

Photocathode Surface Roughness: A Systematic Study

Benjamin L. Rickman, Louis Angeloni, and W. Andreas Schroeder

University of Illinois at Chicago, Department of Physics (m/c 273), 845 W. Taylor Street, Chicago, IL 60607-7059

Introduction

- A detailed systematic study of the effect of surface roughness on the transverse emittance; specifically, the RMS transverse momentum, Δp_T , of electron pulses generated in a 'low-field' laser driven DC gun.
- Agreement between theoretical analyses and experiment confirms that the total Δp_T is dependent upon three factors:
 - The intrinsic RMS transverse momentum of the photocathode, $\Delta p_{T,0}$ – dependent upon $\Delta E = \hbar\omega - \phi$, band structure, etc.
 - The RMS photocathode surface roughness, R_q – introduces local transverse surface fields that increase Δp_T
 - The RMS slope of the photocathode surface, Δ_q – influences electron emission direction through p_T conservation

Approach

Two key components:

- Solenoid scan determination of Δp_T (roughness) with 'low-field' DC gun
 - ⇒ 'Very rough' photocathode surface required
 - ∴ Roughness measurement by optical scattering techniques
- Work function variation with crystal orientation (i.e., ϕ_{ijk})
 - ⇒ Either (i) a polycrystalline photocathode; $\phi_{RMS} \approx$ constant (i.e., crystal facet surface distribution independent of roughness parameters R_q and Δ_q)
 - polycrystalline Mo
 - Or (ii) single-crystal material with $\phi_{ijk} \approx$ constant
 - GaSb(001)

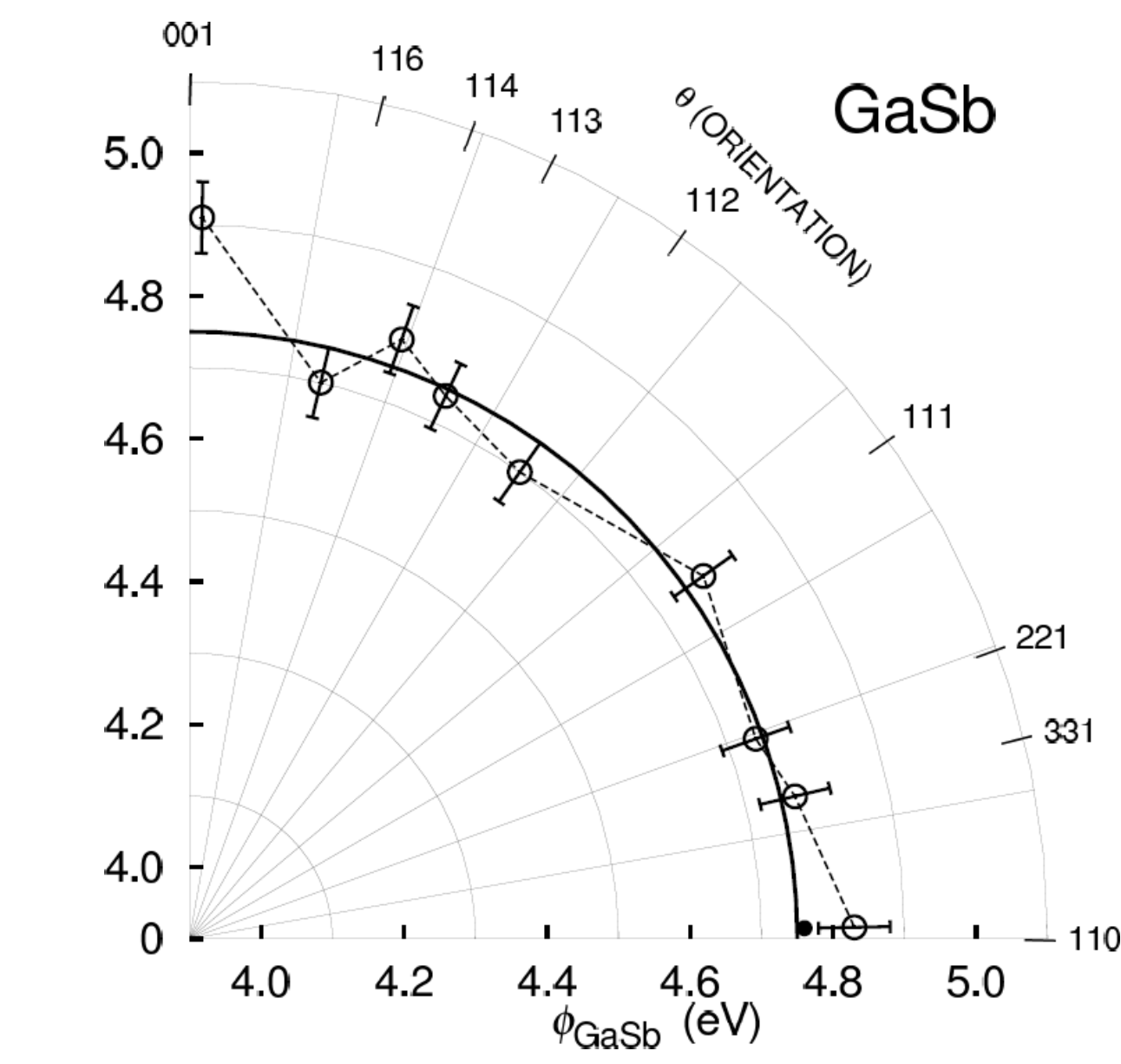


Figure 1
DFT-based thin-slab evaluation of ϕ_{ijk} for GaSb (open circles with ± 0.05 eV uncertainty); $\hbar\omega = 4.75$ eV (solid line); literature value for $\phi_{110} = 4.76$ eV (black circle).

Methods

Optical measurement of R_q and Δ_q

Definitions:

$$\text{RMS surface roughness, } R_q = \sqrt{\frac{1}{L} \int_0^L h^2(x) dx}$$

$$\text{RMS slope, } \Delta_q = \sqrt{\frac{1}{L} \int_0^L \left(\frac{\partial h(x)}{\partial x} \right)^2 dx}$$

$$\dots \text{ for } h(x) = a \cos\left(\frac{2\pi x}{b}\right); R_q = \frac{a}{\sqrt{2}} \text{ and } \Delta_q = \frac{\sqrt{2}\pi a}{b}.$$

Optical scattering measurement:

- Light source: Red ($\lambda = 635$ nm) diode laser beam
- Photocathodes roughened using 5 μ m to 120-grit 'polishing' paper
- Brass and stainless steel 'standards' employed to test technique
- Measurement of $R_{spec}(\theta_i)$ and $I(\theta)$ allows extraction of R_q and Δ_q :

$$(i) \text{ Specular reflection, } R_{spec}(\theta_i) = R_0 \exp\left[-\left(\frac{4\pi R_q \cos \theta_i}{\lambda}\right)^2\right]$$

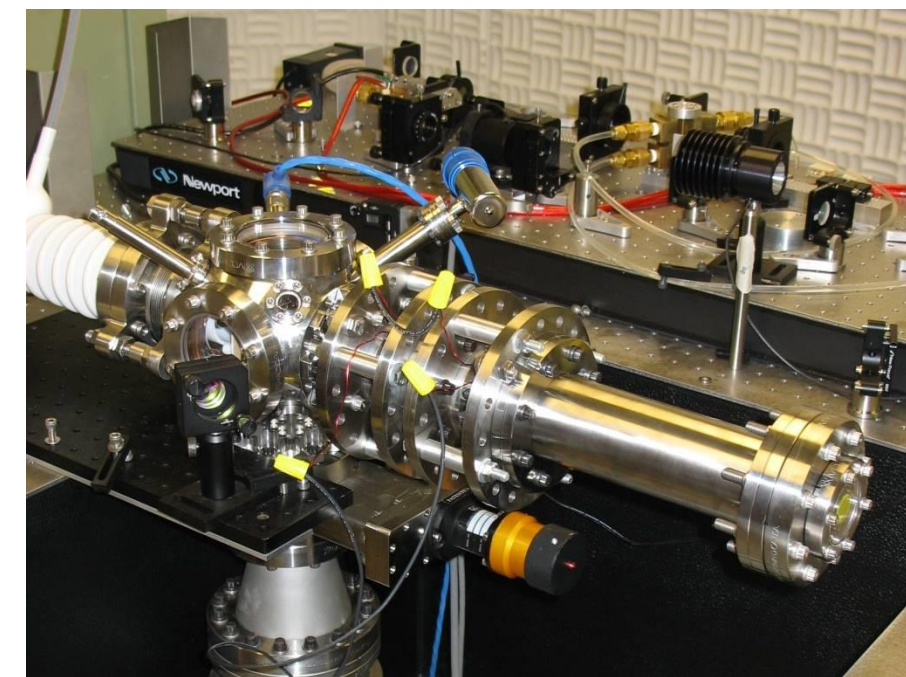
⇒ Measure width of specular reflection peak at different incident angles θ_i to extract R_q

$$(ii) \Delta_q = \frac{1}{2} \sqrt{\frac{1}{I_0} \int_{-\pi/2}^{\pi/2} d\theta (\theta - \theta_r)^2 I(\theta)}$$

$$\text{where } I_0 = \int_{-\pi/2}^{\pi/2} d\theta I(\theta) \text{ and } \theta_r = \frac{1}{I_0} \int_{-\pi/2}^{\pi/2} d\theta \theta I(\theta)$$

⇒ Measure $I(\theta)$ for $\theta_i = 0^\circ$ to extract Δ_q

Solenoid scan technique



- 2W, 250fs, 63MHz, diode-pumped Yb:KGW laser – ~4ps at 261nm ($\hbar\omega = 4.75$ eV)
- YAG scintillator optically coupled to CCD camera
 - Beam size vs. magnetic coil (lens) current measured
 - Analytical Gaussian (AG) pulse propagation model to extract Δp_T

Theory

For $h(x) = a \cos\left(\frac{2\pi x}{b}\right)$, transverse velocity 'kick' due to surface roughness

$$\dot{x}(x_0, \theta) = v_{x,0}(x_0, \theta) + \frac{2\pi a}{b} \left(\frac{qE_0 b}{4m} \right)^{\frac{1}{2}} \exp\left[\frac{\pi m v_{z,0}^2(x_0, \theta)}{qE_0 b} \right] \text{erfc}\left[\left(\frac{\pi m v_{z,0}^2(x_0, \theta)}{qE_0 b} \right)^{\frac{1}{2}} \right] \sin\left(\frac{2\pi x_0}{b} \right)$$

Transverse velocity due to slope angle Transverse velocity due to roughness-induced transverse field components

... E_0 is the acceleration field at the photocathode surface; $v_{x,0}$ and $v_{z,0}$ are the initial transverse and longitudinal electron velocity components at photocathode position x_0

$$\Delta p_T \text{ can then be evaluated using } (\Delta p_T)^2 = \frac{\int_0^b dx_0 \int_{-\pi/2}^{\pi/2} d\theta (m\dot{x}(x_0, \theta))^2 N(\theta)}{\int_0^b dx_0 \int_{-\pi/2}^{\pi/2} d\theta N(\theta)}$$

... where $N(\theta)$ is the angular distribution of emitted electrons for a smooth surface.

Simple analytical expressions result for the 'slope' and 'field' contributions to Δp_T in terms of R_q and Δ_q :

$$\left. \begin{aligned} \Delta p_{T,slope} &= m_0 v \Delta_q \\ \Delta p_{T,field} &= \sqrt{\frac{\pi R_q \Delta_q q E_0 m_0}{2}} \end{aligned} \right\} \text{In limit } \frac{2\pi a}{b} \ll 1$$

The total RMS transverse momentum is then given by

$$\Delta p_T = \sqrt{(\Delta p_{T,0})^2 + (\Delta p_{T,slope})^2 + (\Delta p_{T,field})^2}$$

... where $\Delta p_{T,0}$ is the intrinsic RMS transverse momentum of the flat photocathode

This analytical result is in good agreement with detailed electron trajectory simulations for rough photocathode surfaces:

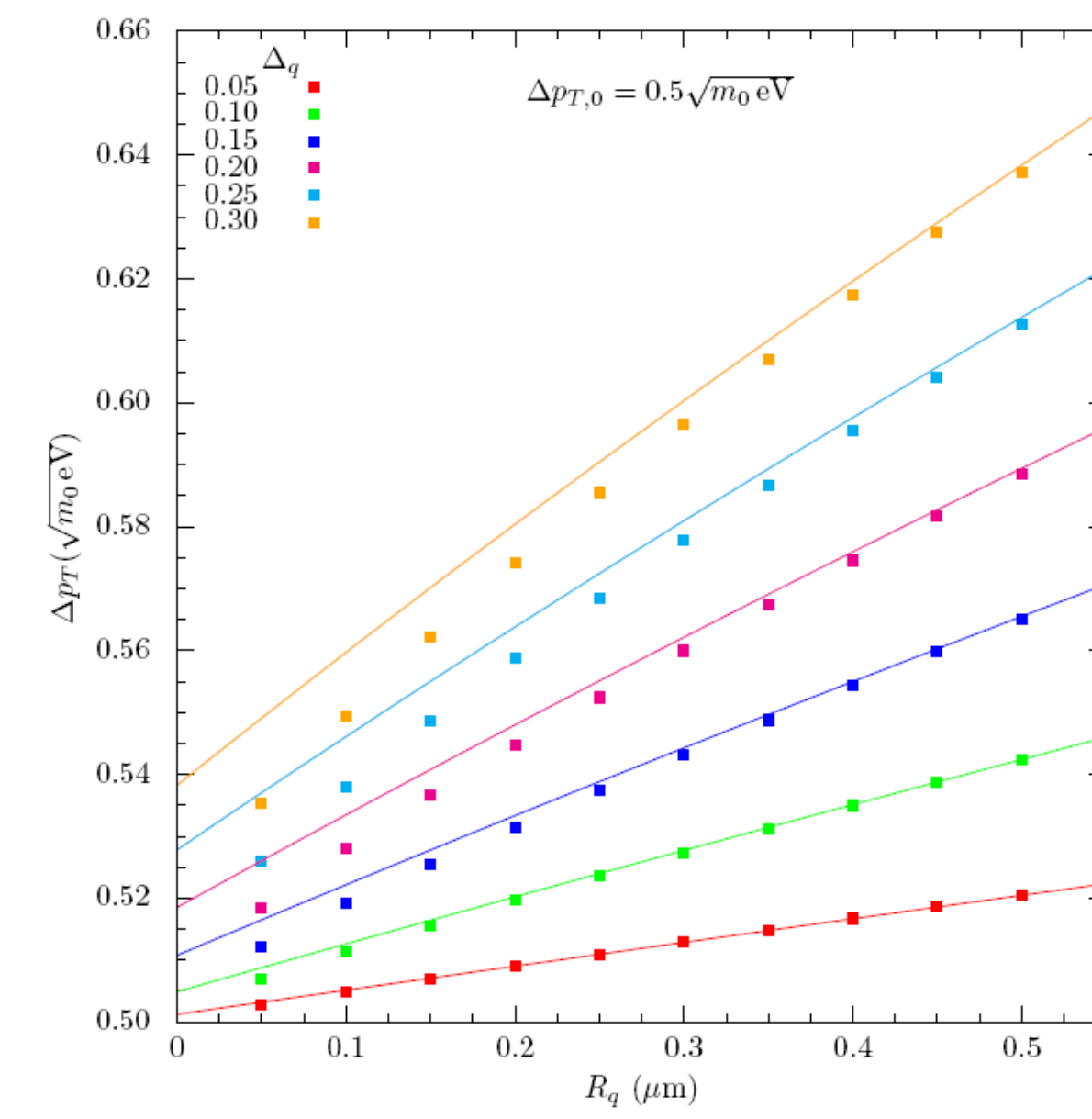
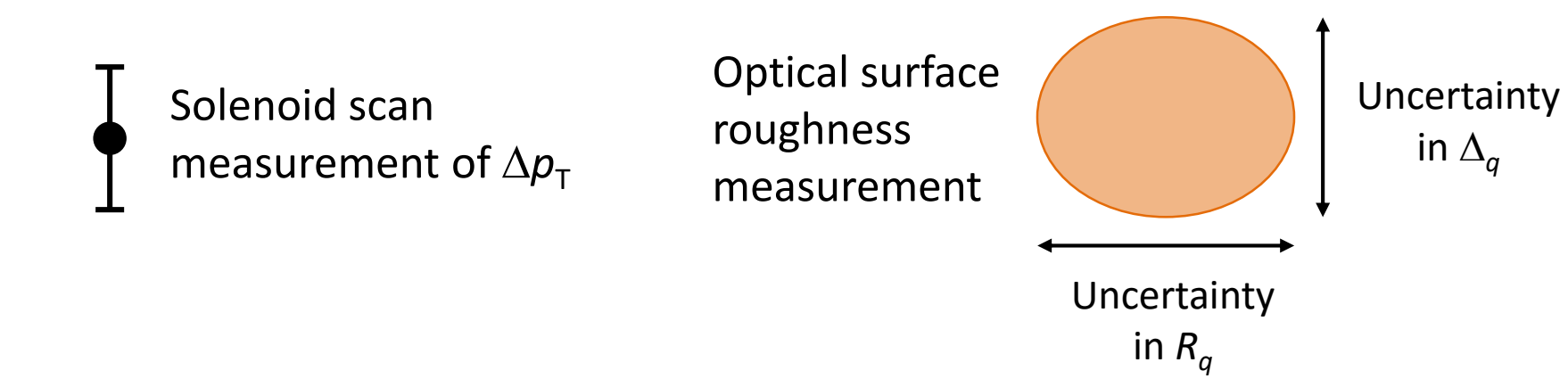
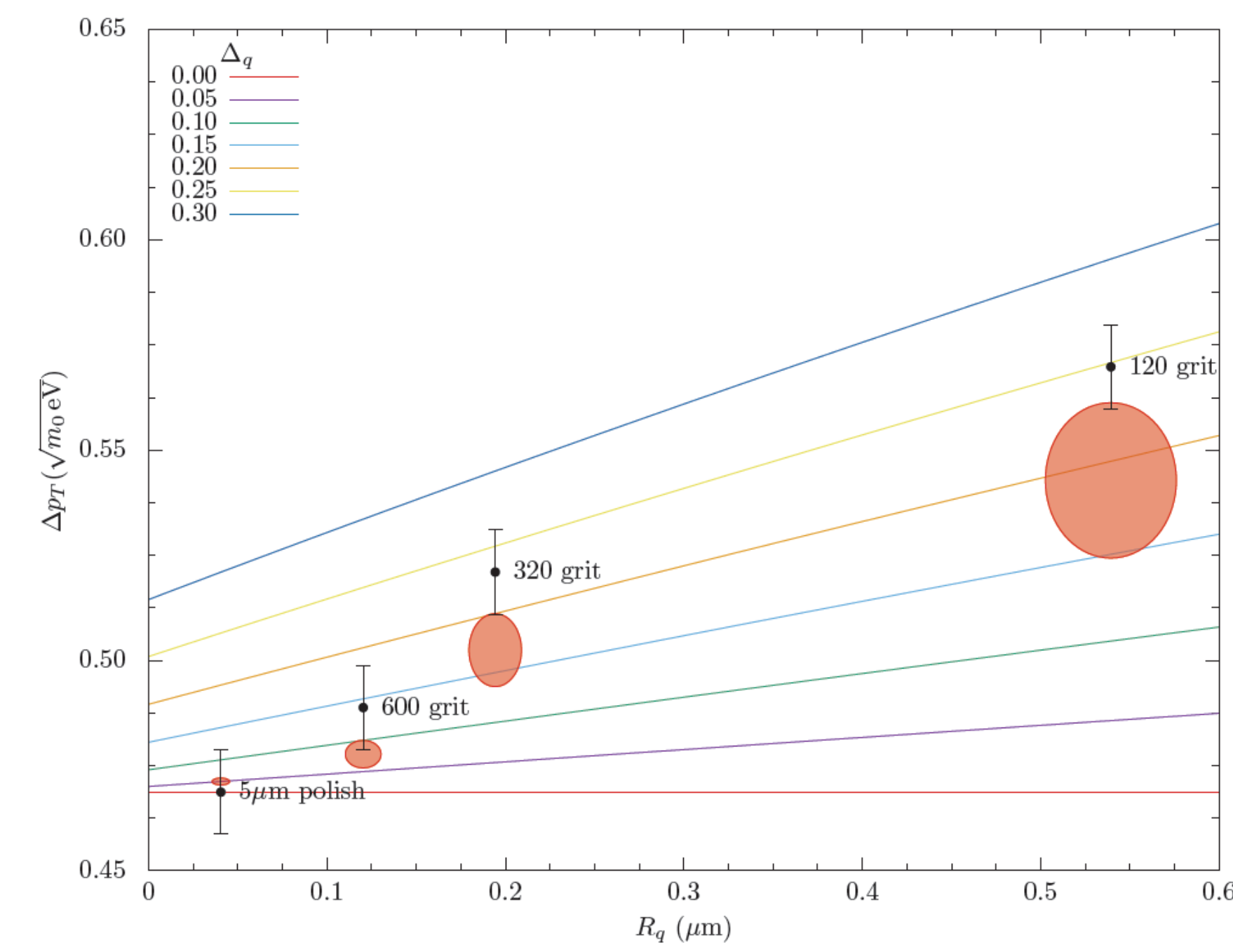


Figure 3
Comparison of analytic approximation (lines) with numeric electron trajectory simulations (points) when $E_0 = 0.5$ MV/m and $\Delta E = 0.25$ eV for a rough photocathode with an intrinsic RMS transverse momentum $\Delta p_{T,0} = 0.5$ ($m_0 \cdot \text{eV}$)^{1/2}.

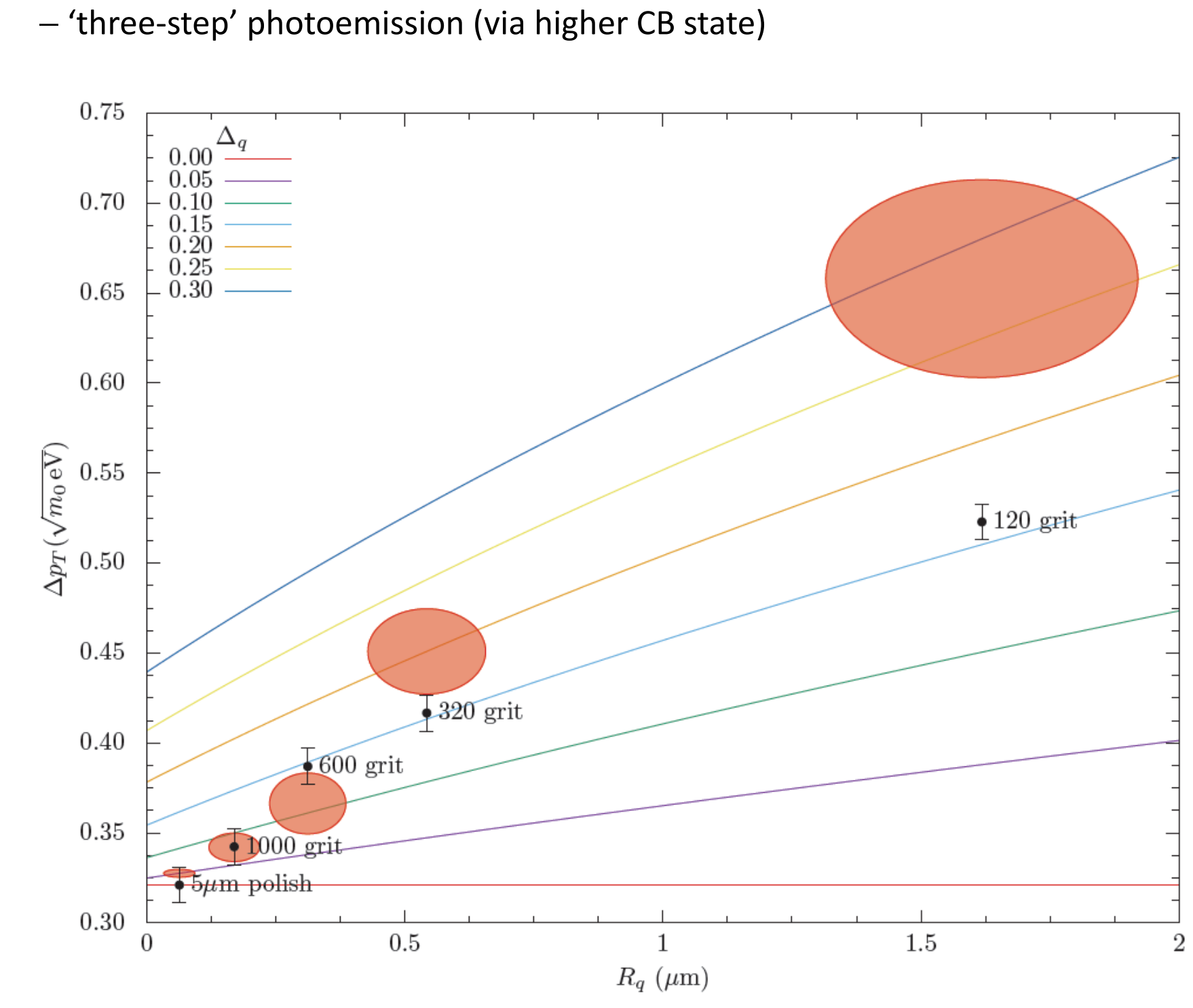
Results



Polycrystalline Mo – 'one-step' photoemission



GaSb(001) – 'three-step' photoemission (via higher CB state)



- Results consistent with theoretical expectations for polycrystalline Mo [a 'hard' material] and at low values of R_q for GaSb(001) [a 'softer' material]
- GaSb(001) shows deviation from theoretical analysis when $R_q > 1\mu\text{m}$

Conclusions

- Emission properties of rough photocathodes are well described by

$$\Delta p_T = \sqrt{(\Delta p_{T,0})^2 + (\Delta p_{T,slope})^2 + (\Delta p_{T,field})^2}$$

where $\Delta p_{T,slope} = m_0 v \Delta_q$
and $\Delta p_{T,field} = \sqrt{\frac{\pi R_q \Delta_q q E_0 m_0}{2}}$

... does **not** include $\Delta\phi$ effects on surface
- For hand-lapped photocathodes (5 μ m polish); $R_q \approx 50$ nm and $\Delta_q \approx 3^\circ$
 - ⇒ Minimal effect on Δp_T in 'low field' DC guns
 - ∴ Prior solenoid scan measurements of Δp_T validated
- Deviation from theory for GaSb(001) at high R_q is likely due to penetration of acceleration field into the dielectric photocathode:

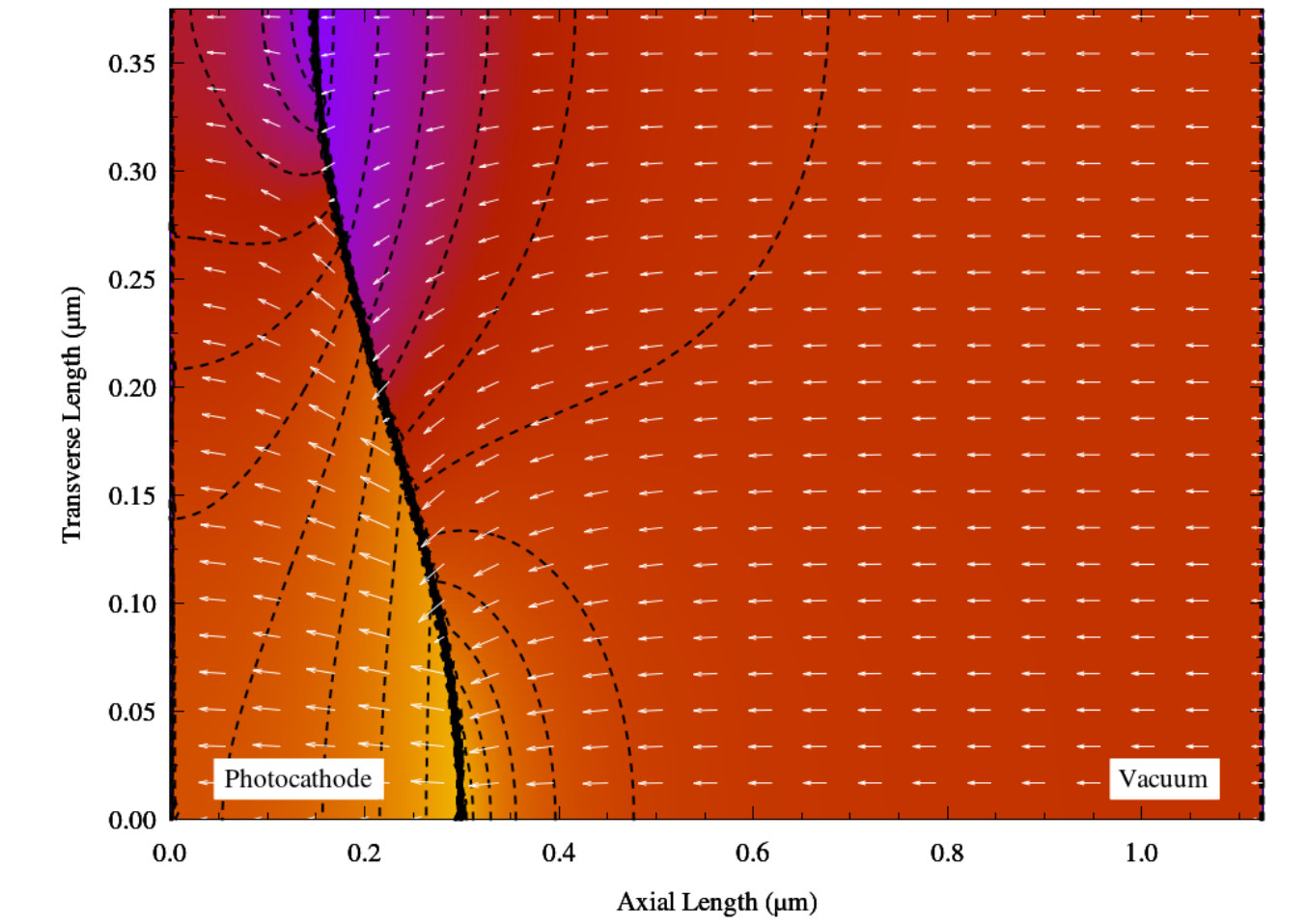


Figure 3
The electric displacement field \mathbf{D} (vector arrows) at the interface of the vacuum with a dielectric (relative permittivity $\epsilon_r = 15$, e.g. GaSb) photocathode surface with a sinusoidal roughness (75 μ m period, 0.75 μ m amplitude); equipotentials indicated by dashed lines.

Lemma

Extension to 3-D with $h(x) = a \cos\left(\frac{2\pi x}{b}\right) \cos\left(\frac{2\pi y}{b}\right)$ gives:

$$\text{Areal RMS roughness, } S_q = \sqrt{\frac{1}{L^2} \int_0^L \int_0^L h^2(x, y) dx dy} = \frac{a}{2}$$

$$\text{Areal RMS slope, } \Gamma_q = \sqrt{\frac{1}{L^2} \int_0^L \int_0^L \left(\frac{\partial h(x, y)}{\partial x} \right)^2 dx dy} = \frac{\pi a}{b}$$

⇒ $\Delta p_{T,slope} = m_0 v \Gamma_q$ ← Same as 2-D

$$\text{and } \Delta p_{T,field} = \sqrt{\frac{\pi S_q \Gamma_q q E_0 m_0}{2\sqrt{2}}} \leftarrow \text{Within } 2^{-1/4} \text{ of 2-D}$$

Acknowledgements

