

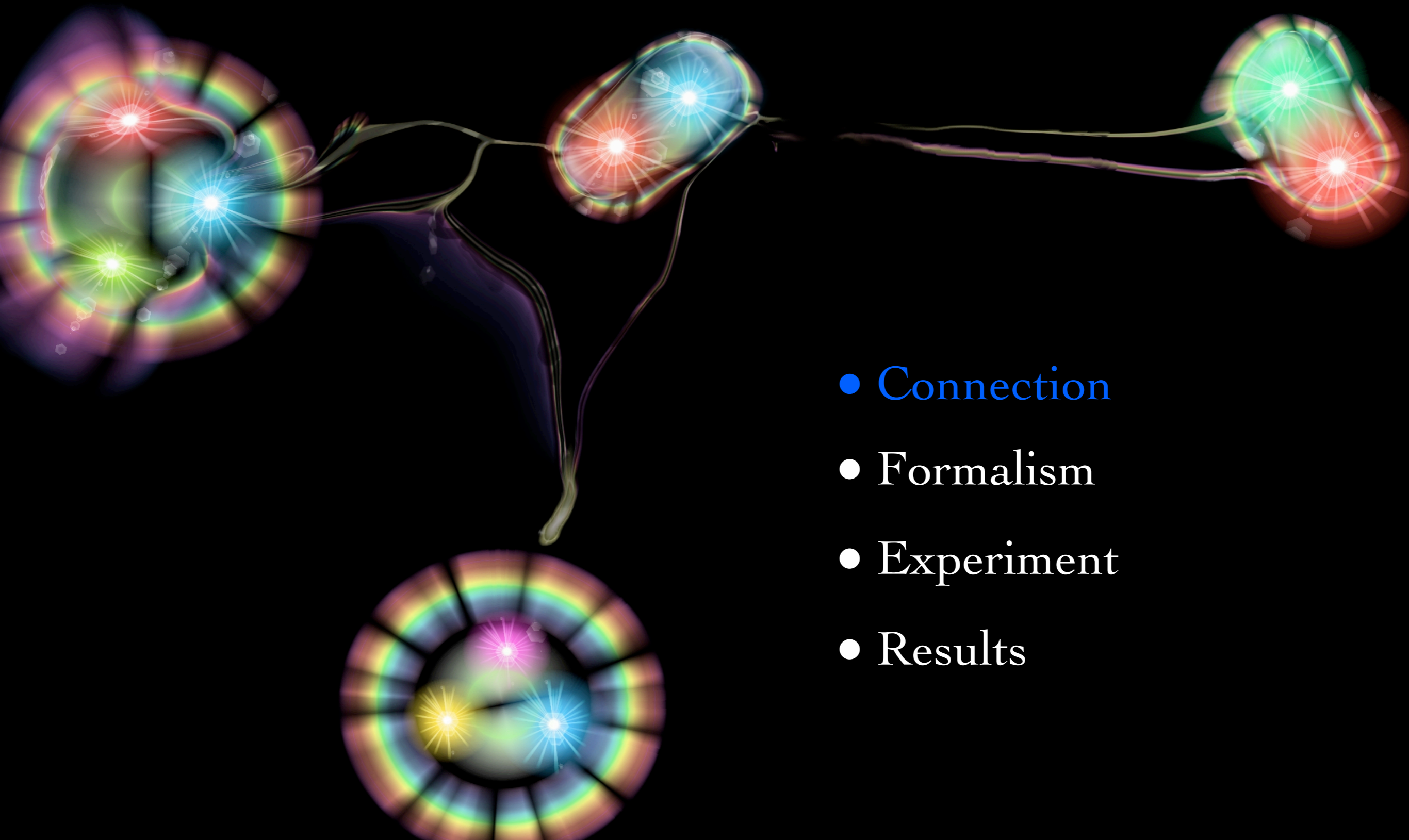
# Pion Electroproduction at CLAS

Taisiya Mineeva



# Synopsis

- Electroproduction measurements were performed off  $^2\text{H}$ ,  $^{12}\text{C}$ ,  $^{56}\text{Fe}$ ,  $^{207}\text{Pb}$  nuclear targets in Jefferson Lab Hall B using incoming electron energy of 5 GeV and CLAS detector
- The goal is to measure SIDIS yield of hadrons along with transverse momentum broadening in large nuclei compared to  $^2\text{H}$
- Such measurements will provide insights onto space-time development of hadronization in cold nuclear medium.



- Connection
- Formalism
- Experiment
- Results

# Nuclear effects in neutrino scattering experiments

## • Incoming $\nu_\mu$ flux

uncertainties on the hadron yield produced off C(Be) targets as a function of  $P_z$  and  $P_T$   
knowledge on  $\pi/K$  ratio gives constraint on  $\nu_\mu$  background

-> MINOS: 9% for  $E_\nu < 6\text{GeV}$  on  $\Delta m^2$

P.Adamson *et al.*, A study of muon neutrino disappearance using the Fermilab Main Injector Neutrino beam 6 arXiv 0711.0769

-> T2K: 7.43% (4.3%) on predicted number of events from SK (w/o oscillations)

model uncertainty on production rate of  $\pi$  is 50%, while for K is 15-100 %

K.Abe *et al.*, First muon-neutrino disappearance study with an off-Axis beam, arXiv:1201.1386

## • Difference in $E_{\text{visible}}$ and $E_\nu$

$$E_\nu = E_\mu + E_{\text{visible}}(\text{hadron}) + E_{\text{missing}}$$

FSI - one of the largest sources of background and inefficiency accounting for  $E_{\text{missing}}$

-> MINOS: uncertainty of neutrino production of hadrons is 8.2%

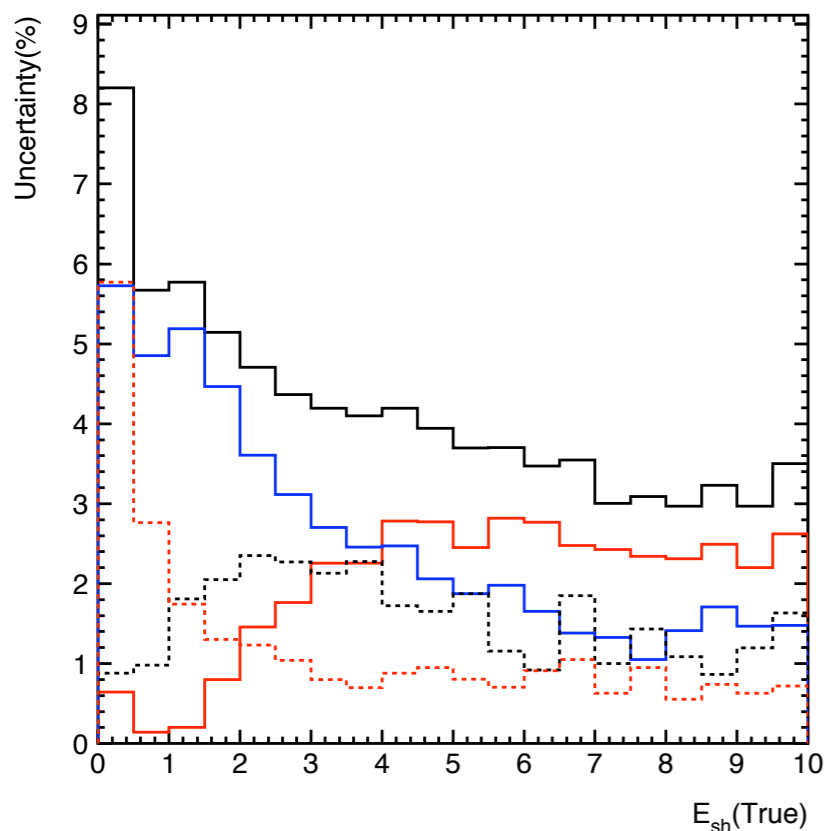
-> T2K: uncertainty due to final-state interactions is 3.2% (5.9%) (w/o oscillations)

## • NC $\pi^0$ production

background in  $\nu_\mu \leftrightarrow \nu_e$ : one of  $\pi^0$  decay products can be misidentified as  $l$  from CC event  
(MiniBooNE, T2K (with dedicated Pi-Zero detector), Icarus, Nova, etc )

# Hadronic shower scale uncertainties

$$E_V = E_\mu + E_{\text{visible}}(\text{hadron}) + E_{\text{missing}}$$

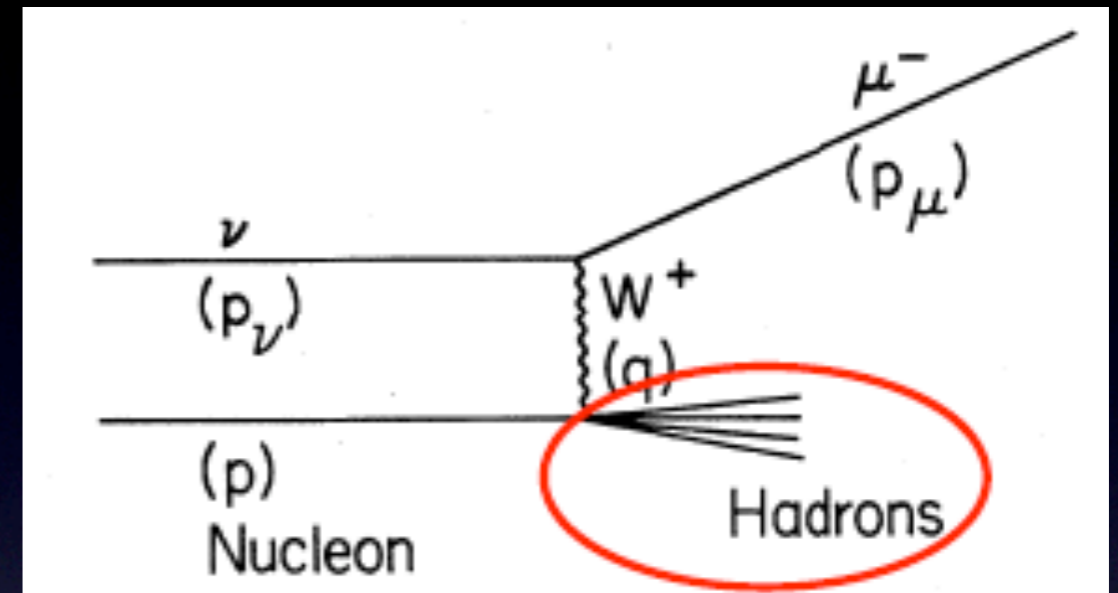
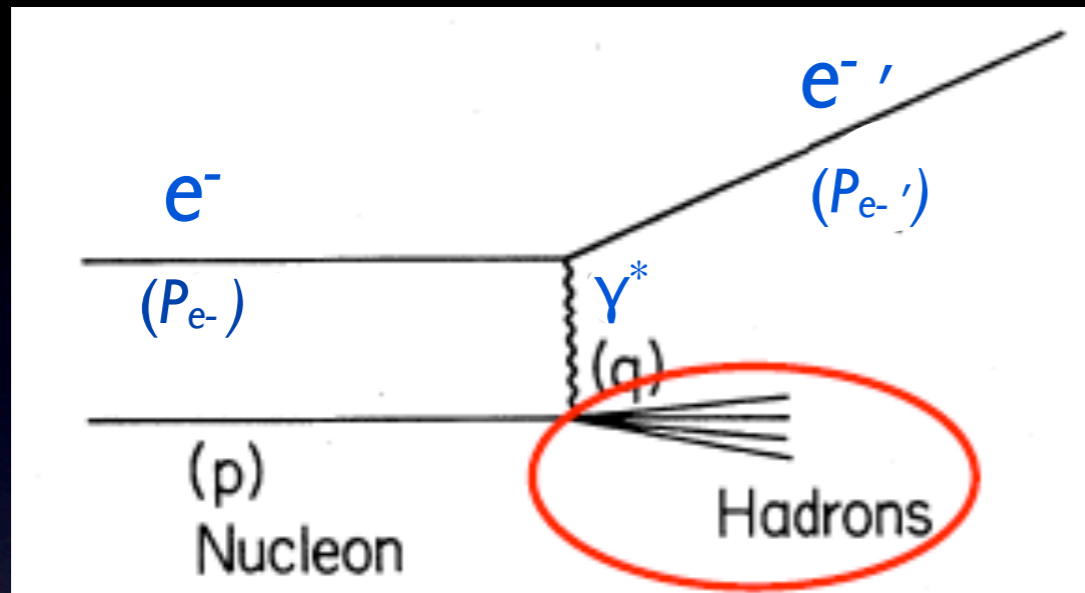


Total uncertainty on the shower energy scale from all the sources of MC parameters (solid black). Illustrated separately are contributions from hadronization model (red solid) uncertainties, formation time (dashed black), external data sources(dashed red), model assumptions (blues).

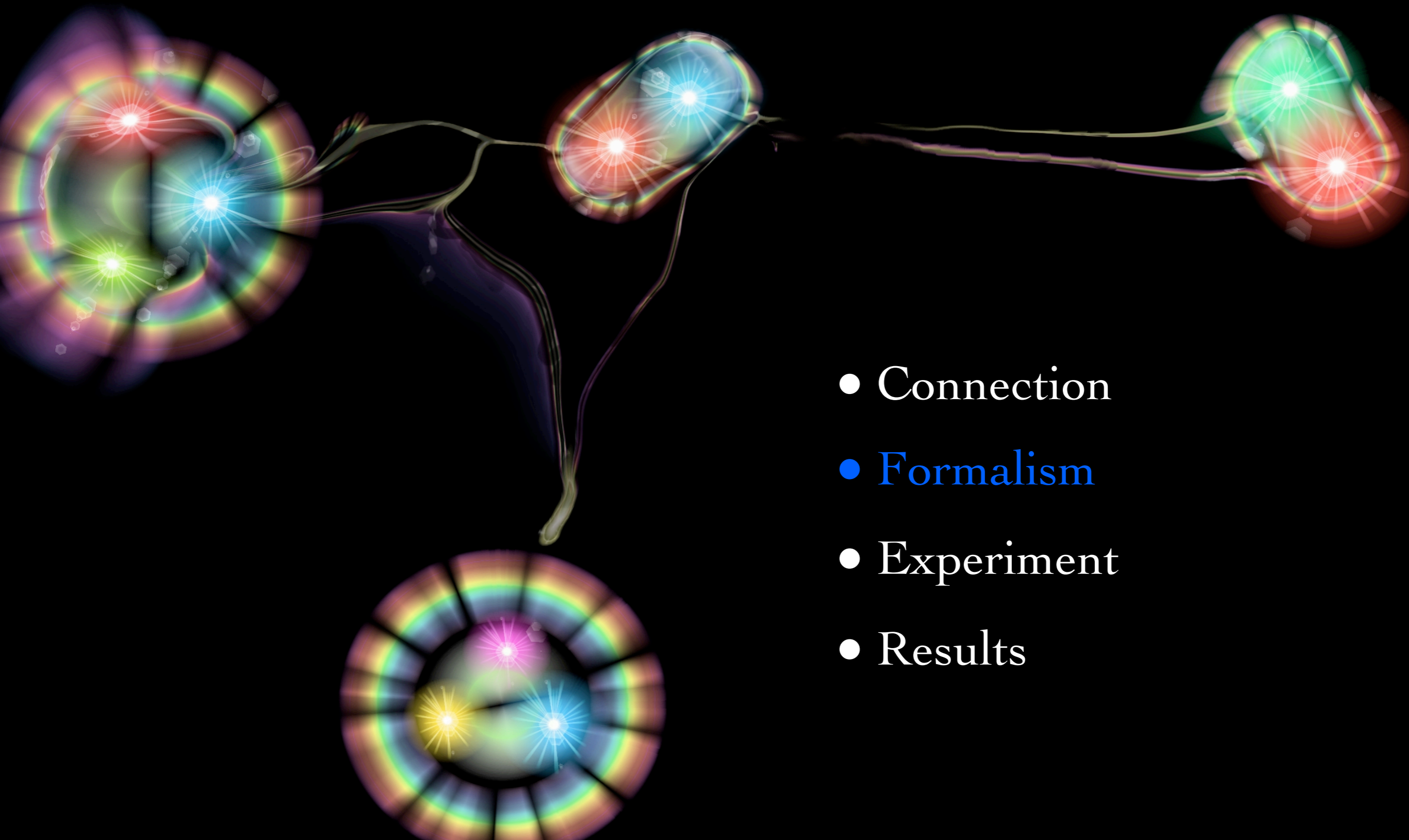
- **Hadronization model** determines set of particles produced from DIS event. modified KNO, JETSET, and linear extrapolation between ( $2.3 < W < 3$  GeV/c)
- **Hadron formation timescale** time to form a hadron during which its interaction cross section in the nucleus is set to zero
- **Intranuclear rescattering** model - how this set of hadrons is altered as by FSI as they exit the target nucleus.
  - $\sigma(h)$  modification on bound N, *etc*
  - $\pi$  absorption/rescattering
  - low momenta secondary nucleon production

S.Dytman, H.Gallagher, M.Kordosky Hadronic shower scale uncertainties in the MINOS experiment, arXiv:0806.2119

# Connection to electron scattering

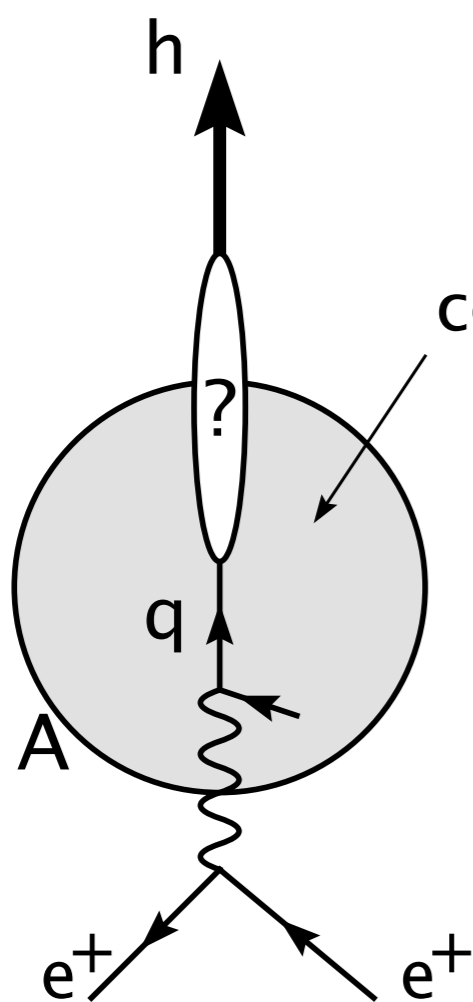


- fragments of the struck parton in  $\nu$  and  $e^-$  scattering should suffer the same FSI: at large ( $W > 2 \text{ GeV}/c^2$ ) invariant masses hadronic systems produced in  $\nu$  systems are similar to those produced in charged lepton and hadron scattering
- neutrino experiment can greatly benefit by tuning their MC based on the improved understanding of parton dynamics and hadronization mechanism constrained by multidimensional data on electroproduction

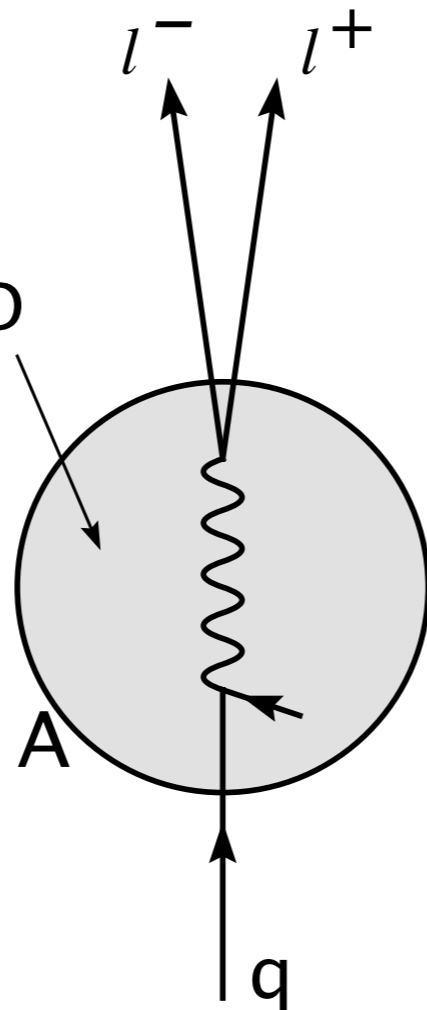


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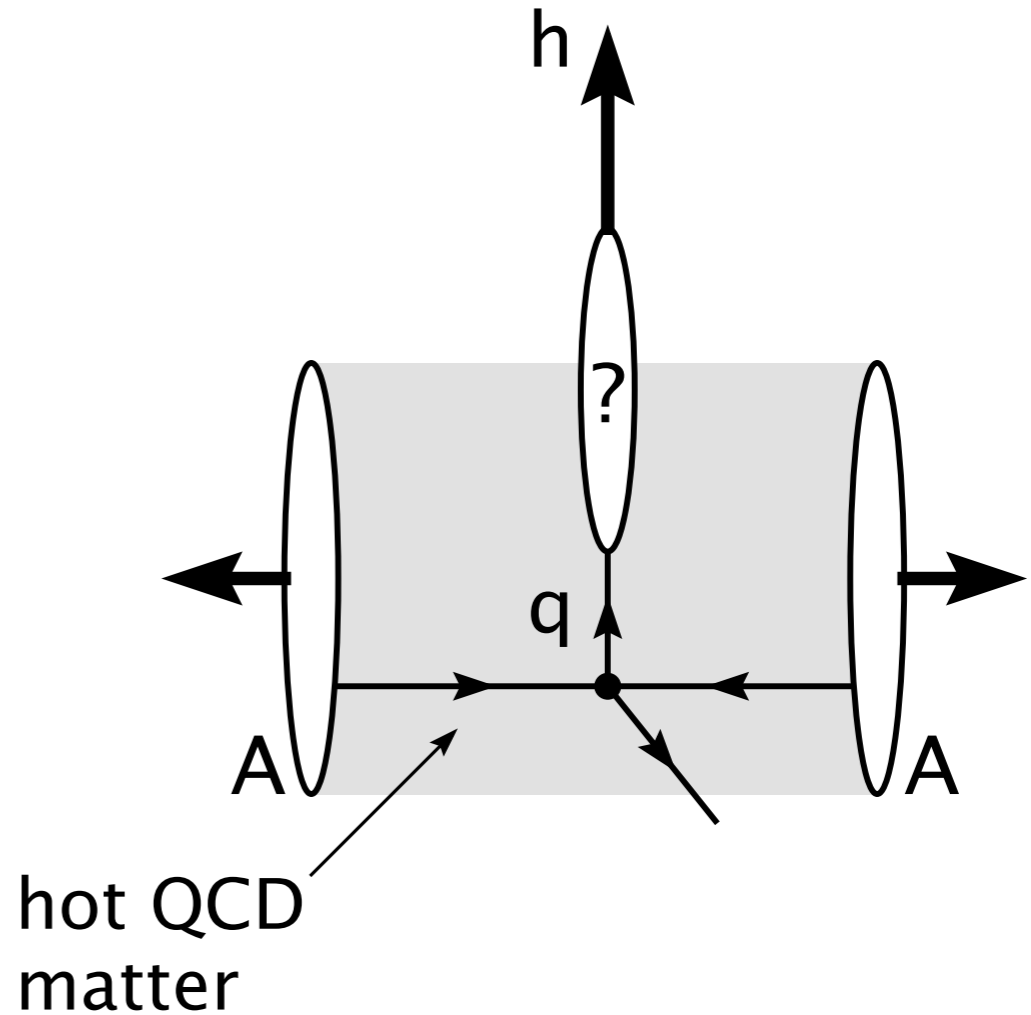
# Reactions



**DIS**  
(DESY, Jefferson Lab)



**Drell-Yan**  
(Fermilab, CERN)

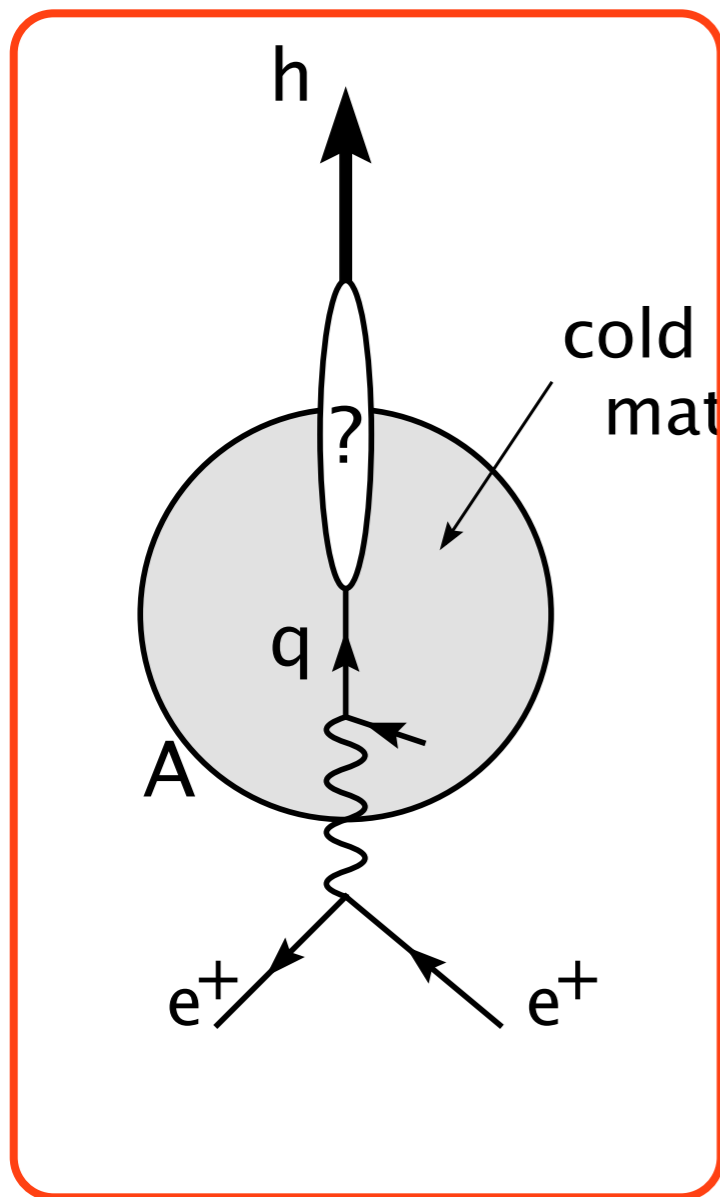


**Heavy-Ion**  
(RHIC, LHC)

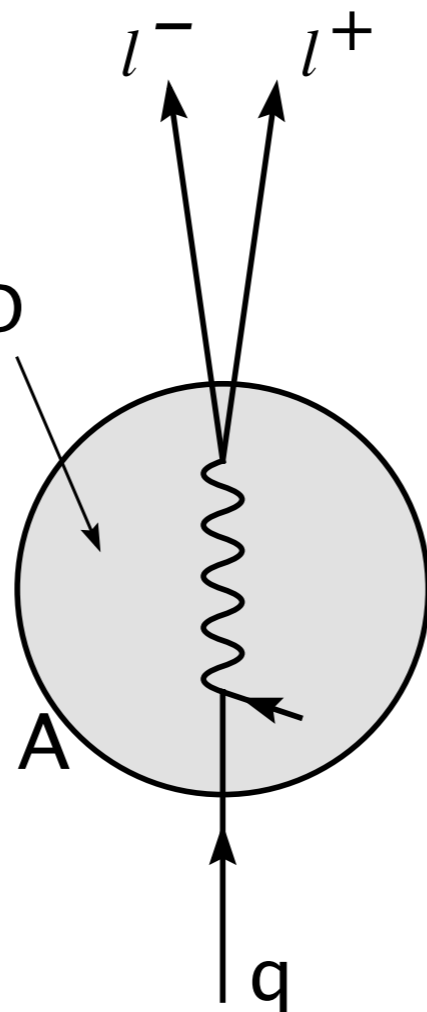
Parton Propagation and Fragmentation in QCD Matter, [A.Accardi](#), [F.Arleo](#), [W.K. Brooks](#), [D. D'Enterria](#), [V.Muccifora](#) [arXiv:0907.3534v1](#) [nucl-th]



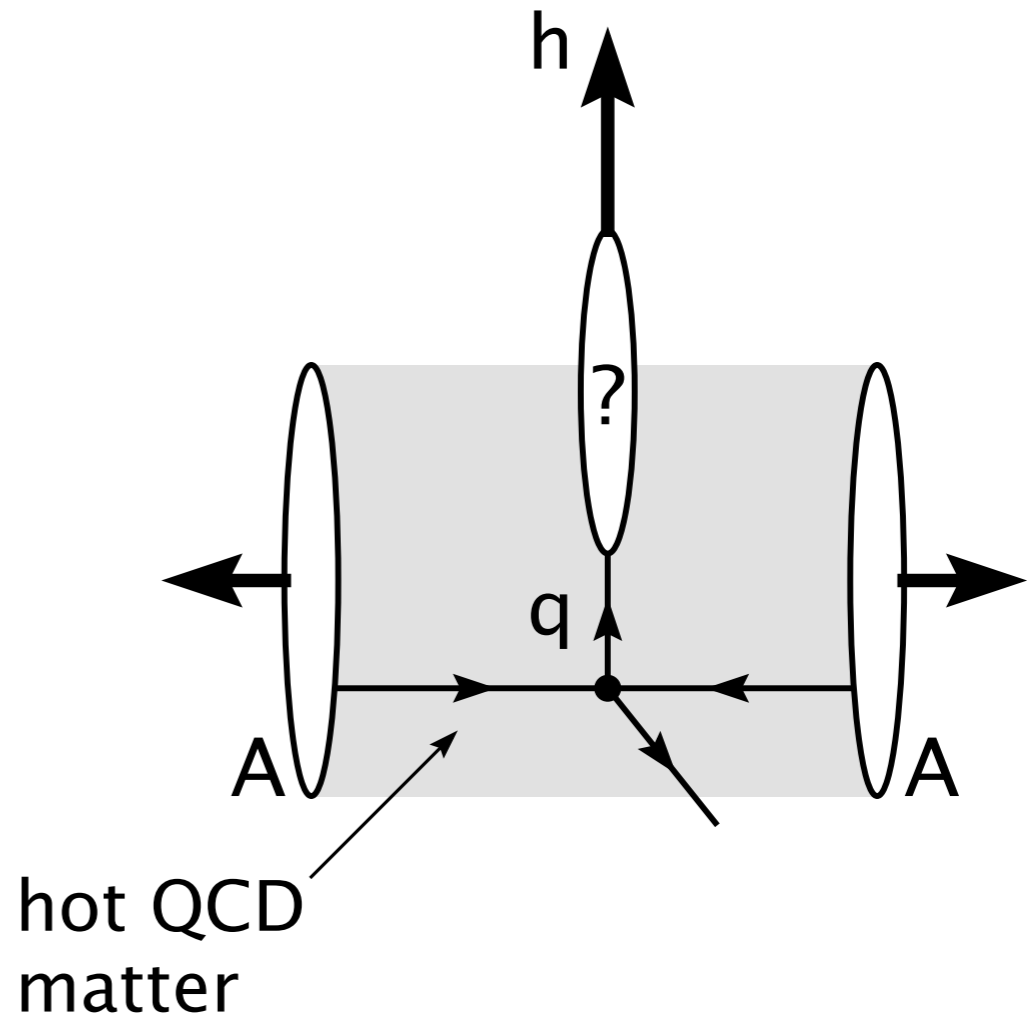
# Reactions



**DIS**  
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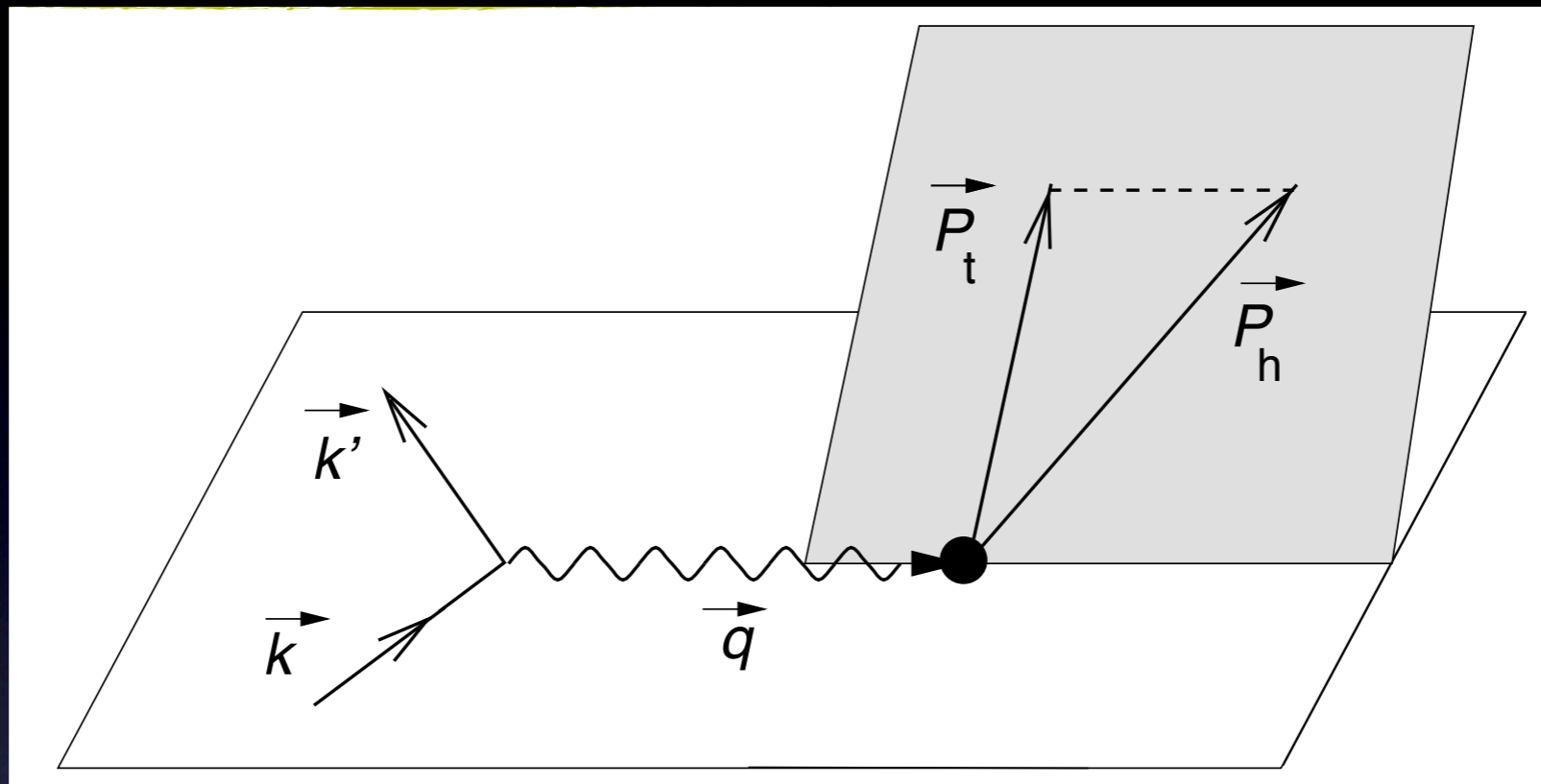
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# SIDIS



$$Q^2 = -q^2$$

four-momentum transferred by the electron;

$$\nu = E - E'$$

energy transferred by the electron, = Initial energy of struck quark;

$$z = E_h / \nu$$

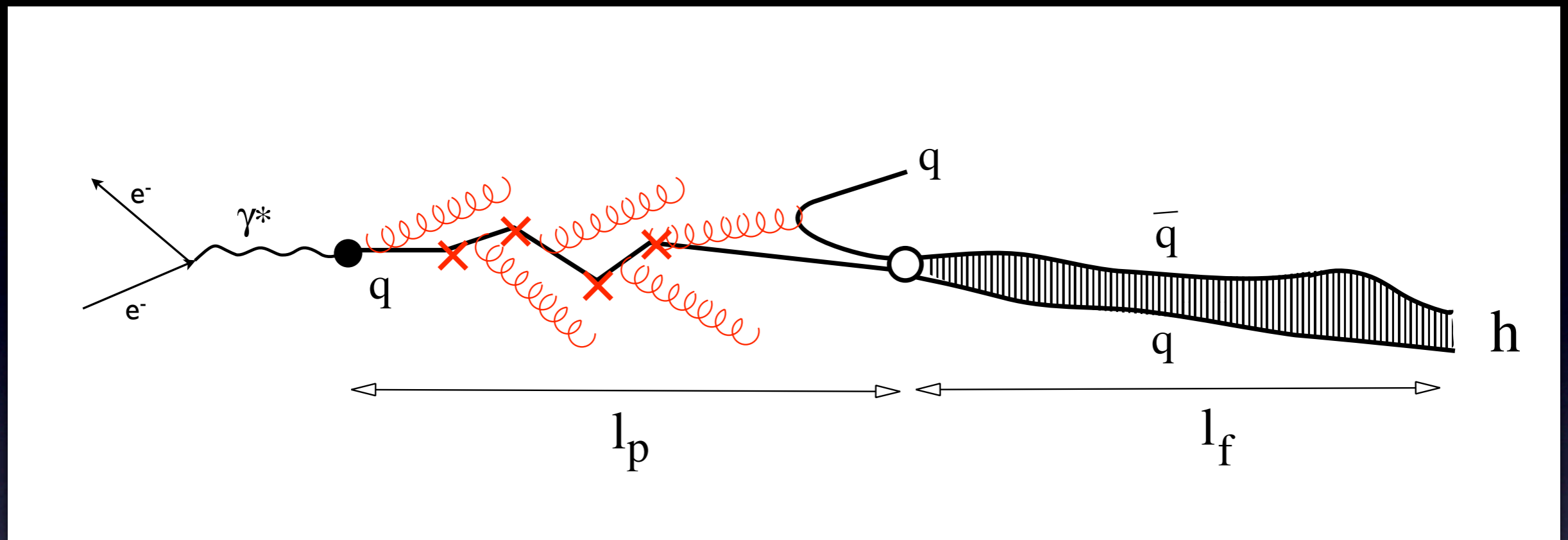
fraction of the struck quark's initial energy that is carried by hadron;

$P_T$

hadron momentum transverse to virtual photon direction;

CLAS kinematics:  $Q^2 > 1 \text{ GeV}^2$ ;  $W > 2 \text{ GeV}$ ;  $0.1 < x_B < 0.55$ ;

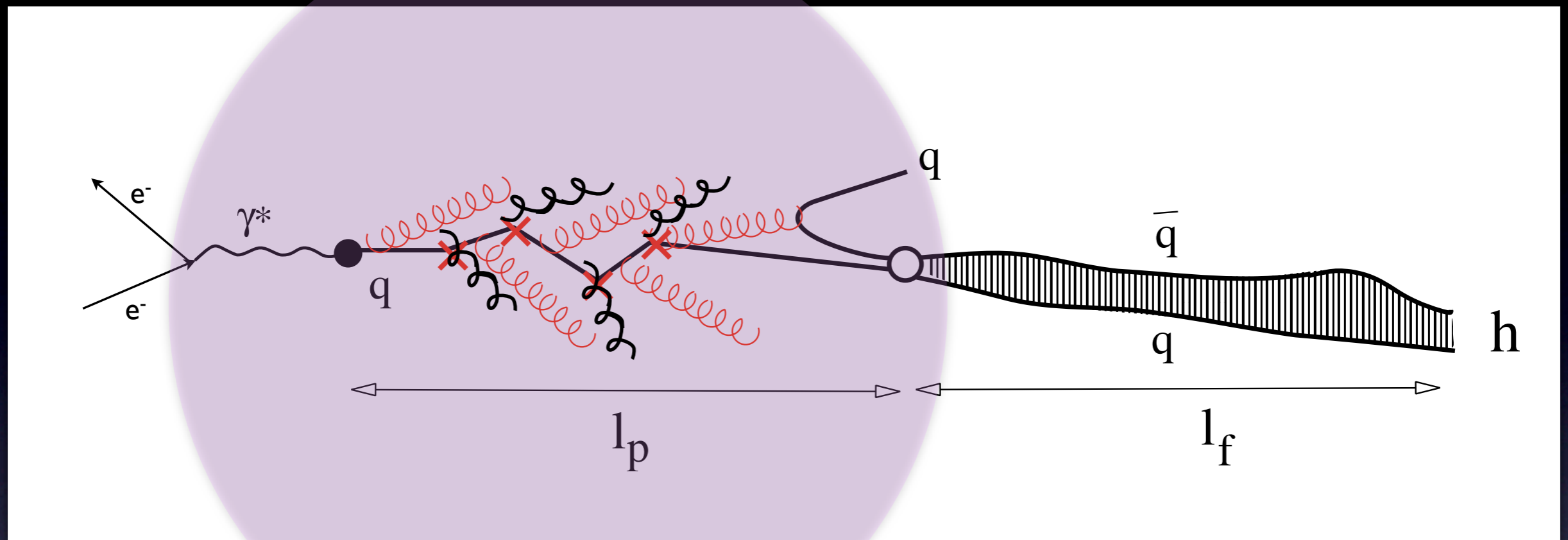
# Hadronization in vacuum



Time Evolution of Jets and perturbative color neutralization B.Z.Kopeliovich, J.Nemchik, I.Schmidt Nucl. Phys A 782 (2007)

- Access to characteristic timescales of parton propagation and hadronization:  
 how long quark remains deconfined?  $\rightarrow$  production time  $\tau_p$   
 how long it takes to form full hadronic wave function?  $\rightarrow$  formation time  $\tau_f$

# Hadronization in medium



- Access to characteristic timescales of parton propagation and hadronization:
  - how long quark remains deconfined?  $\rightarrow$  production time  $\tau_p$
  - how long it takes to form full hadronic wave function?  $\rightarrow$  formation time  $\tau_f$
- Nuclear medium acts as spacial analyzer: the yields and kinematics of detected hadrons are modified by varying the nuclear size which accesses different time scales

# Observables (1)

## Transverse momentum broadening

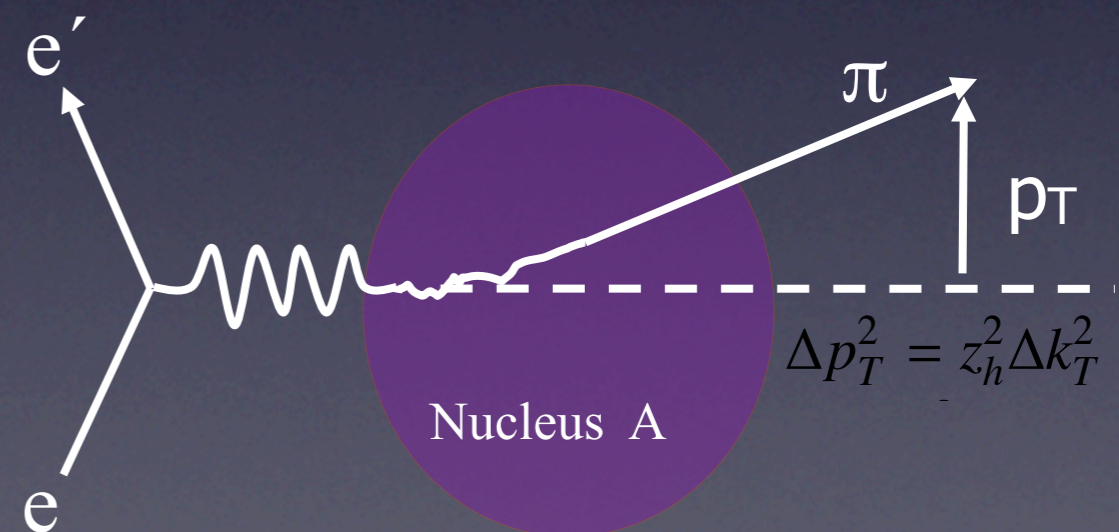
defined with respect to the  $\gamma^*$  direction

$$\Delta p_T^2 = \langle p_T^2 \rangle_A - \langle p_T^2 \rangle_D$$

- propagation of a parton through the medium is accompanied by gluon bremsstrahlung (radiative energy losses) and multiple scattering (collisional energy losses) leading to accumulation of quark transverse momenta.

- *how long can a light quark remain deconfined ?*  
*how much energy it losses ?*

The production time  $\tau_p$  access this via shape and magnitude of  $A$ -dependence.



# Observables

## Hadronic multiplicity ratio

normalized change in production yields of hadrons in nuclei A to D

$$R_A^h(\nu, Q^2, z, p_T, \phi) = \frac{\left. \frac{N_h(\nu, Q^2, z, p_T, \phi)}{N_e(\nu, Q^2)} \right|_{\text{DIS}}}{\left. \frac{N_h(\nu, Q^2, z, p_T, \phi)}{N_e(\nu, Q^2)} \right|_{\text{D}}} \Bigg|_A$$

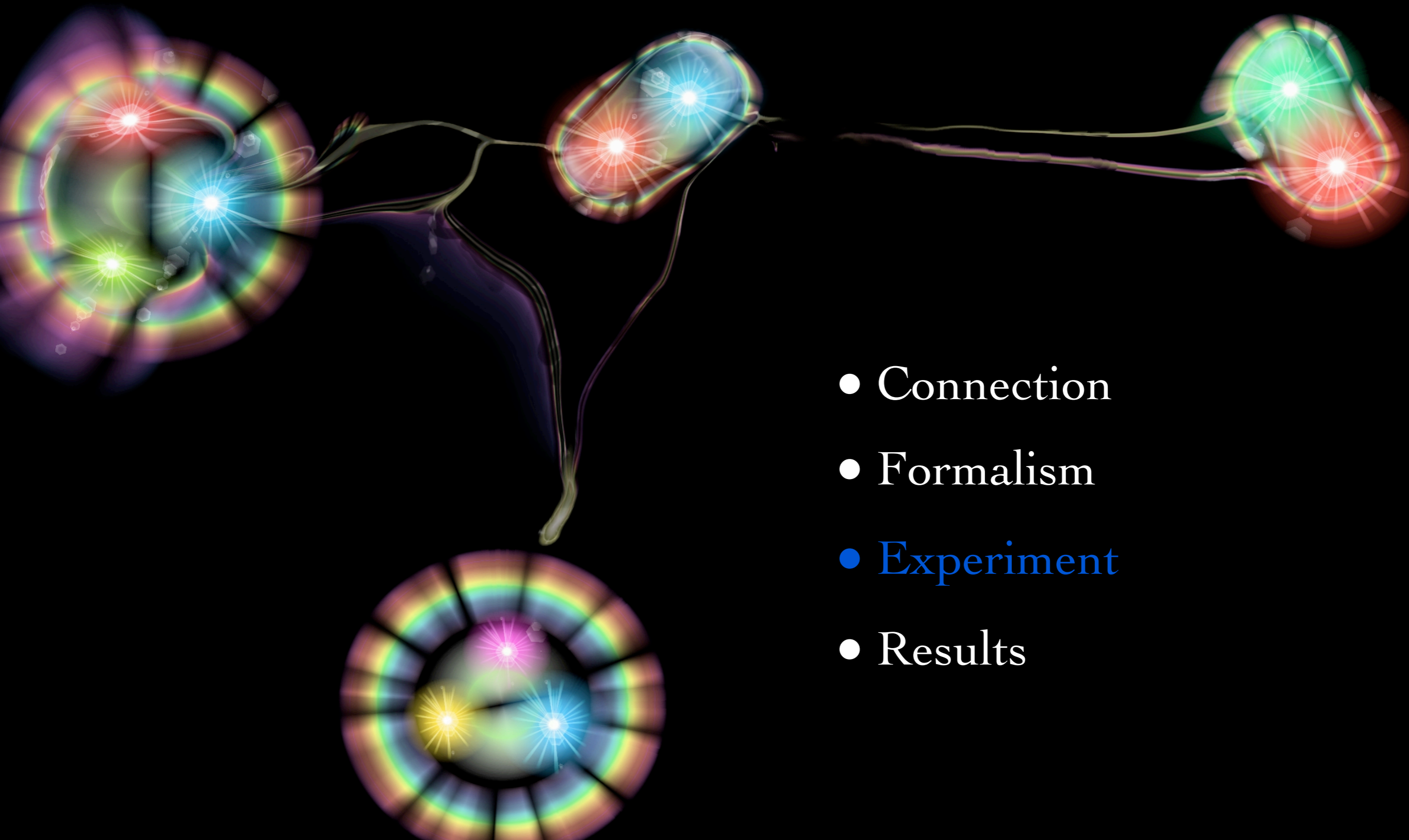
- colorless pre-hadron is created after quark color neutralization
- pre-hadron interaction is dominantly inelastic leading to attenuation
- *how long does it take to form the color field of the hadron ?*

The formation time  $\tau_f$  measures this via  $R_h(Q^2, U, p_T, z_h)$

# Glimpse at available models

Authors	References	Model Description	Multiplicity Results	$p_{\perp}^2$ Broadening Results
Accardi <i>et al.</i>	[Accardi 2003]	$Q^2$ rescaling of FF + hadron absorption	Few comparisons with HERMES & EMC	None
Arleo	[Arleo 2003b]	BDMPS based parton energy loss (quenching weight calculation)	Scarce comparison with HERMES	$\hat{q} = 0.75 \text{ GeV}/\text{fm}^2$ too large for HERMES
Deng <i>et al.</i>	[Deng 2010] [Deng 2011]	Modified DGLAP evolution	Few comparisons with HERMES	$\hat{q} = 0.015 \text{ GeV}^2/\text{fm}$ coherent with HERMES
Falter <i>et al.</i> (GiBUU)	[Falter 2004] [Gallmeister 2005] [Gallmeister 2008]	Pure hadron/prehadron absorption	Extensive comparison with HERMES & EMC	None
Gyulassy and Plümer	[Gyulassy 1990]	Medium modified FF using string-flip model	Comparison with old data (EMC & SLAC)	None
Kopeliovich <i>et al.</i>	[Kopeliovich 2004] [Domdey 2009] [Ciofi degli Atti 2005]	$Q^2$ rescaling of FF, energy loss and prehadron absorption	Extensive comparison with HERMES & EMC	Extensive comparison with HERMES
Salgado and Wiedemann	[Salgado 2002] [Salgado 2003]	BDMPS based parton energy loss (quenching weight calculation)	Few comparisons with HERMES	Extensive comparison with HERMES
Wang and Wang	[Wang 2002]	Pure parton energy loss	Few comparisons with HERMES	None

Based on HERMES data, models do not discriminate between two proposed mechanisms: partonic energy losses vs prehadron absorption. Need higher precision data.

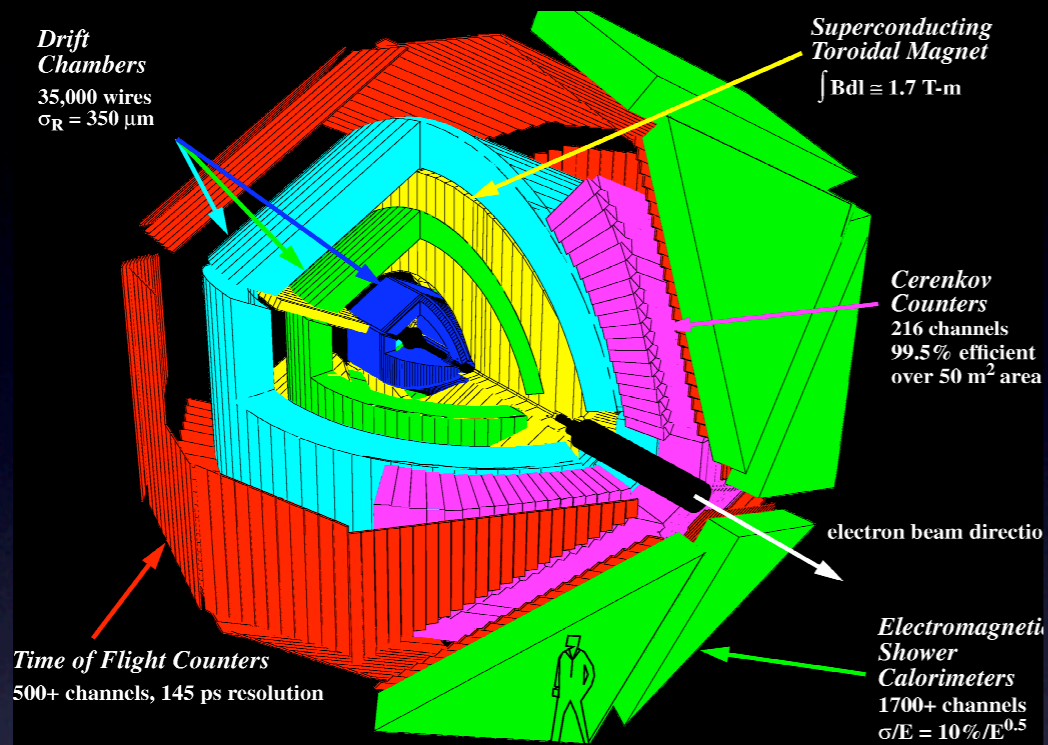


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# Experiment

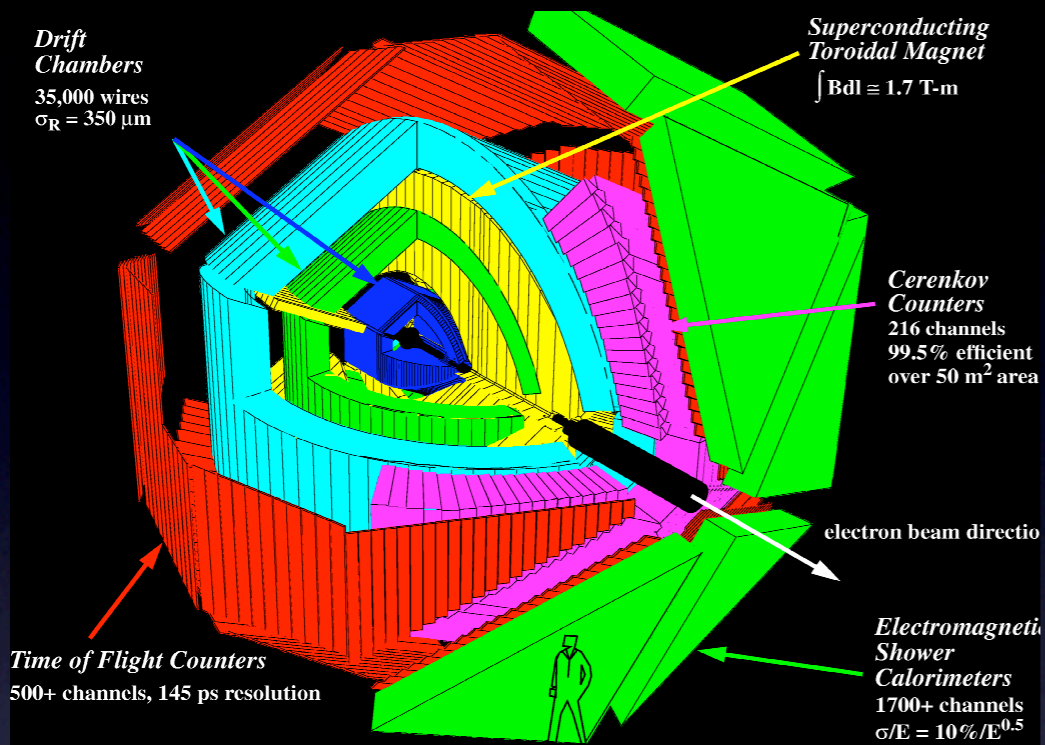
## CEBAF Large Acceptance Spectrometer at Jefferson Lab



- Charged particles polar angles  $8^\circ$ - $144^\circ$   
Momentum resolution 1% down to 0.2 GeV/c
- Neutral particles polar angles  $8^\circ$ - $70^\circ$   
Momentum resolution  $10\%/\sqrt{E}$  down
- Identification of  $e^+/e^-$ ,  $\gamma$ , p, n,  $\pi^\pm$ ,  $K^\pm$

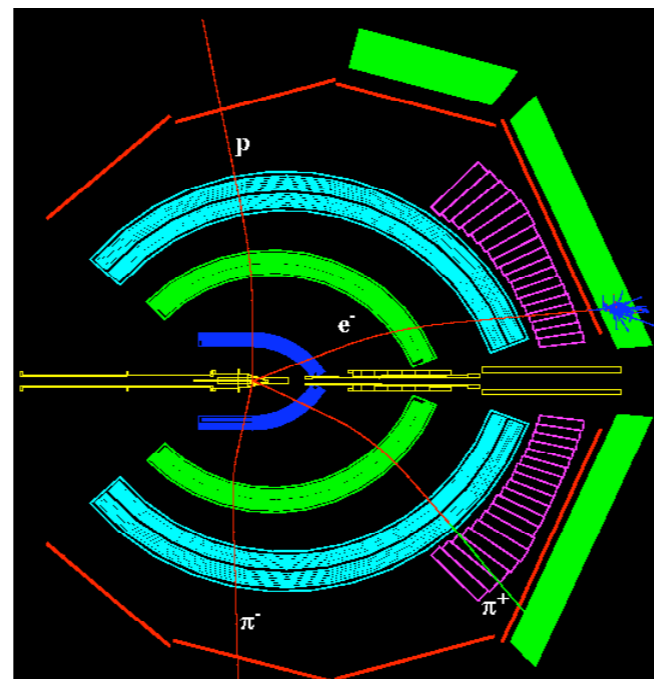
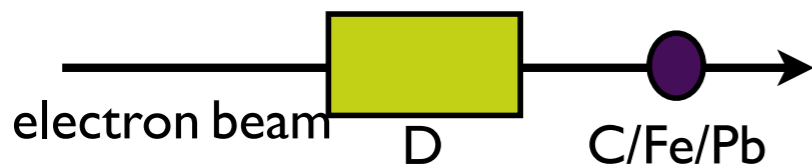
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Momentum resolution  $10\%/\sqrt{E}$  down
- Identification of  $e^+/e^-$ ,  $\gamma$ ,  $p$ ,  $n$ ,  $\pi^\pm$ ,  $K^\pm$

Two targets in the beam simultaneously



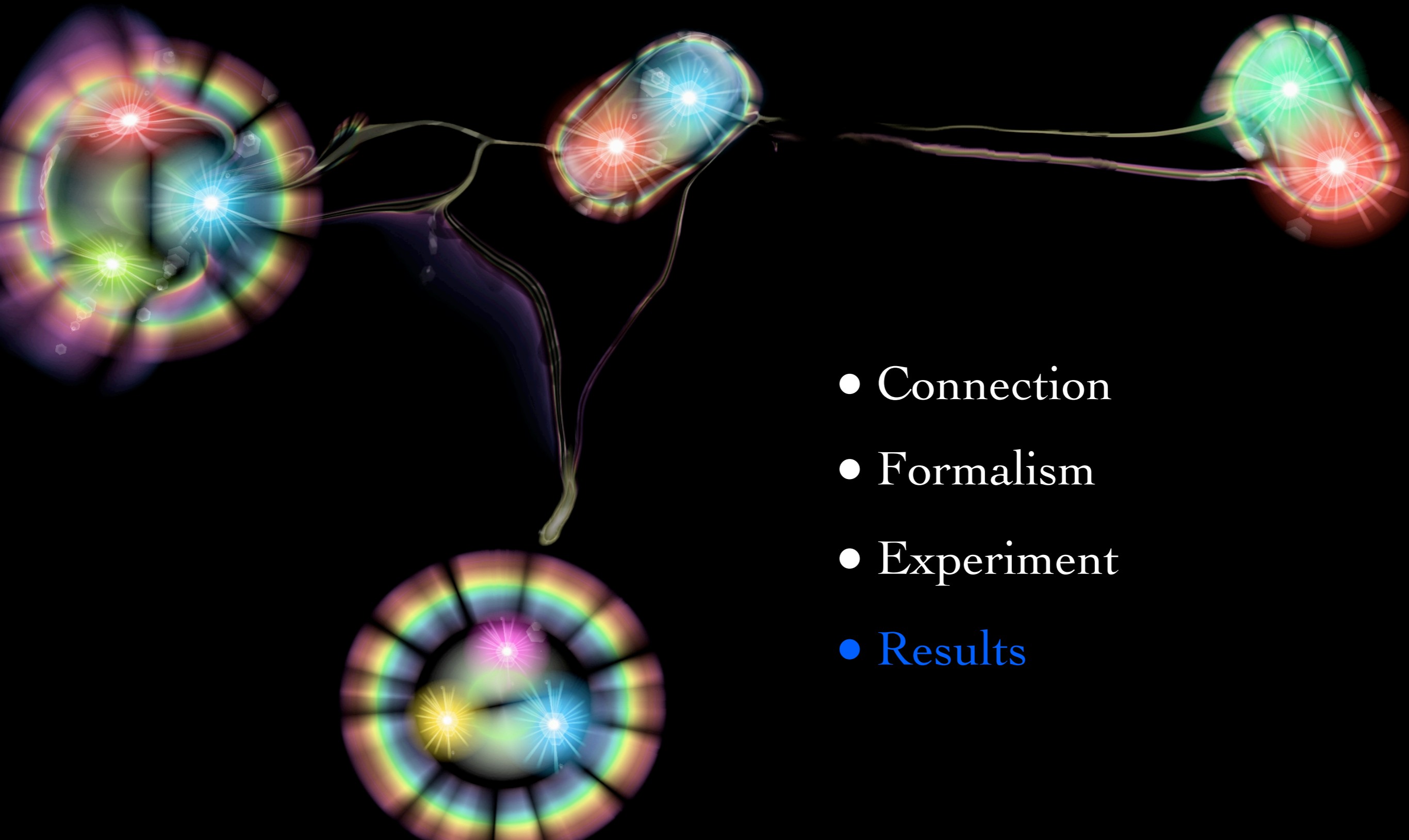
### CLAS EG2

- Electron beam 5.014 GeV
- Targets  $^2\text{H}$ ,  $^{12}\text{C}$ ,  $^{56}\text{Fe}$ ,  $^{207}\text{Pb}$
- Luminosity  $2 \cdot 10^{34} \text{ 1/(s} \cdot \text{cm}^2)$

Statistics:

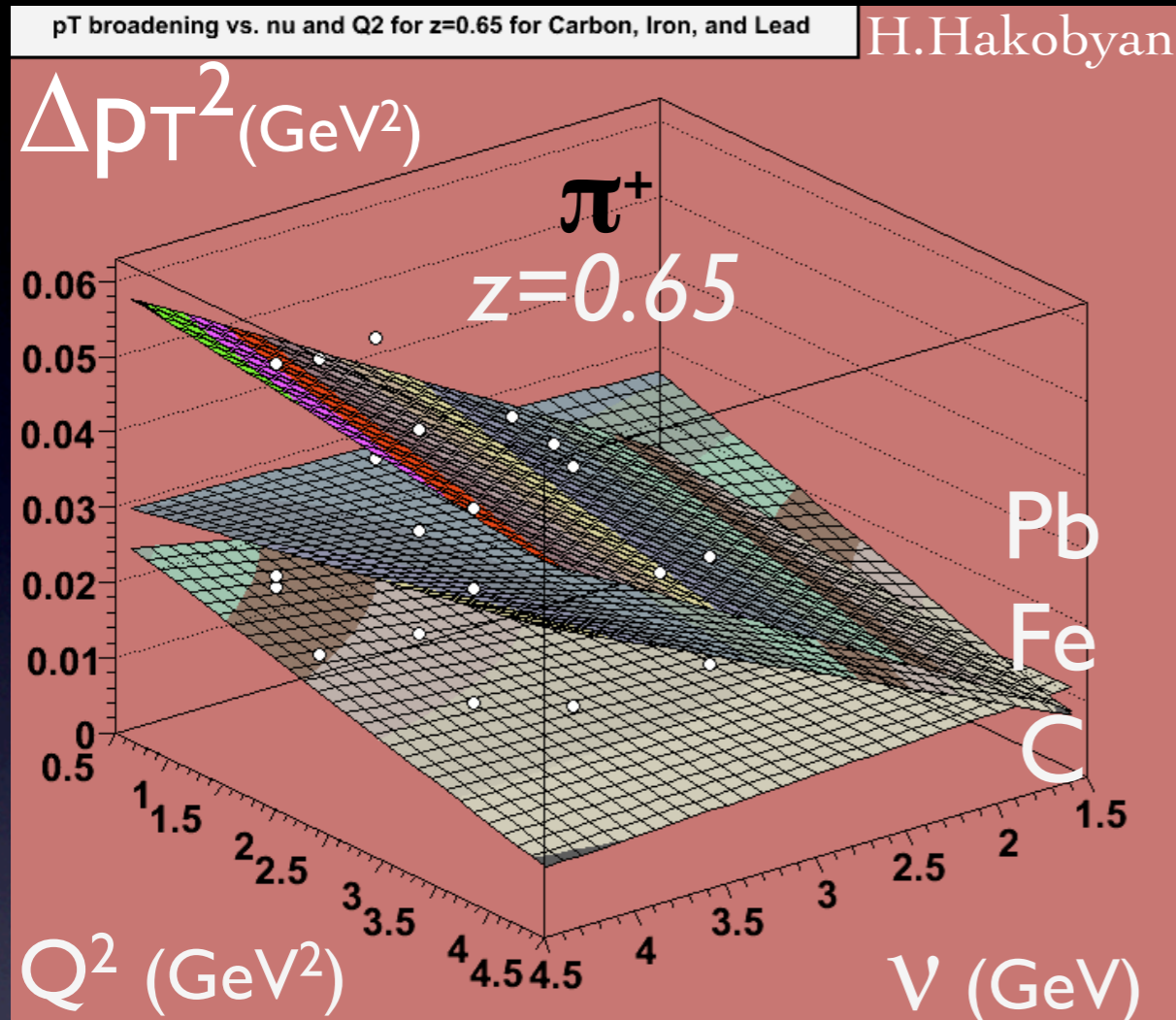
130M DIS electrons

6.6 M  $\pi^+$ , 2.8 M  $\pi^-$ , 2.0 M  $\pi^0$

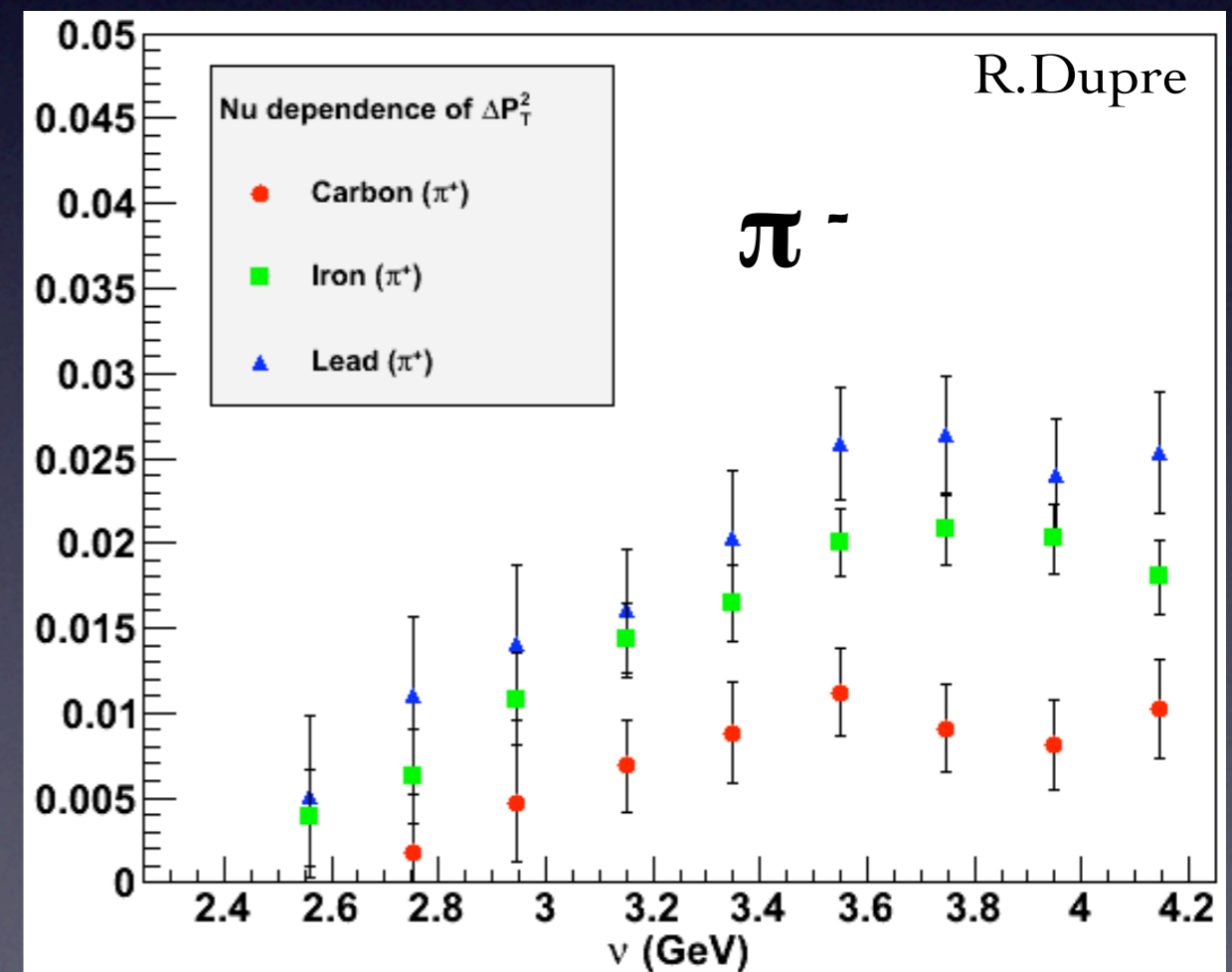


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# Transverse momentum broadening | $\pi^\pm$



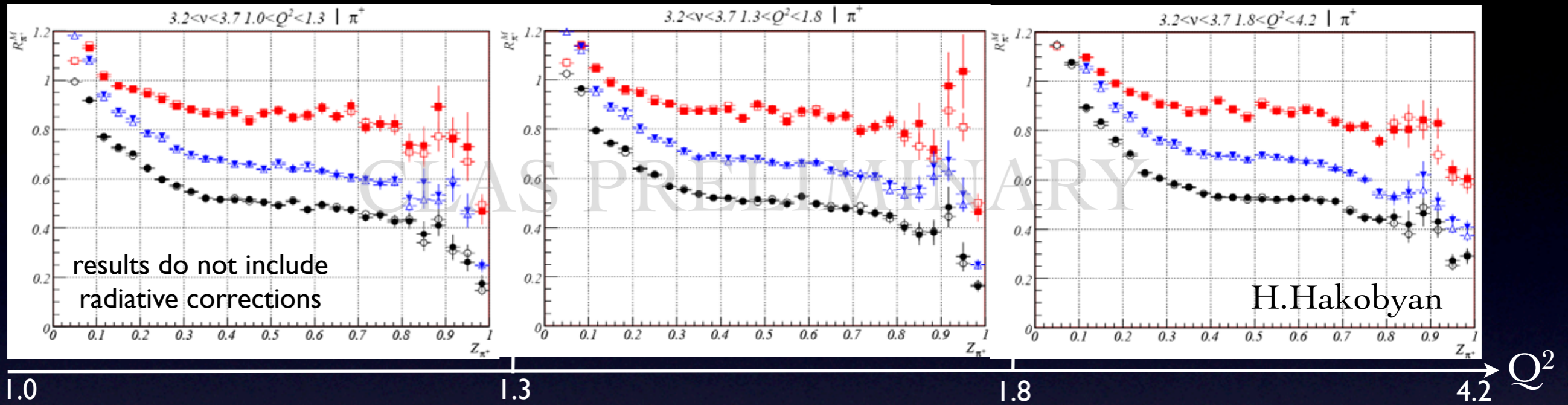
$\pi^+$  :27 bins in  $\nu$ ,  $Q^2$ ,  $z$  for each of 3 targets  
 Radiative corrections for  $\Delta p_T^2$  ( $\pi^+$ ) ~10%  
 Radiative corrections for  $\Delta p_T^2$  ( $\pi^-$ ) are within systematic uncertainties.



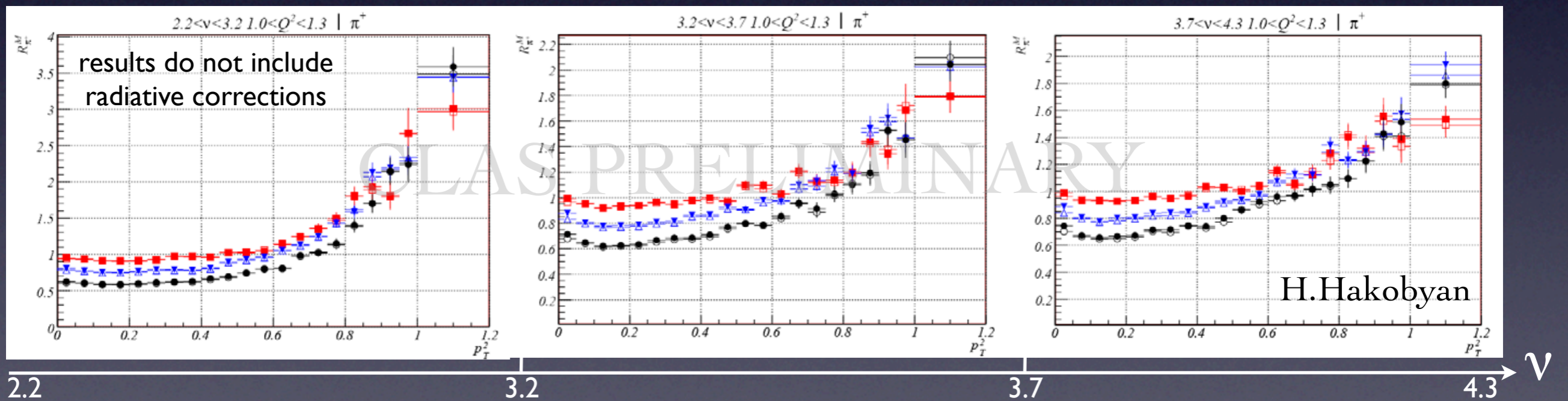
Assuming classical geometrical model:  
 $\Delta p_T^2 \propto l_p$  and  $l_p < R_{Pb}$  (6.5fm)  
 3(4) parameter fit to  $\Delta p_T^2$  on  $\pi^+$  gives  
 range for  $\tau_p$  to be 1.6-2.0 fm

# Multiplicities | $\pi^+$

$R_{\pi^+}$  slice in 3D set of bins ( $Q^2, 3.2 < \nu < 3.7, z$ ) integrated over  $p_T^2$



$R_{\pi^+}$  slice in 3D set of bins ( $1.0 < Q^2 < 1.3, \nu, p_T^2$ ) for  $0.4 < z < 0.7$

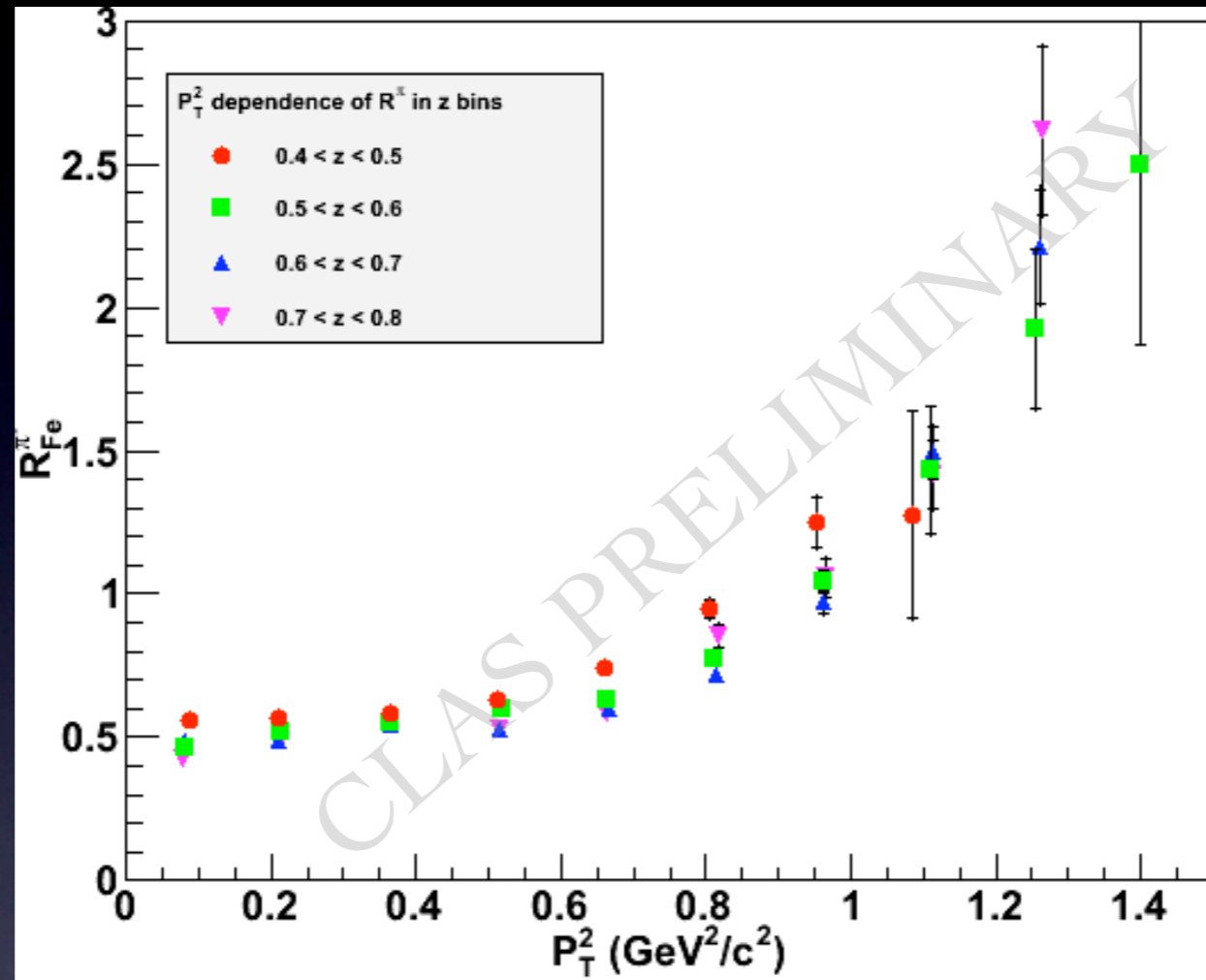


- Attenuation systematically increases for the larger nuclei
- Attenuation of high  $z$  hadrons, enhancement for high  $p_T^2$

# Multiplicities | $\pi^-$

$R_{\pi^-}$  in 2D set in  $(z, p_T^2)$

R.Dupre



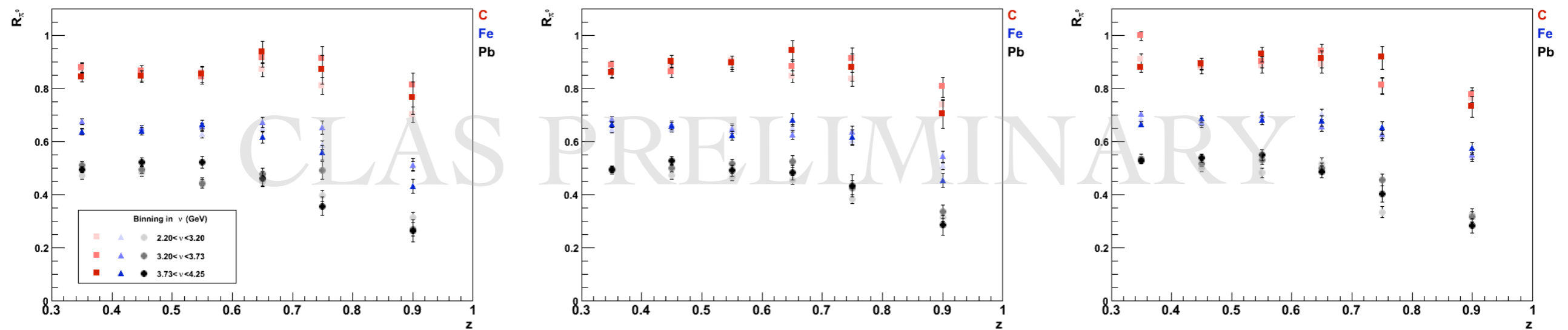
- Attenuation systematically increases for the larger nuclei
- Attenuation of high  $z$  hadrons, enhancement for high  $p_T^2$



# Multiplicities | $\pi^0$

$R_{\pi^0}$  in 3D set of  $(Q^2, \nu, z)$  integrated over  $pT^2$

T.Mineeva



- Attenuation systematically increases for the larger nuclei
- Attenuation of high  $z$  hadrons, enhancement for high  $pT^2$
- Small dependence on  $Q^2$  and  $\nu$



# Future program

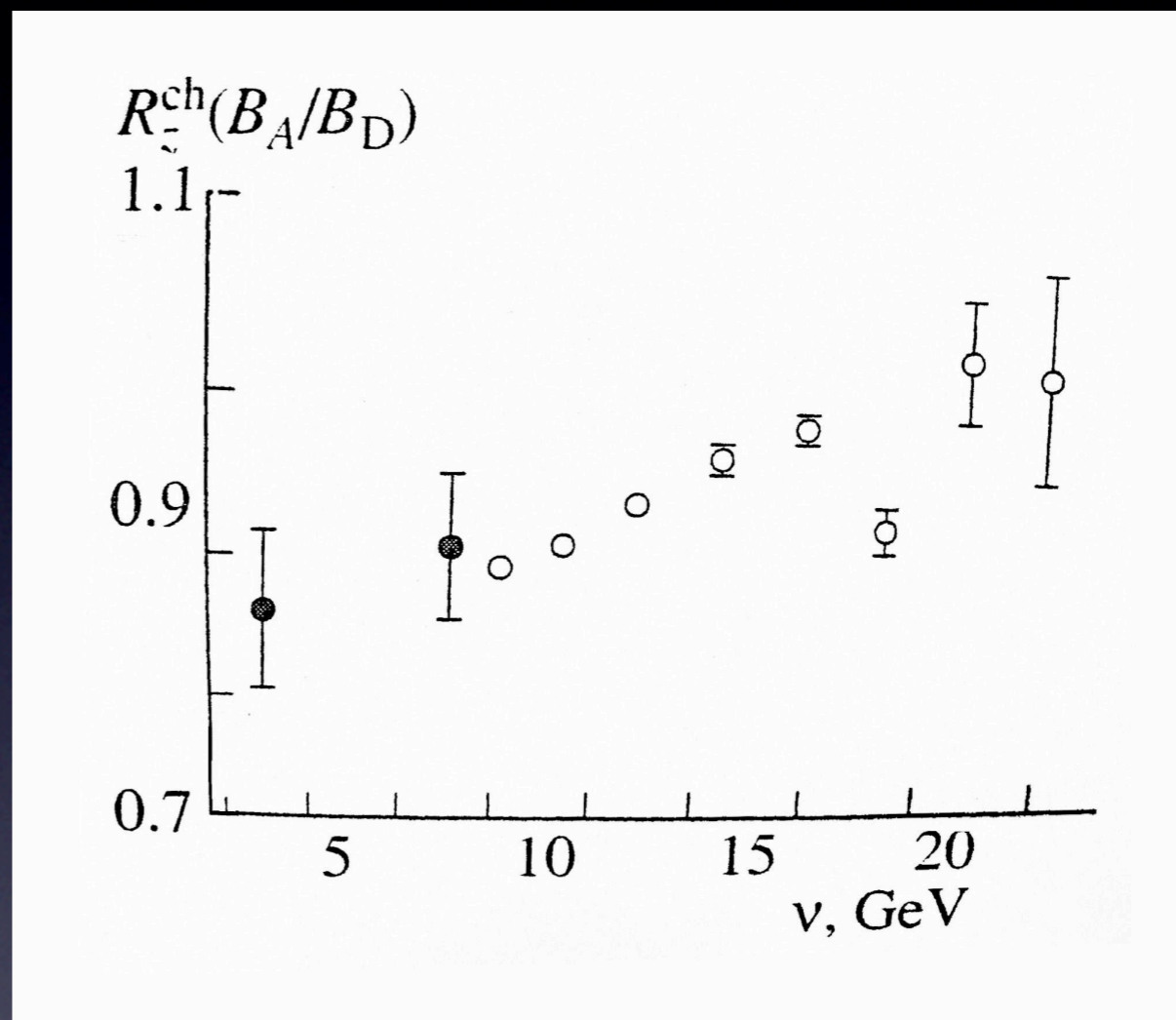
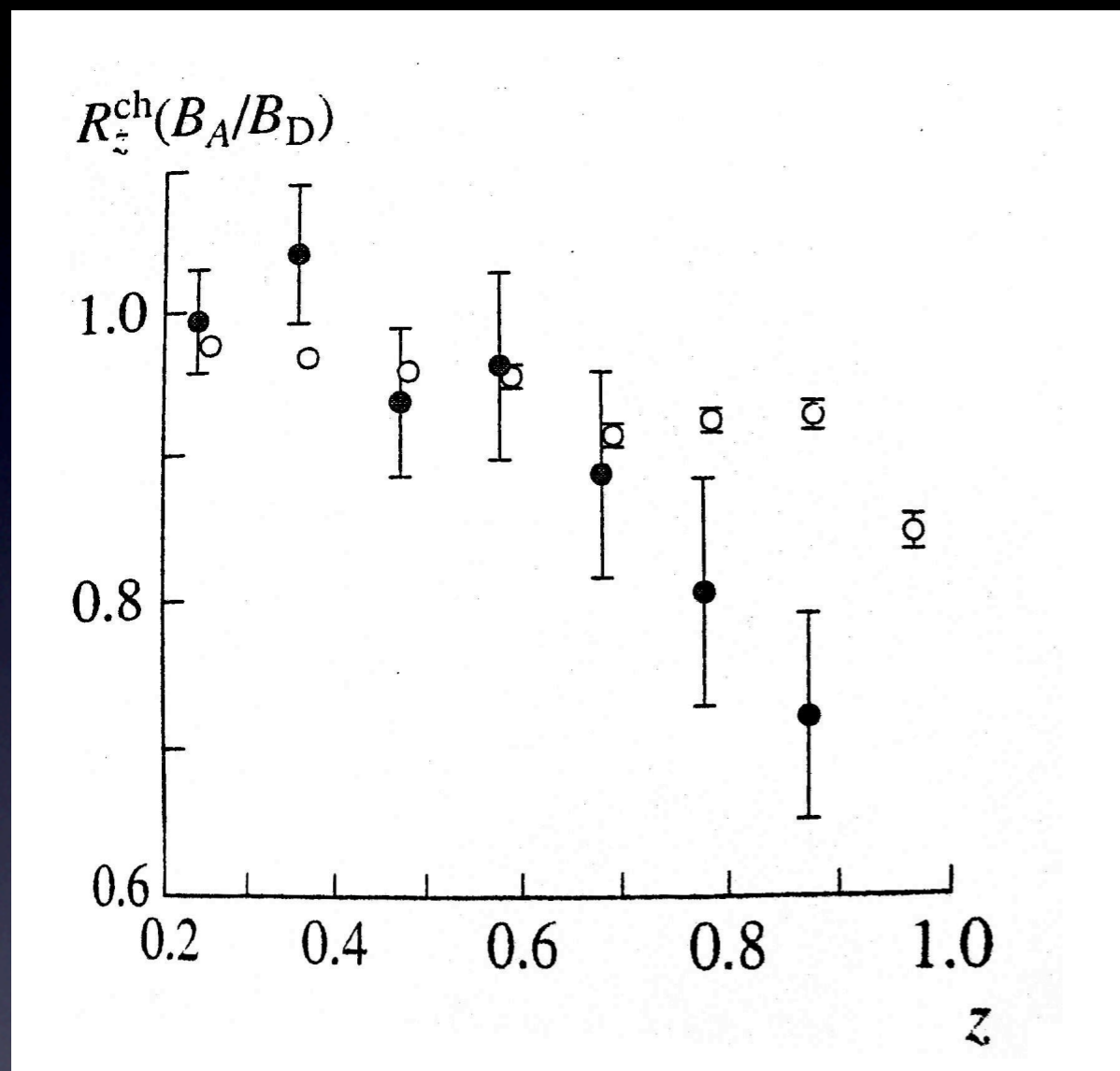
- Extraction of timescales and quark energy losses based on the finalized data analysis
- Measurement from existing EG2 set of  $\pi^0$  vs  $\eta$  suppression (not previously accessed in cold matter), extraction of transverse momentum broadening for both mesons.
- The program to pursue hadronization studies at CLAS12 (E12-06-117) with 11 GeV electron beam has been approved for 120 beam days. It will provide by far the best experimental access to medium-simulated quark energy loss, and enable extraction of 4D multiplicities for wide range of hadrons.
- EIC will offer high energy eA collisions with  $E_{e^-} = 11$  GeV/c and unpolarized heavy nuclei  $E_A = 12-40$  GeV/c per nucleon for  $A > 200$  (Au, Pb).  
Long parton life time, direct access to pQCD  $E_{\text{loss}}$ .

# Summary

- In order to decrease systematic uncertainties from, for example, incoming  $\nu_\mu$  flux, difference in  $E_{\text{visible}}$  and  $E_\nu$ , and NC  $\pi^0$  production neutrino experiment can greatly benefit from improved knowledge on hadroproduction by tuning their MC
- CLAS data set on electroproduction off nuclei is consistent with previous measurements (HERMES) meanwhile providing an access to multidimensional extraction of characteristic time scales of hadronization and quark energy losses
- Program of hadronization studies will continue with CEBAF upgrade @ 11 GeV.

# *Backup*

# Comparison $\nu$ and $e^+$ scattering off nuclei



Ratio of hadrons produced in DIS from  $\nu$  scattering (open circles) on  $A=28$  normalized to D at  $2 < \nu < 15$  GeV (SKAT) compared with  $e^+$  scattered off  $A=14$  in  $7 < \nu < 24$  GeV (HERMES).

Neutrino data show the same trend as observed with electron beam.

# Access to formation time $\tau_f$

## Model dependences on formation time extraction

Model Approach	Formula
single time-scale model	$\tau = \tau_q \frac{E_{quark}}{m_{quark}} \sim \nu$
time dilation of intrinsic size	$\tau = \tau_0 \frac{E_{hadron}}{m_{hadron}} = \tau_0 \frac{z\nu}{m_{hadron}}$
two-timescale parameterization based on the Lund model	$\tau = (1 - \ln(z)) \frac{z\nu}{\kappa c}$
QCD-based model	$\tau \approx \frac{1}{c} \frac{r_{hadron}}{k_{hadron}} E_h \sim z\nu$
gluon-bremsstrahlung model	$\tau = c_{hadron} (1 - z) \nu$

# BACK-OF-ENVELOPE - $\tau_p$

$$t_p = \frac{v}{\left. \frac{dE}{dx} \right|_{\text{vacuum}}} (1 - z_h)$$

Energy  
conservation,  
time dialation

String model

$$\left. \frac{dE}{dx} \right|_{\text{vacuum}} \approx \kappa \approx 1 \text{ GeV} / \text{fm}$$

If take, e.g.,  $z = E_{\text{hadron}}/V = 0.6$ ,  $v = 5 \text{ GeV}$ , then  $t_p \sim 2 \text{ fm}/c$

# BACK-OF-ENVELOPE - $^h\tau_f$

Given hadron of size  $R_h$ , can build color field of hadron in its rest frame in time no less than  $t_0 \sim R_h/c$ . In lab frame this is boosted:

$$t_f \geq \frac{E}{m} R_h$$

If take, e.g., the pion mass, radius 0.66 fm,  
 $E = 4$  GeV, then  $\tau_f \sim 20 \text{ fm}/c$ .





