Analyzing $R_{s}(T, B_{pk})$

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- We would like to understand the linkage between material distribution within the surface and the resulting rf surface resistance.
- For "really good" Nb cavities we take the approximation:

$$R_{s_{eff}}(B_{pk},T) = R_{resid}(B_{pk}) + R_{BCS}(B_{pk},T) \quad \left[=\frac{A(B_{pk})}{T}e^{-\frac{U}{T}}\right]$$

Use multiple data sets with different *T* to make this separation.

Analysis assumptions:

- Effective gap/ T_c is constant, $2\Delta = 1.84$ meV, U = 17.02
- Heat flux is low enough that ΔT between rf surface and LHe bath T is negligible

Surface material determines A and R_{resid}









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- Xiao developed an extension of Mattis-Bardeen theory to non-zero $B: R_s(B,T)$
 - This analysis noted the effects of anisotropies in the distribution function of quasiparticles due to current flow, in the limit of quasiparticle thermal equilibrium.
 - The anisotropic Fermi surface effectively induces a broadening of the peaks in the quasiparticle density of states without significantly modifying the gap.
 - The limit of thermal equilibrium constraint requires that the quasiparticle inelastic scattering time is short compared with the rf cycle.
- A much more general theoretical treatment has recently been proposed by Gurevich
- The presence of thermally annealed surface interstitials lowers the local electron mean-free-path and also significantly affects the low-temperature disposition of available hydrogen.
- Subtle differences in diffusion profiles in Nb appear to significantly affect the quasiparticle scattering, relaxation time, and availability of hydrogen for forming bonds of different types and results in changes in A(B_{pk}).
- Somehow these factors are all interrelated the placement of atoms create all SRF properties.



