

Analysis and Reconstruction Strategies for Exclusive Diffractive eA Physics at ePIC

Exclusive, Diffraction, & Tagging Physics Working Group

Maci Kesler
(on behalf of the ePIC Collaboration)



Towards Improved Hadron Femtography, Edition V
University of Virginia

July 20 - 24, 2026



Exclusive, Diffractive, and Tagging (EDT) Physics Program in eA

Why is eA different from ep ?

- **Nuclear breakup:** spectators, forward nucleons, veto logic
- **Coherent vs incoherent diffraction:** imaging vs breakup background
- **Amplification of gluon density:** access to saturation dynamics
- **Allows 3D imaging:** momentum transfer distribution (t) dependent structure

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Toolbox of reconstruction methods developed to study these processes

- Multiple analysis channels (neutron structure function, diffractive VM production, etc.)
- Strategy depends on **nucleus, final state, and physics goal**



Exclusivity

- Detect full final state
- Scattered nucleus often escapes through beam pipe
- t reconstructed from combinations of final state 4-momenta

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Far-Forward Detector Systems

- **B0:**
 - Forward particles (protons, photons)
- **Roman Pots (RP)/Off-Momentum Detectors (OMD):**
 - **Tag & reconstruct forward protons**
- **Zero Degree Calorimeter (ZDC):**
 - **Tag neutral forward particles**

Will focus on 2 analyses

Neutron SF extraction

Tagged Deep Inelastic Scattering (TDIS)

- Spectator tagging
- Scattered electron reconstruction
- Pole extrapolation
- Obtaining SF

Diffractive ϕ production

- t reconstruction
- Resolving diffractive minima
- Transform to obtain spatial distribution
- Incoherent suppression

Both demonstrate strategies for **extracting nuclear structure** and rely on **exclusive reconstruction & tagging**

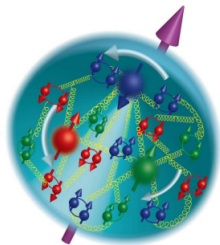
Tagged DIS in eD Collisions

Physics Motivation

Understanding structure of neutrons & light nuclei

Neutron Structure Functions (SF):

- Vital for full **understanding of nucleon structure & QCD**
- Probability distributions of quarks and gluons
- Tells us **how quarks are distributed** inside a nucleon
- Complementary to traditional DIS measurements of proton structure

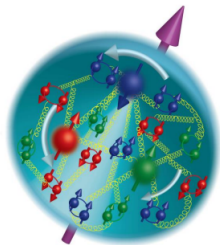


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Main challenges

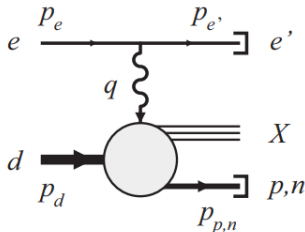
- Neutrons are neutral and cannot be easily accelerated
- Difficult to handle free neutrons

Possible Experimental Solution

Tagged DIS (TDIS) in eD Collisions

$$e + d \rightarrow e' + X + p(n)$$

- DIS on deuteron (d)
- Measure e' and spectator p or n
- Detecting spectator allows us to reconstruct struck nucleons kinematics



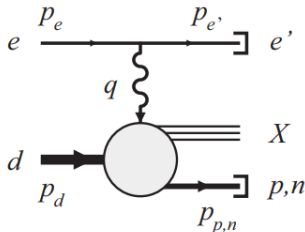
A. Jentsch, Z. Tu, C. Weiss: Phys. Rev. C 104, 065205

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Advantages:

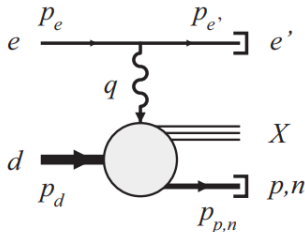
- Straightforward when spectators are successfully detected
- Possible to measure p
- Cross check with standard methods

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Advantages:

- Straightforward when spectators are successfully detected
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Challenges:

- Neutron is bound (not free) $\rightarrow$ **we are interested in free n SF**
- Detection of the spectator p or n

How Do We Overcome The Challenges?

Solution: Pole extrapolation for extraction of free neutron/proton structure

Nucleon reduced cross section obtained by removing pole of deuteron spectral function

$$\bar{\sigma}_{\text{red},n} = \frac{\bar{\sigma}_{\text{red},d}(x, Q^2)}{[2(2\pi)^2] \underbrace{S_d(p_{p,T}, \alpha_p)}_{\text{deuteron spectral function}}[\text{pole}]}$$

- We can measure/calculate this
- External input from model/experiment

Deuteron spectral function

$$S_d(p_{p,T}, \alpha_p)[\text{pole}] = \frac{R}{(p_{p,T}^2 + \alpha_T^2)}$$

- Position of pole:

$$\alpha_T^2 = m_N^2 - \alpha_p(2 - \alpha_p) \frac{M_d^2}{4}$$

- Residue of spectral function:

$$R = 2\alpha_p^2 m_N \Gamma^2 (2 - \alpha_p)$$

$$\Gamma = 0.007885 \text{ GeV}$$

A. Jentsch, Z. Tu, C. Weiss: Phys. Rev. C 104, 065205

M. Strikman and C. Weiss: Phys. Rev. C 97, 035209

W. Cosyn and C. Weiss: Phys. Rev. C 102, 065204

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- $\bar{\sigma}_{\text{red},n}$ at pole corresponds to a free n
 - $p_{pT}^2 \rightarrow -\alpha_T^2$ which means negative (unphysical) p_{pT}^2

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- Same measurement can be done for p
 - Cross check with proton F_2 extracted with traditional method
 - Can be used to select optimal deuteron spectral function

- We can measure/calculate this
- External input from model/experiment

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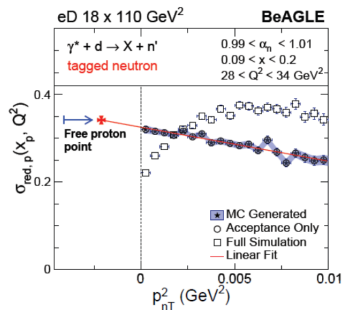
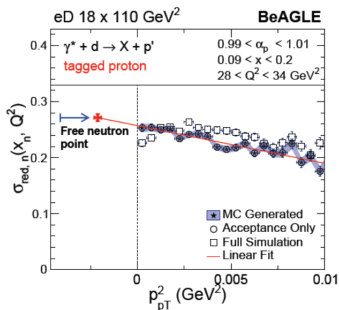
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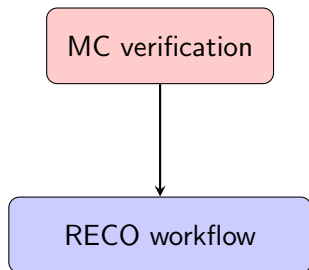


Experimental Method: MC Truth Validation

- Method verified in publication by A. Jentsch, Z. Tu, C. Weiss: Phys. Rev. C 104, 065205
 - Used only MC truth information
 - Manual smearing applied (no EICRecon)
 - RECO not yet incorporated



Experimental Method: This Analysis



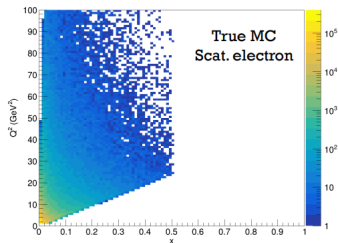
Implement complete TDIS framework:

- Using RECO info from EICRecon
- Full efficiency and acceptance corrections
- More realistic estimate of systematic uncertainties
- Projection of stat. uncertainties to early year EIC luminosity

Scattered Electron Finding

MC

- Find electron with beam electron as parent



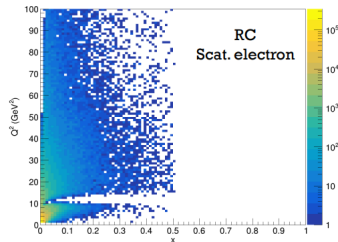
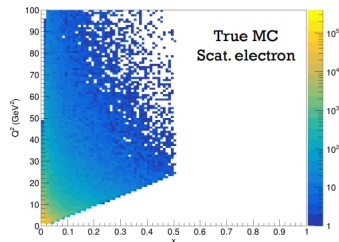
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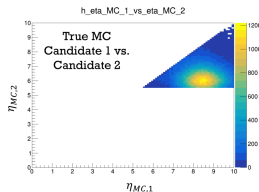
- Find electron with beam electron as parent

RECO (RC)

- Find negative track with well matched nECAL or bECAL cluster with $0.8 < E/p < 1.2$
- Select one with highest momentum
- Correct momentum magnitude using energy deposition in ECAL
- Main source of efficiency drop in final cross sections**



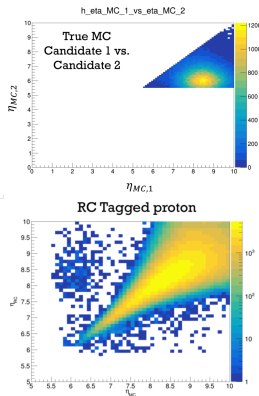
Spectator Selection



MC

- Select final state p, n with $\eta > 5.5$
- If two candidates, take the one with larger η

Spectator Selection



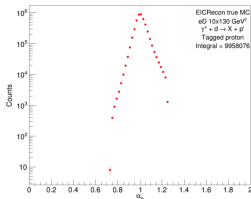
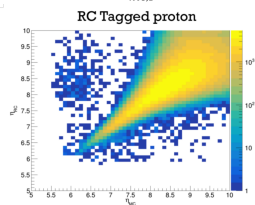
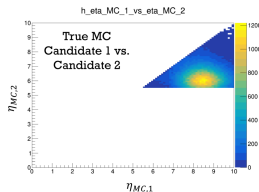
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RECO

- Accept events with candidate in OMD or ZDC
- If two candidates, take the one with larger η

Spectator Selection



MC

- Select final state p, n with $\eta > 5.5$
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RECO

- Accept events with candidate in OMD or ZDC
- If two candidates, take the one with larger η
- Require $0.99 < \alpha_p < 1.01$

$$\alpha_p \equiv \frac{2p_p^+}{p_d^+} = \frac{2(E_p + p_{z,p})}{E_d + p_{z,d}}$$

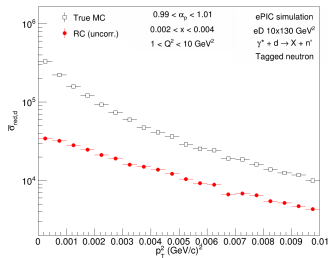
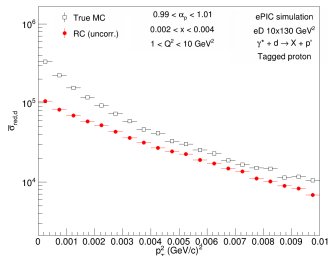
Deuteron Reduced Cross Section

Campaign 26.03.1

Top: Calculated using **tagged protons**

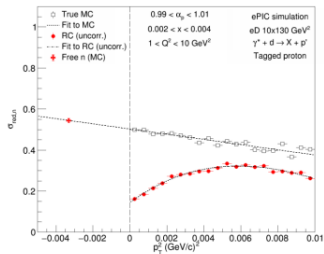
Bottom: Calculated using **tagged neutrons**

- Drop in counts due to efficiency loss by electron finding algorithm
- Change in shape of p_T^2 due to spectator resolution



Nucleon Reduced Cross Section - Uncorrected

Campaign 26.03.1



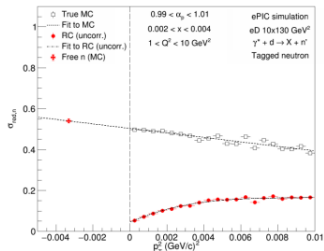
Top: Neutron reduced cross section
Bottom: Proton reduced cross section

MC:

- Fitted with linear function
- Extrapolation uncertainty from 1σ confidence interval of the fit

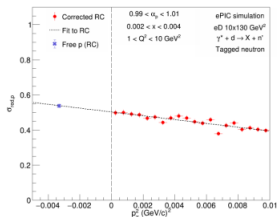
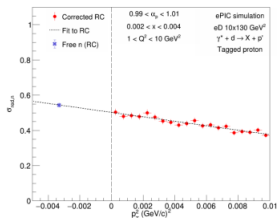
RECO (RC):

- Fitted with 3rd order polynomial



Nucleon Reduced Cross Section - Corrected

Campaign 26.03.1



Top: Fully corrected **neutron** reduced cross section in EICRecon

Bottom: Fully corrected **proton** reduced cross section in EICRecon

Sources of systematic uncertainties

- Luminosity: $\sigma_{\text{lumi}} = 1.5\%$
- Theory/extrapolation: $\sigma_{\text{th}} = 2\%$
- Beam effects: $\sigma_{\text{beam}} = 1.5\%$
- Proton reconstruction: $\sigma_p = 2\%$
- Neutron reconstruction: $\sigma_n = 1\%$

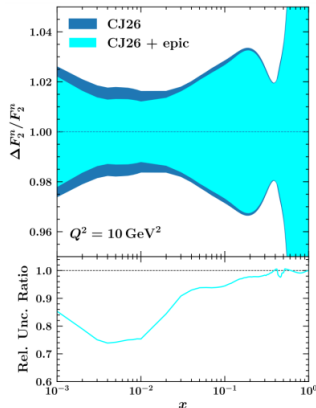
$$\sigma_{\text{sys}} = \sqrt{\sigma_{\text{lumi}}^2 + \sigma_{\text{th}}^2 + \sigma_{\text{beam}}^2 + \sigma_{p/n}^2}$$

$$\sigma_{\text{tot}} = \sqrt{\sigma_{\text{stat}}^2 + \sigma_{\text{sys}}^2}$$

- Error bars are total systematic uncertainties
- Extrapolation uncertainty from 1σ confidence interval of fit to dist. with combined uncertainties

Neutron SF Results

- Results clearly **demonstrate feasibility of measurement of TDIS in eD collisions** in early years of the EIC
- Expected luminosity more than sufficient to have **small or negligible statistical uncertainties**
- Final uncertainty will be dominated by systematic uncertainty
- This **measurement will have significant impact**, even with luminosity in the early EIC years



Calculated by Matteo Cerutti & Alberto Accardi
Analysis by Jan Vanek & Alex Jentsch

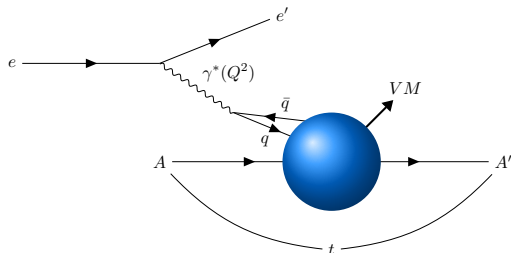
Momentum Transfer from Diffractive ϕ in eAu

Motivation

Critical Measurement:

Exclusive Vector Meson (VM) Production

- Probe to gluon density $\rightarrow$ precisely see structure $\rightarrow$ **saturation**
- Deflection of VM measures **spatial distribution** of gluons

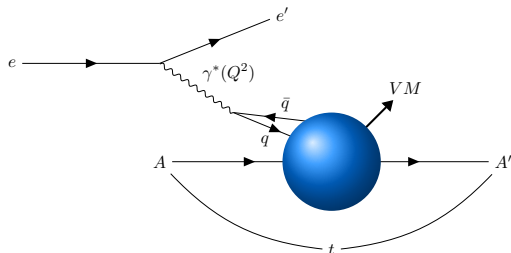


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Distribution of momentum transfer (t)

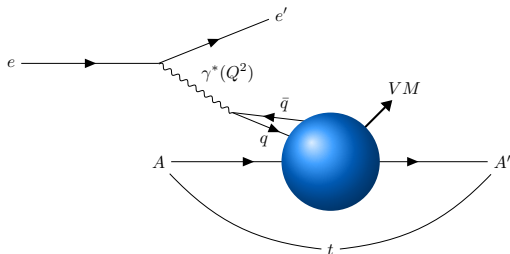
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- Reflects the spatial profile
- **Gluon imaging**

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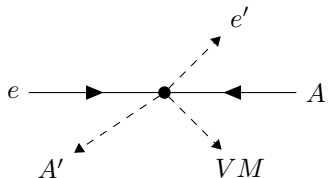
Distribution of momentum transfer (t)

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Saturation models predict at EIC **more than 20% of the cross-section will be diffractive**

EIC YR: <https://doi.org/10.1016/j.nuclphysa.2022.122447>

Reconstruct from exclusive VM production



$$|t| = (P_A - P_{A'})^2$$

To access $|t|$: **need complete final state**

We know 4-momenta of e , A , e' , and VM

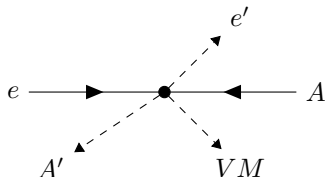
$$e + A \rightarrow e' + A' + VM$$

$$P_{A'} = P_e + P_A - P_{e'} - P_{VM}$$

Cannot measure $P_{A'}$

- For heavy nuclei $\rightarrow$ **stays within beam envelope**
- Tiny momentum change + tiny angular deflection

Reconstruct from exclusive VM production

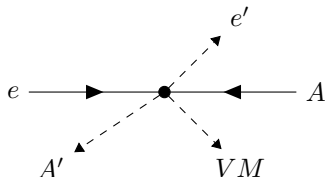


$$|t| = (P_A - P_{A'})^2$$

To access $|t|$: **need complete final state**

- $|t|_{\text{BABE}}$ (method E) MC: $|t| = (P_{VM} + P_{e'} - P_e)^2$
 - Small error/inaccuracy has large effect on $|t| \rightarrow \frac{\sigma_t}{t}$
 - Beam momentum spread

Reconstruct from exclusive VM production

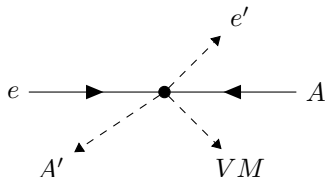


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- $|t|_{\text{eX}}$ (method A) default: $|t| = [\mathbf{p}_T(e') + \mathbf{p}_T(VM)]^2$
 - Underestimates true t
 - Valid only for small $|t|$ and small Q^2

Reconstruct from exclusive VM production



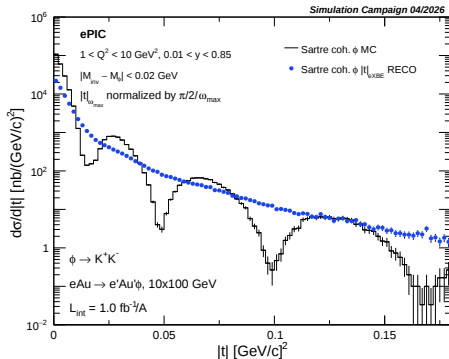
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- $|t|_{\text{eXBE}}$ (method L): $|t| = (P_A - P_{A'}^{\text{corr}})^2$
 - Improvement to $|t|_{\text{eX}}$
 - Only applies to coherent events

Challenges

Measurements of the $|t|$ distribution encounter **2 primary challenges:**



① Limited resolution in measuring $|t|$:

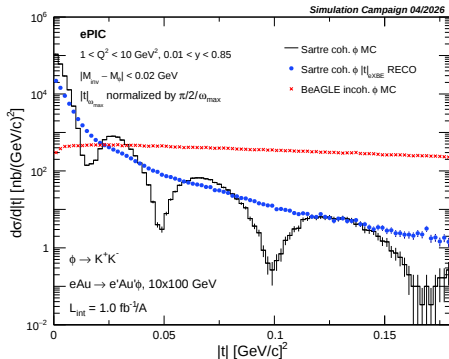
- **Mainly momentum resolution of e'**
 - Tiny angular error $\rightarrow$ large p_T error

Best method of $|t|$ reconstruction $\rightarrow |t|_{\text{eXBE}}$ (blue circles)

EIC YR: <https://doi.org/10.1016/j.nuclphysa.2022.122447>

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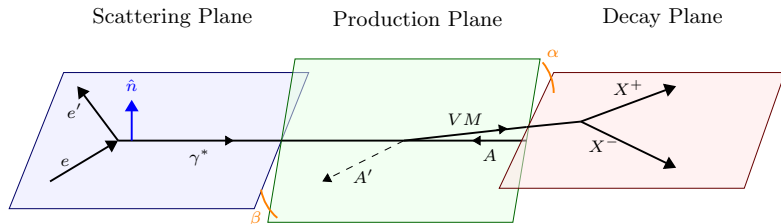
② Overwhelming incoherent background:

- **Red x's**
- Detector **cannot** suppress all incoherent production

Best method of $|t|$ reconstruction $\rightarrow |t|_{\text{eXBE}}$ (**blue circles**)

EIC YR: <https://doi.org/10.1016/j.nuclphysa.2022.122447>

Projected- $|t|$ Method



- Measure **$|t|$ component along normal direction ($\hat{n}$)** of electron scattering plane
- **$|t|\hat{n}$ unaffected by momentum resolution** $\rightarrow$ eliminate momentum resolution contribution from scattered electron

M. Kesler, A. I. Sheikh, et. al., 10.1016/j.physletb.2026.140585

Projection + Wedge Cut

Calculate $\hat{n}$

$$\hat{n} = \hat{p}_e \times \hat{p}_{e'}$$

$|t|_{\text{eXBE}}$

$$|t| = (P_A - P_{A'}^{\text{corr}})^2$$

$|t|_{\text{BABE}}$

$$|t| = (P_{VM} + P_{e'} - P_e)^2$$

$$|t|_{\hat{n}} = [(P_{VM} + P_{e'} - P_e) \cdot \hat{n}]^2$$

$$|t|_{\hat{n}} = (P_{VM} \cdot \hat{n})^2$$

$$|t| = (|t|_x, |t|_{\hat{n}}, |t|_z, |t|_E)$$

Want phase space where $|t|_{\hat{n}}$ dominates

Projection + Wedge Cut

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$$|t| = (|t|_x, |t|_{\hat{n}}, |t|_z, |t|_E)$$

Want phase space where $|t|_{\hat{n}}$ dominates

- e' direction ($\hat{x} = \hat{n} \times \hat{p}_e$)

$$\rightarrow \Delta_x = \sqrt{|t_x|}$$

- $\hat{n}$ direction

$$\rightarrow \Delta_y = \sqrt{|t|_{\hat{n}}}$$

Projection + Wedge Cut

Calculate $\hat{n}$

$$\hat{n} = \hat{p}_e \times \hat{p}_{e'}$$

$|t|_{\text{eXBE}}$

$$|t| = (P_A - P_{A'}^{\text{corr}})^2$$

$|t|_{\text{BABE}}$

$$|t| = (P_{VM} + P_{e'} - P_e)^2$$

$$|t|_{\hat{n}} = [(P_{VM} + P_{e'} - P_e) \cdot \hat{n}]^2$$

$$|t|_{\hat{n}} = (P_{VM} \cdot \hat{n})^2$$

$$|t| = (|t|_x, |t|_{\hat{n}}, |t|_z, |t|_E)$$

Want phase space where $|t|_{\hat{n}}$ dominates

- e' direction ($\hat{x} = \hat{n} \times \hat{p}_e$)

$$\rightarrow \Delta_x = \sqrt{|t_x|}$$

- $\hat{n}$ direction

$$\rightarrow \Delta_y = \sqrt{|t|_{\hat{n}}}$$

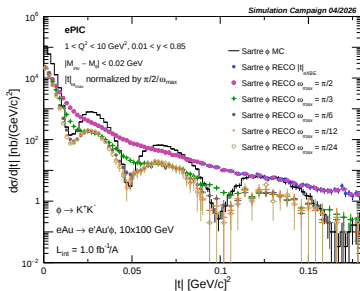
Parameterize $\Delta_{\perp} = \pm \sqrt{|t_{\perp}|}$

$$\Delta_x = \Delta_{\perp} \sin(\omega)$$

$$\Delta_y = \Delta_{\perp} \cos(\omega)$$

$$\omega_{\max} = \tan^{-1} \left(\frac{\Delta_x}{\Delta_y} \right)$$

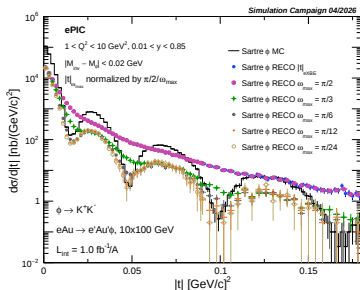
Resulting $|t|$ Distribution and Transformation



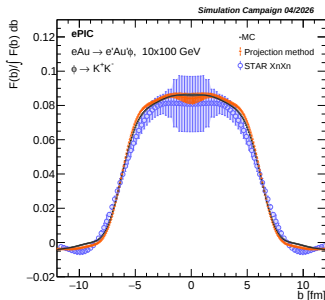
- Increase ω_{max} :
 - **Ability to resolve dips**
 - Lose statistics
- Need **precise location of minima**

Resulting $|t|$ Distribution and Transformation

Significant improvement in both distributions!



- Increase ω_{max} :
 - **Ability to resolve dips**
 - Lose statistics
- Need **precise location of minima**

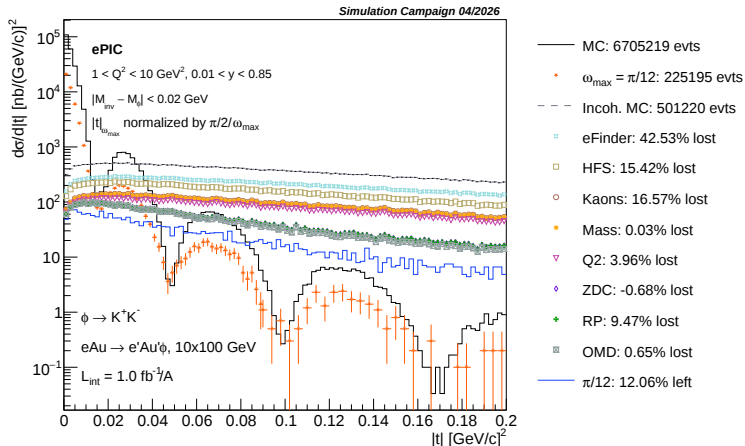


STAR Collaboration, 10.1103/PhysRevC.96.054904

- Small **errors in minima** = large **errors in density**
- Smaller $|b|$ determined by larger $|t|$



Incoherent Production



- **Dominant source of background** → can **only resolve first minima**
- Also studied DIS and ρ background sources and they are found to be negligible

Summary

TDIS: Neutron SF

- Uncertainty reduction on neutron SF
- Outlook: Polarized structure of deuterons

Diffractive ϕ : Gluon Imaging

- Recovery of diffractive minima & transformation
- Outlook: Incoherent suppression



EDT program is delivering early EIC physics!

Backup Slides

Analysis Strategies:

Calculating the distribution of momentum transfer, t

Several methods

- Each with distinct advantages
- Depends on specific reaction under study

t can be calculated using combinations of 4 momenta of final states:

- p_{BA}^μ : outgoing **B**Aryon 4-momentum
- p_{BE}^μ : incoming proton **B**Eam momentum
- $p_{\gamma^*}^\mu$: difference between incoming & scattered electron 4-momenta (i.e. information relating to e & e')
- p_{X}^μ : 4-momentum of the rest of final state, **X**

Tagged DIS in eD Collisions at $10 \times 130 \text{ GeV}^2$

Variables Used in Analysis

Scattered electron:

- $Q^2 = -q^2 = -(p_{e,\text{beam}} - p_{e,\text{scat}})^2$
- $x = \frac{Q^2}{P \cdot q}$
- $x_{\text{nucleon}} = \frac{Q^2}{2P_{\text{nucleon}} \cdot q} = \frac{x}{2 - \alpha_p}$

Kinematic variables:

- $y = \frac{p_d \cdot q}{p_d \cdot p_{e,\text{beam}}}$
- $1 - \epsilon = \frac{y}{1 + (1 - y)^2}$

Light cone momentum fraction:

$$\alpha_p \equiv \frac{2p_p^+}{p_d^+} = \frac{2(E_p + p_{z,p})}{E_d + p_{z,d}}$$

Proton transfer momentum:

- $p_{T,p} = \sqrt{p_{x,p}^2 + p_{y,p}^2}$

Fine structure constant:

- $\alpha_{\text{em}} = \frac{1}{137}$

Luminosity and cross section:

- $L_{\text{int}} = \frac{N_{\text{events}}}{\sigma_{\text{tot}}}$
 - $\text{mb} = 2.568 \text{ GeV}^{-2}$
 - $\sigma_{\text{tot}} = 3.538 \cdot 10^{-4} \text{ mb (en)}$
 - $\sigma_{\text{tot}} = 3.707 \cdot 10^{-4} \text{ mb (ep)}$

Tagged DIS in eD Collisions at $10 \times 130 \text{ GeV}^2$

Deuteron Reduced Cross Section

- Differential cross section on d can be written in terms of deuteron reduced cross section ($\sigma_{\text{red},d}$) and photon flux A. Jentsch, Z. Tu, C. Weiss: Phys. Rev.

C 104, 065205

M. Strikman and C. Weiss: Phys. Rev. C 97, 035209

W. Cosyn and C. Weiss: Phys. Rev. C 102, 065204

$$d\sigma_d = \underbrace{\frac{2\pi\alpha_{\text{em}}^2[1+(1-y)^2]}{Q^4x}}_{\text{photon flux}} \times \sigma_{\text{red},d} \times \underbrace{dx dQ^2 \frac{d\phi_{e'}}{2\pi}}_{\text{scat. electron}} \underbrace{[2(2\pi)^3]^{-1} \frac{d\alpha_p}{\alpha_p} \frac{dp_{T,p}^2}{2} d\phi_p}_{\text{spectator}}$$

- Measured reduced cross section (integrated over $d\phi_p$)

$$\bar{\sigma}_{\text{red},d} = \underbrace{\left(\frac{2.568}{L_{\text{int}}}\right)}_{\text{lumi + unit conversion}} \frac{Q^4x}{2\pi\alpha_{\text{em}}^2[1+(1-y)^2]} \frac{[4(2\pi)^3]\alpha_p}{\Delta\alpha_p} \frac{dN}{\Delta x \Delta Q^2 \Delta p_{T,p}^2 2\pi}$$

Full azimuthal coverage for electron: $\Delta\phi_{e'}/2\pi = 2\pi/2\pi = 1$

Full azimuthal coverage for spectator: $\Delta\phi_p = 2\pi$

Tagged DIS in eD Collisions at $10 \times 130 \text{ GeV}^2$

Simulation: BeAGLE

Overview:

- Campaign 26.03.1
- 20M events of eD $10 \times 130 \text{ GeV}^2$
 - 10M en and 10M ep

Luminosity and cross section:

- Total integrated luminosity (EICRecon):

$$L_{\text{int}} = 0.056 \text{ fb}^{-1}$$

- Expected luminosity for year 2:

$$L_{\text{int}} = 1.5 \text{ fb}^{-1}$$

- All statistical uncertainties projected to expected luminosity

- en : $\sigma_{\text{tot}} = 3.538 \cdot 10^{-4} \text{ mb}$, $L_{\text{int}} = N_{\text{evt}}/\sigma_{\text{tot}} = 10\text{M}/(3.538 \cdot 10^{-4} \text{ mb}) = 10\text{M}/(3.538 \cdot 10^8 \text{ fb}) = L_{\text{int}} = 0.028 \text{ fb}^{-1}$
- ep : $\sigma_{\text{tot}} = 3.707 \cdot 10^{-4} \text{ mb}$, $L_{\text{int}} = 0.028 \text{ fb}^{-1}$

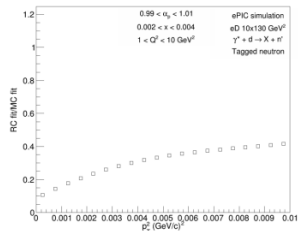
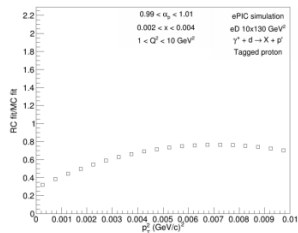


<https://eic.github.io/software/beagle.html>

Tagged DIS in eD Collisions at $10 \times 130 \text{ GeV}^2$

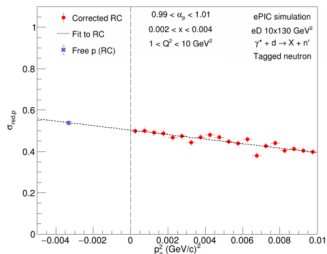
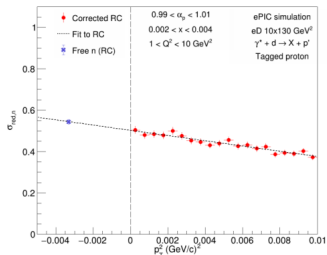
Efficiency and Acceptance

- Efficiency and bin migration correction done using ratio of MC and RC fit functions to nucleon reduced cross section
 - $\text{Eff.} = \text{RC fit}/\text{MC fit}$
 - Functions evaluated at histogram bin centers
- Efficiency then applied back to RC reduced cross section point by point



Tagged DIS in eD Collisions at $10 \times 130 \text{ GeV}^2$

Systematic Uncertainties (1/2)



Luminosity: $\sigma_{\text{lumi}} = 1.5\%$

- From YR and current discussion

Theory/extapolation: $\sigma_{\text{th}} = 2\%$

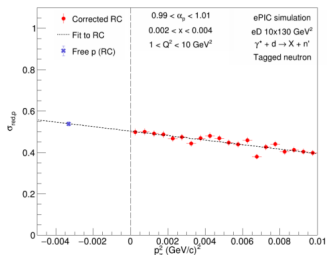
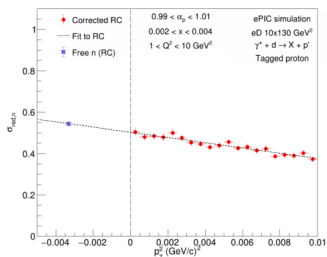
- From original paper
- A. Jentsch, Z. Tu, C. Weiss: Phys. Rev. C 104, 065205

Beam effects: $\sigma_{\text{beam}} = 1.5\%$

- Crossing angle, angular divergence, beam momentum spread
- Based on values cited by ATLAS or HERA
- ATLAS Collaboration: The European Physical Journal C 83 (2023)
- HERA Collaboration: Nuclear Physics B 816, 1–61 (2009)

Tagged DIS in eD Collisions at $10 \times 130 \text{ GeV}^2$

Systematic Uncertainties (2/2)



Proton reconstruction: $\sigma_p = 2\%$

- Uncertainty of transport matrix
- Based on values cited by ATLAS
- ATLAS Collaboration, The European Physical Journal C 83 (2023)

Neutron reconstruction: $\sigma_n = 1\%$

- Uncertainty of the primary vertex resolution
- Based on values cited by HERA and ePIC Roman Pot studies
- HERA Collaboration, Nuclear Physics B 816, 1-61 (2009)

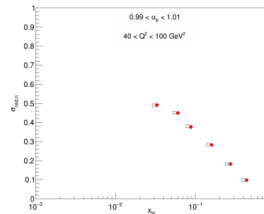
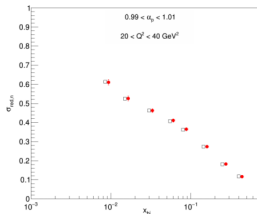
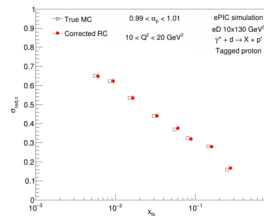
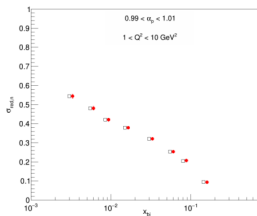
Total systematic uncertainty:

$$\sigma_{\text{sys}} = \sqrt{\sigma_{\text{lumi}}^2 + \sigma_{\text{th}}^2 + \sigma_{\text{beam}}^2 + \sigma_{n/p}^2}$$

Tagged DIS in eD Collisions at 10×130 GeV²

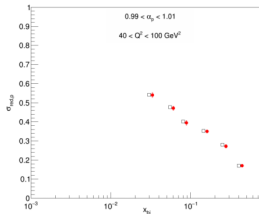
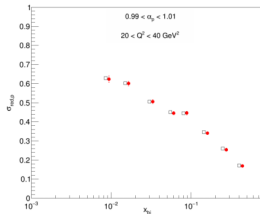
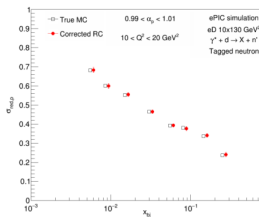
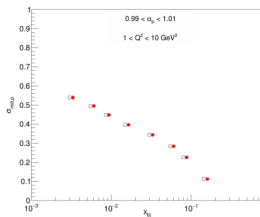
Neutron Reduced Cross Section vs x_{bj}

- Comparison of **MC** and **RECO** neutron reduced cross sections in bins of (x_{bj}, Q^2)
- Uncertainties dominated by systematics
- **Statistical uncertainties negligible** even in early years of running the EIC



Tagged DIS in eD Collisions at $10 \times 130 \text{ GeV}^2$

Proton Reduced Cross Section vs x_{bj}

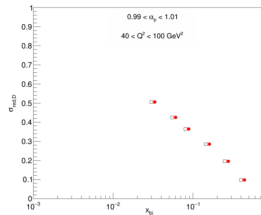
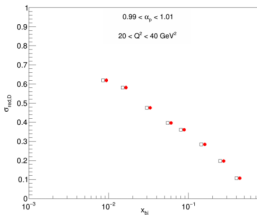
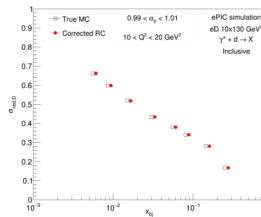
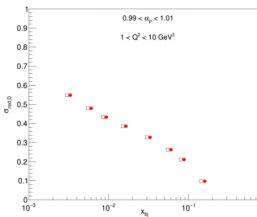


- Comparison of **MC** and **RECO** proton reduced cross sections in bins of (x_{bj}, Q^2)
- Uncertainties dominated by systematics
- **Statistical uncertainties negligible** even in early years of running the EIC

Tagged DIS in eD Collisions at 10×130 GeV²

Dueteron Reduced Cross Section vs x_{bj}

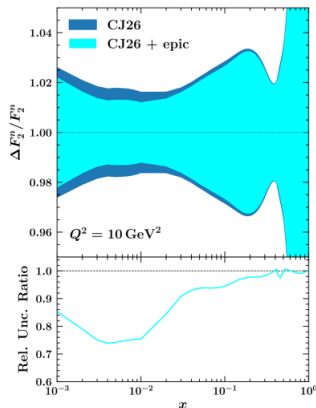
- Comparison of **MC** and **RECO** inclusive deuteron reduced cross sections in bins of (x_{bj} , Q^2)
- Used primarily for calculations by our theory colleagues for ESR



Tagged DIS in eD Collisions at $10 \times 130 \text{ GeV}^2$

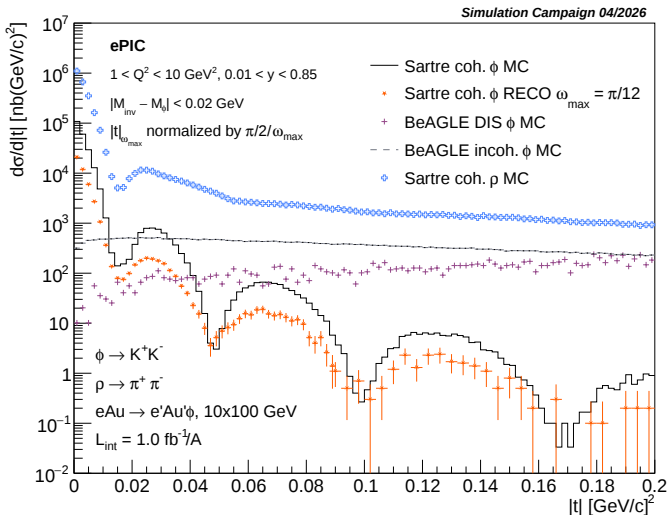
Result for ESR

- Impact study demonstrating uncertainty reduction on neutron SF F_2^n after using TDIS pseudo-data from ePIC
- Expected 20% uncertainty reduction for $10^{-3} < x < 10^{-2}$
- Calculation by Matteo Cerutti and Alberto Accardi



Diffractive ϕ Production

Background Contributions

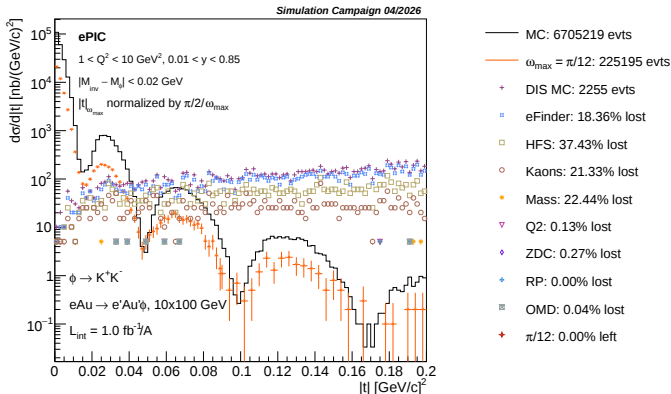


- 1 DIS background
- 2 Incoherent production
- 3 ρ contamination

Note: will implement more sources of uncertainty in future

Diffractive ϕ Production

DIS Background

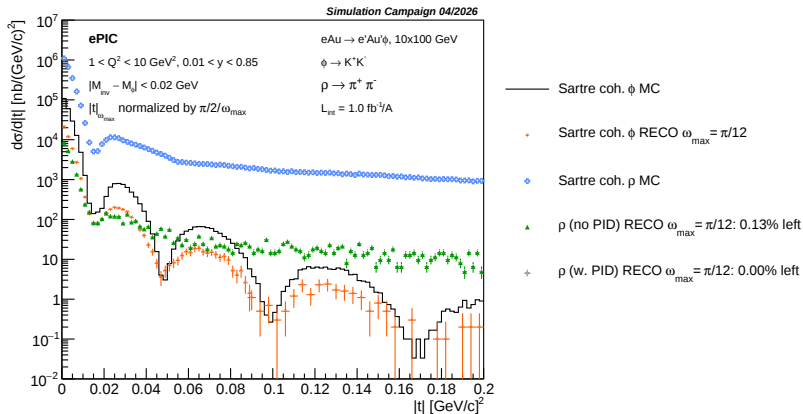


Distribution is dominated by DIS background contributions
Applying all vetoes and cuts **removes DIS background**



Diffractive ϕ Production

ρ Contamination



VM daughters can be misidentified
Applying PID removes ρ contribution