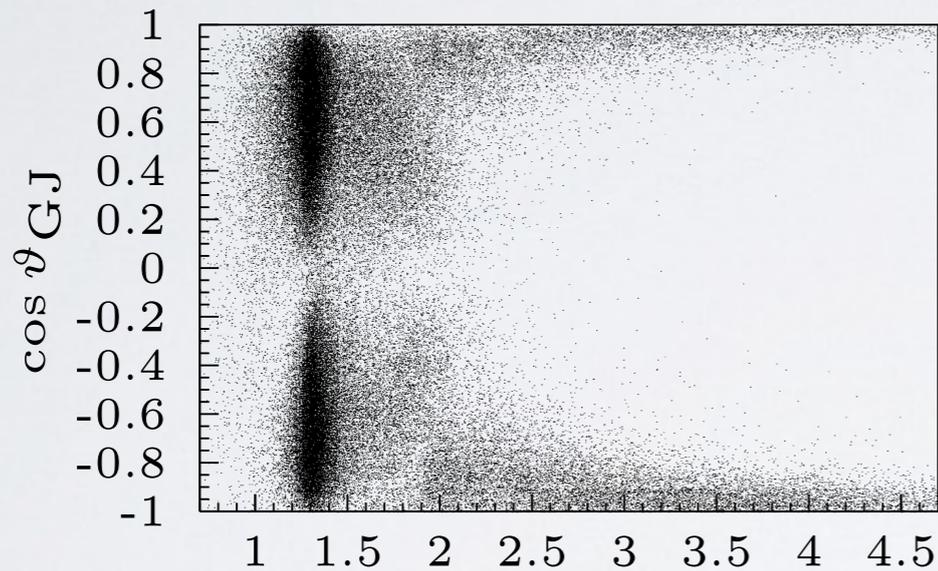
# $\eta\pi$ production at COMPASS

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$$J^{PC} = 1^{-+}, 3^{-+}, 5^{-+}, \dots$$

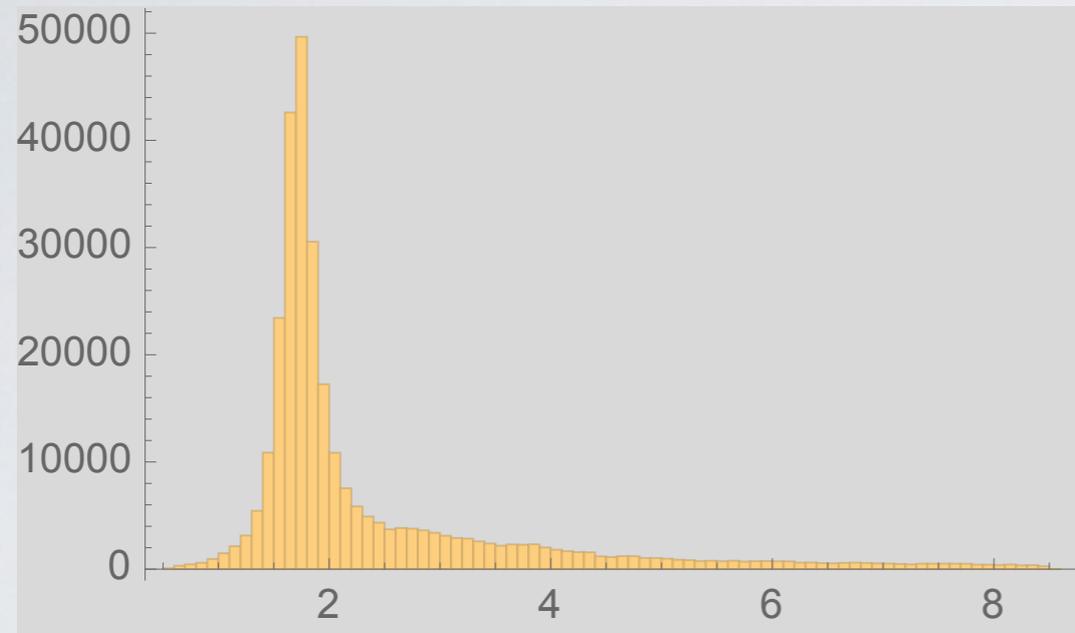
exotic quantum numbers



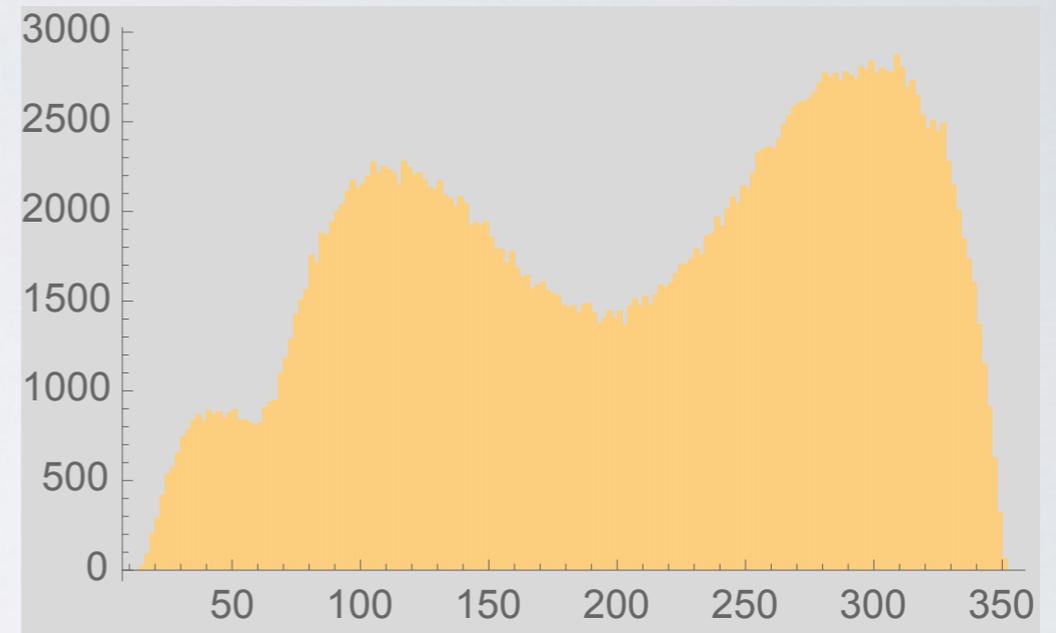
# Kinematic ranges

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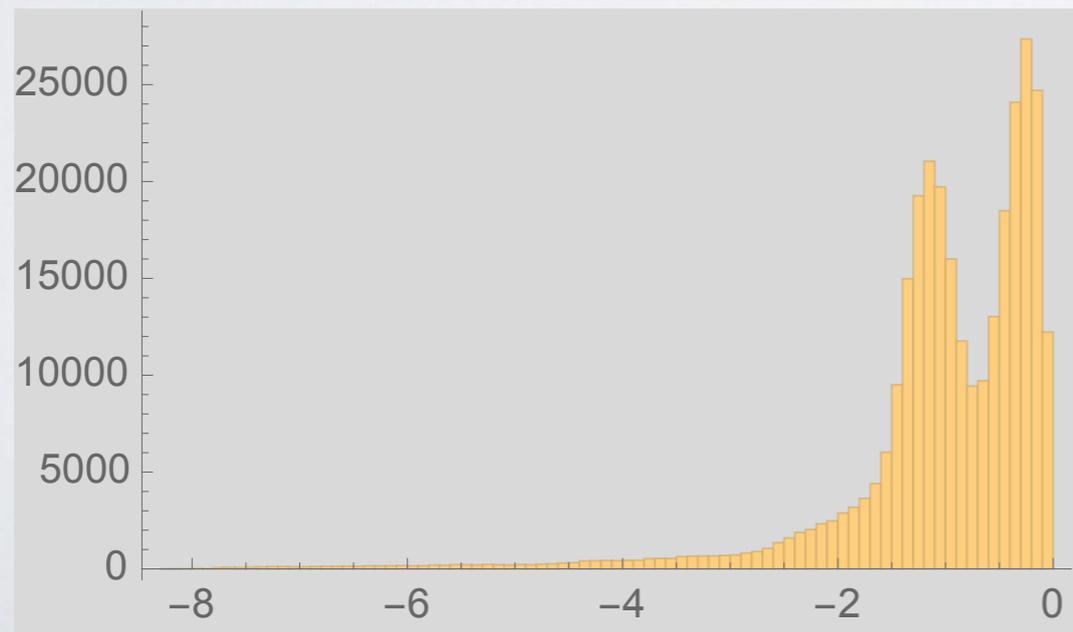
$$0.5 \text{ GeV} < s_1 = (p_1 + p_2)^2 < 5 - 6 \text{ GeV}$$



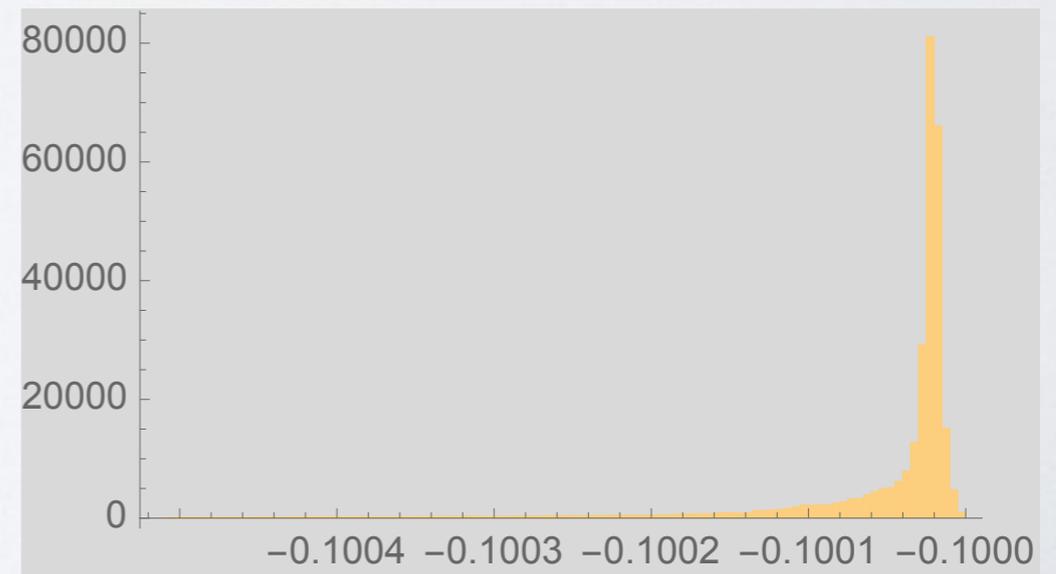
$$15 \text{ GeV} < s_2 = (p_2 + p_3)^2 < 350 \text{ GeV}$$



$$0 \text{ GeV} > t_1 = (p_a - p_1)^2 > -2 \text{ GeV}$$

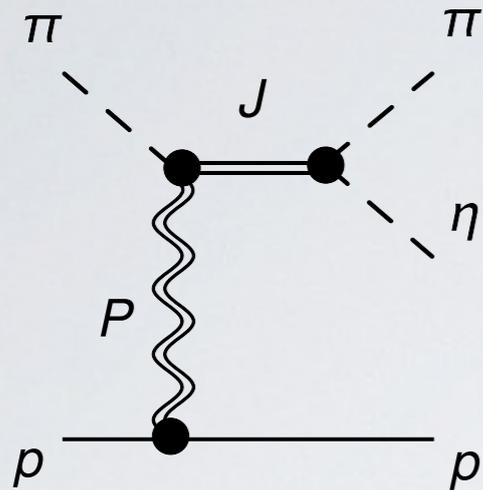


$$t_2 = (p_b - p_3)^2 \approx -0.1 \text{ GeV}$$

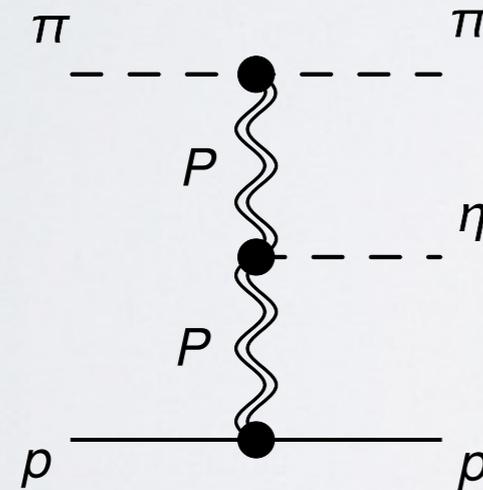


# Production mechanism

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resonance production



central production

partial-wave expansion

$$A_{\mu\mu'}(s_1, s_2, t_1, t_2) = \sum_{\lambda=-J}^J A_{\lambda}^J(s_1, t_2) \frac{g_J(s_1)}{m_J^2 - im_J \Gamma(s_1) - s_1} d_{\lambda 0}^J(z) e^{i\lambda\phi} \times \eta(\alpha_2(t_2)) \gamma_{2\mu\mu'}(t_2) s_2^{\alpha_2(t_2)}$$

Regge behavior

**FESR**

$$A_{\mu\mu'}(s_1, s_2, t_1, t_2) = \eta(\alpha_1(t_1)) \gamma_1(t_1) s_1^{\alpha_1(t_1)} \sum_{\lambda} e^{i\lambda\omega} \gamma_{\lambda}(t_1, t_2) \times \eta(\alpha_2(t_2)) \gamma_2^{\mu\mu'}(t_2) s_2^{\alpha_2(t_2)}$$

Toller angle

$$\cos \omega = \frac{(\mathbf{p}_a \times \mathbf{p}_1) \cdot (\mathbf{p}_b \times \mathbf{p}_3)}{|\mathbf{p}_a \times \mathbf{p}_1| |\mathbf{p}_b \times \mathbf{p}_3|} \Big|_{\mathbf{p}_2=0}$$

$$A_{H_s}(s, t) = 16\pi \sum_{J=M}^{\infty} (2J+1) A_{HJ}(s) \mathcal{D}_{\mu\mu'}^{J*}(\phi_s, \theta_s, -\phi_s)$$

$$\zeta_{\mu\mu'}(z_s) = \left(\frac{1-z_s}{2}\right)^{\frac{1}{2}|\mu-\mu'|} \left(\frac{1+z_s}{2}\right)^{\frac{1}{2}|\mu+\mu'|} + \text{crossing symmetry or threshold behavior}$$

## low-t dependence

$$\gamma^{\lambda_1\lambda_2}(t) \sim (-t)^{|\lambda_1-\lambda_2|/2}$$

$$\gamma_{\lambda\lambda_2-\lambda}^{\lambda_2} \sim (-t_1)^{|\lambda_1|/2} (-t_2)^{|\lambda_2|/2}$$

## general form of the double-Regge residue

parity

$$\gamma_{\lambda_1\lambda_2}^{\lambda_h} = -\gamma_{-\lambda_1-\lambda_2}^{-\lambda_h}$$

$$\gamma_0(t_1, t_2, \omega) = 0$$

$$\beta(t_1, t_2, \omega) \sim 2i(-t_1)^{1/2}(-t_2)^{1/2} \sin \omega \sum_{\lambda=0}^{\infty} \sin^{\lambda} \omega \tilde{\gamma}_{\lambda}(t_1, t_2)$$

$$\beta(t_1, t_2, \omega) = \sum_{\lambda=-\infty}^{\infty} e^{i\lambda\omega} \gamma_{\lambda}(t_1, t_2)$$

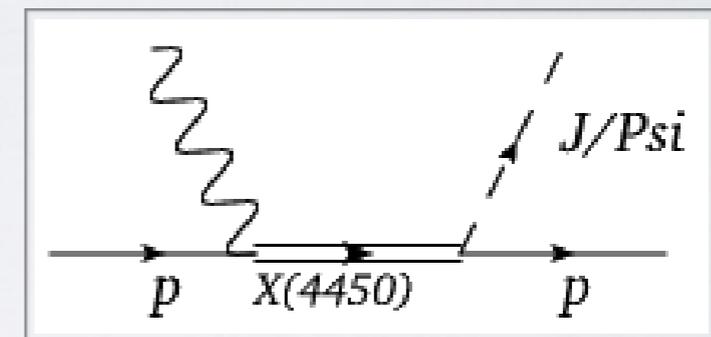
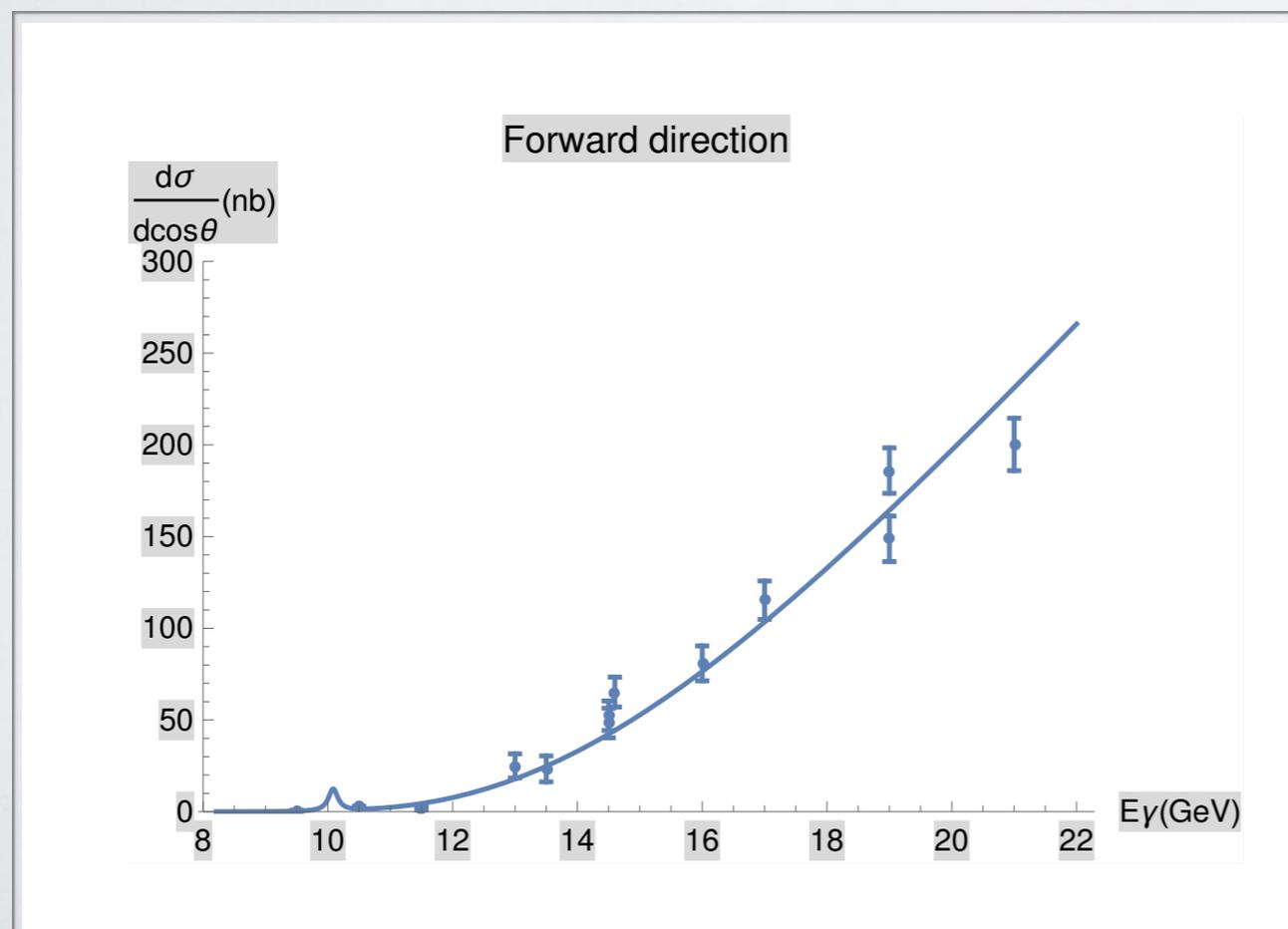
# Studying the pentaquark in the $J/\Psi$ photoproduction =9=

- Recent discovery of a narrow (39 MeV) exotic resonance compatible with a pentaquark  $P_c(4459)$  in the  $J/\Psi p$  channel.

LHCb collaboration (2015) arXiv:1507.03414

- Proposed as an excellent candidate for  $J/\Psi$  photoproduction off a proton target.

arXiv:1508.01496, arXiv:1508.00339, arXiv:1508.00888

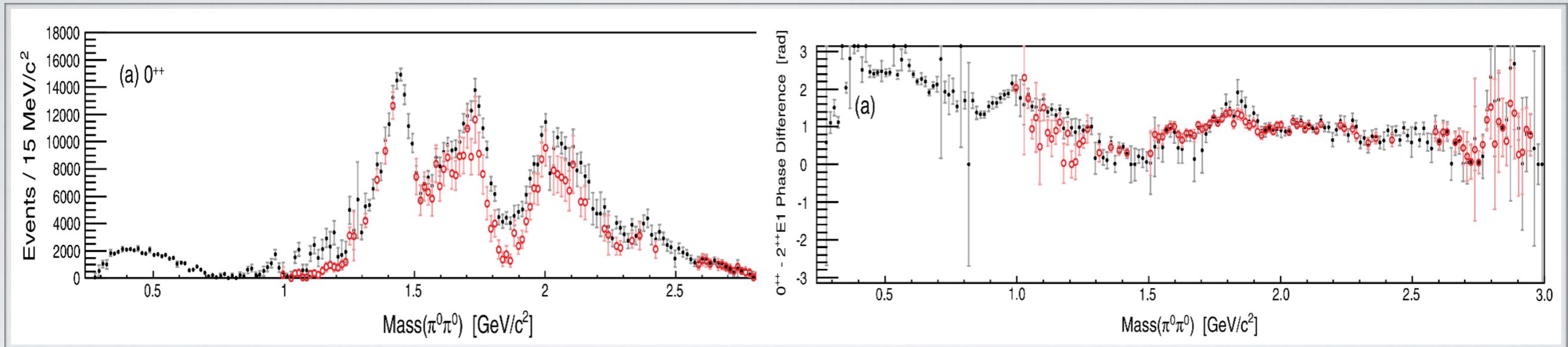


Testing these predictions is within the capabilities of **JLAB's CLAS** detector, also in a wider range of scattering angles!

# $J/\Psi \rightarrow \gamma \pi \pi$

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Good channel for the search of ground-state scalar glueball  
Data provided by BESIII both for the charged and neutral channel



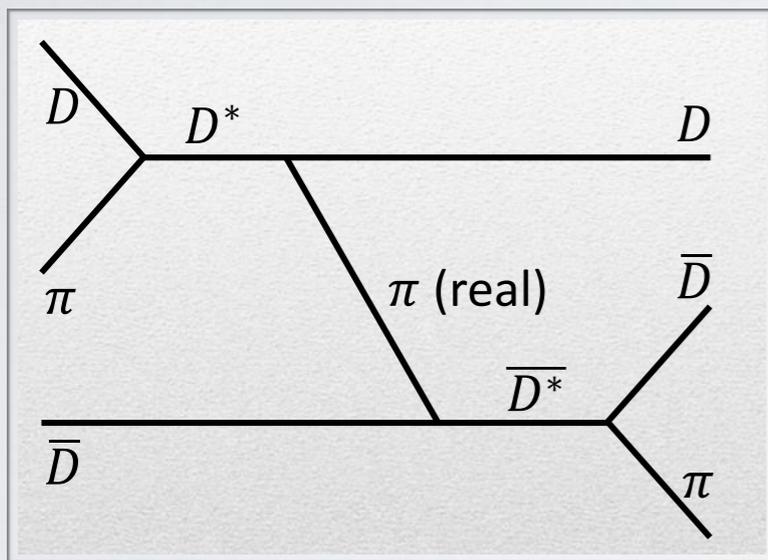
$t$ -channel dominated by  $\rho$  exchange  $\rightarrow$  model for LHC

$$f_J(s) = \frac{1}{\pi} \int \frac{\text{disc } f_J^\rho(s')}{s' - s} + \frac{1}{\pi} \int_{4m_\pi^2}^{\infty} \frac{\eta \sin \delta_J e^{-i\delta_J} f_J(s')}{s' - s} ds'$$

Solved in terms of the Omnès function up to  $\sqrt{s} \sim 1.2$  GeV  
new parametrizations tested in the  $1.2 \text{ GeV} < \sqrt{s} < M_\psi$

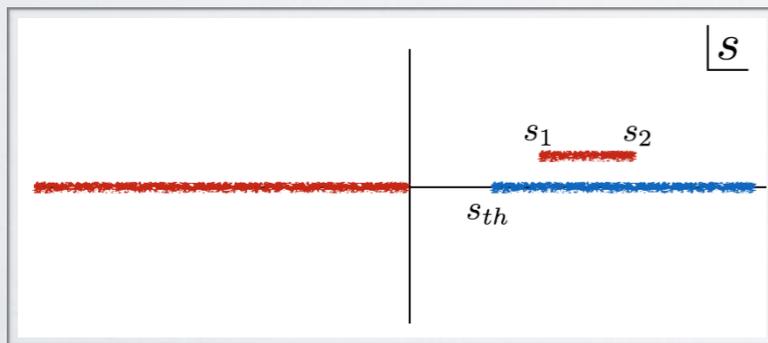
# 3-body scattering: $DD\pi \rightarrow DD\pi$ ( $X(3872)$ ) = || =

The dominant binding mechanism is expected to be the exchange of one pion in the u channel, but in the literature, this has been evaluated in the static limit only (virtual pion)



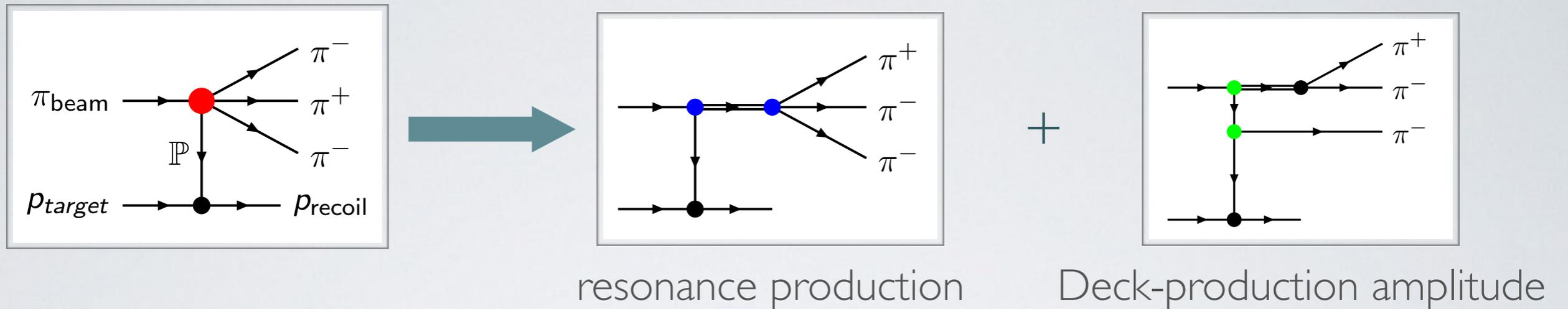
However, the  $\pi$  can happen to be on shell: this creates another cut, which might spoil the binding mechanism of  $DD^*$

Cusp effect if the branch points pinch the real axis



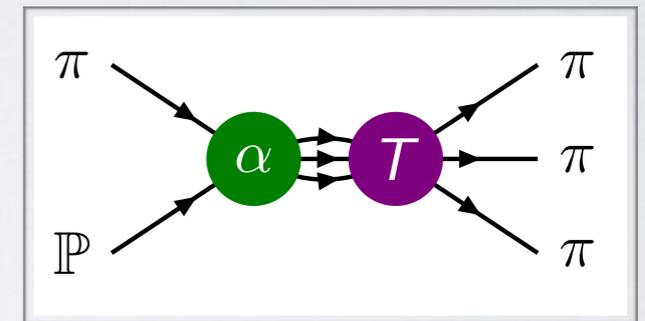
Once developed, the formalism can be extended to other  $3 \rightarrow 3$  channels, like  $\rho \pi \rightarrow \rho \pi \dots$

# 3 $\pi$ production at COMPASS experiment = 12=



resonance production

$$A = \langle \pi \mathbb{P} | \hat{T} | 3\pi \rangle =$$



- **Unitarity** has to be respected  $\hat{S} \times \hat{S}^\dagger = \mathbb{I}$

$$2\text{Im}T = iT\rho T^\dagger \quad \hat{S} = \mathbb{I} + iT$$

- Convenient parameterization of  $T$  is a **K-matrix**

- **Non-resonance processes** is a physical background?

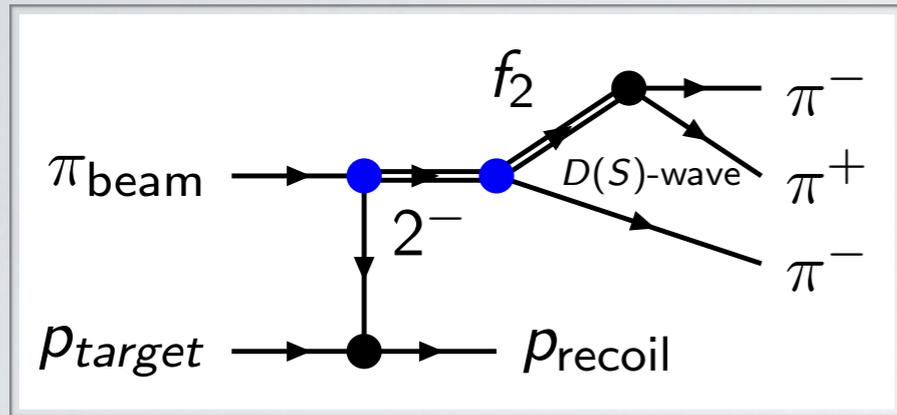
$$2\text{Im}A = iA\rho T^\dagger \quad A(m_{3\pi}) = \alpha(m_{3\pi}) \times T(m_{3\pi})$$

- **Quasi-two-body** phase space ( $f_2$  is  $(\pi\pi\pi)_D$ -state)

M. Mikhashenko, A. Jackura

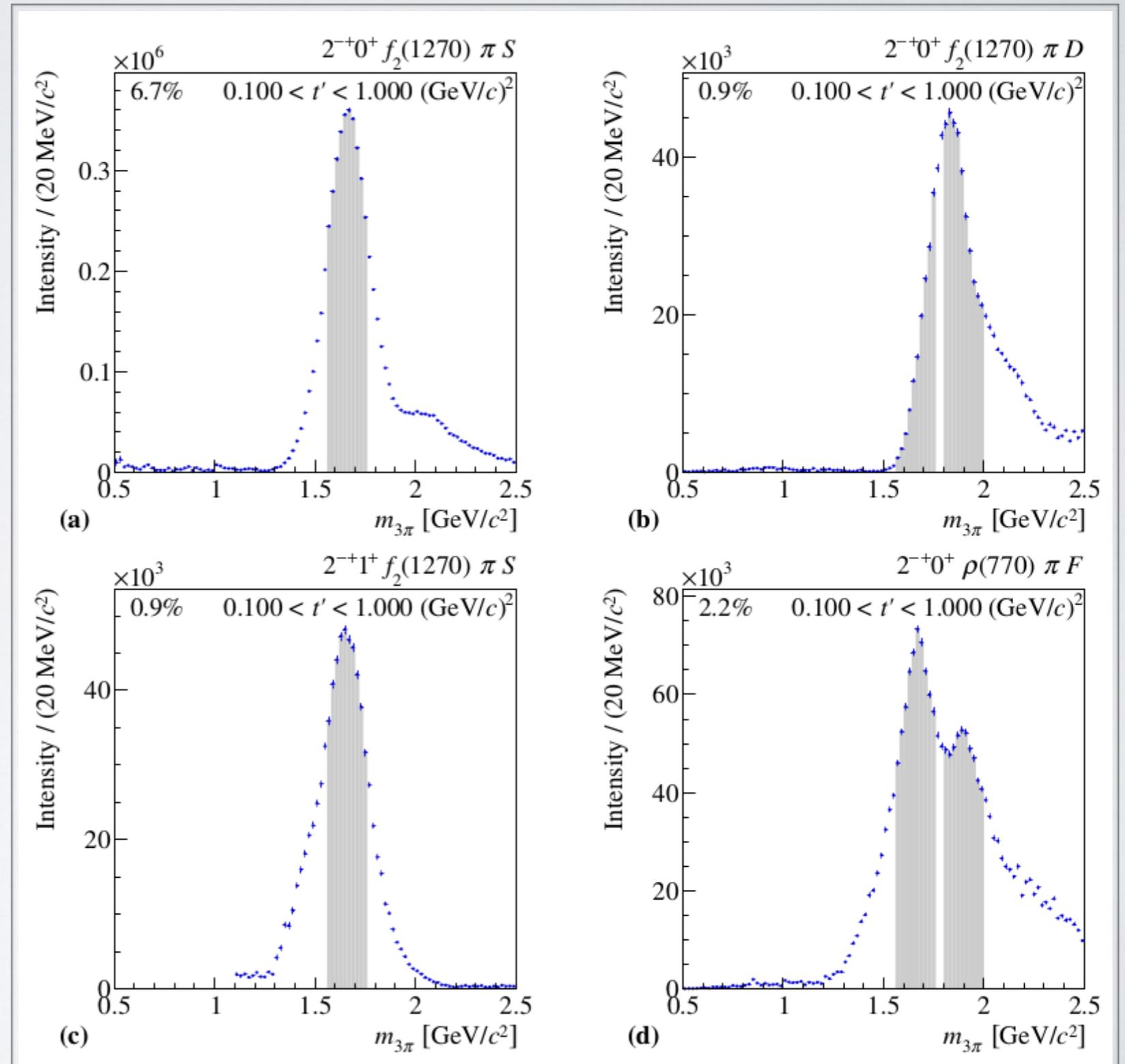
# $2^{-+}$ sector of the $3\pi$ final state

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Long standing puzzle about  $\pi_2(1670) - \pi_2(1880)$  interplay

The goal is to **develop** a method of the analysis and **demonstrate** an applicability.



# Lepton pair production on a proton target = |4=

Measure the ratio of the e vs  $\mu$  cross sections

$$R_{\mu/e} \equiv \frac{d\sigma(\mu^- \mu^+ + e^- e^+)}{d\sigma(e^- e^+)} - 1$$

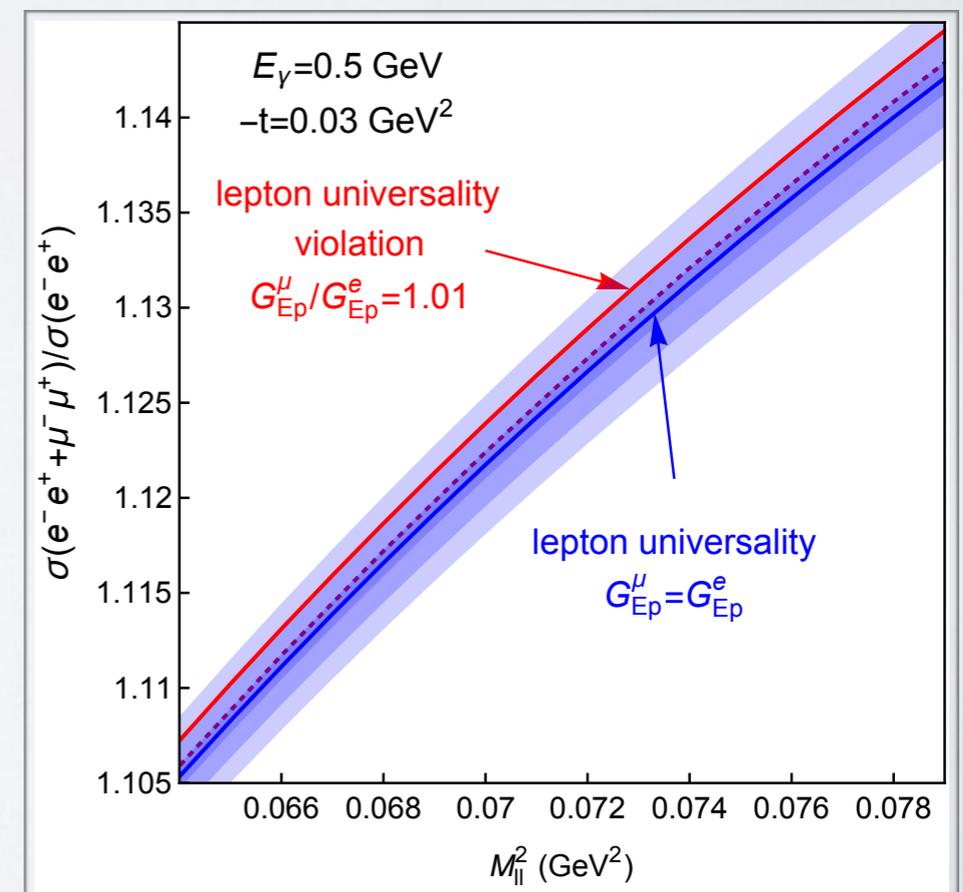
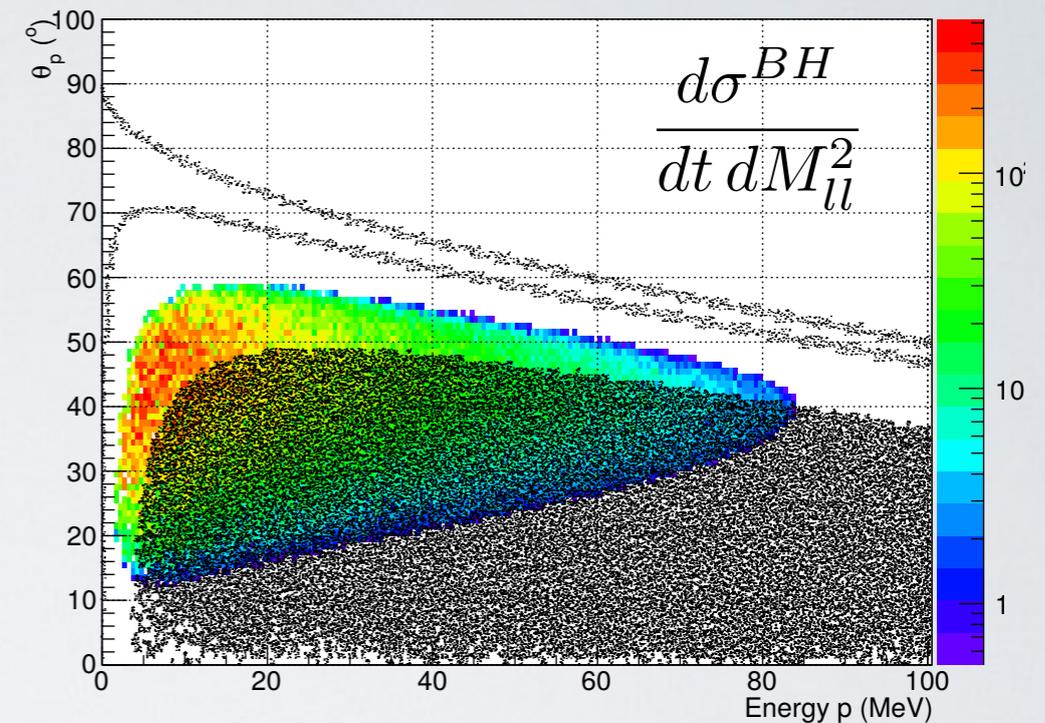
VP, Vanderhaeghen, Phys. Rev. Lett. 115, 221804

at small  $t$  the ratio  $R_{\mu/e}$  gives **direct access** to the ratio of the **proton electric form factor** in the  $\mu p$  versus  $ep$  scattering

the deviation from the unity will be a sign of **violation of the lepton universality**

**JLab**

access the proton form factor by analyzing angular distributions of the lepton pairs



# Overview

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$$\pi N \rightarrow \pi N$$

V. Mathieu, et al.

arXiv:1506.01764  
PRD92 7 074004

$$\gamma p \rightarrow \pi^0 p$$

V. Mathieu, et al.

arXiv:1505.02321  
PRD92 7 074013

$$\eta \rightarrow \pi^+ \pi^- \pi^0$$

P. Guo, et al.

arXiv:1505.01715  
PRD92 5 054016

$$\omega, \phi \rightarrow \pi^+ \pi^- \pi^0$$

I. Danilkin, et al.

arXiv:1409.7708  
PRD91 9 094029

$$\gamma p \rightarrow K^+ K^- p$$

M. Shi, et al.

arXiv:1411.6237  
PRD91 3 034007

$$KN \rightarrow KN$$

C. Fernandez-Ramirez, et al.

arXiv:1510.07065