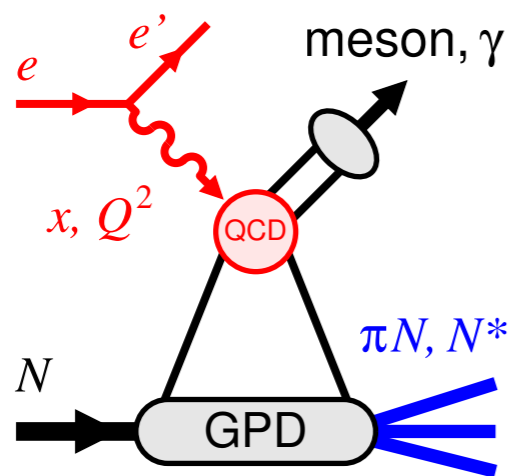


Transition GPDs and mechanical properties of excited baryons

C. Weiss (JLab), Center for Nuclear Femtography Workshop, JLab, 26-27 Mar 2026



Emerging physics program: 2025 White Paper

S. Diehl et al., EPJA 61, 131 (2025) [\[INSPIRE\]](#)

Recent developments in theory

J-Y. Kim, C. Weiss, PLB 870 (2025) 139924 [\[INSPIRE\]](#)

S. Son, K. Semenov-Tian-Shansky, JHEP 01, 119 (2025) [\[INSPIRE\]](#)

J-Y. Kim et al. PRD 111, 114010 (2025) [\[INSPIRE\]](#)

First results from JLab12 experiments

S. Diehl et al. PRL 131 (2023) 021901 [\[INSPIRE\]](#)

Transition GPDs

Exclusive processes $N \rightarrow \pi N, N^*$

Resonance transition GPDs

Excited baryon structure

Transition light-front densities

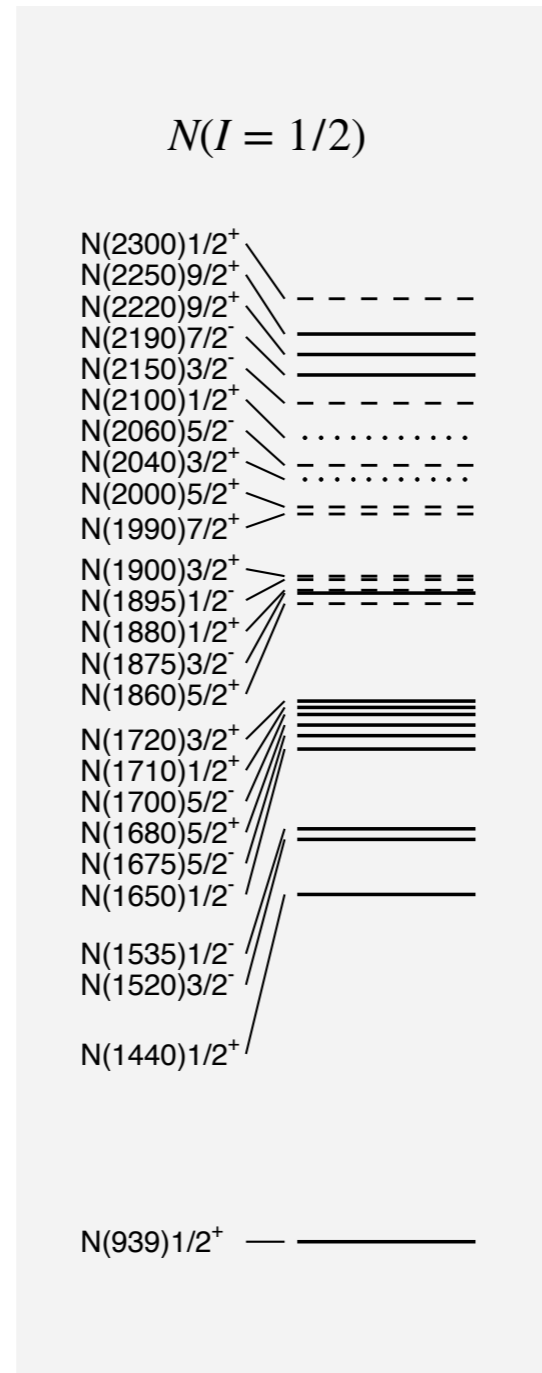
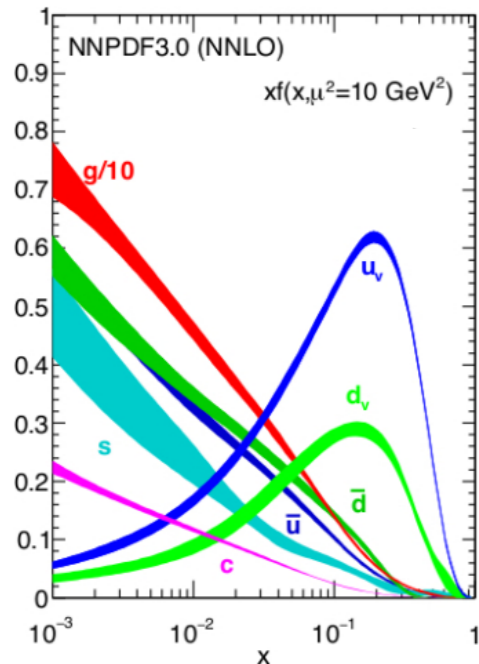
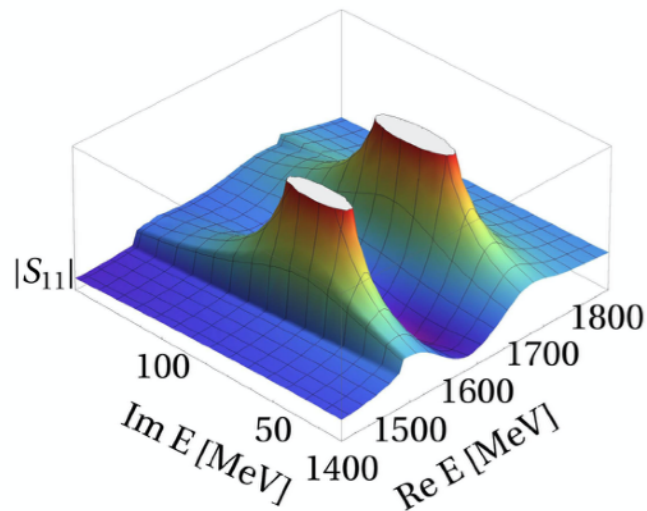
Transition angular momentum: $N \rightarrow \Delta$

Transition PDFs: $N \rightarrow \Delta$ tensor polarized PDF

[Processes and experiments]

$N \rightarrow \Delta$ in pion production and DVCS at JLab12

Baryon transitions with EIC far-forward detectors



[Image credits: NNPDF 3.0, PDG 2016, Roenchen - FZ Jülich]

Structure of ground-state nucleon

High-momentum-transfer processes:
Short-distance probe, “microscope”

Quark/gluon distributions 1D → 3D

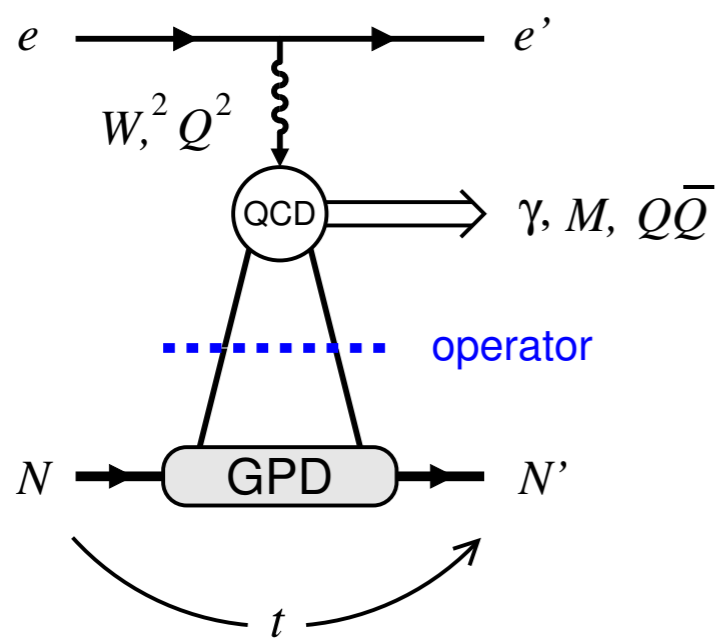
Structure of excited baryons?

Rich spectrum of excited baryons N^* , Δ :
Resonances, unstable particles

Fundamentals of parton picture:
Structure ↔ excitation spectrum

Relevant for dense matter, astrophysics,
 νN interactions

Need short-distance probe suitable for
baryon resonances!



Process with $Q^2, W^2 \gg \mu_{\text{had}}^2 \sim 1 \text{ GeV}^2, \quad |t| \sim \mu_{\text{had}}^2$

Scattering takes place on single quark/gluon in nucleon

Amplitude expressed as matrix element of QCD operator between incoming/outgoing nucleon states

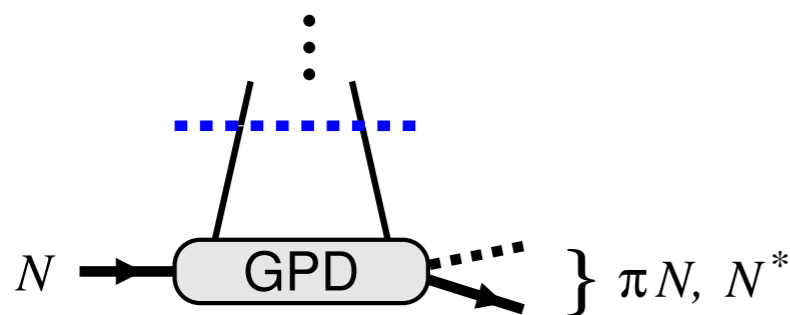
$$\langle N | \bar{\psi}(z) \dots \psi(0) | N \rangle \leftrightarrow \text{nucleon GPDs}$$

Transition GPDs

Same factorization works for scattering processes with $N \rightarrow \pi N, N^*$ transitions

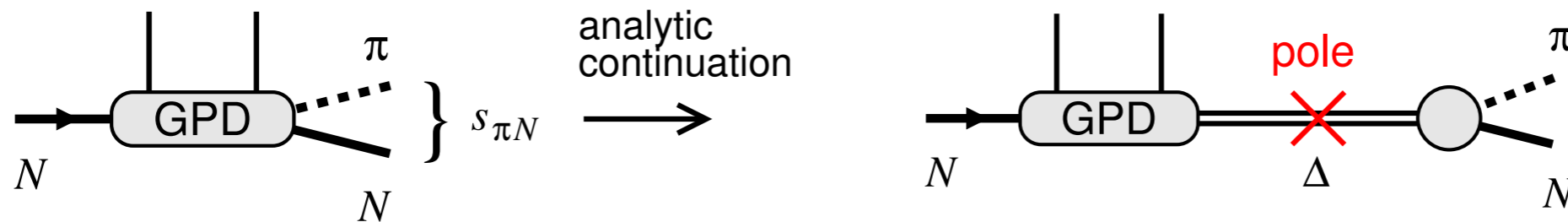
$$\langle \pi N | \bar{\psi}(z) \dots \psi(0) | N \rangle \leftrightarrow \text{transition GPDs}$$

N^*



Resonance excitation with defined QCD operator, rich set of quantum numbers

Probes quark/gluon structure of N^*



$$\langle \pi N | \mathcal{O} | N \rangle = \frac{\langle \pi N | \Delta \rangle \langle \Delta | \mathcal{O} | N \rangle}{s_{\pi N} - M_{\Delta}^2} + \text{less singular}$$

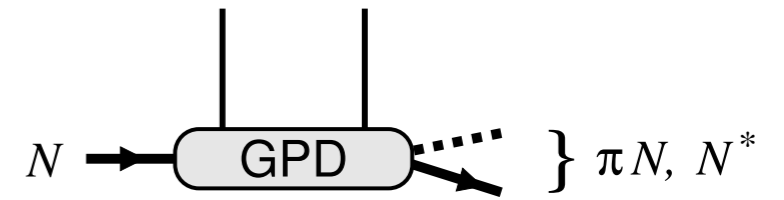
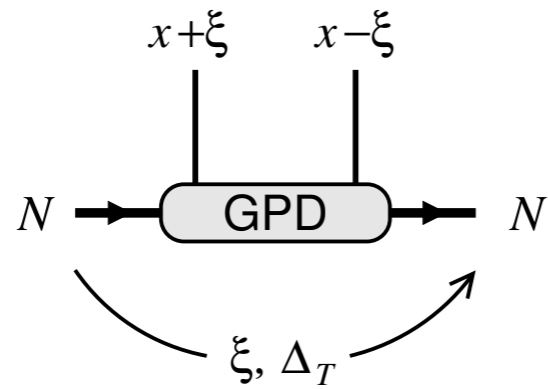
Definition of resonance GPDs

Transition matrix element to multihadron final state, e.g. πN

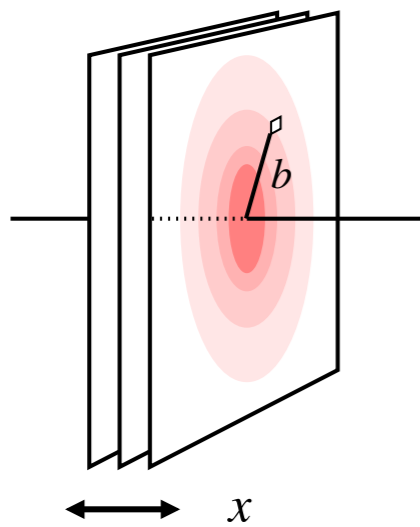
Analytic continuation in invariant mass $s_{\pi N}$: Pole at $s_{\pi N} = M_{\Delta}^2$, resonance structure defined at pole, residue factorizes

Rigorous definition of “resonance GPDs” using methods of S-matrix theory

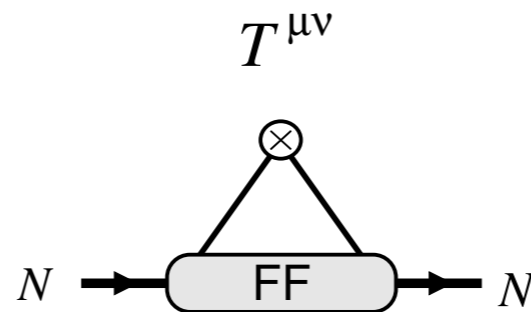
In physical region: Resonant + non-resonant contributions, needs theory



Concepts can be extended to $N \rightarrow N^*$ transitions!



Spatial distribution of quarks/gluons
“3D imaging”



QCD energy-momentum tensor:
mass, angular momentum, forces
“Mechanical properties”

$$|\mathbf{X}_T\rangle = \int \frac{d^2P_T}{(2\pi)^2} e^{i\mathbf{P}_T\mathbf{X}_T} |\mathbf{P}_T\rangle$$

$$\rho_{N \rightarrow N^*}(\mathbf{b}) = \int \frac{d^2\Delta_T}{(2\pi)^2} e^{-i\Delta_T\mathbf{b}} \langle N^*(-\Delta_T/2) | \mathcal{O} | N(\Delta_T/2) \rangle$$

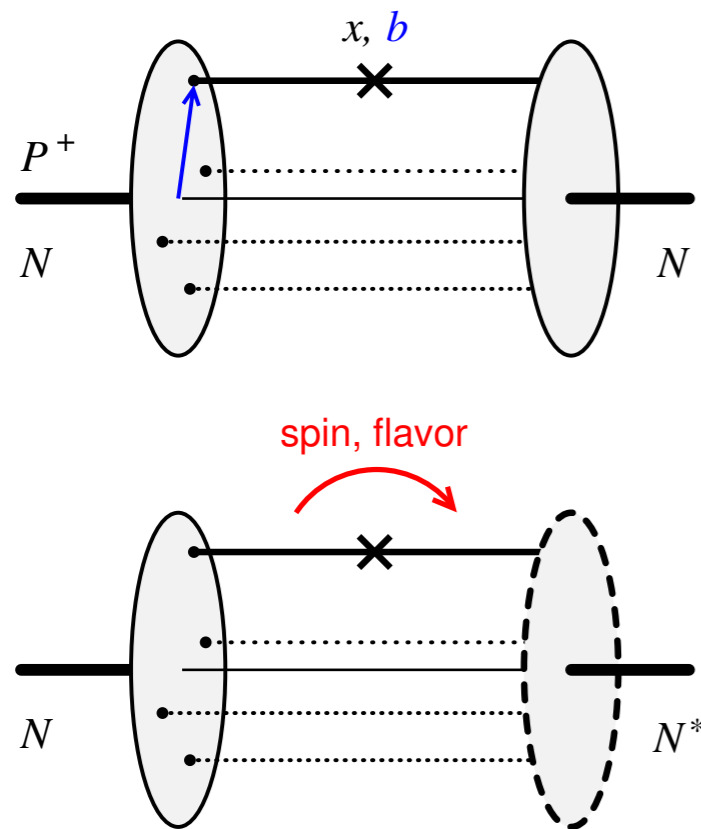
same P^+ , momentum transfer $\Delta_T, \Delta^- \neq 0$

LF dynamics: Transverse motion independent of mass of state. Galilean, same as nonrelativistic

Transversely localized states defined independently of mass

Transition densities $N \rightarrow N^*$

Cf. instant form: Motion depends on mass of state, definition of transition density ambiguous



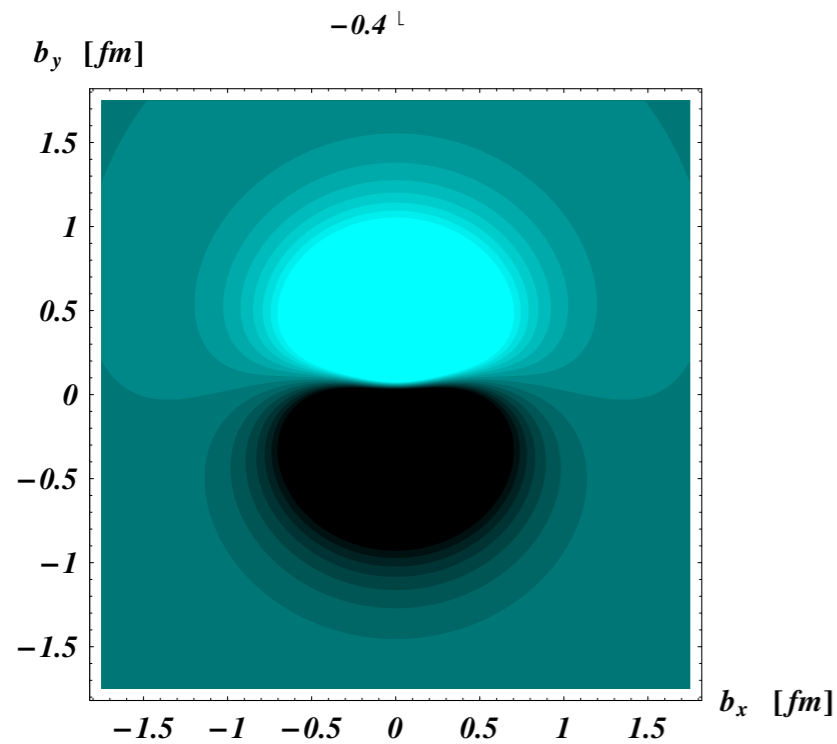
Interpretation

LF wave function overlap $N \rightarrow N^*$

Operator localized in transverse position, measures charge/current of constituents

Hadron undergoes excitation $N \rightarrow N^*$ (LF energy transfer $\Delta^- \neq 0$)

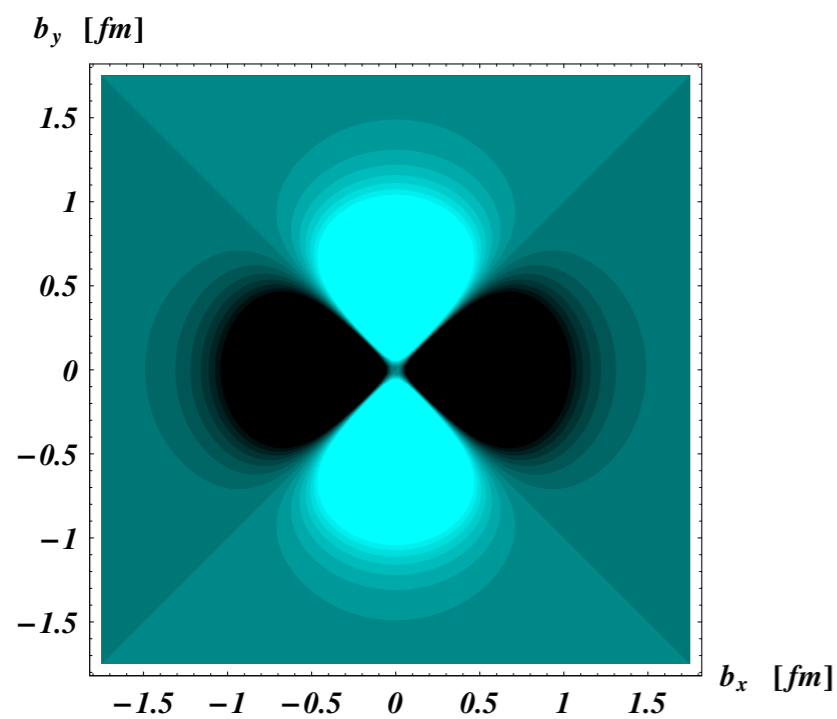
Operator can transfer quantum numbers: Spin, flavor/isospin



Example: Transverse charge density in $p \rightarrow \Delta^+$ transition (J^+ current density) from empirical transition form factors

Carlson, Vanderhaeghen 2007

Dipole and quadrupole patterns



$$T^{\mu\nu} = \sum_q T_q^{\mu\nu} + T_g^{\mu\nu} \quad \partial_\mu T^{\mu\nu}(x) = 0$$

Total EMT of QCD, conserved current (Noether thm)

$$T_u^{\mu\nu}, T_d^{\mu\nu} \quad T_{u\pm d}^{\mu\nu}$$

Individual quark flavor or gluon terms, not conserved

$$\rho_{N \rightarrow N^*}(\mathbf{b}) \quad \leftrightarrow \quad \langle N^*(-\Delta_T/2) | T_q^{++} | N(\Delta_T/2) \rangle$$

$$T_q^{+T}$$

LF transition densities from transition matrix elements of EMT terms

Transitions must be allowed by quantum numbers, e.g. $\langle \Delta | T_{u-d} | N \rangle$ by isospin

Conserved currents can have nonzero transition densities: see EM current

Conserved charges (= integrals of EMT) leave hadronic states unchanged and have no transitions

$$\hat{Q}|N\rangle = q|N\rangle \quad \langle N^*|\hat{Q}|N\rangle = q\langle N^*|N\rangle = 0$$

e.g. $\hat{Q} = P^+ \equiv \int dx^- d^2x_T T^{++}(x)$

total LF momentum

$$\mathbf{T}_{N \rightarrow N^*}^{+T}(\mathbf{b}) \leftrightarrow \langle N^*(-\Delta_T/2) | \mathbf{T}^{+T} | N(\Delta_T/2) \rangle$$

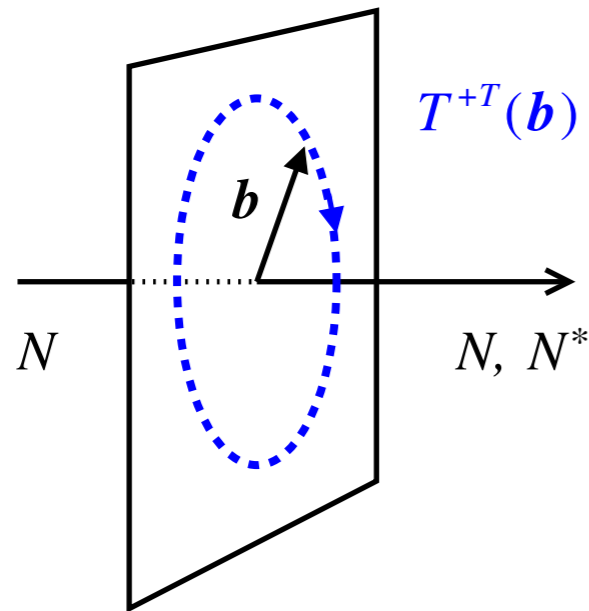
Transverse density of T momentum in hadron

$$J_{N \rightarrow N^*}^z = C_{\text{spin}} \int d^2b [\mathbf{b} \times \mathbf{T}^{+T}(\mathbf{b})]$$

Angular momentum as integral of transverse density

Kim, Won, Goity, Weiss, PLB 844,138083 (2023) [INSPIRE]

For AM definitions see e.g. Lorce, Mantovani, Pasquini 2017



$N \rightarrow \Delta$ transition angular momentum

Isovector quark angular momentum $u - d$

$$1/N_c \text{ expansion: } J_{p \rightarrow p}^{u-d} = \frac{1}{\sqrt{2}} J_{p \rightarrow \Delta^+}^{u-d}$$

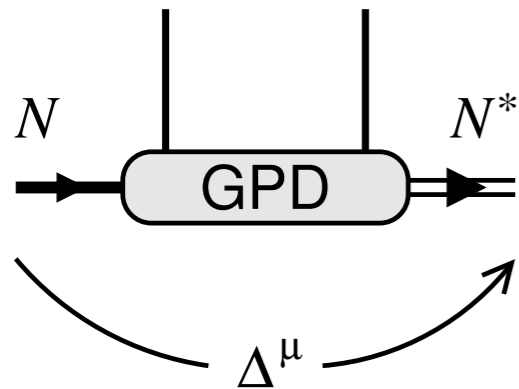
$p \rightarrow \Delta$ estimated using LQCD results for $p \rightarrow p$

Other gravitational form factors in $N \rightarrow \Delta$

Kim 2022; Alharazin et al. 2024

Lattice QCD	$J_{p \rightarrow p}^S$	$J_{\Delta^+ \rightarrow \Delta^+}^S$	$J_{p \rightarrow p}^V$	$J_{p \rightarrow \Delta^+}^V$	$J_{\Delta^+ \rightarrow \Delta^+}^V$
[9] $\mu^2 = 4 \text{ GeV}^2$	0.33*	0.33	0.41*	0.58	0.08
[10] $\mu^2 = 4 \text{ GeV}^2$	0.21*	0.21	0.22*	0.30	0.04
[11] $\mu^2 = 4 \text{ GeV}^2$	0.24*	0.24	0.23*	0.33	0.05
[12] $\mu^2 = 1 \text{ GeV}^2$	–	–	0.23*	0.33	0.05
[13] $\mu^2 = 4 \text{ GeV}^2$	–	–	0.17*	0.24	0.03

[9] Gökeler 2004. [10] Hägler 2008. [11] Bratt 2010. [12] Bali 2019. [13] Alexandrou 2020



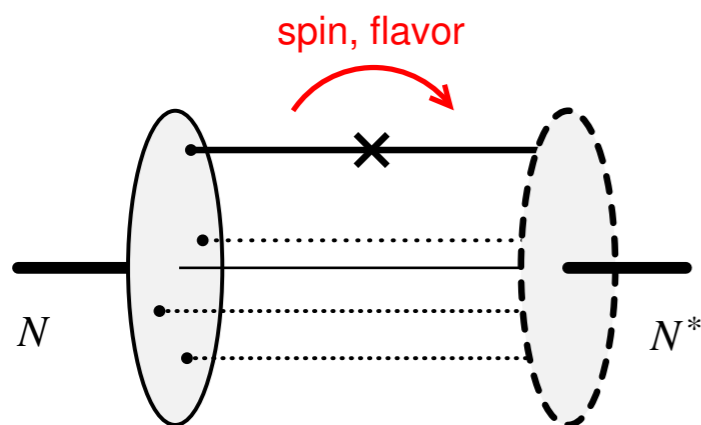
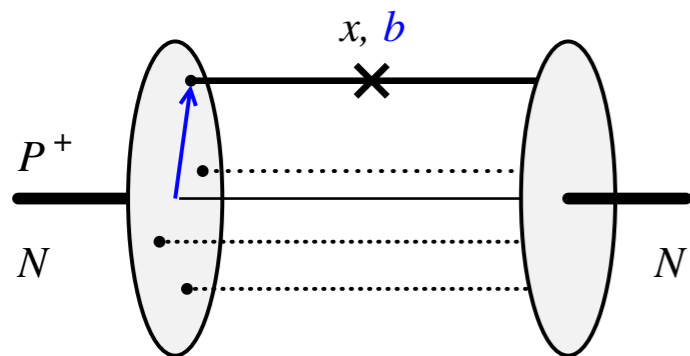
Forward limit of transition GPDs

$$\Delta^+ = -2\xi P^+ \rightarrow 0, \quad \Delta_T \rightarrow 0 \quad \text{light-front momentum transfer}$$

$$\Delta^- = \frac{m_{N^*}^2 - m_N^2}{2P^+} \neq 0 \quad \text{light-front energy transfer}$$

Zero light-front momentum transfer, only energy transfer

J-Y. Kim, C. Weiss, PLB 870 (2025) 139924 [INSPIRE]



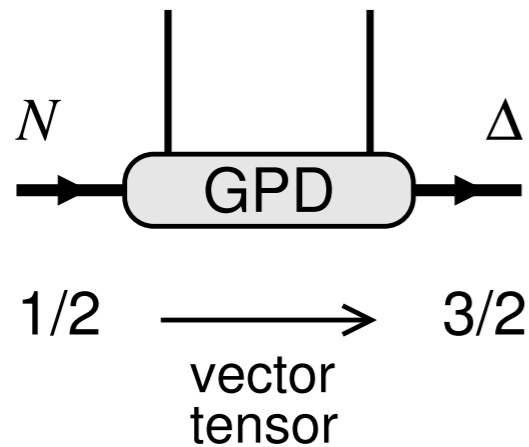
Transition PDFs

Describe excitation of baryon at rest (in LF momentum)

Represented as LF wave function overlap $N \rightarrow N^*$

Realize spin-isospin quantum numbers not accessible in $N \rightarrow N$

Also: Impact parameter dependence: $\Delta_T \neq 0 \rightarrow b$



$N \rightarrow \Delta$ spin transition

Transition vector and tensor

$$\psi_{3/2}^* [\dots] \psi_{1/2} \rightarrow V^i, Q^{ij} \quad \text{in 3D rest frame}$$

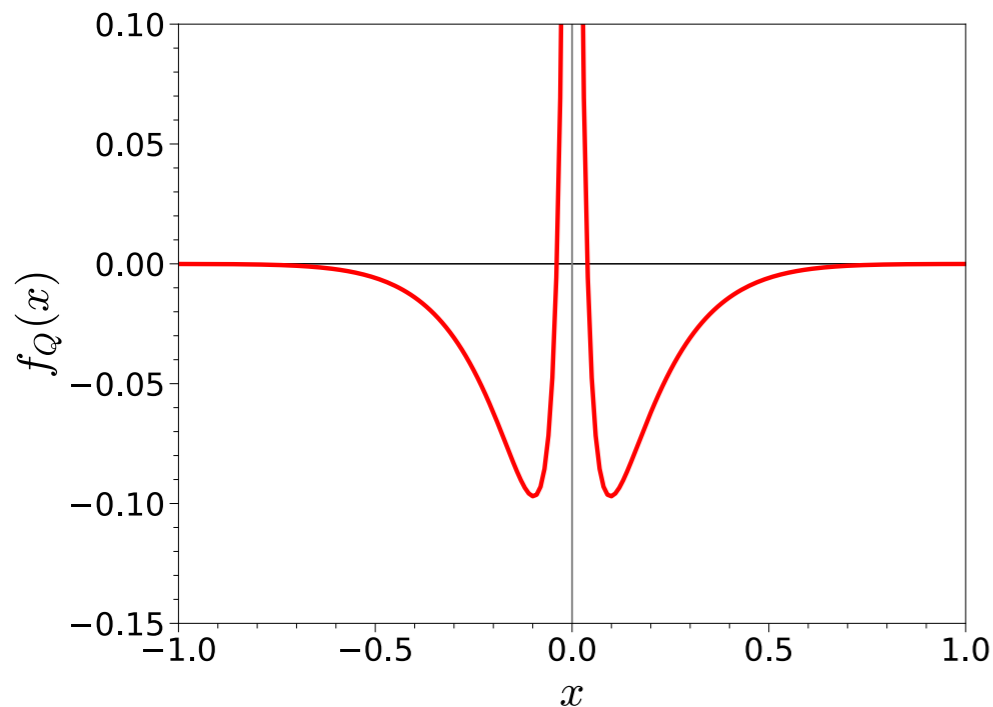
$N \rightarrow \Delta$ spin tensor PDF

Structure $\sim Q^{ij} n^i n^j$ 3D part of light-like 4-vector n

Unique to $1/2 \rightarrow 3/2$ transition matrix element of nonlocal partonic operator, absent in local vector current

$$\int_{-1}^1 dx f_Q(x) = 0 \quad \text{zero sum rule}$$

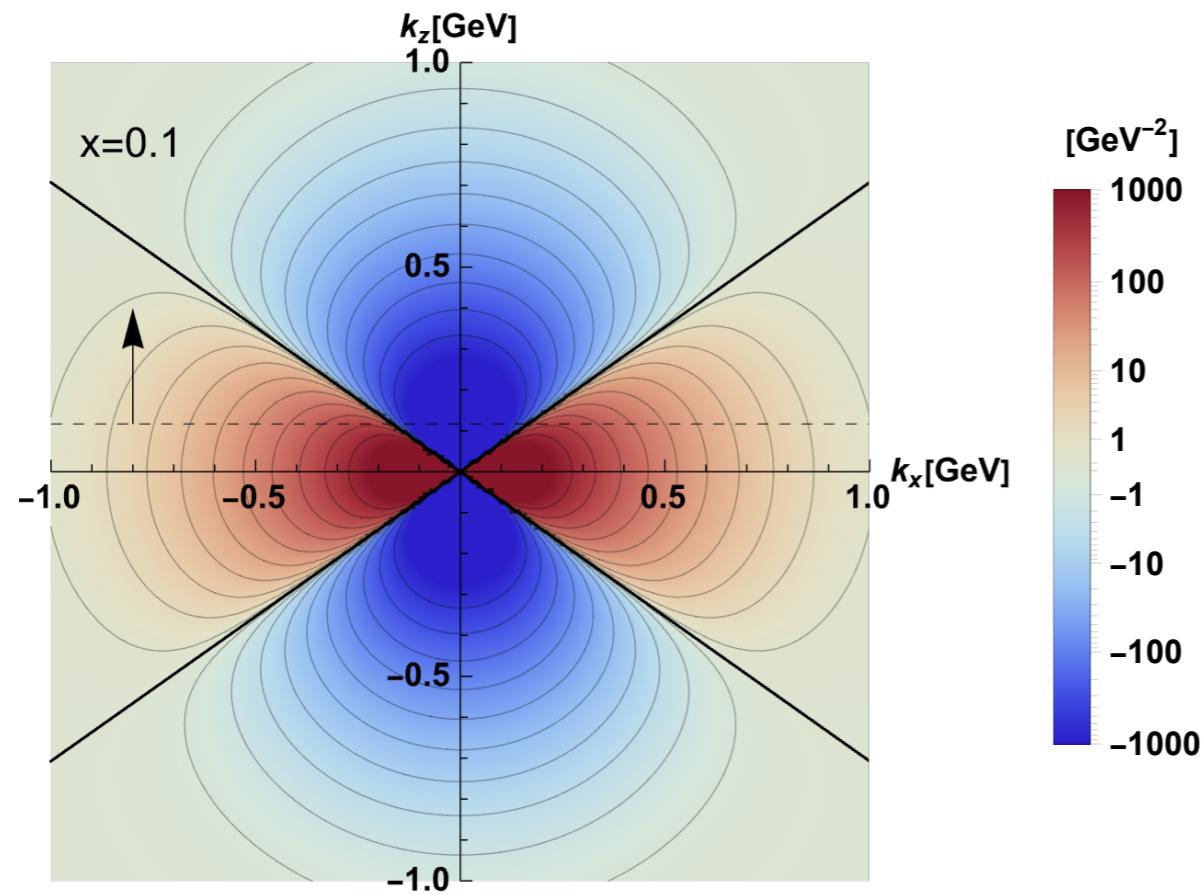
Subleading in $1/N_c$ expansion
 Estimated in chiral quark-soliton model



Complete spin structure of $N \rightarrow \Delta$ transition GPDs

J-Y. Kim et al. PRD 111, 114010 (2025) [\[INSPIRE\]](#)

J-Y. Kim, C. Weiss, PLB 870 (2025) 139924 [\[INSPIRE\]](#)



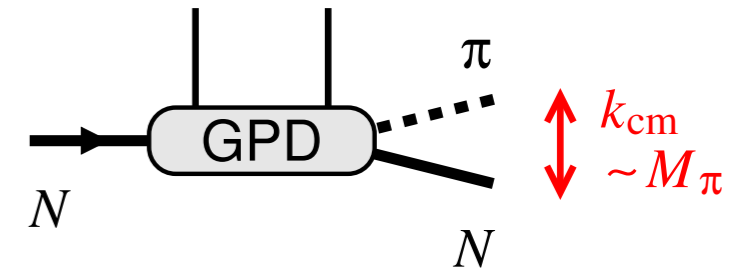
Rest-frame momentum distribution of quarks giving rise to $N \rightarrow \Delta$ spin tensor PDF
(chiral quark-soliton model, large N_c -limit)

J-Y. Kim, C. Weiss, PLB 870 (2025) 139924 [\[INSPIRE\]](#)

Chiral dynamics

Soft-pion theorems relate $N \rightarrow \pi N$ and $N \rightarrow N$ matrix elements

Pobylitsa, Polyakov, Strikman 2001; Guichon, Mossé, Vanderhaeghen 2003; Chen, Savage 2004; Birse 2004

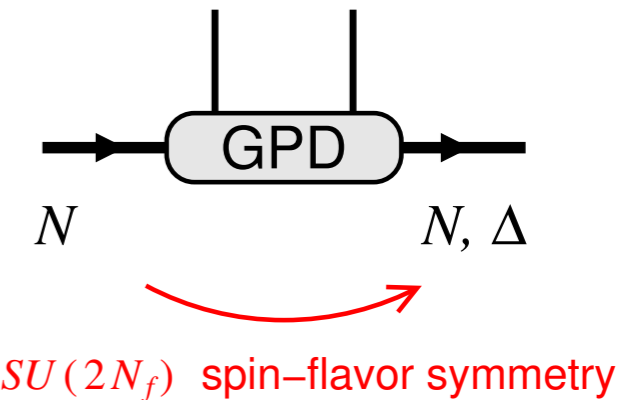


$1/N_c$ expansion of QCD

Spin-flavor symmetry relates $N \rightarrow N$ and $N \rightarrow \Delta$ transitions:

$$\langle \Delta | \mathcal{O} | N \rangle = [\text{symmetry factor}] \times \langle N | \mathcal{O} | N \rangle$$

Frankfurt, Polyakov, Strikman 1998. FPS, Vanderhaeghen 2000; Kim, Won, Goity, Weiss 2023



Effective degrees of freedom

Chiral soliton model, light-front quark models, holographic models, instanton vacuum

→ Talk Mamo

Lattice QCD

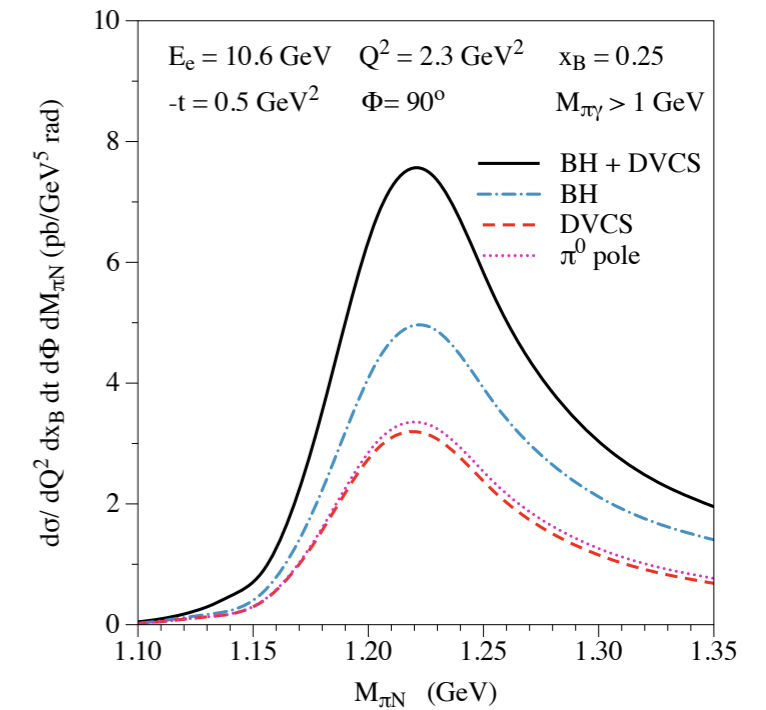
Partonic structure from Euclidean correlation functions

Deeply-virtual Compton scattering

$$e + p \rightarrow e' + \gamma + \Delta^+ \quad (\rightarrow \pi^0 p, \pi^+ n) \quad \text{also higher } N^*$$

Probes chiral-even GPDs

Cross section predictions: Semenov-Tian-Shansky, Vanderhaeghen 2023



Pseudoscalar meson production

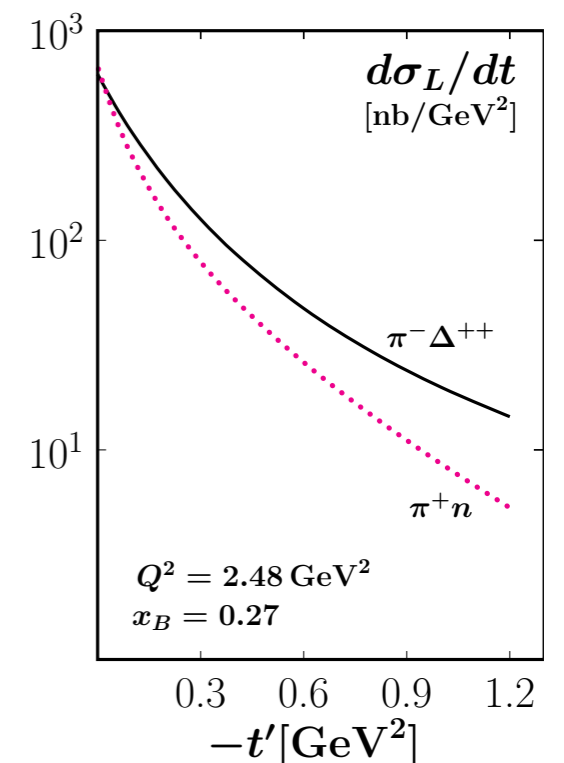
$$e + p \rightarrow e' + \pi^+ + \Delta^0 \quad \text{also } \eta, K \text{ mesons}$$

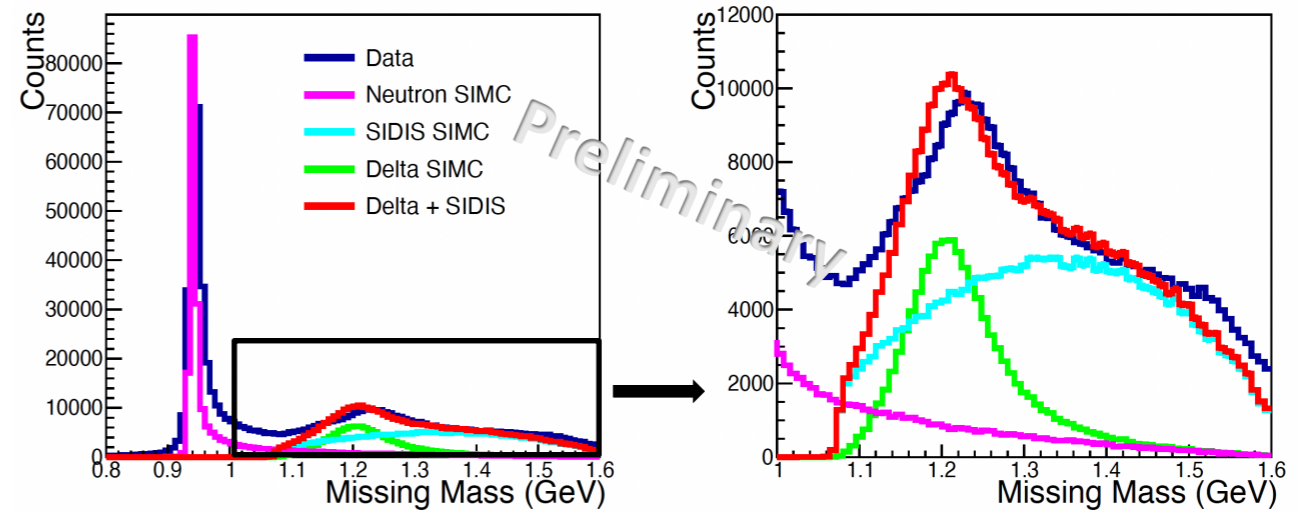
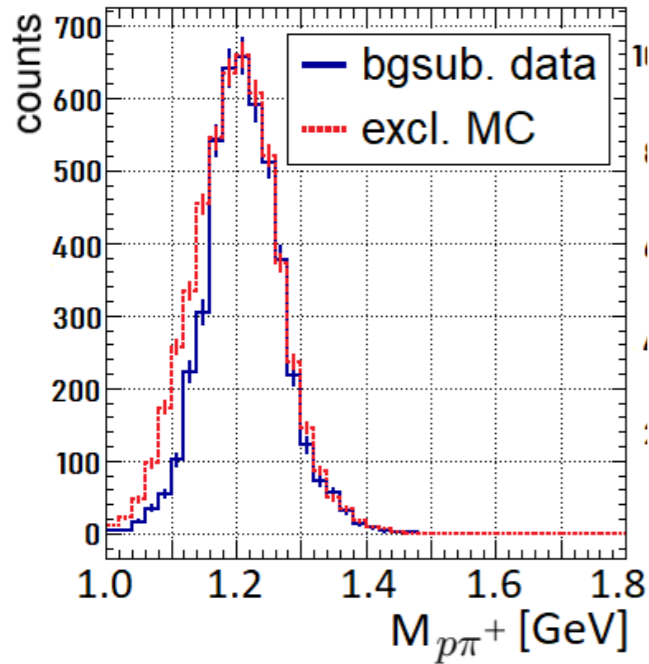
$$\pi^0 + \Delta^+$$

$$\pi^- + \Delta^{++}$$

Probes chiral-odd GPDs ($x \gtrsim 0.1$), mechanism tested in $p \rightarrow p$

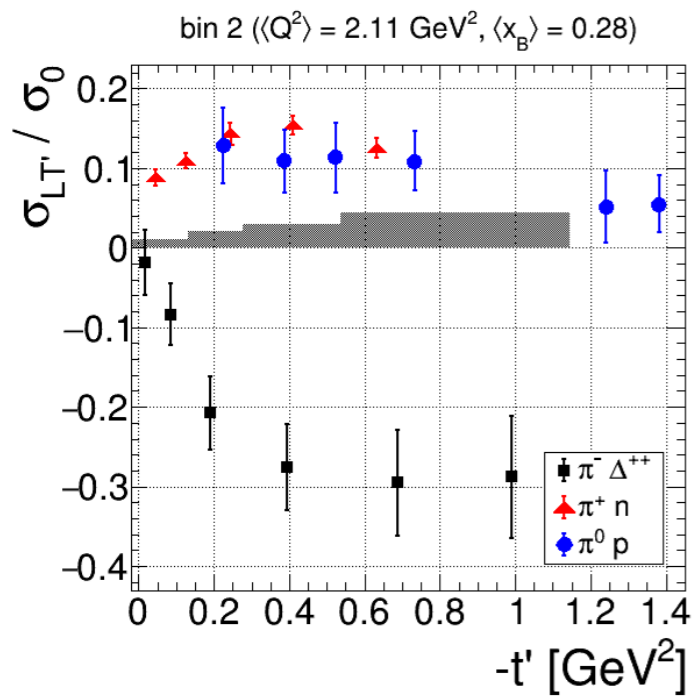
Cross section predictions: Kroll, Passek-Kumericki 2023





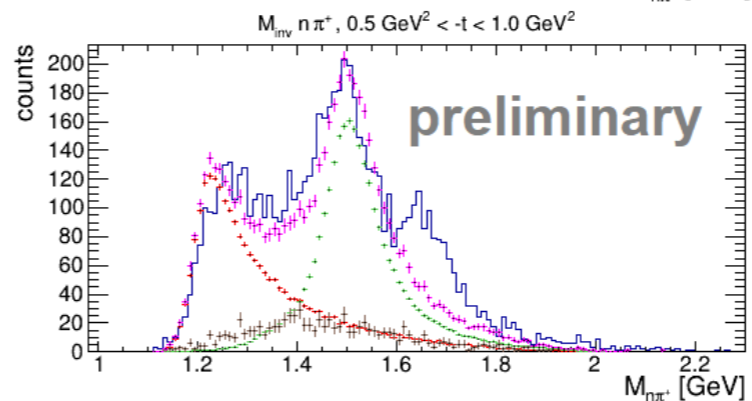
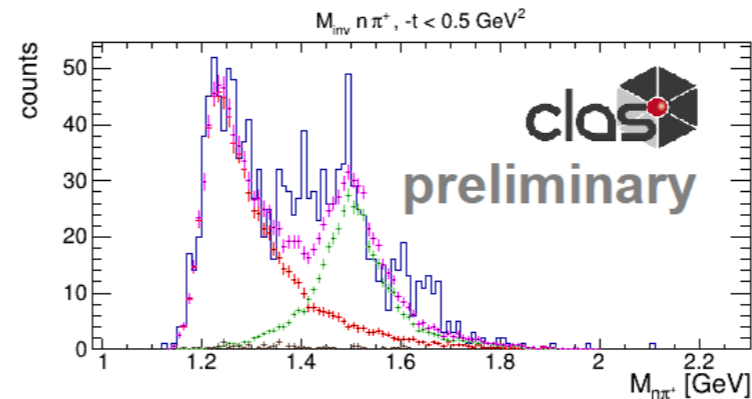
Hall C $ep \rightarrow e' \pi^+ \Delta^0$

A. Usman, ECT* Trento Workshop Aug 2023

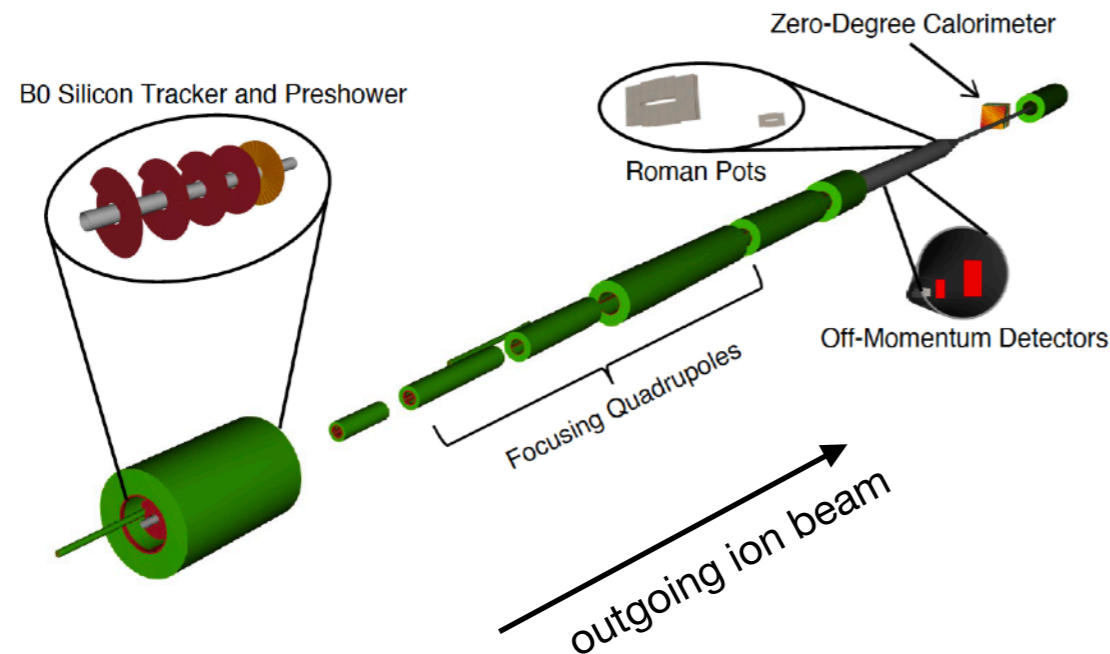


CLAS12 $ep \rightarrow e' \pi^- \Delta^{++}$

S. Diehl et al. PRL 131 (2023) 021901 [INSPIRE]



CLAS12 $ep \rightarrow e' \gamma n \pi^+$
DVCS



Charged hadrons: Forward spectrometer
 Neutral hadrons: Zero-Degree Calorimeter

Transition GPDs present “new” final states, complement/extend elastic channels

E.g. forward π^0 , forward π^\pm rigidity \ll beam

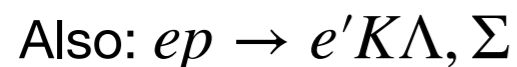
Transition channels that should be simulated



Strong decay, happens at vertex



Different decay modes of same Δ activate different detectors — charged-neutral, neutral-neutral, charged-charged. Could be used for tests and calibration besides physics interest



Can we reconstruct forward baryon resonances at EIC?

- Transition GPDs present new probe of excited baryon structure: Well-defined QCD operator, various quantum numbers accessible
- Transverse densities can be extended to $N \rightarrow N^*$ transitions with unequal masses
- EMT densities and mechanical properties can be formulated for $N \rightarrow N^*$ transitions
- Transition PDFs defined as forward limit of GPDs: Simple interpretation, new quantum numbers
- Theoretical methods for nonperturbative dynamics, experimental results
- Emerging field of study... many opportunities

Related topics

Transition GPDs in diffractive vector meson production $N \rightarrow X$ at small x

Transition GPDs in hadron-induced exclusive processes

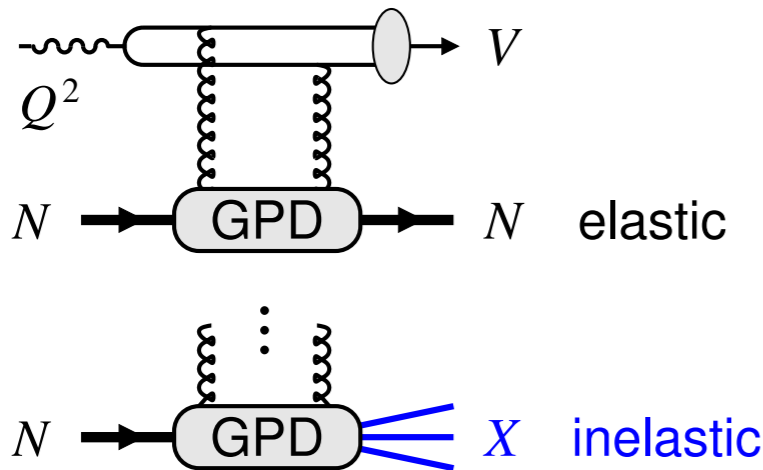
→ S. Diehl et al., EPJA 61, 131 (2025) [[INSPIRE](#)]

Transition GPDs with JLab 22 GeV

→ A. Accardi et al. EPJA 60, 173 (2024) [[INSPIRE](#)]

Supplemental material

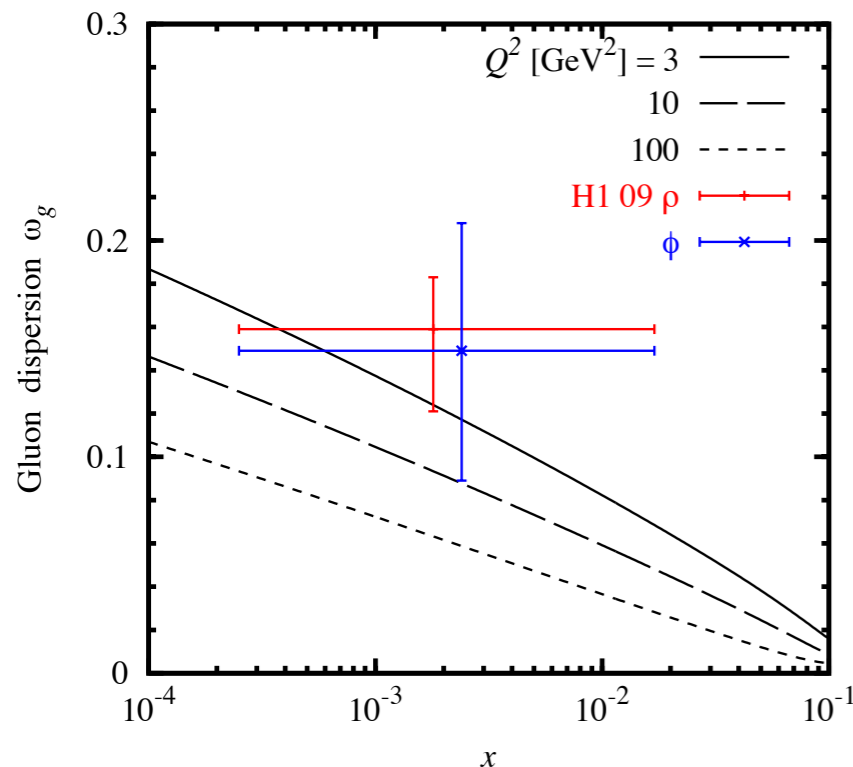
Transition GPDs: Vector meson production at small x 19



Diffractive vector meson production ($V = J/\psi, \phi, \rho^0$) with $N \rightarrow X(\text{low-mass})$ transitions

Probes quantum fluctuations of gluon density in nucleon: Frankfurt, Strikman, Treleani, Weiss PRL 101, 202003 (2008) [INSPIRE]

$$\omega_g \equiv \frac{\langle G^2 \rangle - \langle G \rangle^2}{\langle G \rangle^2} = \frac{d\sigma/dt (\gamma^* N \rightarrow VX)}{d\sigma/dt (\gamma^* N \rightarrow VN)} \Big|_{t=0}$$



Fluctuations formulated in context of collinear factorization and transition GPDs. Alt formulation in dipole model Schlichting, Schenke 2014; Mäntisaari, Schenke 2016

Discussed as part of diffraction at HERA and EIC: Inelastic diffraction

Connection: Diffraction \leftrightarrow transition GPDs