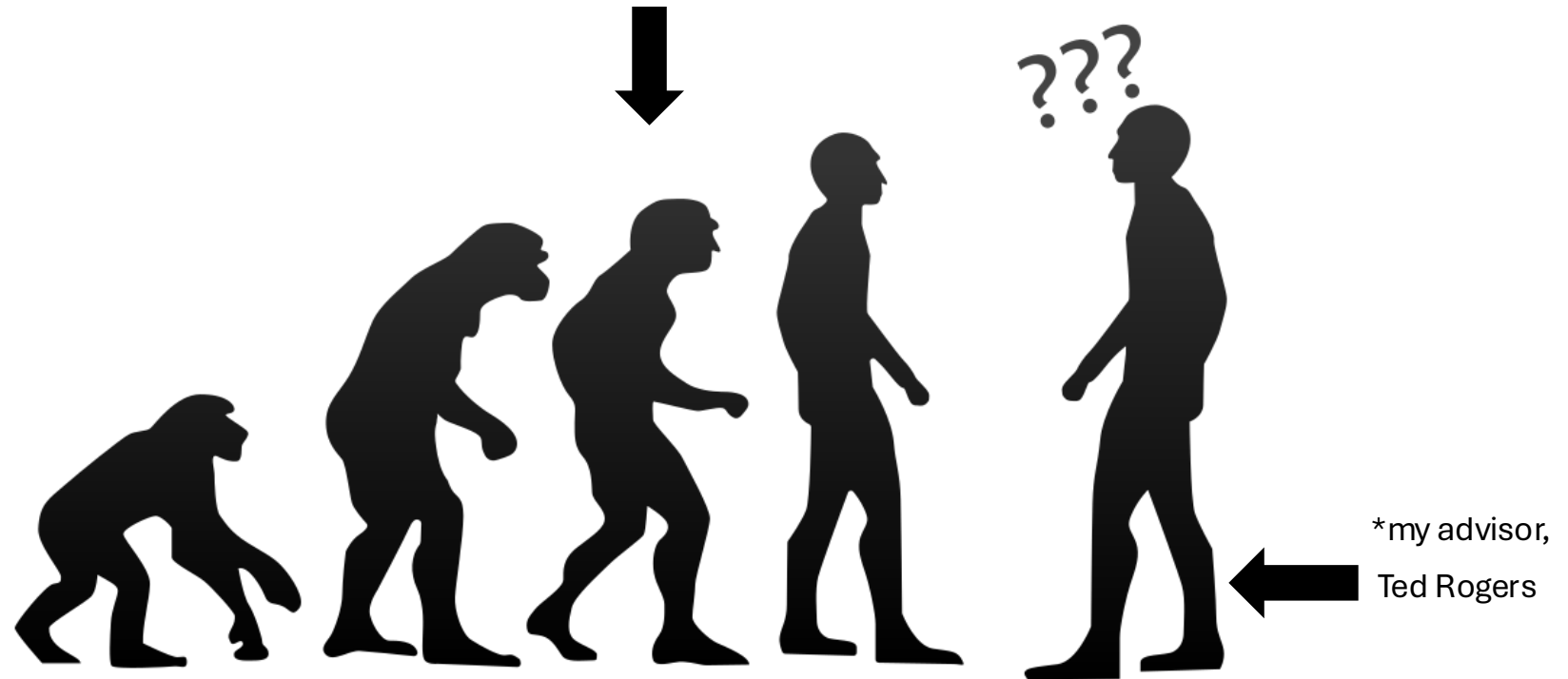


***Evolution* of Dihadron Fragmentation function**

**Kazuki Makino\***

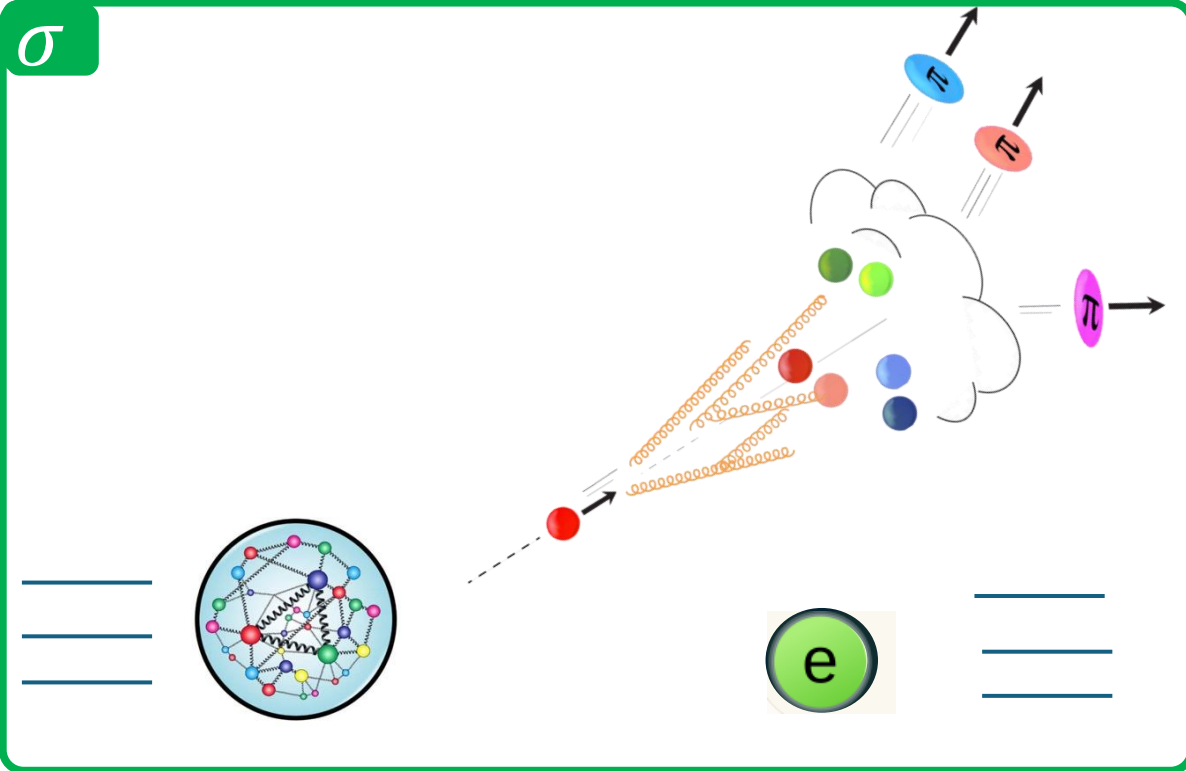


# Factorization

Isolate perturbative & nonperturbative QCD

Factorization formula

$$\sigma \sim \hat{\sigma} \otimes d \otimes f$$

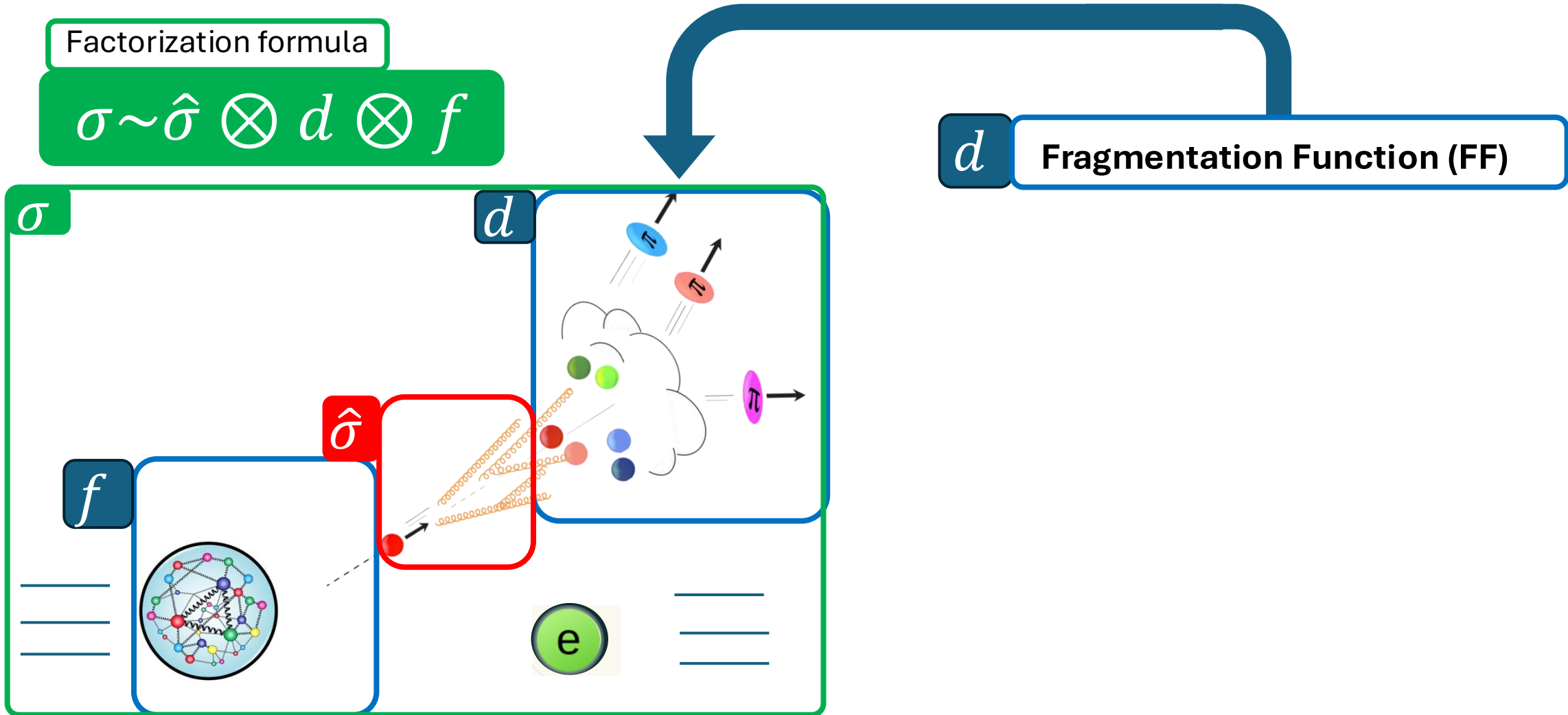


# Factorization

Isolate perturbative & nonperturbative QCD

Factorization formula

$$\sigma \sim \hat{\sigma} \otimes d \otimes f$$



# Fragmentation function (FFs)

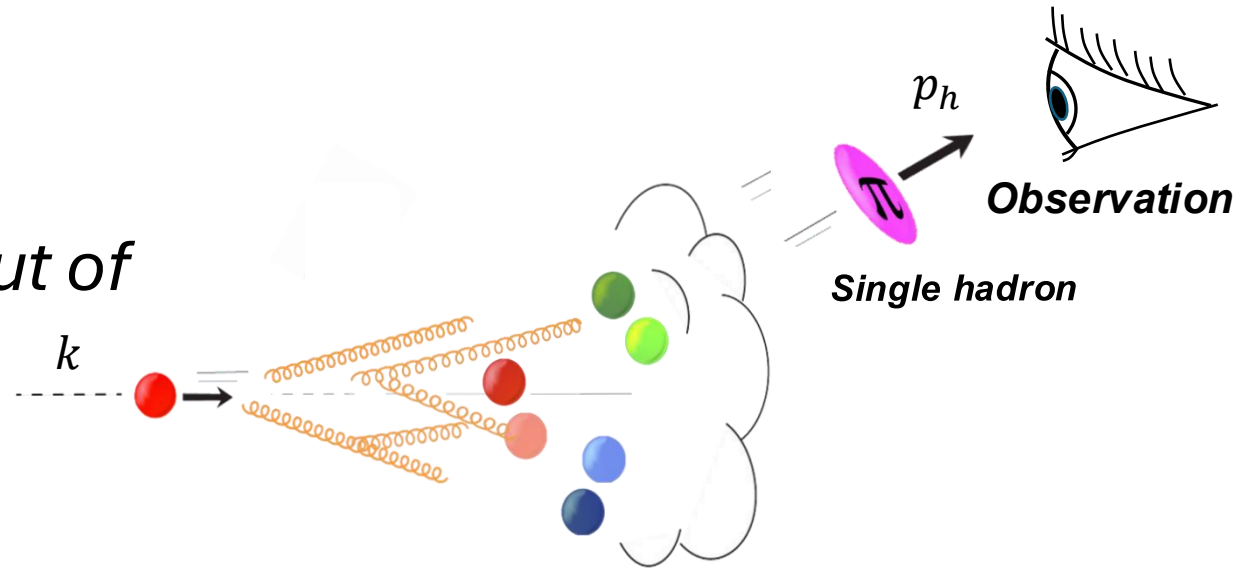
How does a quark turn into the hadron?

$d(\xi)$

Fragmentation function (FFs)

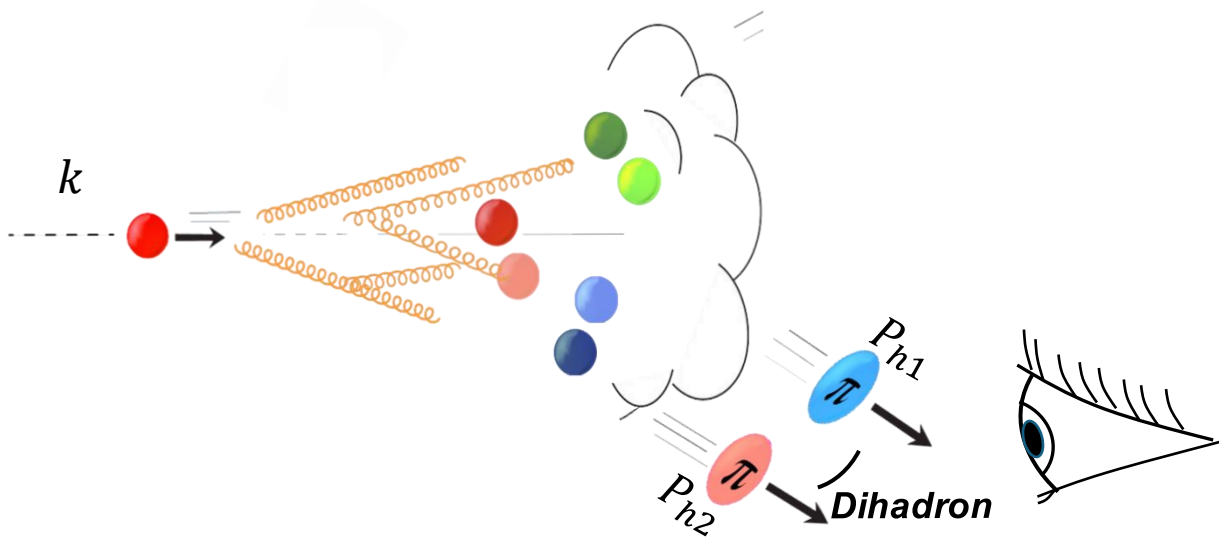
Probability of finding a hadron in a quark as a function of  $\xi = p_h^+ / k^+$

Quark coming out of a collision



# **(Dihadron)** Fragmentation function (FFs)

How does a quark turn into the ***pair of*** hadron?



$d(\xi)$

**Fragmentation function (FFs)**

Probability of finding a hadron in a quark as a function of  $\xi = p_h^+/k^+$

$d(\xi)$

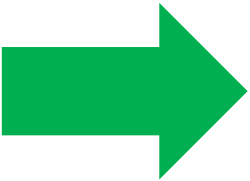
**Dihadron Fragmentation function (DiFFs)**

Probability of finding a pair of hadrons in a quark as a function of  $\xi = p_h^+/k^+$

$$P_h = P_{h1} + P_{h2}$$

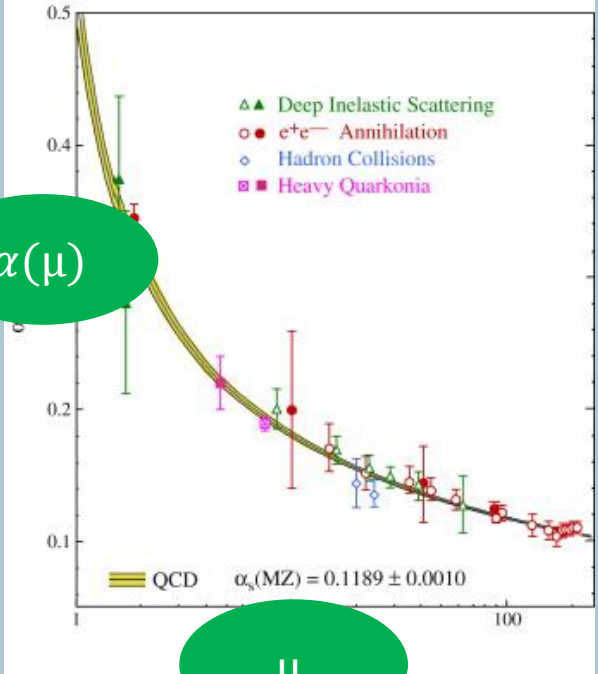
# Evolution

$\alpha(\mu)$  evolves in (energy) scale  $\mu$



Fragmentation functions evolves in energy scale  $\mu$

## Asymptotic freedom



Coupling

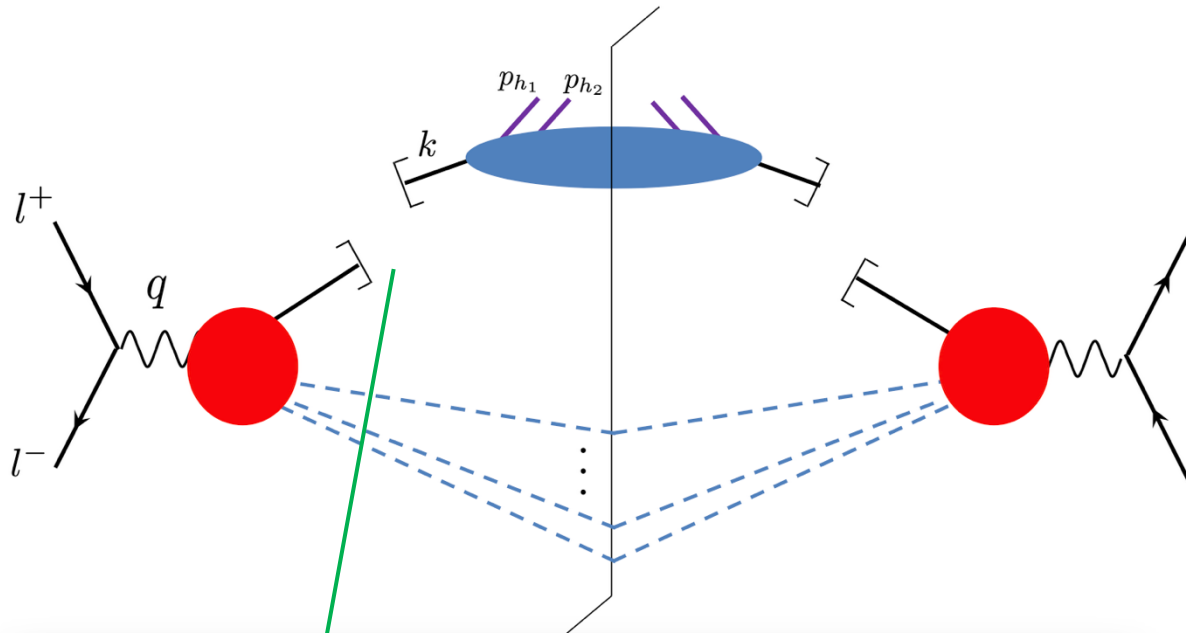
$$\frac{d}{d \ln \mu} g_s(\mu) = \beta g_s(\mu)$$



Fragmentation function (DGLAP)

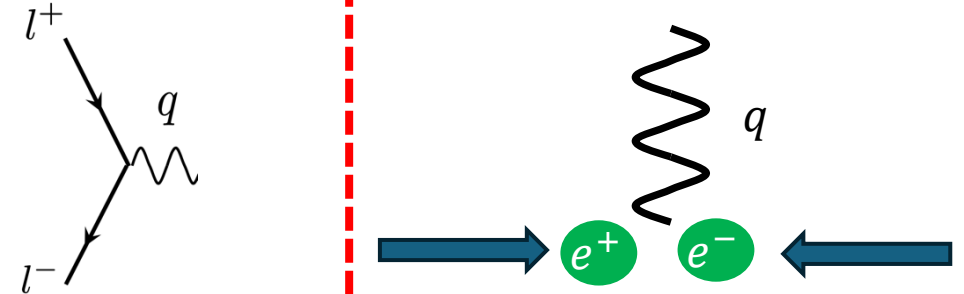
$$\frac{d}{d \ln \mu} d(\xi; \mu) = [P \otimes d](\xi; \mu)$$

# Dihadron Fragmentation function (DiFF)



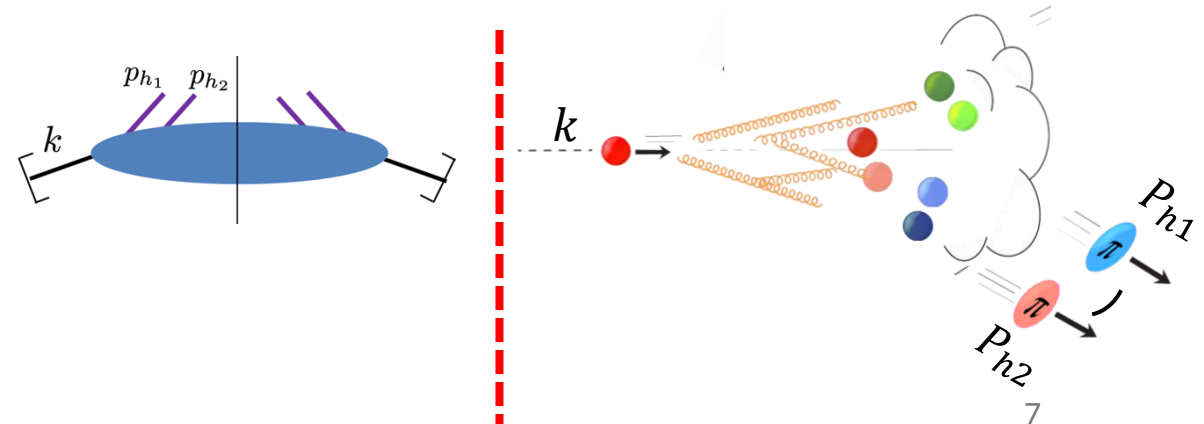
Hooks  $[[$  to indicates factorization

$Q^2$  controls how much energy is put into the system

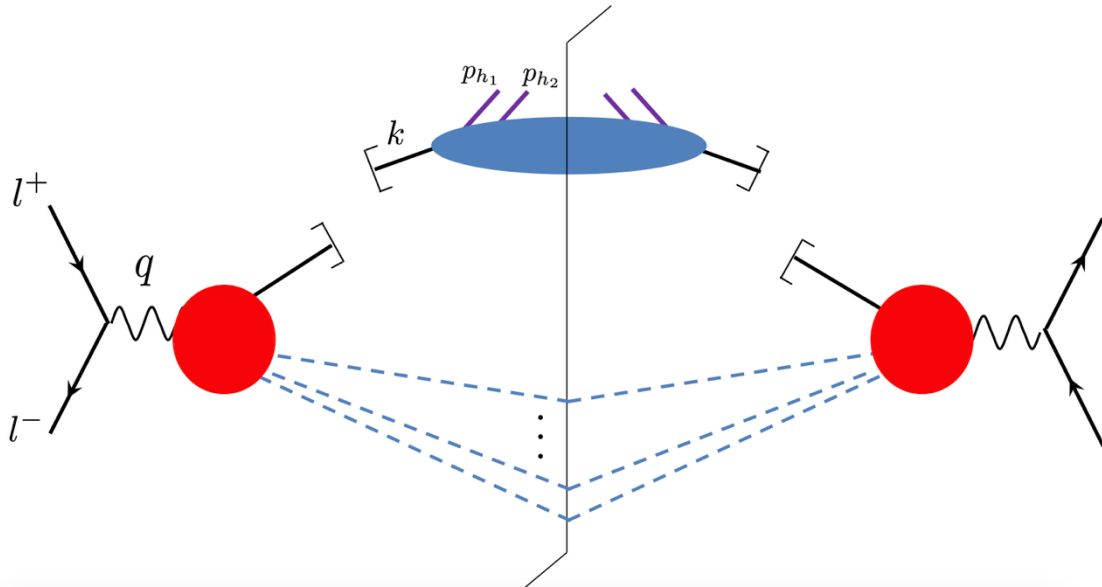


For DiFFs, invariant mass:

$$M_h^2 = (P_{h1} + P_{h2})^2$$

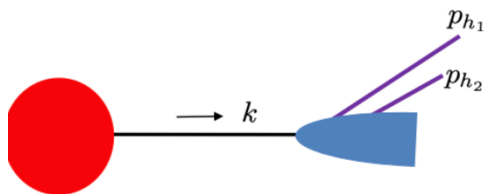


# Evolution of Dihadron Fragmentation function (DiFF) in $M_h^2$



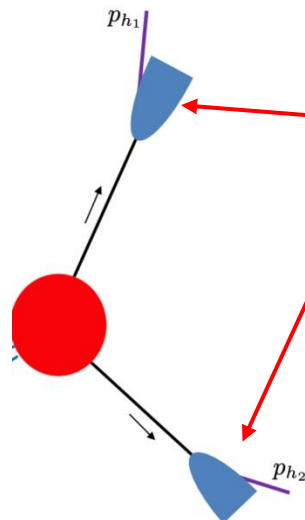
Evolution in  $M_h^2$  and  $Q^2$   
to cover full phase space

Dihadron fragmentation  
function

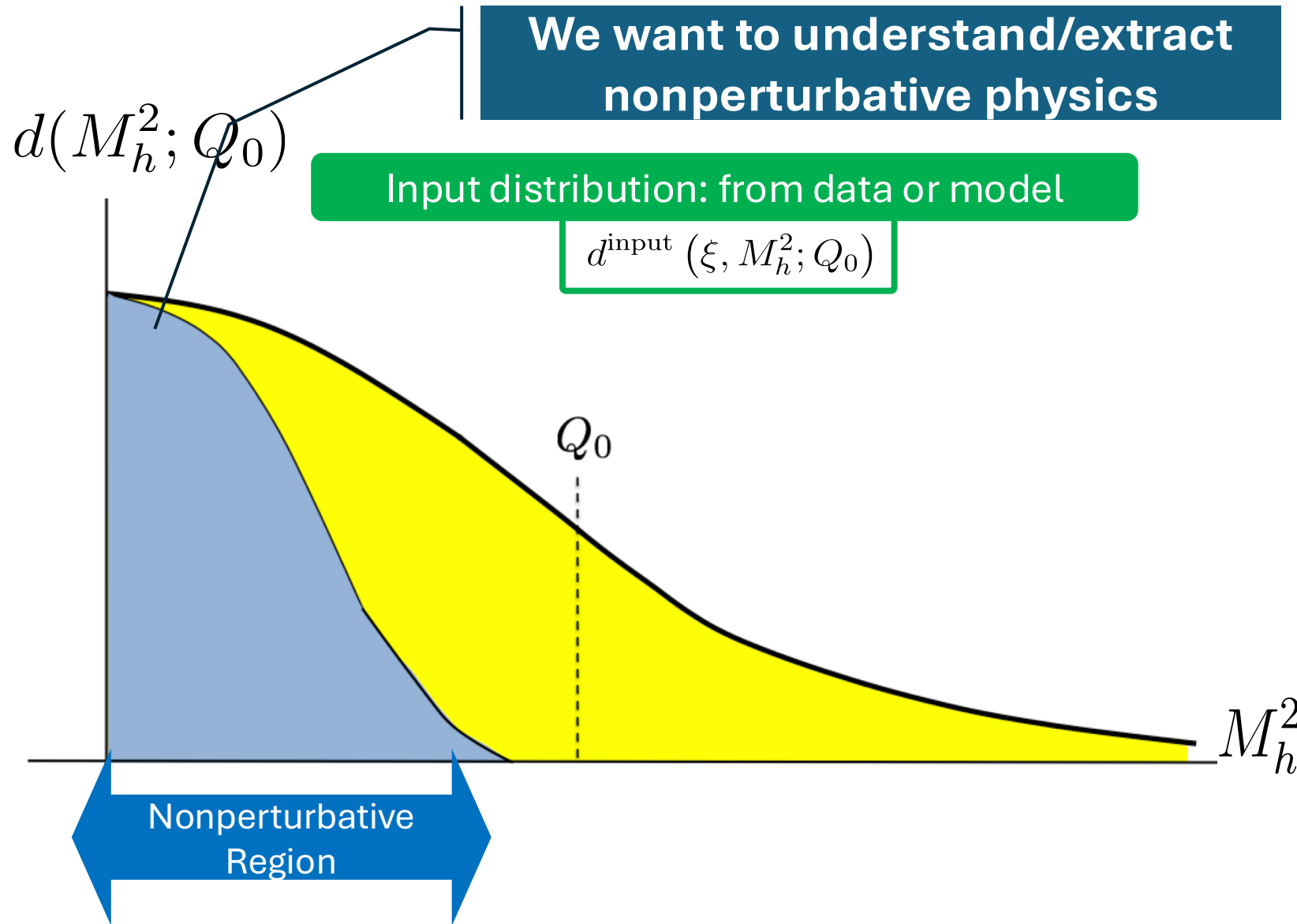


Larger  $M_h$

two single hadron  
fragmentation functions



# Evolution from small to large $M_h^2$



# Evolution from small to large $M_h^2$

We want to understand/extract nonperturbative physics

$d(M_h^2; Q_0)$

Input distribution: from data or model

$d^{\text{input}}(\xi, M_h^2; Q_0)$

Evolve *bottom-up* (Hadron structure oriented)

$\Delta d^{\text{evol}}(\xi, M_h^2; Q_0)$

high  $M_h^2$ : perturbative correction evolution

$$M_h^2 \gg Q_0^2, \alpha_s(Q_0) \ln \frac{M_h}{Q_0} \rightarrow \infty$$

Nonperturbative Region

Ultra-perturbative Region

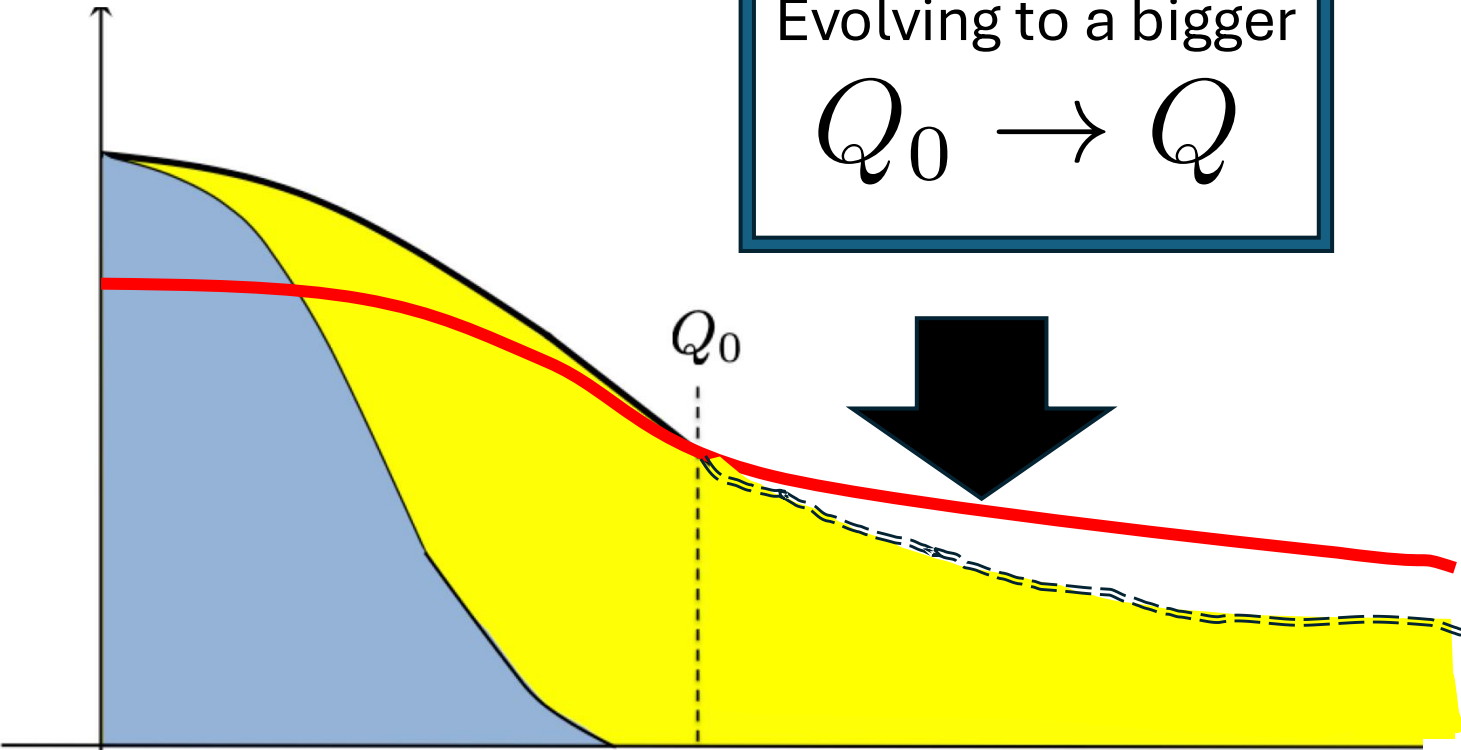
# Final DIFF

## Evolution into high $Q^2$

Needed to evolve in full range of  $M_h$  to get evolution to larger  $Q$  !

Evolving to a bigger  $Q_0 \rightarrow Q$

$d(M_h^2; Q_0)$



Fit with experiments

$M_h^2$

# DIFF evolution in a nutshell

Coupling evolution

$$\frac{d}{d \ln \mu} g_s(\mu) = \beta g_s(\mu) \quad \text{Recall}$$

Normal FF evolution in  $Q$

$$\frac{d}{d \ln \mu} d(\xi, M_h^2; \mu) = P(\xi, \alpha(\mu)) \otimes d(\xi, M_h^2; \mu) \quad \text{Recall}$$

Independent of  $M_h$

Separate evolution  
with additional scale  $M_h$

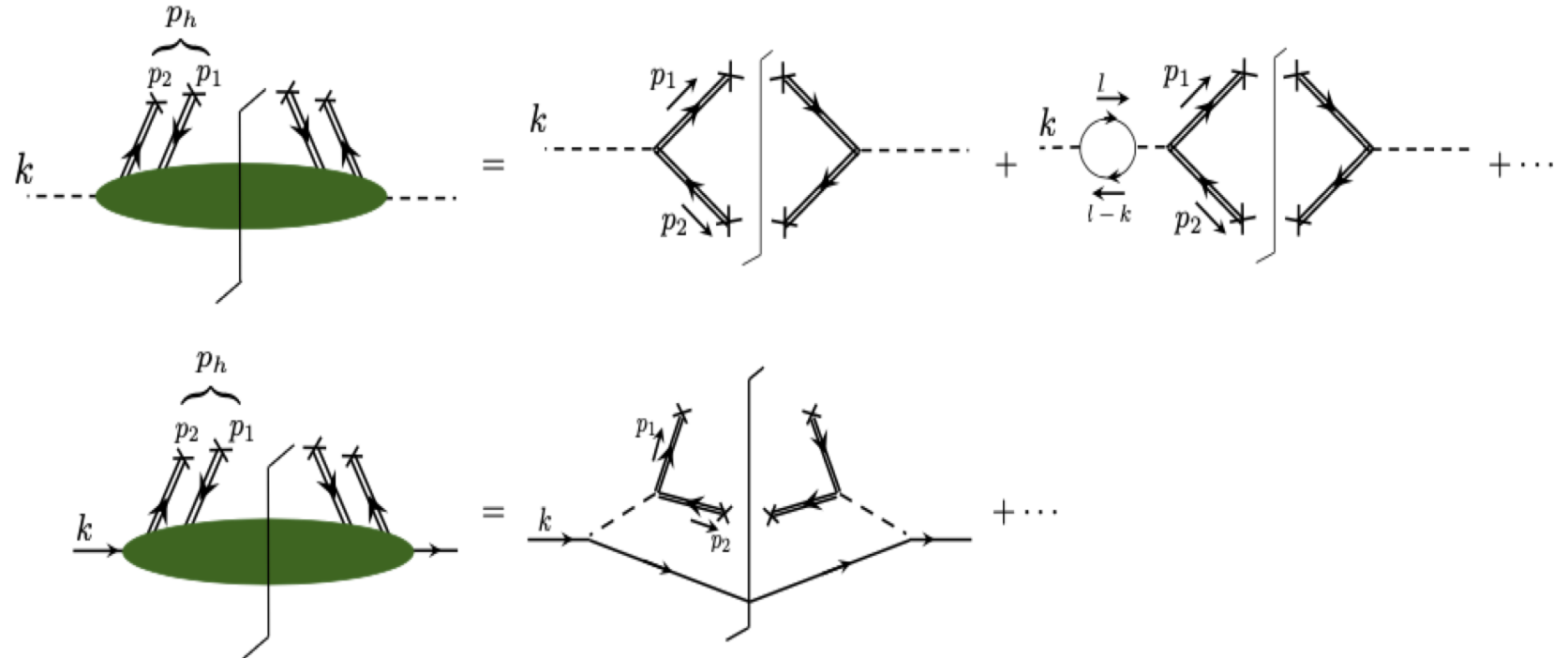
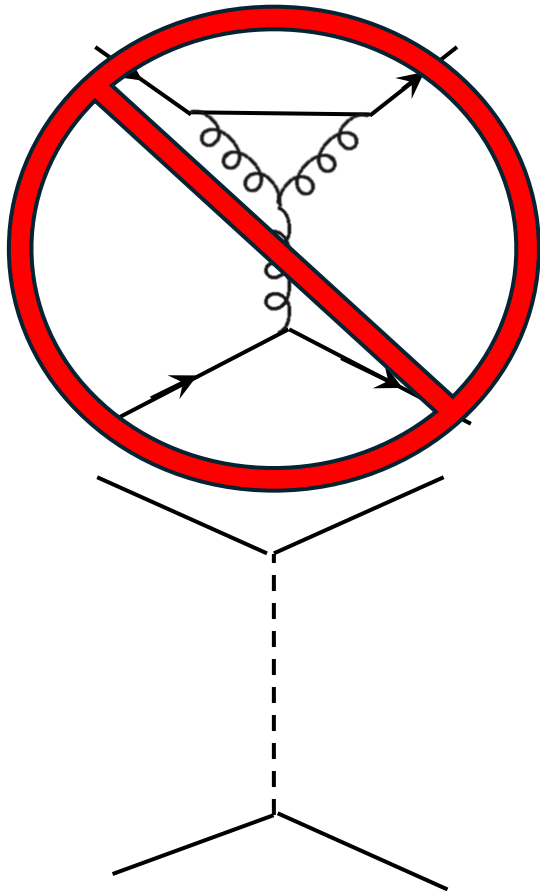
**New!**

$$\frac{\partial}{\partial \ln M_h^2} d(\xi, M_h^2; \mu) = K(\xi, M_h^2) \otimes d(\xi, M_h^2; \mu)$$

Independent of  $\mu$

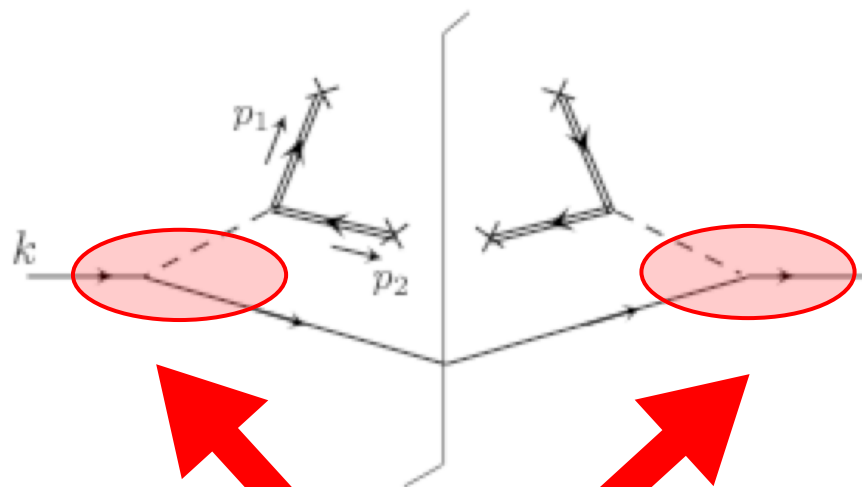
# (DiFFs) in Yukawa theory

Proof of principle in a theory where everything is easily calculable



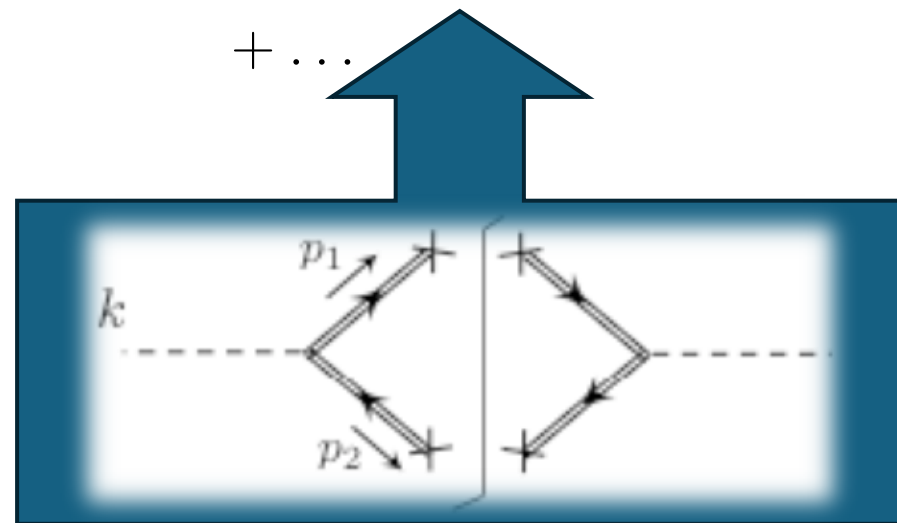
# (DiFFs) in Yukawa theory

## Evolution in $M_h$



$$\frac{\partial}{\partial \ln M_h^2} d_{q \rightarrow h\bar{h}}^{\text{evol}}(\xi, M_h^2; Q_0) = d_{s \rightarrow h\bar{h}}^{\text{evol}}(\xi, M_h^2; Q_0) \otimes K_{q \rightarrow s}^{\text{pert}}(\xi, M_h^2)$$

+ ...



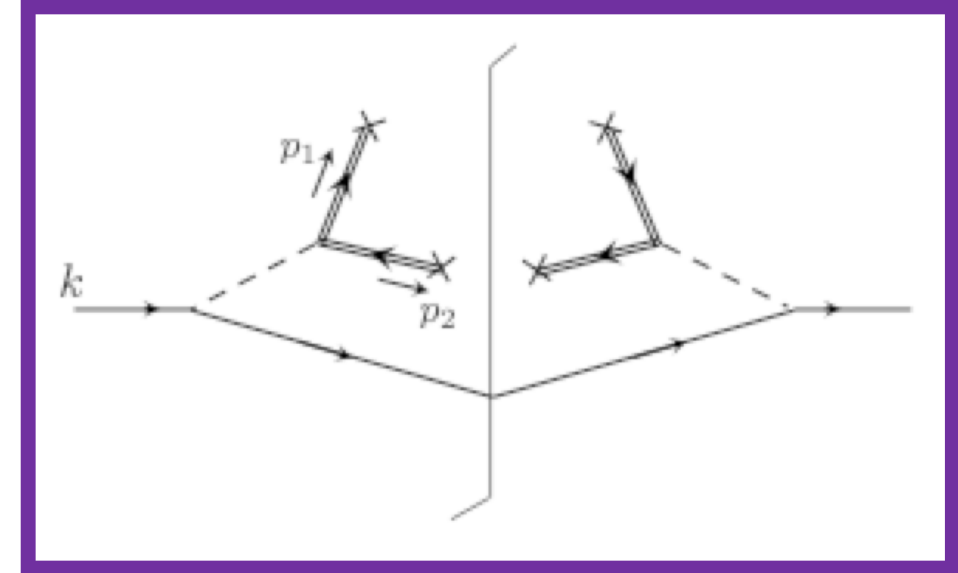
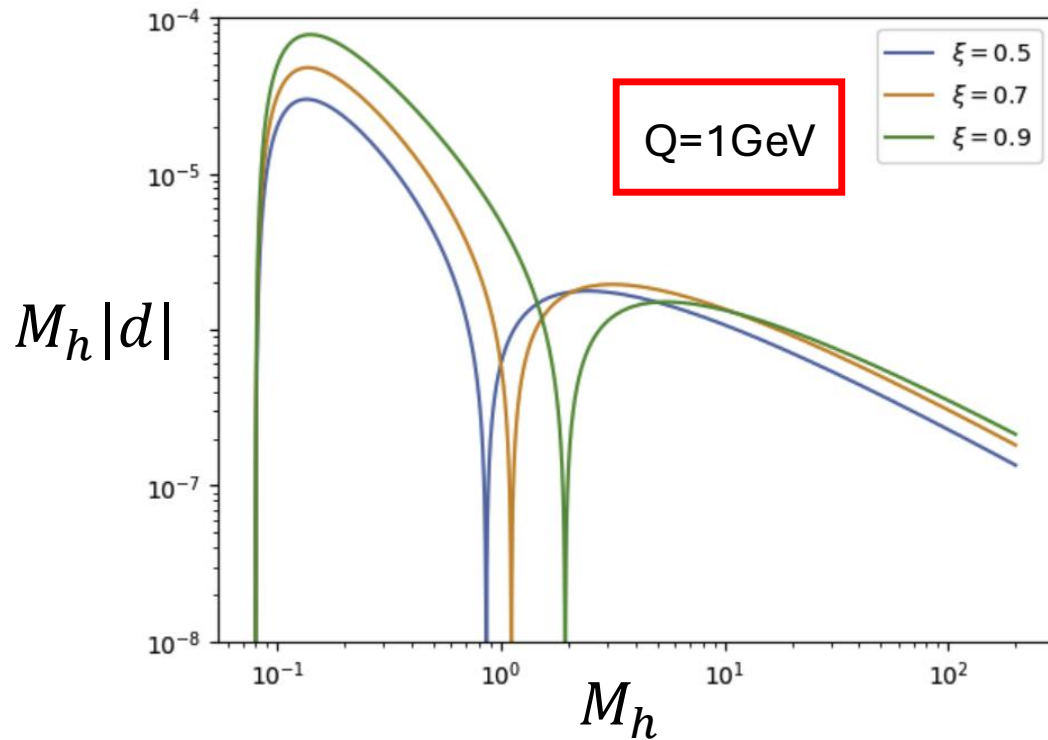
**Perturbative calculable**

- at large  $M_h^2$  in QCD
- at all  $M_h^2$  in Yukawa theory

# evolution in $M_h$

## DiFFs in Yukawa theory

$$d(\xi, M_h^2; Q_0) = d^{\text{input}}(\xi, M_h^2; Q_0) + \Delta d^{\text{evol}}(\xi, M_h^2; Q_0)$$

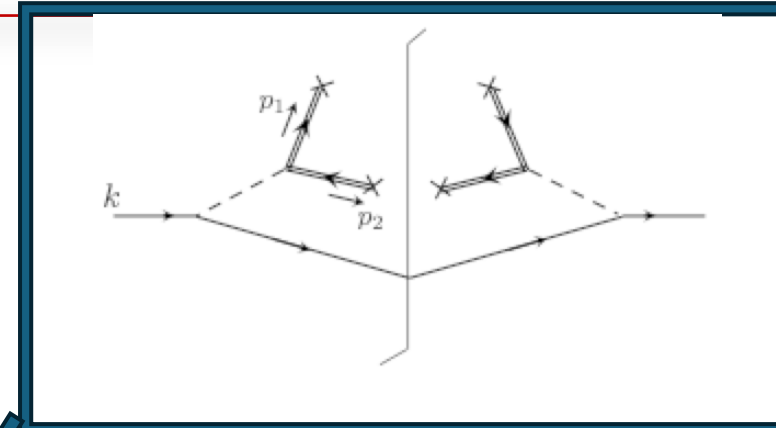


# Numerical evolution in Q

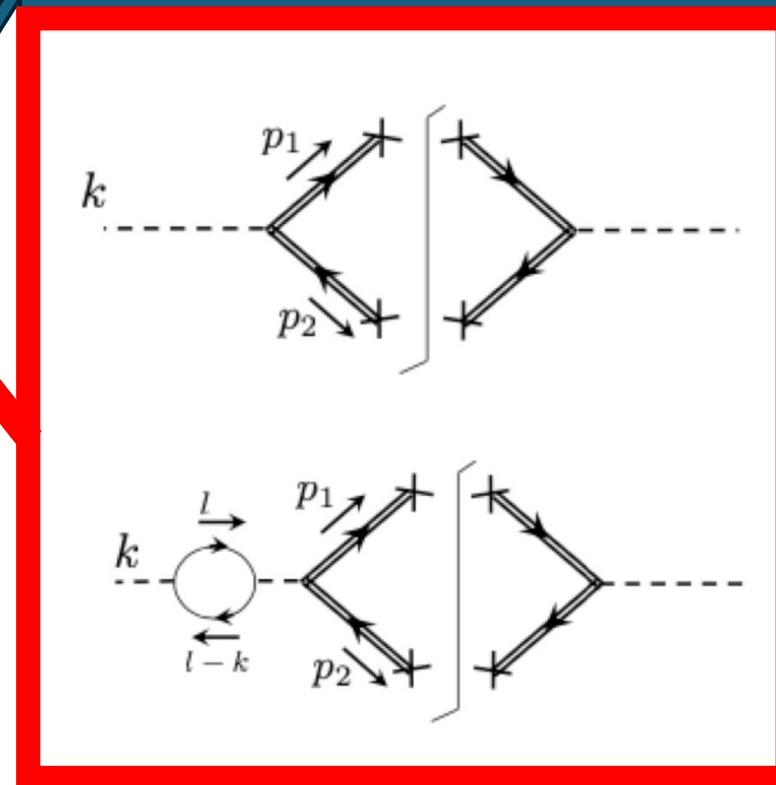
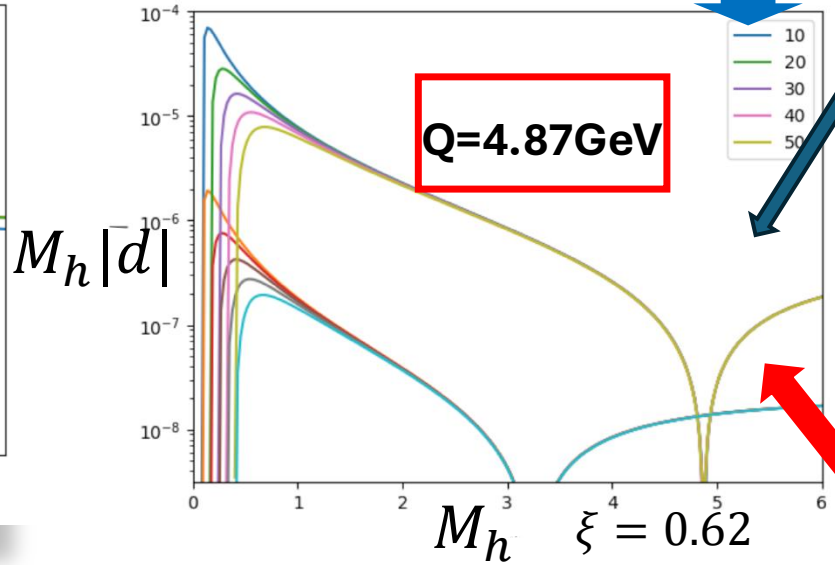
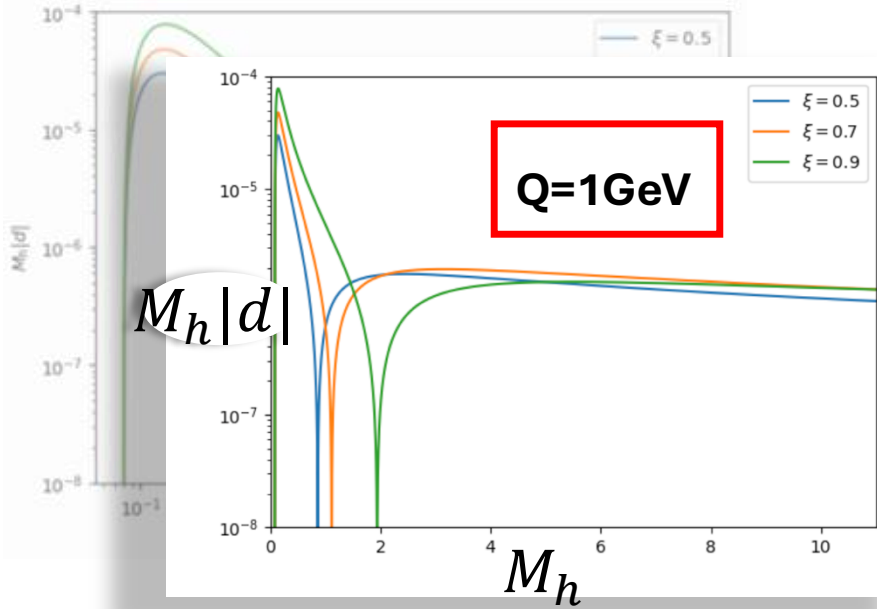
## DiFF in Yukawa theory

$$\frac{d}{d \ln Q} d_{s \rightarrow h\bar{h}}(\xi, M_h^2; Q) = \int_{\xi}^1 \frac{dz}{z} d_{f \rightarrow h\bar{h}}(\xi/z, M_h^2; Q) P_{sf}(z) + \int_{\xi}^1 \frac{dz}{z} d_{s \rightarrow h\bar{h}}(\xi/z, M_h^2; Q) P_{ss}(z)$$

$$\frac{d}{d \ln Q} d_{f \rightarrow h\bar{h}}(\xi, M_h^2; Q) = \int_{\xi}^1 \frac{dz}{z} d_{f \rightarrow h\bar{h}}(\xi/z, M_h^2; Q) P_{sf}(z) + \int_{\xi}^1 \frac{dz}{z} d_{s \rightarrow h\bar{h}}(\xi/z, M_h^2; Q) P_{fs}(z)$$

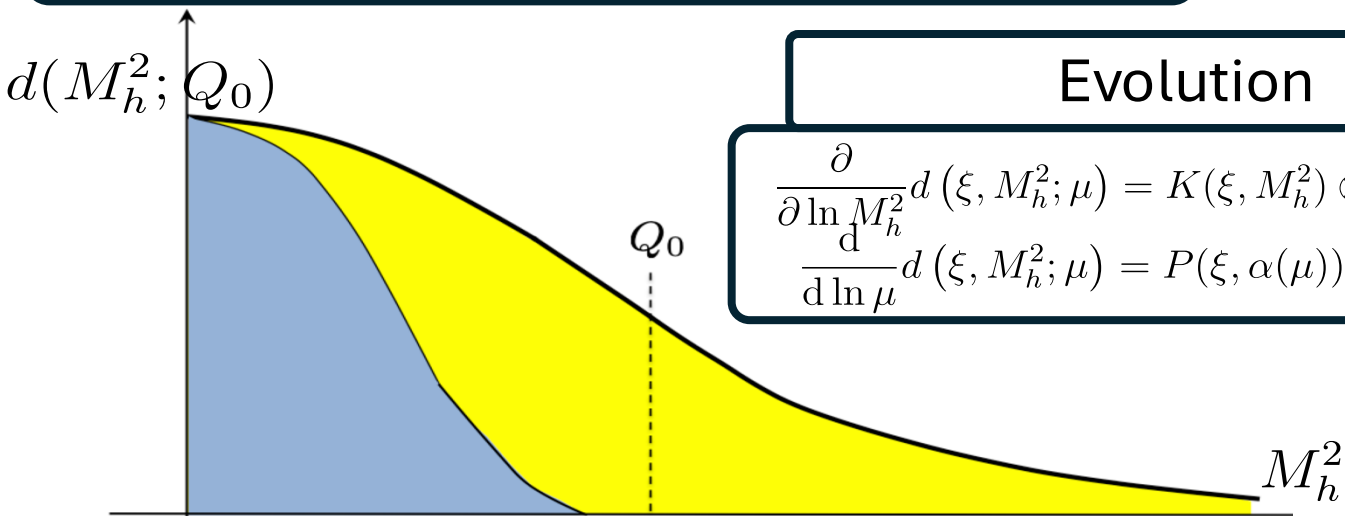
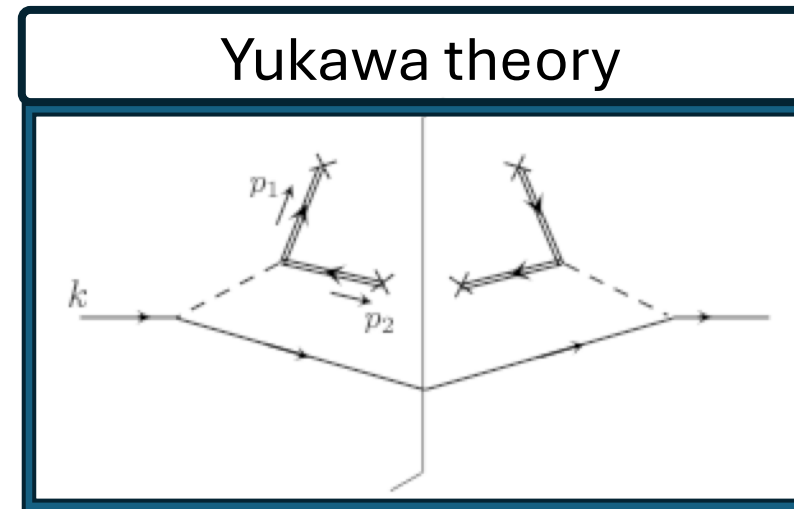
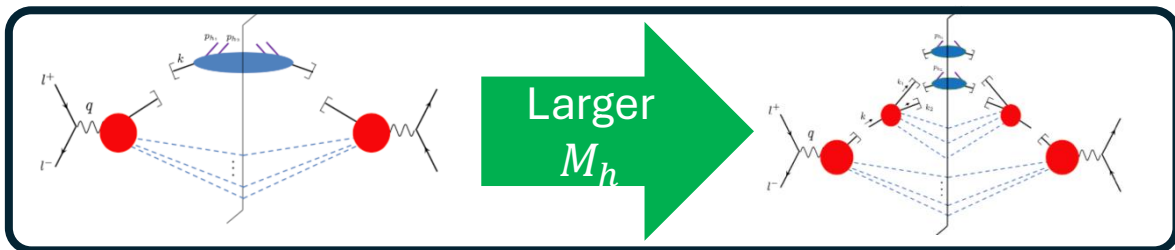


Different particle mass scales



# DiFFs: summary

- Defining  $d$  in all region of  $M_h^2$  and  $Q^2$ 
  - Using separate evolution for  $M_h^2$  and  $Q^2$
- Proof of principle in Yukawa theory



Evolution

$$\frac{\partial}{\partial \ln M_h^2} d(\xi, M_h^2; \mu) = K(\xi, M_h^2) \otimes d(\xi, M_h^2; \mu)$$

$$\frac{d}{d \ln \mu} d(\xi, M_h^2; \mu) = P(\xi, \alpha(\mu)) \otimes d(\xi, M_h^2; \mu)$$

*The End*