

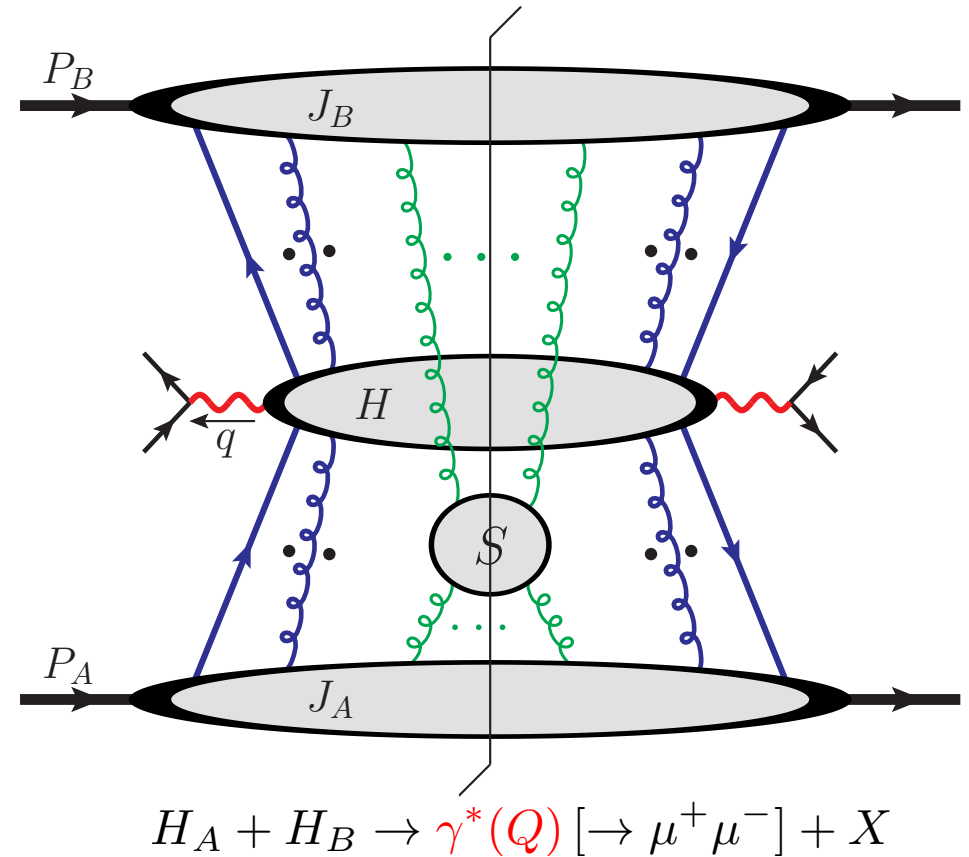
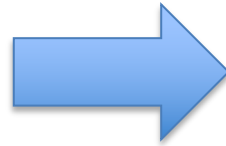
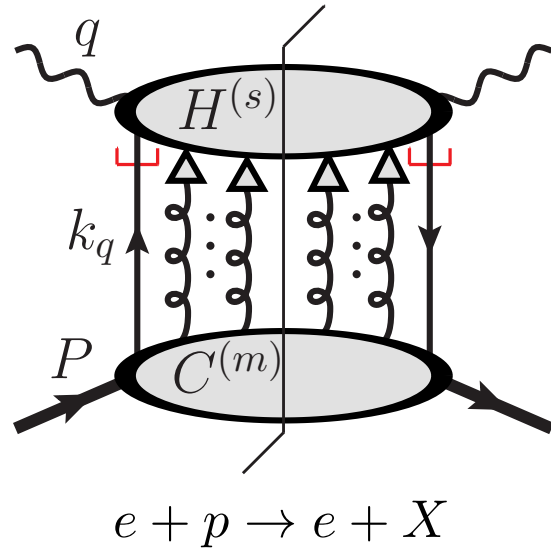


QCD factorization and global analysis

Zhite Yu (BNL, High Energy Physics)

- Introduction to QCD factorization
- Parton distribution functions and global analysis
- From inclusive to exclusive processes
- Generalized parton distributions and global analysis

From DIS to Drell-Yan

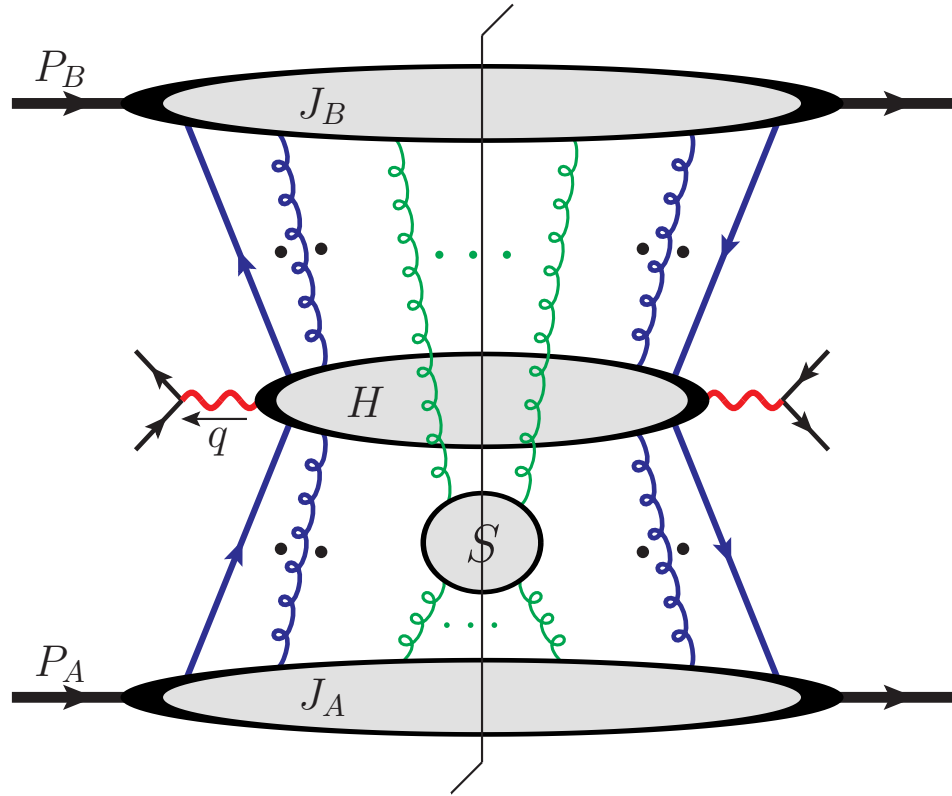


➤ **DIS factorization:** factorize collinear pinches.

➤ **Goal of Drell-Yan factorization:**

- Collinear sectors from two hadrons are separated **from each other**
- They are also factorized **from the hard part**
- Connected to the PDF in DIS

Drell-Yan factorization



$$= \sum_{a,b} \int_0^1 dx_a \int_0^1 dx_b f_{a/H_A}(x_a, \mu) f_{b/H_B}(x_b, \mu) C_{ab}(x_a, x_b; Q/\mu)$$

- Key: inclusiveness \Rightarrow unitarity cancels soft gluons
- PDFs from Drell-Yan factorization are the same as in DIS
- Hard part involves subtraction for collinear pinches

Global analysis of PDFs

➤ Parametrize PDFs at a given scale μ_0

E.g., JAM20-SIDIS $f_i(x; \mu_0) = T_i(x; \mathbf{a}_i) = M_i x^{\alpha_i} (1 - x)^{\beta_i} (1 + \gamma_i \sqrt{x} + \delta_i x)$

$$i = g, u, d, s, c, \bar{u}, \bar{d}, \bar{s}, \bar{c}$$

JAM, PRD 104, 016015 (2021)

9 independent flavors. Typically, μ_0 is just above m_c

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➤ Solve DGLAP evolution to obtain PDFs at scale Q

$$\frac{d}{d \ln \mu^2} f_i(x, \mu) = \sum_j \int_x^1 \frac{dz}{z} P_{ij}(x/z) f_j(z, \mu)$$

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➤ Use factorization to make prediction for observables at scale Q

E.g., DIS structure function $F_1^h(x_B, Q^2) = \sum_i \int_{x_B}^1 \frac{dx}{x} f_{i/h}(x, \mu) \hat{F}_1^i\left(\frac{x_B}{x}, \frac{Q^2}{\mu^2}\right)$

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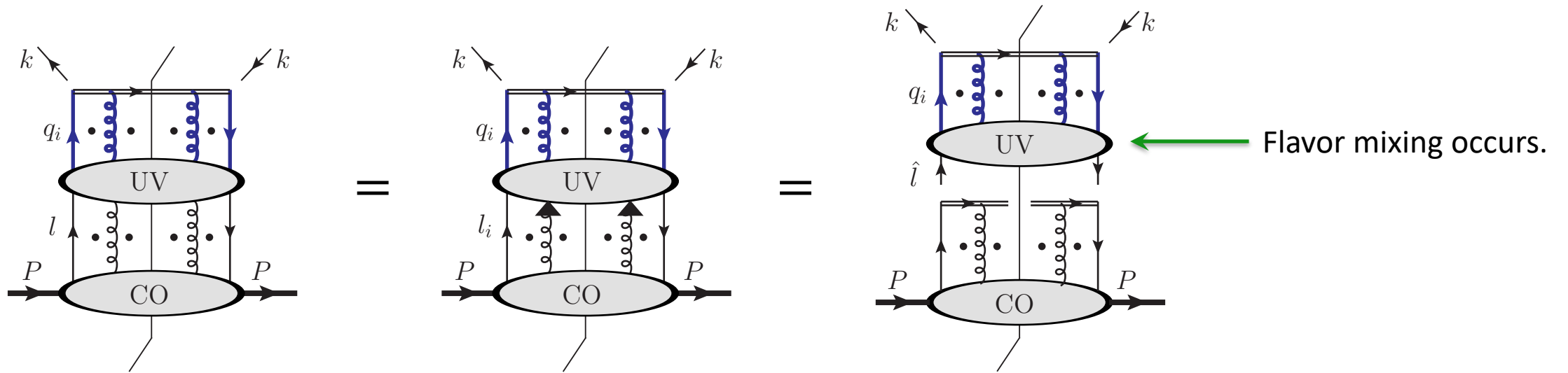
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➤ Compare with experimental measurements

- Design loss function, e.g., $\chi^2(\vec{a})$, and minimize it to find the best fit
- Quantify uncertainties

DGLAP equation

Recall: DGLAP arises from **UV renormalization** of bare PDFs



$$f_{i/H}(x, \mu) = \sum_j \int_x^1 \frac{dz}{z} Z_{ij}(x/z, 1/\epsilon; \mu) f_{j/H}^{\text{bare}}(z, 1/\epsilon)$$

- Z_{ij} collects all UV poles. Scheme dependent. Typically, $\overline{\text{MS}}$ scheme.
- **UV property of operator.** Does not depend on target state or **parton mass**.
- Perturbatively calculated, order by order.

Symmetry of DGLAP kernel

Replace H by a parton state.

$$f_{i/k}(x, \mu) = \sum_j \int_x^1 \frac{dz}{z} Z_{ij}(x/z, 1/\epsilon; \mu) f_{j/k}^{\text{bare}}(z, 1/\epsilon) \quad i, j, k = q, \bar{q}, g$$

❖ Charge conjugation symmetry

$$f_{i/k}^{\text{bare}} = f_{\bar{i}/\bar{k}}^{\text{bare}} \longrightarrow \text{UV} \left[f_{i/k}^{\text{bare}} \right] = \text{UV} \left[f_{\bar{i}/\bar{k}}^{\text{bare}} \right] \longrightarrow Z_{ij} = Z_{\bar{i}\bar{j}} \longrightarrow f_{i/k} = f_{\bar{i}/\bar{k}}$$

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❖ Together with $\text{SU}(n_f)$ flavor symmetry

$$P_{q_i g} = P_{\bar{q}_i g} \equiv P_{qg}, \quad P_{gq_i} = P_{g\bar{q}_i} \equiv P_{gq}$$

$$P_{ud} = P_{\bar{u}\bar{d}} = P_{us} = P_{\bar{u}\bar{s}} = \dots \quad P_{uu} = P_{\bar{u}\bar{u}} = P_{dd} = P_{\bar{d}\bar{d}} = \dots$$

$$P_{u\bar{d}} = P_{\bar{u}d} = P_{u\bar{s}} = P_{\bar{u}s} = \dots \quad P_{u\bar{u}} = P_{\bar{u}u} = P_{d\bar{d}} = P_{\bar{d}d} = \dots$$

Seven independent DGLAP kernels

$$P_{q_i g} = P_{\bar{q}_i g} \equiv P_{qg}, \quad P_{gq_i} = P_{g\bar{q}_i} \equiv P_{gq}$$

$$P_{ud} = P_{\bar{u}\bar{d}} = P_{us} = P_{\bar{u}\bar{s}} = \dots \quad P_{uu} = P_{\bar{u}\bar{u}} = P_{dd} = P_{\bar{d}\bar{d}} = \dots$$

$$P_{u\bar{d}} = P_{\bar{u}d} = P_{u\bar{s}} = P_{\bar{u}s} = \dots \quad P_{u\bar{u}} = P_{\bar{u}u} = P_{d\bar{d}} = P_{\bar{d}d} = \dots$$

$$P_{q_i q_j} = P_{\bar{q}_i \bar{q}_j} = P_{qq}^v \delta_{ij} + P_{qq}^s = \begin{cases} P_{qq}^v + P_{qq}^s & \text{if } i = j \\ P_{qq}^s & \text{if } i \neq j \end{cases}$$

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$$P_{q_i g} = P_{\bar{q}_i g} = P_{qg},$$

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$$P_{gg}$$

Flavor decomposition of DGLAP equation

Notation: Momentum convolution $(A \otimes B)(x) \equiv \int_x^1 \frac{dz}{z} A(x/z) B(z)$

Flavor index $i, j = u, d, s, \dots$ & g
 $\bar{i}, \bar{j} = \bar{u}, \bar{d}, \bar{s}, \dots$

$$\begin{aligned} P_{q_i q_j} &= P_{\bar{q}_i \bar{q}_j} = P_{qq}^v \delta_{ij} + P_{qq}^s \\ P_{q_i \bar{q}_j} &= P_{\bar{q}_i q_j} = P_{q\bar{q}}^v \delta_{ij} + P_{q\bar{q}}^s \\ P_{q_i g} &= P_{\bar{q}_i g} = P_{qg} \\ P_{g q_i} &= P_{g \bar{q}_i} = P_{gq} \end{aligned} \quad P_{gg}$$

$$\frac{d f_i}{d \ln \mu^2} = \sum_j P_{ij} \otimes f_j + \sum_{\bar{j}} P_{i\bar{j}} \otimes f_{\bar{j}} + P_{ig} \otimes f_g$$

$$\frac{d f_{\bar{i}}}{d \ln \mu^2} = \sum_j P_{i\bar{j}} \otimes f_j + \sum_{\bar{j}} P_{\bar{i}\bar{j}} \otimes f_{\bar{j}} + P_{\bar{i}g} \otimes f_g$$

$$\frac{d f_g}{d \ln \mu^2} = \sum_j P_{gj} \otimes f_j + \sum_{\bar{j}} P_{g\bar{j}} \otimes f_{\bar{j}} + P_{gg} \otimes f_g$$

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How to disentangle flavors?

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How to disentangle flavors?



f_i and $f_{\bar{i}}$ mix with f_g in the same way



Charge conjugation symmetry!



Define $f_i^\pm \equiv f_i \pm f_{\bar{i}}$

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 P_{q_i g} &= P_{\bar{q}_i g} = P_{qg} & P_{gg} \\
 P_{g q_i} &= P_{g \bar{q}_i} = P_{gq}
 \end{aligned}$$

Task (3mins): Work out the evolution equation of f_i^\pm

$$\frac{d f_i}{d \ln \mu^2} = P_{qq}^v \otimes f_i + P_{q\bar{q}}^v \otimes f_{\bar{i}} + P_{qq}^s \otimes \sum_j f_j + P_{q\bar{q}}^s \otimes \sum_{\bar{j}} f_{\bar{j}} + P_{qg} \otimes f_g,$$

$$\frac{d f_{\bar{i}}}{d \ln \mu^2} = P_{qq}^v \otimes f_{\bar{i}} + P_{q\bar{q}}^v \otimes f_i + P_{qq}^s \otimes \sum_{\bar{j}} f_{\bar{j}} + P_{q\bar{q}}^s \otimes \sum_j f_j + P_{qg} \otimes f_g,$$

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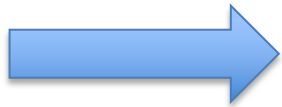
Flavor decomposition of DGLAP equation

$$\frac{d f_i}{d \ln \mu^2} = P_{qq}^v \otimes f_i + P_{q\bar{q}}^v \otimes f_{\bar{i}} + P_{qq}^s \otimes \sum_j f_j + P_{q\bar{q}}^s \otimes \sum_{\bar{j}} f_{\bar{j}} + P_{qg} \otimes f_g,$$

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$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes \sum_j (f_j + f_{\bar{j}}).$$

$$f_i^\pm \equiv f_i \pm f_{\bar{i}}$$



$$\frac{d f_i^+}{d \ln \mu^2} = (P_{qq}^v + P_{q\bar{q}}^v) \otimes f_i^+ + (P_{qq}^s + P_{q\bar{q}}^s) \otimes \sum_j f_j^+ + 2P_{qg} \otimes f_g,$$

$$\frac{d f_i^-}{d \ln \mu^2} = (P_{qq}^v - P_{q\bar{q}}^v) \otimes f_i^- + (P_{qq}^s - P_{q\bar{q}}^s) \otimes \sum_j f_j^-,$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes \sum_j f_j^+.$$

- f_i^+ and f_i^- don't mix
- f_i^+ mixes with their flavor sum, and gluon
- f_i^- only mixes with their flavor sum

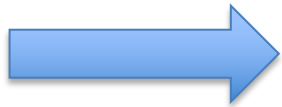
Flavor singlet and non-singlet

$$\frac{d f_i}{d \ln \mu^2} = P_{qq}^v \otimes f_i + P_{q\bar{q}}^v \otimes f_{\bar{i}} + P_{qq}^s \otimes \sum_j f_j + P_{q\bar{q}}^s \otimes \sum_{\bar{j}} f_{\bar{j}} + P_{qg} \otimes f_g,$$

$$\frac{d f_{\bar{i}}}{d \ln \mu^2} = P_{qq}^v \otimes f_{\bar{i}} + P_{q\bar{q}}^v \otimes f_i + P_{qq}^s \otimes \sum_{\bar{j}} f_{\bar{j}} + P_{q\bar{q}}^s \otimes \sum_j f_j + P_{qg} \otimes f_g,$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes \sum_j (f_j + f_{\bar{j}}).$$

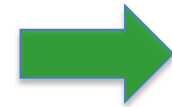
$$f_i^\pm \equiv f_i \pm f_{\bar{i}}$$



$$\frac{d f_i^+}{d \ln \mu^2} = (P_{qq}^v + P_{q\bar{q}}^v) \otimes f_i^+ + (P_{qq}^s + P_{q\bar{q}}^s) \otimes \sum_j f_j^+ + 2P_{qg} \otimes f_g,$$

$$\frac{d f_i^-}{d \ln \mu^2} = (P_{qq}^v - P_{q\bar{q}}^v) \otimes f_i^- + (P_{qq}^s - P_{q\bar{q}}^s) \otimes \sum_j f_j^-,$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes \sum_j f_j^+.$$



$$f_s \equiv \sum_j f_j^+$$

$$f_{ns} \equiv \sum_j f_j^-$$

Task (2mins): Work out the evolution equation of f_s and f_{ns} ,
assuming there are n_f active quark flavors.

Flavor singlet and non-singlet

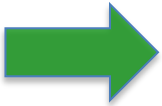
$$\frac{d f_i^+}{d \ln \mu^2} = (P_{qq}^v + P_{q\bar{q}}^v) \otimes f_i^+ + (P_{qq}^s + P_{q\bar{q}}^s) \otimes f_s + 2P_{qg} \otimes f_g,$$

$$\frac{d f_i^-}{d \ln \mu^2} = (P_{qq}^v - P_{q\bar{q}}^v) \otimes f_i^- + (P_{qq}^s - P_{q\bar{q}}^s) \otimes f_{ns},$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes f_s.$$

$$f_s \equiv \sum_j f_j^+$$

$$f_{ns} \equiv \sum_j f_j^-$$



3 equations

$$\frac{d f_{ns}}{d \ln \mu^2} = [P_{qq}^v - P_{q\bar{q}}^v + n_f (P_{qq}^s - P_{q\bar{q}}^s)] \otimes f_{ns}$$

$$\frac{d f_s}{d \ln \mu^2} = [P_{qq}^v + P_{q\bar{q}}^v + n_f (P_{qq}^s + P_{q\bar{q}}^s)] \otimes f_s + 2n_f P_{qg} \otimes f_g$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes f_s$$

 Evolves by itself

 Coupled together

Flavor singlet and non-singlet

$$\frac{d f_i^+}{d \ln \mu^2} = (P_{qq}^v + P_{q\bar{q}}^v) \otimes f_i^+ + (P_{qq}^s + P_{q\bar{q}}^s) \otimes f_s + 2P_{qg} \otimes f_g,$$

$$\frac{d f_i^-}{d \ln \mu^2} = (P_{qq}^v - P_{q\bar{q}}^v) \otimes f_i^- + (P_{qq}^s - P_{q\bar{q}}^s) \otimes f_{ns},$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes f_s.$$

$$f_s \equiv \sum_j f_j^+$$

$$f_{ns} \equiv \sum_j f_j^-$$

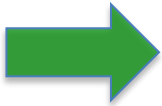
$$\frac{d f_{ns}}{d \ln \mu^2} = [P_{qq}^v - P_{q\bar{q}}^v + n_f (P_{qq}^s - P_{q\bar{q}}^s)] \otimes f_{ns}$$

$$\frac{d f_s}{d \ln \mu^2} = [P_{qq}^v + P_{q\bar{q}}^v + n_f (P_{qq}^s + P_{q\bar{q}}^s)] \otimes f_s + 2n_f P_{qg} \otimes f_g$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes f_s$$

→ Evolves by itself

} Coupled together


3 equations

Remaining $2(n_f - 1)$ equations are for relative differences among f_i^+ and f_i^-

E.g., $f_{ij}^\pm \equiv f_i^\pm - f_j^\pm$ evolve by themselves $\frac{d f_{ij}^\pm}{d \ln \mu^2} = (P_{qq}^v \pm P_{q\bar{q}}^v) \otimes f_{ij}^\pm$

Flavor singlet and non-singlet

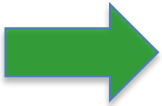
$$\frac{d f_i^+}{d \ln \mu^2} = (P_{qq}^v + P_{q\bar{q}}^v) \otimes f_i^+ + (P_{qq}^s + P_{q\bar{q}}^s) \otimes f_s + 2P_{qg} \otimes f_g,$$

$$\frac{d f_i^-}{d \ln \mu^2} = (P_{qq}^v - P_{q\bar{q}}^v) \otimes f_i^- + (P_{qq}^s - P_{q\bar{q}}^s) \otimes f_{ns},$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes f_s.$$

$$f_s \equiv \sum_j f_j^+$$

$$f_{ns} \equiv \sum_j f_j^-$$



3 equations

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$$\frac{d f_s}{d \ln \mu^2} = [P_{qq}^v + P_{q\bar{q}}^v + n_f (P_{qq}^s + P_{q\bar{q}}^s)] \otimes f_s + 2n_f P_{qg} \otimes f_g$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes f_s$$

 Evolves by itself

 Coupled together

Remaining $2(n_f - 1)$ equations are for relative differences among f_i^+ and f_i^-

Or, $\Delta f_i^+ = f_s - n_f f_i^+$
 $\Delta f_i^- = f_{ns} - n_f f_i^-$ evolve by themselves, too, by the same kernels

$$\frac{d \Delta f_i^\pm}{d \ln \mu^2} = (P_{qq}^v \pm P_{q\bar{q}}^v) \otimes \Delta f_i^\pm$$

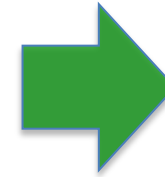
Flavor decomposed DGLAP equations

$$\frac{d f_{\text{ns}}}{d \ln \mu^2} = [P_{qq}^v - P_{q\bar{q}}^v + n_f (P_{qq}^s - P_{q\bar{q}}^s)] \otimes f_{\text{ns}}$$

$$\frac{d f_s}{d \ln \mu^2} = [P_{qq}^v + P_{q\bar{q}}^v + n_f (P_{qq}^s + P_{q\bar{q}}^s)] \otimes f_s + 2n_f P_{qg} \otimes f_g$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes f_s$$

3 equations



$$f_s \equiv \sum_j f_j^+$$

$$f_{\text{ns}} \equiv \sum_j f_j^-$$

$$f_g$$

$$\Delta f_i^+ = f_s - n_f f_i^+$$

$$\Delta f_i^- = f_{\text{ns}} - n_f f_i^-$$

$$\frac{d \Delta f_i^\pm}{d \ln \mu^2} = (P_{qq}^v \pm P_{q\bar{q}}^v) \otimes \Delta f_i^\pm \quad 2(n_f - 1) \text{ independent equations}$$



$$f_i^+ = \frac{1}{n_f} (f_s - \Delta f_i^+)$$

$$f_i^- = \frac{1}{n_f} (f_{\text{ns}} - \Delta f_i^-)$$

Proper matching at quark thresholds is needed as μ goes up.

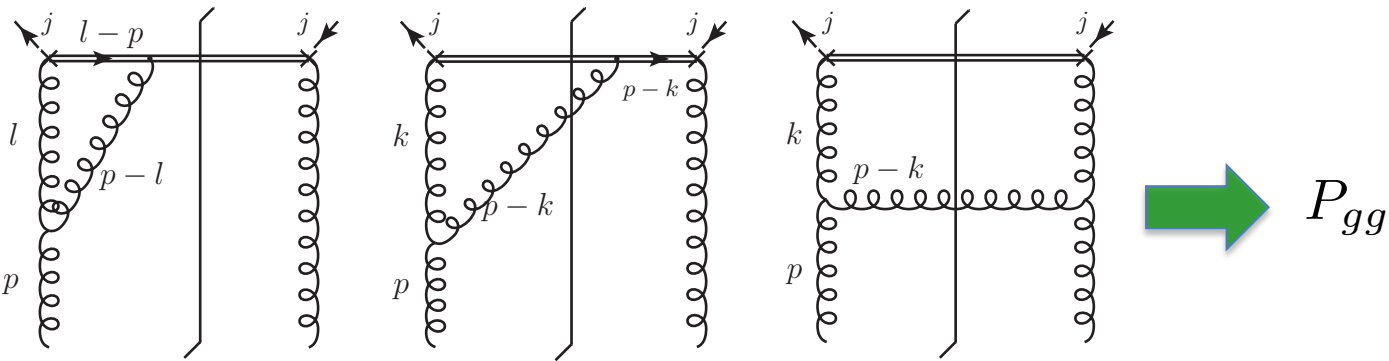
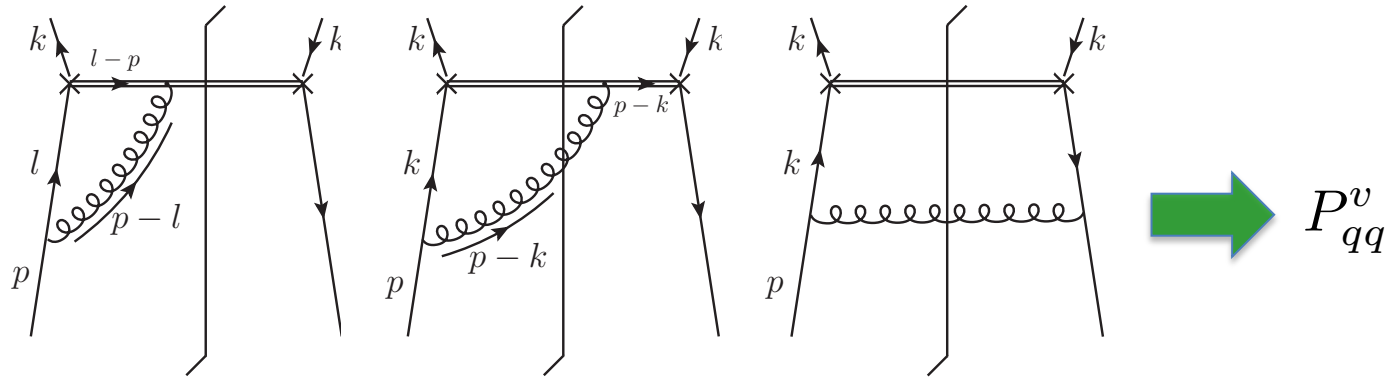
DGLAP kernels: Are they all nonzero and independent?

Question (1min): Which kernels are nonzero at LO? $\mathcal{O}(\alpha_s)$

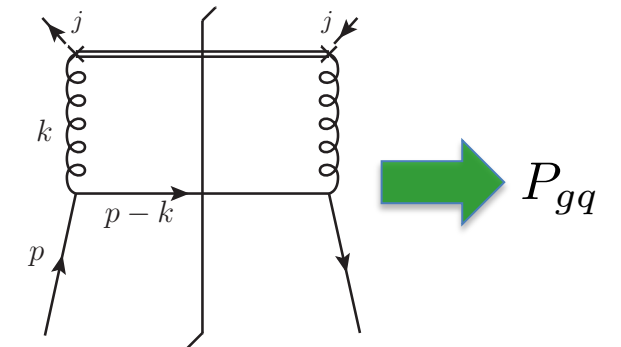
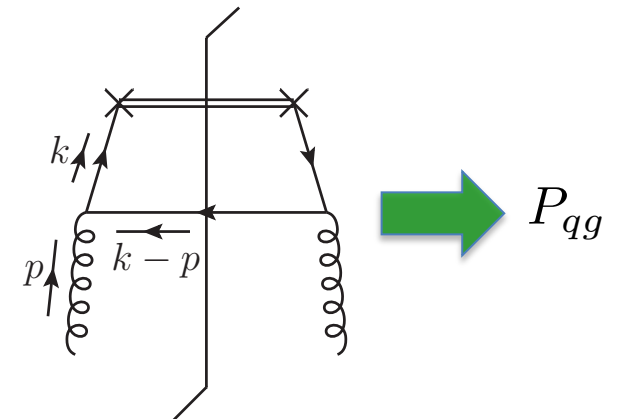
$$\begin{aligned}P_{q_i q_j} &= P_{\bar{q}_i \bar{q}_j} = P_{qq}^v \delta_{ij} + P_{qq}^s \\P_{q_i \bar{q}_j} &= P_{\bar{q}_i q_j} = P_{q\bar{q}}^v \delta_{ij} + P_{q\bar{q}}^s \\P_{q_i g} &= P_{\bar{q}_i g} = P_{qg} \\P_{g q_i} &= P_{g \bar{q}_i} = P_{gq} \quad P_{gg}\end{aligned}$$

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 P_{g q_i} &= P_{g \bar{q}_i} = P_{gq}
 \end{aligned}$$



But $P_{q\bar{q}}^v = P_{qq}^s = P_{q\bar{q}}^s = 0$ at LO

DGLAP kernels: Are they all nonzero and independent?

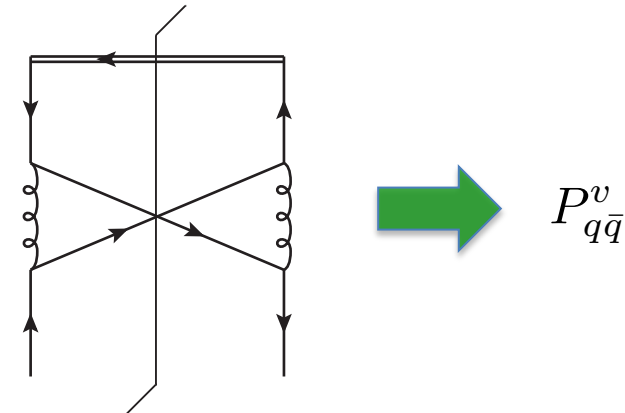
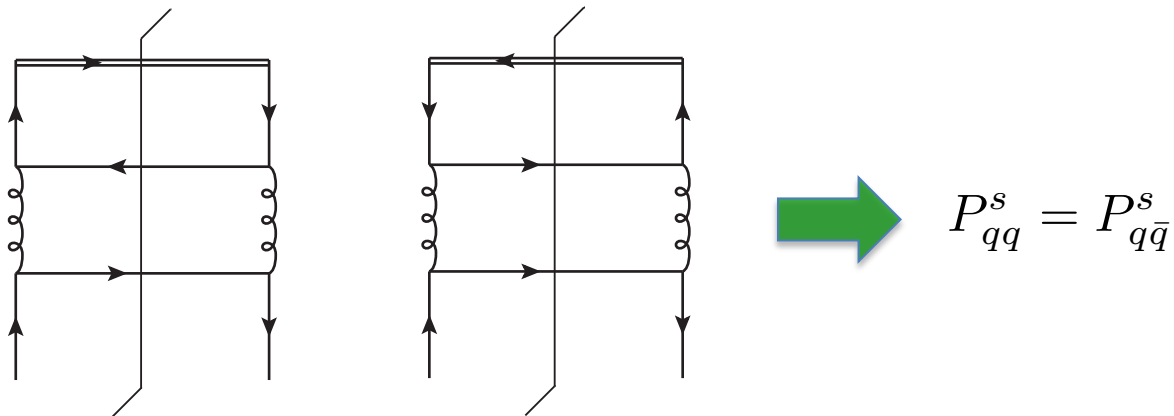
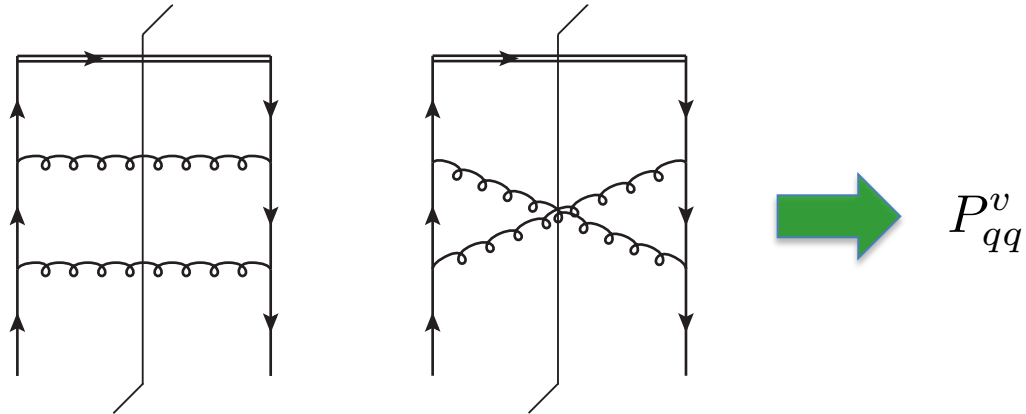
Question (3mins): Any *new* kernels occur at NLO? $\mathcal{O}(\alpha_s^2)$

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 P_{g q_i} &= P_{g \bar{q}_i} = P_{gq}
 \end{aligned}$$



- All the seven kernels are nonzero.
- But there is a degeneracy.

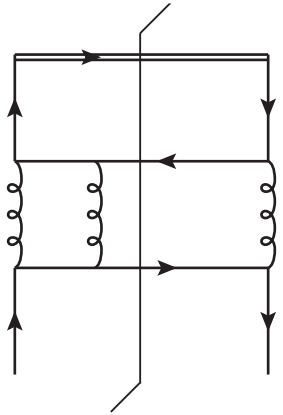
DGLAP kernels: Are they all nonzero and independent?

Question (3mins): Does the degeneracy retain at NNLO? $\mathcal{O}(\alpha_s^3)$

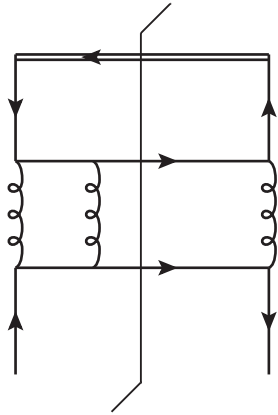
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$\rightarrow P_{qq}^s$



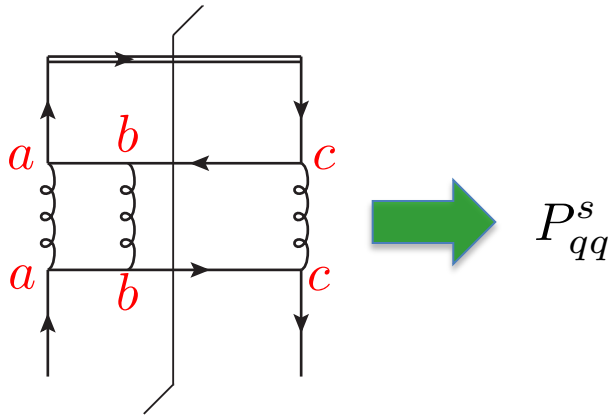
$\rightarrow P_{q\bar{q}}^s$

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 P_{q_i g} &= P_{\bar{q}_i g} = P_{qg} \\
 P_{g q_i} &= P_{g \bar{q}_i} = P_{gq} \\
 & \qquad \qquad \qquad P_{gg}
 \end{aligned}$$

DGLAP kernels: They **ARE** all nonzero and independent!

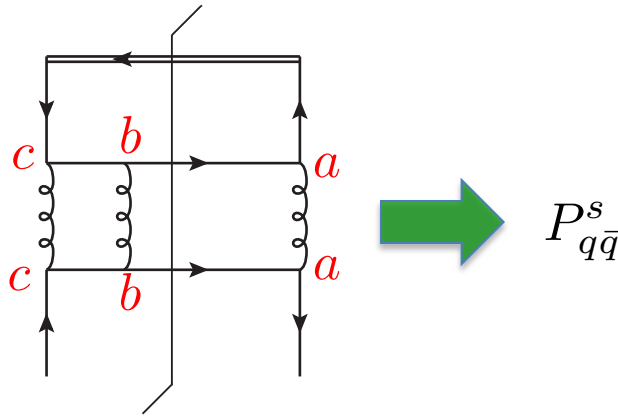
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 \end{aligned}$$



P_{qq}^s

$$\begin{aligned}
 &-\frac{1}{N_c} \text{tr}\{t^a t^b t^c\} \text{tr}\{t^c t^b t^a\} \\
 &= \frac{-1}{16N_c} (d_{abc}^2 + f_{abc}^2)
 \end{aligned}$$



$P_{q\bar{q}}^s$

$$\begin{aligned}
 &\frac{1}{N_c} \text{tr}\{t^a t^b t^c\} \text{tr}\{t^a t^b t^c\} \\
 &= \frac{1}{16N_c} (d_{abc}^2 - f_{abc}^2)
 \end{aligned}$$



$$P_{qq}^s \neq P_{q\bar{q}}^s$$

starting at NNLO, due to **Abelian quantum interference** effect!

What shall we expect?

At LO:

$$\frac{d f_i}{d \ln \mu^2} = P_{qq}^v \otimes f_i + P_{qg} \otimes f_g,$$

$$\frac{d f_{\bar{i}}}{d \ln \mu^2} = P_{qq}^v \otimes f_{\bar{i}} + P_{qg} \otimes f_g,$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes f_s.$$

$$P_{qg}(z) = \frac{\alpha_s}{2\pi} T_F [z^2 + (1-z)^2] > 0$$

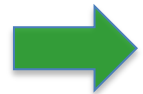
$$P_{gq}(z) = \frac{\alpha_s}{2\pi} C_F \left[\frac{1 + (1-z)^2}{z} \right] > 0$$

$$P_{gg}(z) = \frac{\alpha_s}{2\pi} \left[2C_A \left(\frac{z}{(1-z)_+} + z(1-z) + \frac{1-z}{z} \right) + \left(\frac{11}{6} C_A - \frac{2}{3} n_f T_F \right) \delta(1-z) \right]$$

$$P_{qq}^v(z) = \frac{\alpha_s}{2\pi} C_F \left[\frac{2}{(1-z)_+} - 1 - z + \frac{3}{2} \delta(1-z) \right]$$

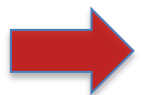
Property of the convolution:

$$(P \otimes f)(x) \equiv \int_x^1 \frac{dz}{z} P(x/z) f(z) = \int_0^1 dz_1 \int_0^1 dz_2 \delta(x - z_1 z_2) P(z_1) f(z_2),$$



Change in $f(x)$ is determined by the product $P(z_1) f(z_2)$ with $z_1, z_2 \in [x, 1]$

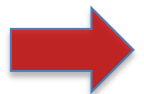
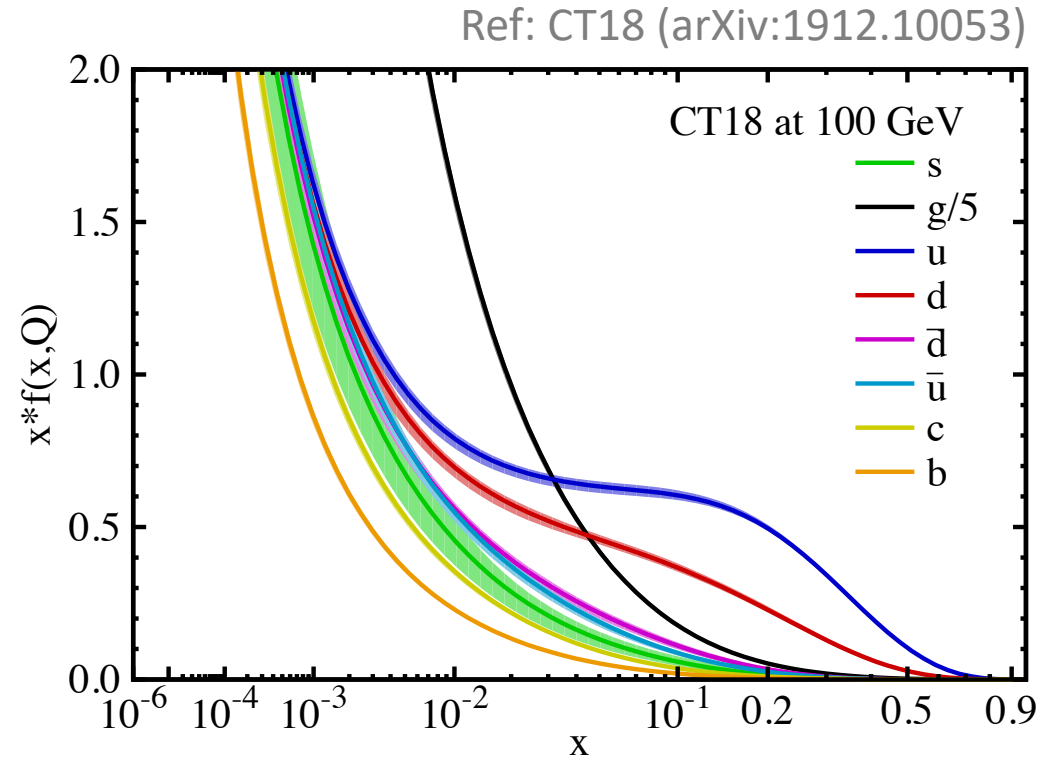
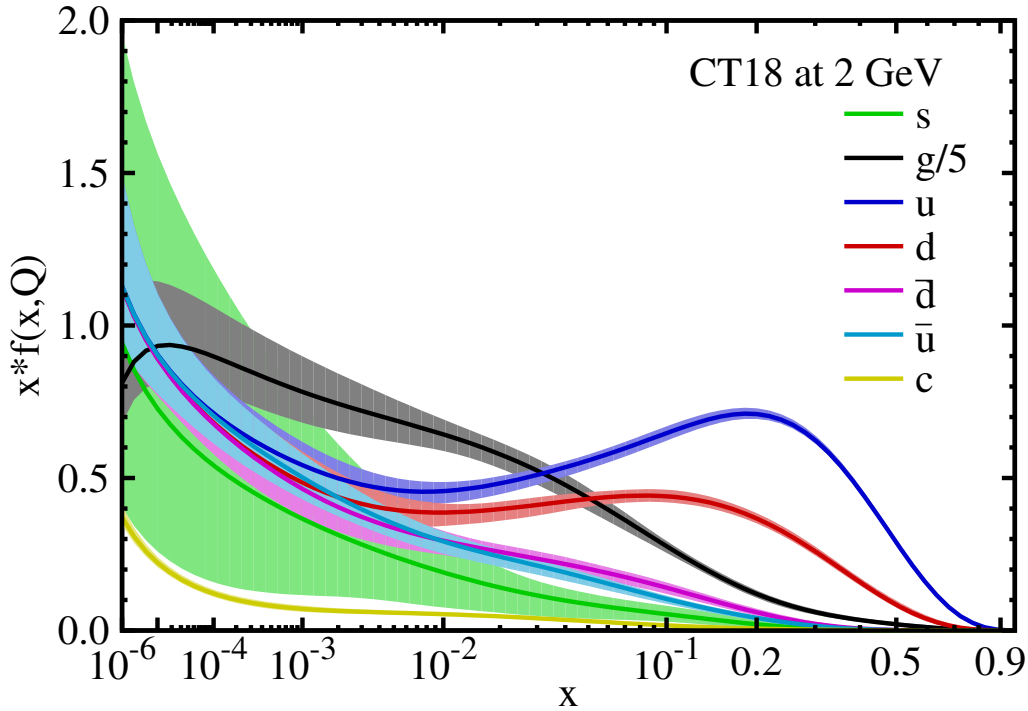
- **At large x :** evolution is dominated by singular pieces in P_{qq}^v and P_{gg} , which are **negative**.
- **At small x :** evolution is dominated by $1/z$ pieces from gluon evolution, which are **positive**.



Evolution generally brings PDFs from large x down to small x .

What shall we expect?

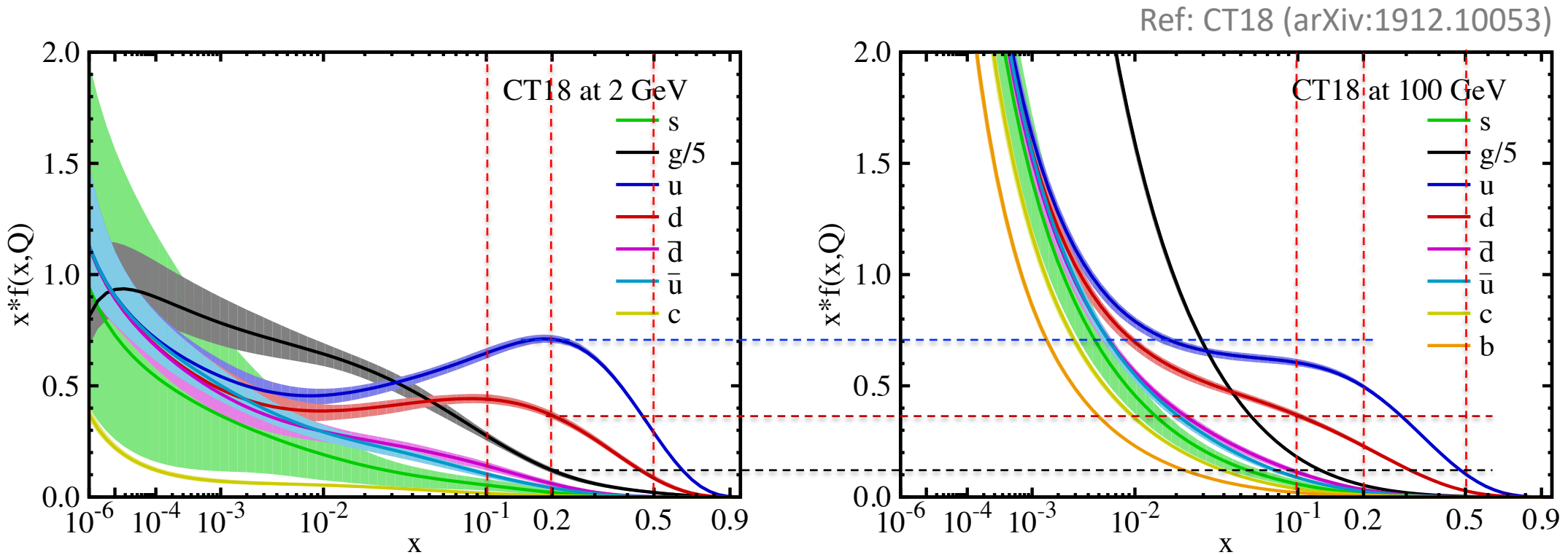
At LO:



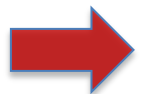
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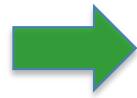
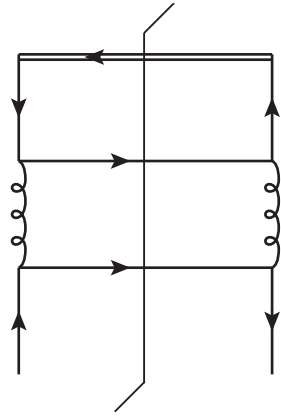
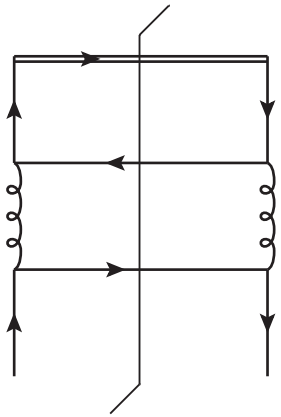
This overall property should **NOT** be altered by high-order corrections.



Evolution generally brings PDFs from large x down to small x .

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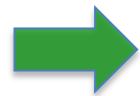
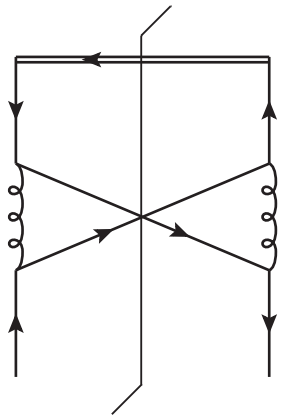
At NLO: 3 new kernels



$$P_{qq}^s = P_{q\bar{q}}^s > 0$$



u has equal probability to split into d/\bar{d} , s/\bar{s} , c/\bar{c} , etc.



$$P_{q\bar{q}}^v \propto \frac{1}{N_c} \text{tr}\{t^a t^b t^a t^b\} = C_F^2 - \frac{1}{2} C_F C_A = -\frac{2}{9} < 0$$

negative due to color factor



$u \rightarrow \bar{u}$ is **suppressed** as compared to $u \rightarrow \bar{d}$, \bar{s} , etc.



Flavor difference between antiquarks.

What shall we expect?

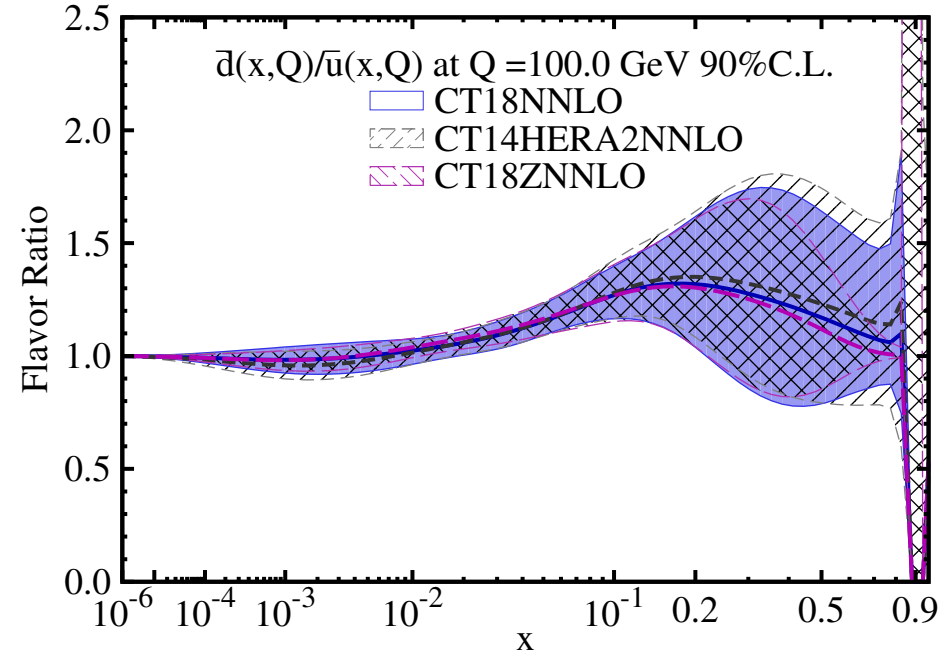
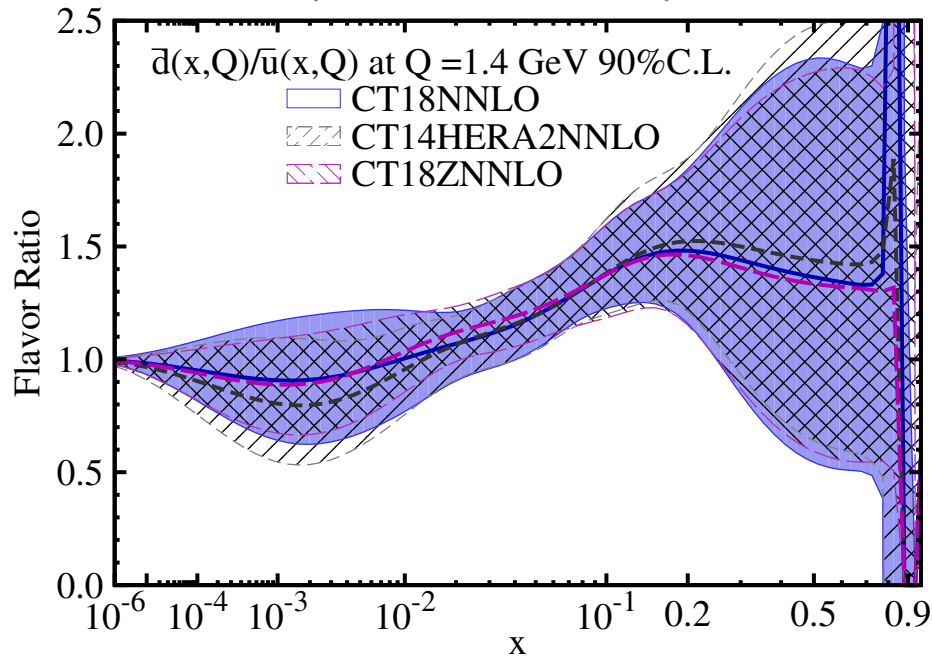
At NLO: Antiquark flavor asymmetry

$$\frac{d}{d \ln \mu^2} (u - d) = P_{qq}^v \otimes (u - d) + P_{q\bar{q}}^v \otimes (\bar{u} - \bar{d}),$$

$$\frac{d}{d \ln \mu^2} (\bar{u} - \bar{d}) = P_{qq}^v \otimes (\bar{u} - \bar{d}) + \boxed{P_{q\bar{q}}^v \otimes (u - d)}$$

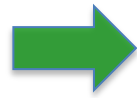
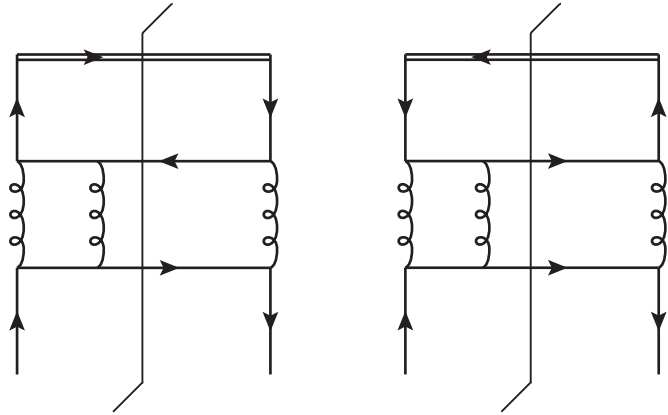
➔ \bar{d} tends to $> \bar{u}$ at large x .

Ref: CT18 (arXiv:1912.10053)



What shall we expect?

At NNLO: Sea quark-antiquark asymmetry



$$P_{qq}^s \neq P_{q\bar{q}}^s, \quad \text{but} \quad \int_0^1 dz \left[P_{qq}^s(z) - P_{q\bar{q}}^s(z) \right] = 0$$



Equal splitting between $u \rightarrow s$ and $u \rightarrow \bar{s}$ is broken

$$\frac{d}{d \ln \mu^2} (s - \bar{s}) = (P_{qq}^v - P_{q\bar{q}}^v) \otimes (s - \bar{s}) + (P_{qq}^s - P_{q\bar{q}}^s) \otimes f_{ns}$$

→ $s(x) \neq \bar{s}(x)$ generally, although $\int_0^1 dx (s - \bar{s})(x) = 0$

→ Similarly, $c(x) \neq \bar{c}(x)$ and $b(x) \neq \bar{b}(x)$

S. Catani, et al., PRL 93, 152003 (2004)

So, all “naïve flavor symmetries” of PDFs are broken

Since they are broken perturbatively, no reason they are reserved at their initial scale μ_0 .

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➤ How to separate u from d ?

DIS $F_{1,2}$ structure functions $\propto e_u^2 (u + \bar{u}) + e_d^2 (d + \bar{d}) + \dots$ (in proton target by default)

With a neutron target (from deuteron) $\rightarrow e_u^2 (d + \bar{d}) + e_d^2 (u + \bar{u}) + \dots$ (isospin symmetry)

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Drell-Yan production cross section, charge asymmetry, A_{FB} , etc. $q(x_1) \bar{q}'(x_2) \pm \bar{q}'(x_1) q(x_2)$

DIS F_3 structure function $\propto e_u^2 (u - \bar{u}) + e_d^2 (d - \bar{d}) + \dots$

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➤ How to constrain gluon? [1min]

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DIS $F_{1,2}$ structure functions $\propto e_u^2 (u + \bar{u}) + e_d^2 (d + \bar{d}) + \dots$ (in proton target by default)

With a neutron target (from deuteron) $\rightarrow e_u^2 (d + \bar{d}) + e_d^2 (u + \bar{u}) + \dots$ (isospin symmetry)

➤ How to separate \bar{q} from q ?

Drell-Yan production cross section, charge asymmetry, A_{FB} , etc. $q(x_1) \bar{q}'(x_2) \pm \bar{q}'(x_1) q(x_2)$

DIS F_3 structure function $\propto e_u^2 (u - \bar{u}) + e_d^2 (d - \bar{d}) + \dots$

➤ How to constrain gluon?

$u, d, \bar{u}, \bar{d}, g$

Jet production at hadron colliders $q + g \rightarrow q + g, \dots$

So, all “naïve flavor symmetries” of PDFs are broken

Since they are broken perturbatively, no reason they are reserved at their initial scale μ_0 .

But in an actual fit, one needs data to constrain and separate different flavors!

➤ **How to probe strange quark? (2mins)**

So, all “naïve flavor symmetries” of PDFs are broken

Since they are broken perturbatively, no reason they are reserved at their initial scale μ_0 .

But in an actual fit, one needs data to constrain and separate different flavors!

➤ How to probe strange quark?

- It's all about precision ...
- By flavor tagging, e.g., dimuon data from CCFR NuTeV SIDIS

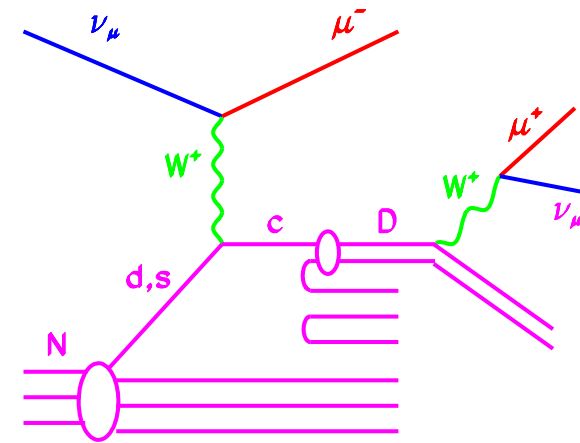
$$\nu_\mu + s \rightarrow \mu^- + c (\rightarrow \mu^+ + \nu_\mu) + X$$

$$\bar{\nu}_\mu + \bar{s} \rightarrow \mu^+ + \bar{c} (\rightarrow \mu^- + \bar{\nu}_\mu) + X$$

- $W + c$ production with charm jet tagging at hadron collider

$$s + g \rightarrow W^- + c, \quad \bar{s} + g \rightarrow W^+ + \bar{c}$$

- Simultaneous fit of PDF and fragmentation function, by tagging kaon production JAM, PRD 101, 074020 (2020)
- Lattice QCD constraint ...



CCFR, PRD 64 (2001), 112006

Summary

➤ DGLAP kernels and flavor decomposition

$$\frac{d f_{\text{ns}}}{d \ln \mu^2} = [P_{qq}^v - P_{q\bar{q}}^v + n_f (P_{qq}^s - P_{q\bar{q}}^s)] \otimes f_{\text{ns}}$$

$$\frac{d f_s}{d \ln \mu^2} = [P_{qq}^v + P_{q\bar{q}}^v + n_f (P_{qq}^s + P_{q\bar{q}}^s)] \otimes f_s + 2n_f P_{qg} \otimes f_g$$

$$\frac{d f_g}{d \ln \mu^2} = P_{gg} \otimes f_g + P_{gq} \otimes f_s$$

$$\frac{d \Delta f_i^\pm}{d \ln \mu^2} = (P_{qq}^v \pm P_{q\bar{q}}^v) \otimes \Delta f_i^\pm$$

$$P_{q_i q_j} = P_{\bar{q}_i \bar{q}_j} = P_{qq}^v \delta_{ij} + P_{qq}^s$$

$$P_{q_i \bar{q}_j} = P_{\bar{q}_i q_j} = P_{q\bar{q}}^v \delta_{ij} + P_{q\bar{q}}^s$$

$$P_{q_i g} = P_{\bar{q}_i g} = P_{qg} \quad P_{gg}$$

$$P_{gq_i} = P_{g\bar{q}_i} = P_{gq}$$

➤ This is a great example of

- How high-order corrections bring qualitative impacts;
- How perturbation theory guides nonperturbative “intuition” or modeling.

Further high order corrections only give quantitative changes.

➤ How this plays a role in actual global analysis