

Spectator tagging with polarized deuteron

Wim Cosyn

JLab tensor SIDIS workshop

in collaboration with Ch. Weiss

1906.11119, 2006.03033, 2603.23699, 2603.23700

JLab LDRD project on EIC spectator tagging

1409.5768, 1601.06665



Supported by



Deuteron tensor polarized observables

- Many structures appear in cross section etc.
→ Challenging to disentangle
- Direction of polarization relative to scattering plays a role

$$\sigma = \sigma_U(1 + P A^V(N_{\text{pol}}) + Q A^T(N_{\text{pol}}))$$

→ Effective polarizations

[See also 2410.12764]

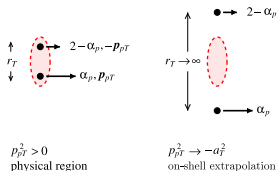
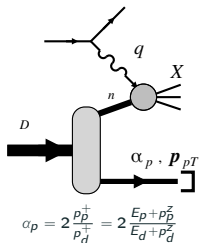
- Tensor polarization needs deuteron D -wave ($L = 2$)

$$\sim \left(2f_0(k) + \frac{f_2(k)}{\sqrt{2}} \right) \frac{f_2(k)}{\sqrt{2}}$$

→ **Small** when averaged over all internal momenta

- In inclusive measurement nuclear correction is an important systematic uncertainty

Deuteron Spectator tagging



- Detection of nucleon in target fragmentation region: **“spectator”**
- Interaction w photon can be elastic, DIS, hard exclusive, SIDIS, etc.
- **Control** over your initial nuclear state
→ Active nucleon identified, create effective targets
- Spectator can **reinteract** with other final-state hadrons
→ “Simple” for deuteron
- Spectator kinematics gives handle on initial state
→ **On-shell extrapolation** for free nucleon
[Sargsian, Strikman, PLB05]
→ Larger momenta for medium modifications
- Measurements with fixed target [BONuS, Deeps, LAD, BAND, Alert,..] and collider [EIC] with large acceptance far forward detectors

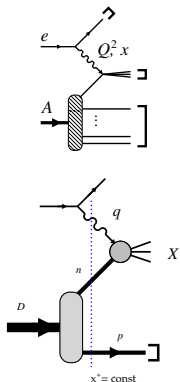
- General expression of SIDIS for a polarized spin 1 target
 - ▶ Tagged spectator DIS is SIDIS in the target fragmentation region

$$\vec{e} + \vec{T} \rightarrow e' + X + h$$

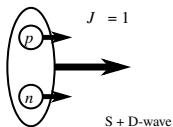
- ▶ 41 (18 + 23) structure functions
- Light-front structure of the deuteron
 - ▶ Natural for high-energy reactions as **off-shellness of nucleons** in LF quantization remains **finite**
- Dynamical model to express structure functions of the reaction
 - ▶ First step: impulse approximation (IA) model
 - ▶ Results for all SF: deuteron LF distributions \times nucleon collinear pdfs
 - ▶ FSI corrections

High-energy scattering with nuclei [Frankfurt, Strikman 80s+]

- Interplay of two scales: **high-energy scattering** and **low-energy nuclear structure**
- Virtual photon probes nucleus at fixed lightcone time $x^+ = x^0 + x^3$
- Scales can be separated using methods of light-front quantization and QCD factorization
- Tools for high-energy scattering known from ep
- Nuclear input: light-front momentum densities, spectral functions, overlaps with specific final states in breakup/tagging reactions
 - ▶ framework known for deuteron, can be extended to ^3He but challenges
 - ▶ still **low-energy** nuclear physics, just formulated differently



Deuteron light-front wave function



- Up to momenta of a few 100 MeV: dominated by NN
- Can be evaluated in LFQM [Berestetsky, Frankfurt, Strikman, Terentev]
- Overlap with **on-shell** free two-nucleon state

- Schrödinger (non-rel) like eq. for the wf components, rotational invariance recovered

$$\Psi_{\lambda}(\mathbf{k}, \lambda_p, \lambda_n) = \sqrt{E_k} \sum_{\lambda'_p \lambda'_n} \mathcal{D}_{\lambda_p \lambda'_p}^{\frac{1}{2}}[R_{fc}(k_1^{\mu}/m)] \mathcal{D}_{\lambda_n \lambda'_n}^{\frac{1}{2}}[R_{fc}(k_2^{\mu}/m)] \Phi_{\lambda}(\mathbf{k}, \lambda'_p, \lambda'_n)$$

- Differences with non-rel wave function:

- ▶ **Melosh rotations** to account for light-front spin
- ▶ \mathbf{k} is the rel. 3-momentum in the rest frame of the on-shell NN state



- Allows for the definition of nucleon LF momentum distributions \sim nucleon TMDs in deuteron

Unpolarized	$P_{[U]}(\alpha_p, \mathbf{p}_{pT} U, S_d, T_d)$	γ^+	deut: unpol, tensor
L polarized	$P_{[S_L]}(\alpha_p, \mathbf{p}_{pT} (U, S_d, T_d)$	$\gamma^+ \gamma_5$	deut: vector
T polarized	$P_{[S_T]}(\alpha_p, \mathbf{p}_{pT} U, S_d, T_d)$	$i\sigma^{i+} \gamma_5$	deut: vector

Nucleon LF momentum distributions

- Unpolarized, helicity independent

$$P_{[U,U]} = \frac{f_0^2(k) + f_2^2(k)}{2 + \alpha_p}$$

- Tensor polarized, helicity independent: needs D-wave, $\sim Y_{2m}(\theta, 0)$

$$P_{[T_{LL},U]}(\alpha_p, \mathbf{p}_{pT}) = -\frac{1}{2 + \alpha_p} \left(2f_0 + \frac{f_2}{\sqrt{2}} \right) \frac{f_2}{\sqrt{2}} \frac{3}{2} (3 \cos^2 \theta_k - 1),$$

$$P_{[T_{LT},U]}(\alpha_p, \mathbf{p}_{pT}) = -\frac{1}{2 + \alpha_p} \left(2f_0 + \frac{f_2}{\sqrt{2}} \right) \frac{f_2}{\sqrt{2}} 6 \cos \theta_k \sin \theta_k,$$

$$P_{[T_{TT},U]}(\alpha_p, \mathbf{p}_{pT}) = -\frac{1}{2 + \alpha_p} \left(2f_0 + \frac{f_2}{\sqrt{2}} \right) \frac{f_2}{\sqrt{2}} 3 \sin^2 \theta_k,$$

- L/T nucleon polarized

$$P_{[S_L,S_L]}(\alpha_p, \mathbf{p}_{pT}) = \frac{1}{2 + \alpha_p} \left(f_0 - \frac{f_2}{\sqrt{2}} \right) \left[A_{|\neq|}(\mathbf{k}) \left(f_0 - \frac{f_2}{\sqrt{2}} \right) + B_{|\neq|}(\mathbf{k}) \left(f_0 + \sqrt{2}f_2 \right) \right]$$

$$P_{[S_T,S_L]}(\alpha_p, \mathbf{p}_{pT}) = \frac{1}{2 + \alpha_p} \left(f_0 - \frac{f_2}{\sqrt{2}} \right) \left[A_{|\neq|}(\mathbf{k}) \left(f_0 - \frac{f_2}{\sqrt{2}} \right) - B_{|\neq|}(\mathbf{k}) \left(f_0 + \sqrt{2}f_2 \right) \right]$$

$$P_{[S_L,S_T]}(\alpha_p, \mathbf{p}_{pT}) = \frac{1}{2 + \alpha_p} \left(f_0 - \frac{f_2}{\sqrt{2}} \right) \left[B_{|\neq|}(\mathbf{k}) \left(f_0 - \frac{f_2}{\sqrt{2}} \right) - A_{|\neq|}(\mathbf{k}) \left(f_0 + \sqrt{2}f_2 \right) \right]$$

$$P_{[S_T,S_T]}^{\parallel}(\alpha_p, \mathbf{p}_{pT}) = \frac{1}{2 + \alpha_p} \left(f_0 - \frac{f_2}{\sqrt{2}} \right) \left[B_{|\neq|}(\mathbf{k}) \left(f_0 - \frac{f_2}{\sqrt{2}} \right) + A_{|\neq|}(\mathbf{k}) \left(f_0 + \sqrt{2}f_2 \right) \right]$$

- Remarkable symmetry between helicity and transversity
→ Rotational invariance + relativistic spin effects

Nucleon LF momentum distributions: Sum Rules

■ baryon

$$\int \frac{d\alpha_p}{\alpha_p} d^2 p_{pT} P_{[U,U]}(\alpha_p, \mathbf{p}_{pT}) = 1,$$

$$\int \frac{d\alpha_p}{\alpha_p} d^2 p_{pT} P_{[T,U]}(\alpha_p, \mathbf{p}_{pT}) = 0,$$

■ momentum

$$\int \frac{d\alpha_p}{\alpha_p} d^2 p_{pT} (2 - \alpha_p) P_{[U,U]}(\alpha_p, \mathbf{p}_{pT}) = 1,$$

$$\int \frac{d\alpha_p}{\alpha_p} d^2 p_{pT} (2 - \alpha_p) P_{[T,U]}(\alpha_p, \mathbf{p}_{pT}) = 0$$

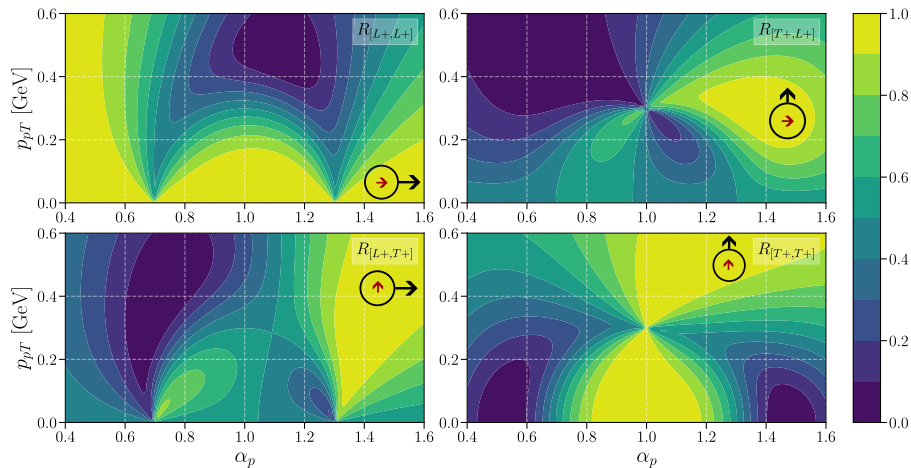
■ axial

$$\int \frac{d\alpha_p}{\alpha_p} d^2 p_{pT} P_{[S_L, S_L]}(\alpha_p, \mathbf{p}_{pT}) = S_d^z \frac{g_{Ad}}{2g_A}, \quad \int \frac{d\alpha_p}{\alpha_p} d^2 p_{pT} P_{[S_L, S_T]}(\alpha_p, \mathbf{p}_{pT}) = \epsilon_{S_L},$$

$$\int \frac{d\alpha_p}{\alpha_p} d^2 p_{pT} P_{[S_T, S_L]}^{\parallel}(\alpha_p, \mathbf{p}_{pT}) = \epsilon_{S_T} \approx -\epsilon_{S_L}, \quad \int \frac{d\alpha_p}{\alpha_p} d^2 p_{pT} P_{[S_T, S_T]}(\alpha_p, \mathbf{p}_{pT}) = S_d^T \frac{g_{Ad}}{2g_A},$$

$$1 - \frac{3}{2} \omega_2 = \frac{g_{Ad}}{2g_A}. \quad \rightarrow \text{cfr correction in inclusive polarized ed DIS}$$

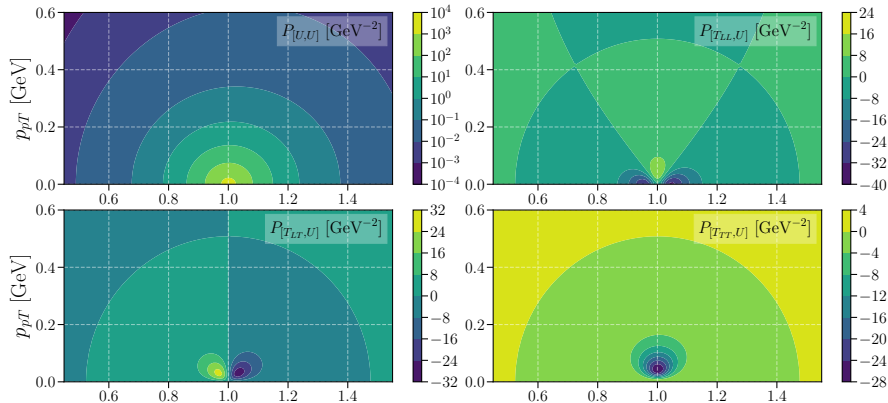
Probabilistic Distributions for pure states



Probabilistic Distributions for pure states

- At small momenta S -wave dominates
→ polarization of nucleon is transferred from deuteron one
- Depolarization at larger momenta
→ D -wave
- Limits at $p_{dT} = 0$
→ on-shell extrapolation
- Relativistic spin effects
→ sideways polarization
- Positivity remains for any **mixed deuteron state**

Tensor polarized distributions

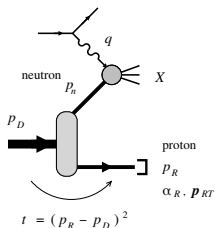


$$P_{[LL,U]}(\alpha_p, \mathbf{p}_{pT}) = -\frac{1}{2 + \alpha_p} \left(2f_0 + \frac{f_2}{\sqrt{2}} \right) \frac{f_2}{\sqrt{2}} \frac{3}{2} (3 \cos^2 \theta_k - 1),$$

$$P_{[TL,U]}(\alpha_p, \mathbf{p}_{pT}) = -\frac{1}{2 + \alpha_p} \left(2f_0 + \frac{f_2}{\sqrt{2}} \right) \frac{f_2}{\sqrt{2}} 6 \cos \theta_k \sin \theta_k,$$

$$P_{[TT,U]}(\alpha_p, \mathbf{p}_{pT}) = -\frac{1}{2 + \alpha_p} \left(2f_0 + \frac{f_2}{\sqrt{2}} \right) \frac{f_2}{\sqrt{2}} 3 \sin^2 \theta_k,$$

Tagged DIS with deuteron: model for the IA



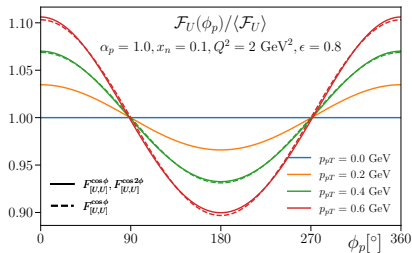
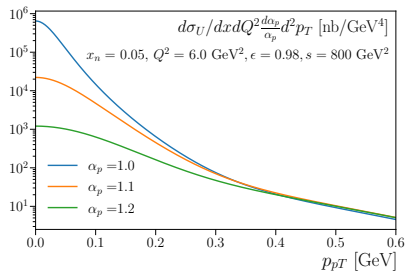
- Hadronic tensor can be written as a product of nucleon hadronic tensor with deuteron light-front densities

All SF can be written as

$$F_{ij}^k = \{\text{kin. factors}\} \times \{F_{1,2}(\tilde{x}, Q^2) \text{ or } g_{1,2}(\tilde{x}, Q^2)\} \times \{P_{[U,U]}, P_{[T,U]} \text{ or } P_{[S,S]}\}$$

- In the IA the following structure functions are **zero** \rightarrow sensitive to FSI
 - ▶ beam spin asymmetry [$F_{LU}^{\sin \phi_h}$]
 - ▶ target vector polarized single-spin asymmetry [8 SFs]
 - ▶ target tensor polarized double-spin asymmetry [7 SFs]
- Applied in studies of F_2 and tensor and vector polarized asymmetries [1906.11119, 2006.03033, 2108.08314, 2603.23699, 2603.23700, etc.]

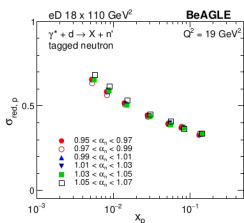
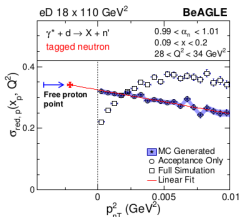
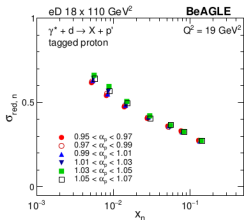
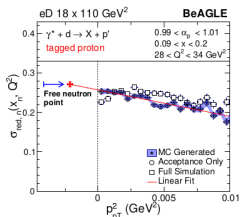
Unpolarized observables



- Sizable rates
- Drops fast with α_p , p_{pT}
- ϕ -modulations are power suppressed (p_{pT}/Q)
- Verification of reaction dynamics
- Pole extrapolation: unpolarized neutron structure

On-shell extrapolation of F_{2N}

$$F_{2d} = [2(2\pi)^3] S_d(\alpha_p, \mathbf{p}_{pT}) [\text{unpol}] F_{2n}(\tilde{x}, Q^2)$$

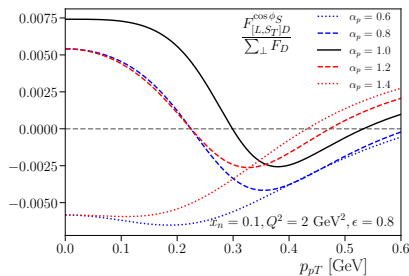
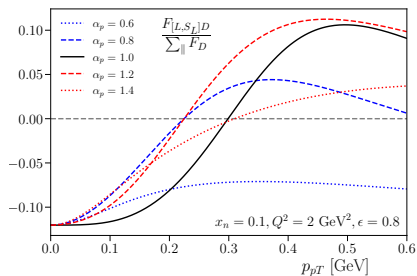


Detailed simulations for EIC [Jentsch, Tu, Weiss, PRC 21]

Vector pol double asymmetries (ϕ_p -integrated)

■ $A^V = D_{[S_L]} \frac{F_{[LS_L]D}}{\Sigma F} \sim \frac{g_{1n}}{F_{1n}}$

$A_{\perp}^V = D_{\perp[S_L]} \frac{F_{[LS_L]D}}{\Sigma_{\perp} F} + D_{\perp[S_T]} \frac{F_{[LS_T]S}^{\cos \phi_S}}{\Sigma_{\perp} F}$



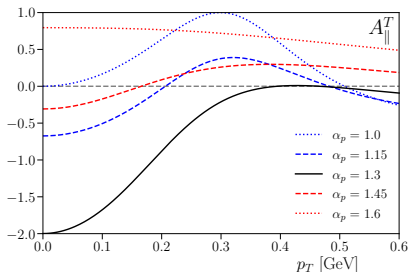
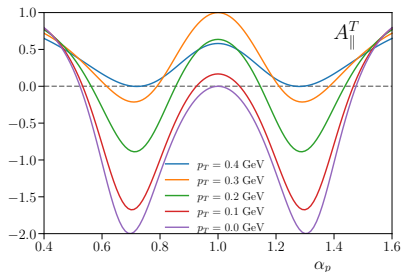
- Pole extrapolation of polarized neutron structure g_{1n}, g_{2n}
- At larger momenta: polarized (tagged) EMC effect

Tagged Tensor asymmetry

- No more averaging \rightarrow nucleon (unpol.) DIS part cancels!

$$A_{\parallel}^T = -2A_{\perp}^T \stackrel{LP}{=} \frac{\left(2f_0 + \frac{1}{\sqrt{2}}f_2\right) \frac{1}{\sqrt{2}}f_2}{f_0^2 + f_2^2} (1 - 3\cos^2\theta_k).$$

- Maximal $A^T = \{-2, 1\}$ can be reached at D/S ratio $f_2/f_0 = \{\sqrt{2}, -\frac{1}{\sqrt{2}}\}$
- Constrains deuteron wf
- Needs FSI corrections (can be sizeable)

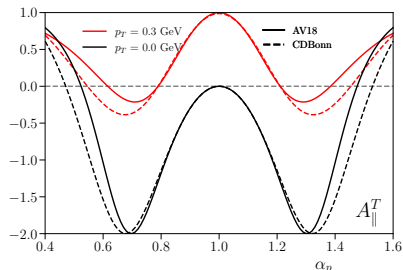
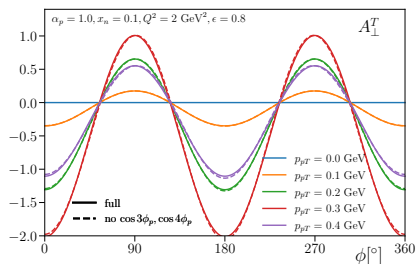


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Outlook & Conclusions

- Tagged reactions on deuteron have unique capabilities
 - control over nuclear effects
- Light-front is natural picture for high-energy nuclear scattering
- Deuteron structure can be quantified using polarized LF distributions
 - “our” TMDs (of nucleons in deuteron)
- Interesting relativistic spin effects
- Creates phi-dependent structures, spin-orbit effects etc.
- In proper reaction theory FSI need to be accounted for.
 - Distorted momentum distributions
- Extensions: tagged SIDIS/exclusive, tagging beyond proton, $A > 2$ etc.
- Lots of interesting physics opportunities and work to be done...