Robustness of theoretical model predictions

Michael Ratz



NuFact, Virginia, July 27, 2012

Based on:

- T. Araki, T. Kobayashi, J. Kubo, S. Ramos–Sánchez, M.R. & P. Vaudrevange Nucl. Phys. B 805, 124–147 (2008)
- M.-C. Chen, M. Fallbacher, M.R. & C. Staudt, to appear

Origin of mass hierarchies and mixing

ilder Standard model: ${\it O}(20)$ physical quantities which seem unrelated

Origin of mass hierarchies and mixing

- ? Are there testable relations?

Origin of mass hierarchies and mixing

- ? Are there testable relations?
- Symmetries may relate the parameters and thus reduce their number

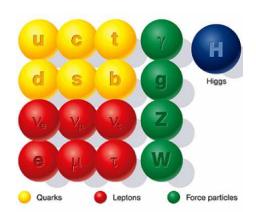
 ${\mathscr G}$ Grand Unified Theories : $G_{\rm SM} = {\rm SU}(3) \times {\rm SU}(2) \times {\rm U}(1) \subset G_{\rm GUT}$

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- $G_{GUT} \in \{SU(5), SO(10), \ldots\}$
- GUT symmetry relates quarks and leptons



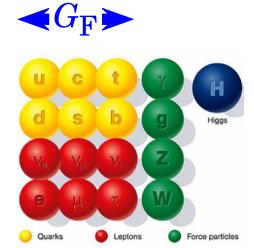


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- Quarks & leptons reside in the same GUT multiplets
- One set of Yukawa coupling for a given GUT multiplet
 intra-family relations

Family symmetries $G_{\rm F}$ relate different families \sim inter-family relations





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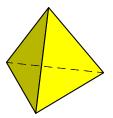
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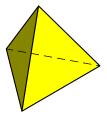


- Recently strong model building activities based on discrete family symmetry groups:
 - A_4 (tetrahedron)

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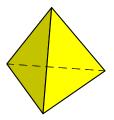
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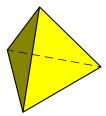




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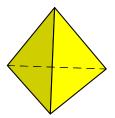
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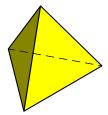




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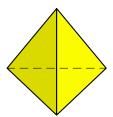
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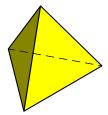




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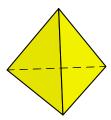
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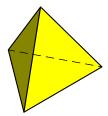






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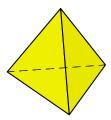


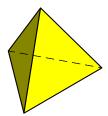




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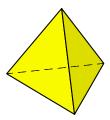


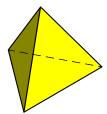




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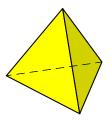






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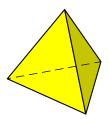
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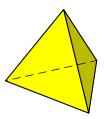
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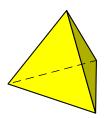
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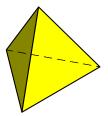




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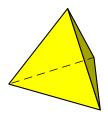
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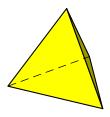
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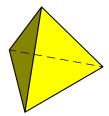




Non-Abelian discrete flavor symmetries

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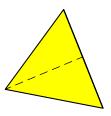


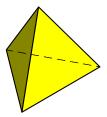




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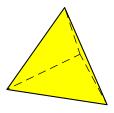
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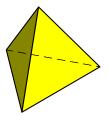






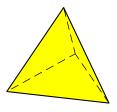
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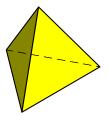




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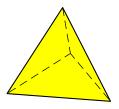
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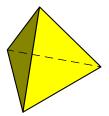






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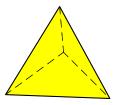


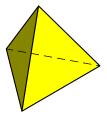




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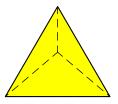
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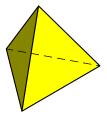




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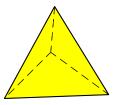
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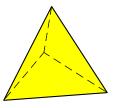
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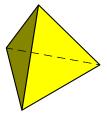






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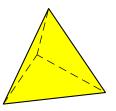


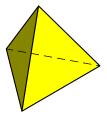
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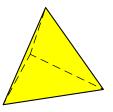


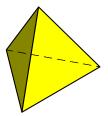




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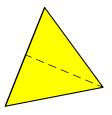
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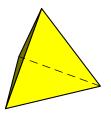


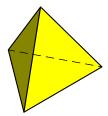




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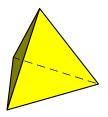
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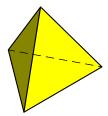






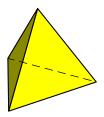
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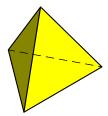






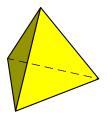
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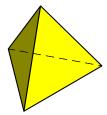






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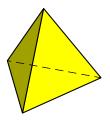


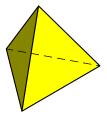




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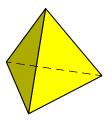


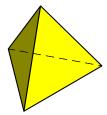




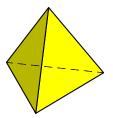
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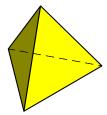
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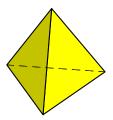


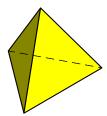




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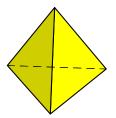
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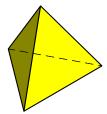




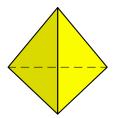
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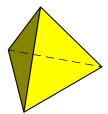
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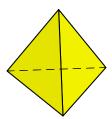


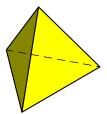




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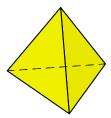
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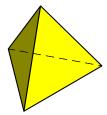






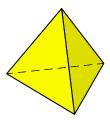
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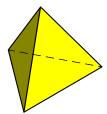






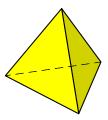
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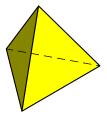






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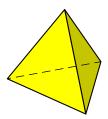






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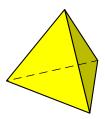


Family symmetries

Non-Abelian discrete flavor symmetries

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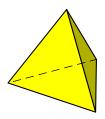
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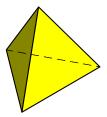




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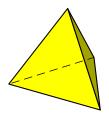
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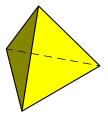






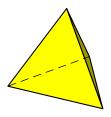
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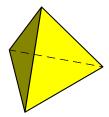






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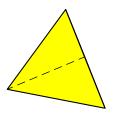


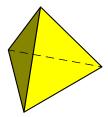




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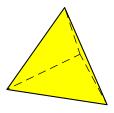
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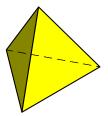




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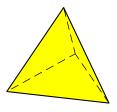
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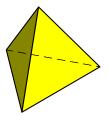




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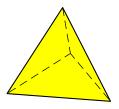
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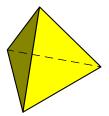






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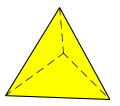




Family symmetries



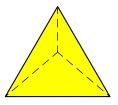
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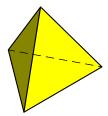






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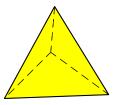






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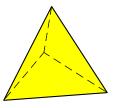
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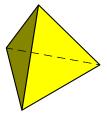






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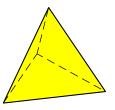




Family symmetries

Non-Abelian discrete flavor symmetries

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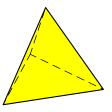


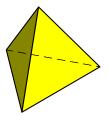




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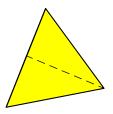
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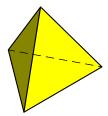




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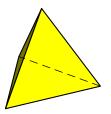
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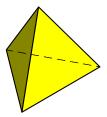






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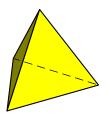


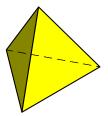




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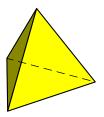
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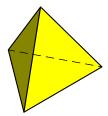




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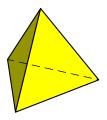
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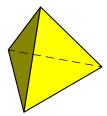




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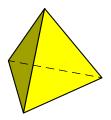
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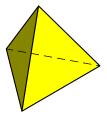






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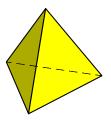


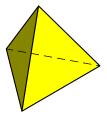




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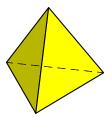
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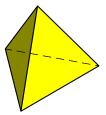




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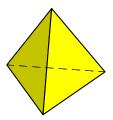
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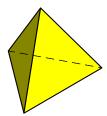




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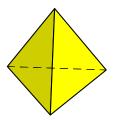
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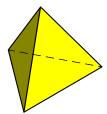




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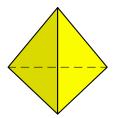
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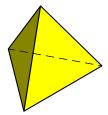




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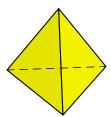


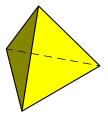




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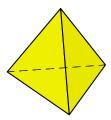
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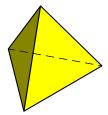




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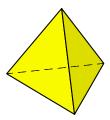


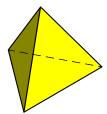




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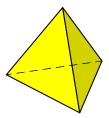


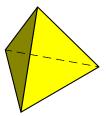




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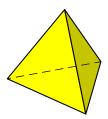
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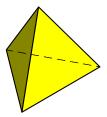






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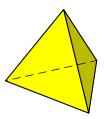






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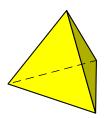
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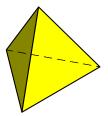




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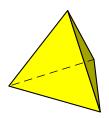
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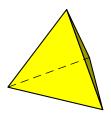


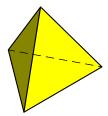




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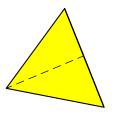
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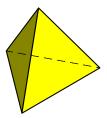




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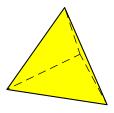


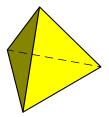




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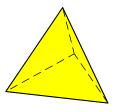
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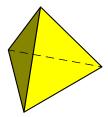






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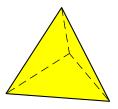






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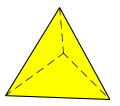






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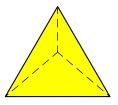
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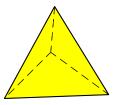
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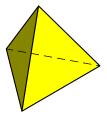




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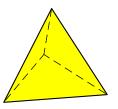


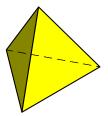
Family symmetries

Non-Abelian discrete flavor symmetries

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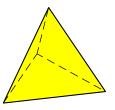


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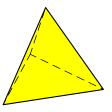


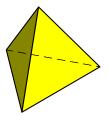




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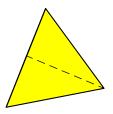
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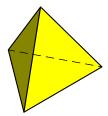




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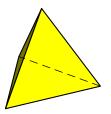


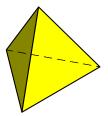




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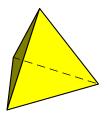
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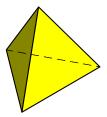




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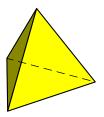


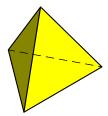






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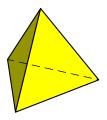


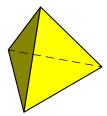




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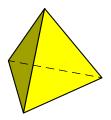
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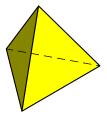






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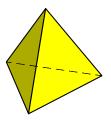


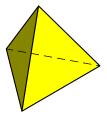




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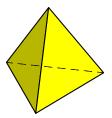
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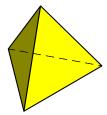




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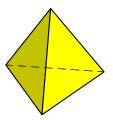
- Recently strong model building activities based on discrete family symmetry groups:
 - A_4 (tetrahedron)

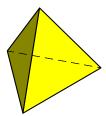




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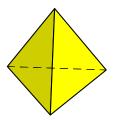
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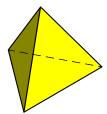




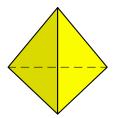
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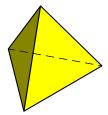
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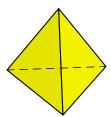


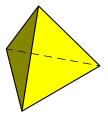




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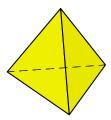
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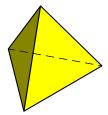




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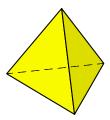
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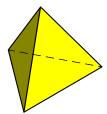






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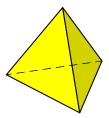


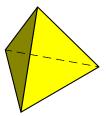




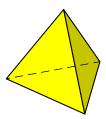
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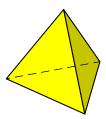
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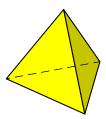
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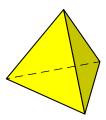


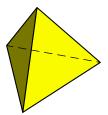




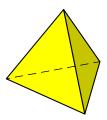


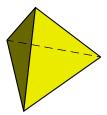
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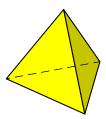


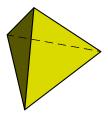




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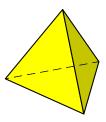
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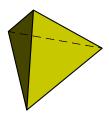






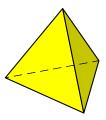
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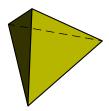




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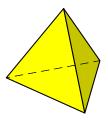
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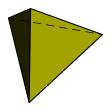




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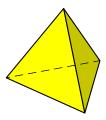
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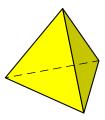
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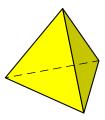
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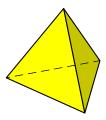
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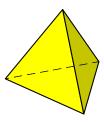
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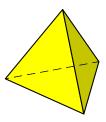
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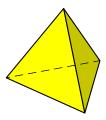
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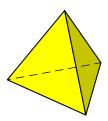






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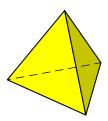
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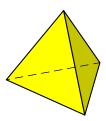
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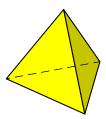
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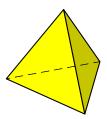
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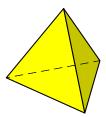
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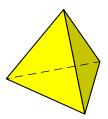






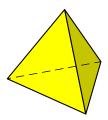
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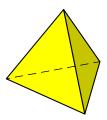






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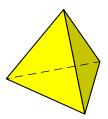
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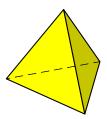
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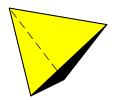
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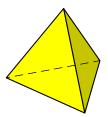
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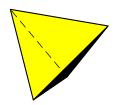






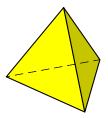
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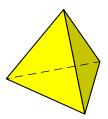






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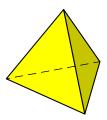
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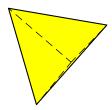






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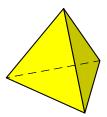


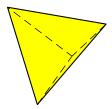


Family symmetries

Non-Abelian discrete flavor symmetries

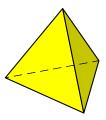
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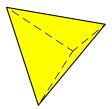




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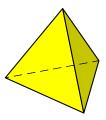
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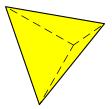




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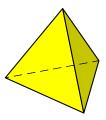
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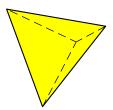




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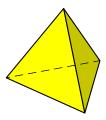
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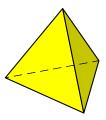
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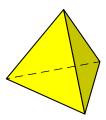


Family symmetries

Non-Abelian discrete flavor symmetries

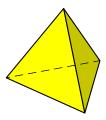
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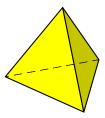


Family symmetries

Non-Abelian discrete flavor symmetries



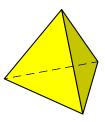
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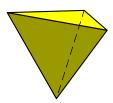






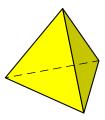
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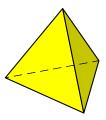
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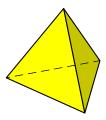
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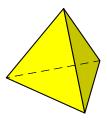
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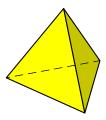
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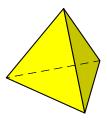
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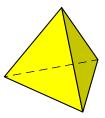
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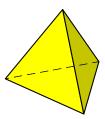
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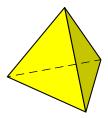


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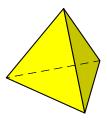
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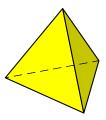




Family symmetries

Non-Abelian discrete flavor symmetries

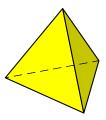
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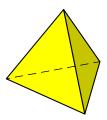
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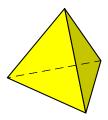


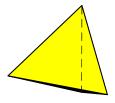
- Recently strong model building activities based on discrete family symmetry groups:
 - A_4 (tetrahedron)





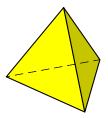
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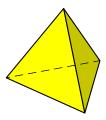
Family symmetries

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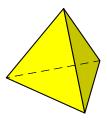


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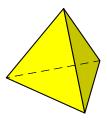


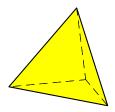






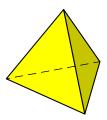
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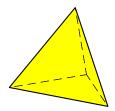






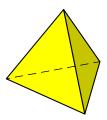
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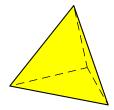






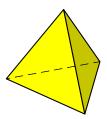
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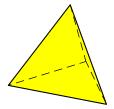






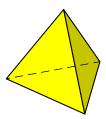
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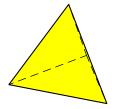






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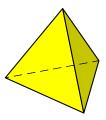


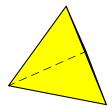




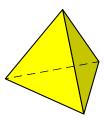


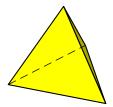
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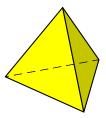
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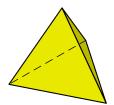






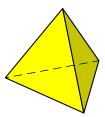
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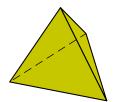




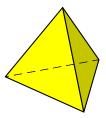


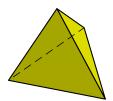
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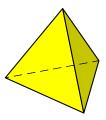
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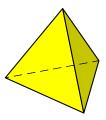
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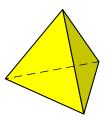
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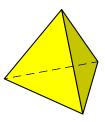
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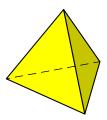
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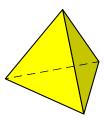
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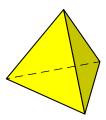


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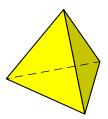








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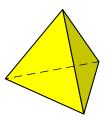








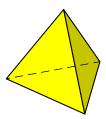
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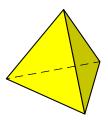
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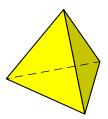
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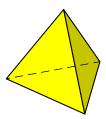
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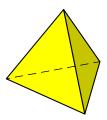
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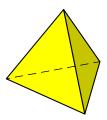






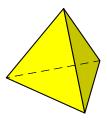


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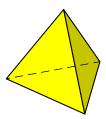
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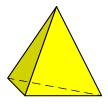






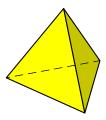
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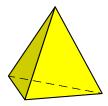






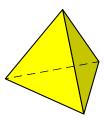
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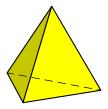






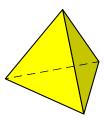
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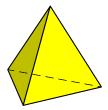




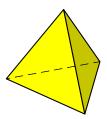


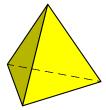
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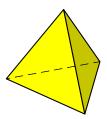
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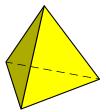




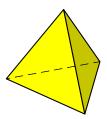


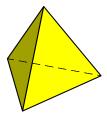
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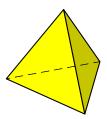
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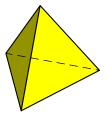




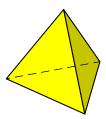


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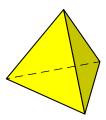


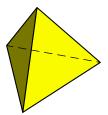




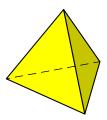


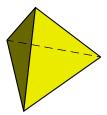
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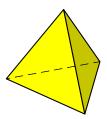


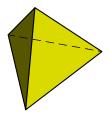




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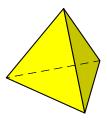
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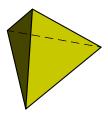






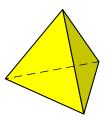
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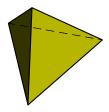




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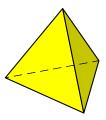


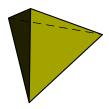


Family symmetries

Non-Abelian discrete flavor symmetries

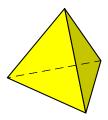
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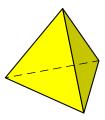
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- Recently strong model building activities based on discrete family symmetry groups:
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- Recently strong model building activities based on discrete family symmetry groups:
 - A_4 (tetrahedron)
 - T' (double tetrahedron)

Family symmetries



- Recently strong model building activities based on discrete family symmetry groups:
 - A_4 (tetrahedron)

 - S_3 (equilateral triangle)

Family symmetries Non-Abelian discrete flavor symmetries



- Recently strong model building activities based on discrete family symmetry groups:
 - A_4 (tetrahedron)
 - T' (double tetrahedron)
 - S₃ (equilateral triangle)
 - S_4 (octahedron, cube)



- Recently strong model building activities based on discrete family symmetry groups:
 - A_4 (tetrahedron)

 - S_3 (equilateral triangle)
 - S_4 (octahedron, cube)
 - A_5 (icosahedron, dodecahedron)



- Recently strong model building activities based on discrete family symmetry groups:
 - A_4 (tetrahedron)

 - S_3 (equilateral triangle)
 - S_4 (octahedron, cube)
 - A_5 (icosahedron, dodecahedron)
 - Q₄



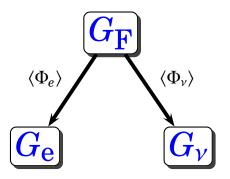
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 - . . .

Interplay between the symmetry breaking patterns in two sectors lead to lepton mixing (BM, TBM, . . .)

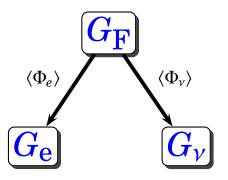
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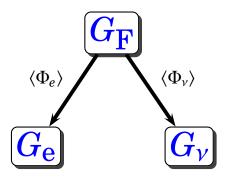
Some sub-sector(s) may preserve different residual symmetry

- Interplay between the symmetry breaking patterns in two sectors lead to lepton mixing (BM, TBM, . . .)
- Symmetry breaking achieved through flavon VEVs



- Some sub-sector(s) may preserve different residual symmetry
- However, the full Lagrangean does not have these residual symmetries

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- Symmetry breaking achieved through flavon VEVs

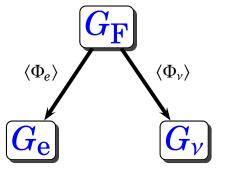


Some sub-sector(s) may preserve different residual symmetry

However, the full Lagrangean does not have these residual symmetries

Quantum corrections

- Interplay between the symmetry breaking patterns in two sectors lead to lepton mixing (BM, TBM, ...)
- Symmetry breaking achieved through flavon VEVs



- Some sub-sector(s) may preserve different residual symmetry
- However, the full Lagrangean does not have these residual symmetries
- Quantum corrections

question:

how robust are such predictions?

Altarelli and Feruglio (2005)

Superpotential couplings

$$\mathcal{W}_{v} = \frac{\lambda_{1}}{\Lambda \Lambda_{v}} \left\{ \left[\left(LH_{u} \right) \times \left(LH_{u} \right) \right]_{3} \times \Phi_{v} \right\}_{1} + \frac{\lambda_{2}}{\Lambda \Lambda_{v}} \left[\left(LH_{u} \right) \times \left(LH_{u} \right) \right]_{1} \xi$$

$$\begin{array}{c} \text{left-handed} \\ \text{lepton doublets} \\ \text{transform as } A_{4} \text{ triplet} \\ L = \left(L_{e}, L_{\mu}, L_{\tau} \right)^{T} \end{array} \right.$$

Altarelli and Feruglio (2005)

Superpotential couplings

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$$\text{cut-off} \qquad \text{see-saw} \qquad \text{triplet} \qquad \text{singlet} \qquad \text{flavon}$$

Attarelli and Feruglio (2005)

Superpotential couplings

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$$\begin{array}{c} \text{triplet} \\ \text{contraction} \end{array}$$

Altarelli and Feruglio (2005)

Superpotential couplings

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Altarelli and Feruglio (2005)

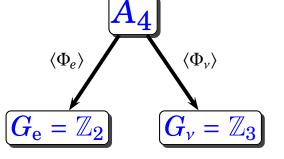
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A₄ symmetry broken by VEVs of flavons

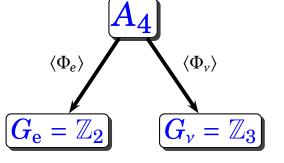
$$\langle \Phi_{v} \rangle = (v, v, v)$$

 $\langle \Phi_{e} \rangle = (v', 0, 0)$

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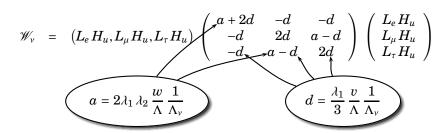
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Tri-bi-maximal mixing (TBM)

Structure lepton masses

After inserting the flavon VEVs

Altarelli and Feruglio (2005)



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After inserting the flavon VEVs

$$\begin{split} \mathscr{W}_{\nu} & = & \left(L_{e} H_{u}, L_{\mu} H_{u}, L_{\tau} H_{u} \right) \left(\begin{array}{ccc} a + 2d & -d & -d \\ -d & 2d & a - d \\ -d & a - d & 2d \end{array} \right) \left(\begin{array}{c} L_{e} H_{u} \\ L_{\mu} H_{u} \\ L_{\tau} H_{u} \end{array} \right) \\ \mathscr{W}_{e} & = & \left(L_{e}, L_{\mu}, L_{\tau} \right) \left(\begin{array}{c} y_{e} & 0 & 0 \\ 0 & y_{\mu} & 0 \\ 0 & 0 & v_{\tau} \end{array} \right) \left(\begin{array}{c} e_{\mathrm{R}} \\ \mu_{\mathrm{R}} \\ \tau_{\mathrm{R}} \end{array} \right) H_{d} \end{aligned}$$

After inserting the electroweak VEVs

$$\mathscr{W}_{\scriptscriptstyle V} \stackrel{H_u o (0, v_u)^T}{\longrightarrow} \stackrel{v_u^2}{\longrightarrow} \left(\begin{matrix} v_e, v_\mu, v_\tau \end{matrix} \right) \left(egin{array}{ccc} a + 2d & -d & -d \\ -d & 2d & a -d \\ -d & a -d & 2d \end{array} \right) \left(egin{array}{c} v_e \\ v_\mu \\ v_\tau \end{array} \right)$$

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Harrison et al. (2002)

Structure of neutrino masses (in the basis in which the charged lepton masses are diagonal)

$$m_{\scriptscriptstyle Y} \propto \left(egin{array}{ccc} a+2d & -d & -d \ -d & 2d & a-d \ -d & a-d & 2d \end{array}
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Tri-bi-maximal (P)MNS

mixing matrix
$$U_{(\mathrm{P)MNS}}^{\mathrm{TBM}} = \begin{pmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix}$$
 • Mixing angles:
$$\begin{cases} \theta_{12} \simeq 35^{\circ} \\ \theta_{13} = 0 \\ \theta_{23} = 45^{\circ} \end{cases}$$
 • δ undefined for $\theta_{13} = 0$

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Unrealistic <a>¬`corrections' required

Many analyses: include high order terms in holomorphic superpotential

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- However: possible to construct models where higher order holomorphic superpotential terms vanish to all orders
- Also popular: contribution from right-handed sector (may be determined by symmetries as well)
- ? Conceptual question: how predictive are such models?

Superpotential: holomorphic, e.g.

$$\mathscr{W}_{v} = \frac{1}{2} \left(L H_{u} \right)^{T} \kappa_{v} L H_{u}$$

e.g. Leurer et al. (1994)

Superpotential: holomorphic, e.g.

$$\mathscr{W}_{v} = \frac{1}{2} \left(L H_{u} \right)^{T} \kappa_{v} L H_{u}$$

Kähler potential: non-holomorphic (real analytic)

$$K = K_{\text{canonical}} + \Delta K$$

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Canonical K\u00e4hler potential

$$K_{
m canonical} \supset \sum_f \left[\left(L_f
ight)^\dagger L_f + \left(R_f
ight)^\dagger R_f
ight]$$
 charged lepton singlets $R = (e_{
m R}, \mu_{
m R}, au_{
m R})$

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ight)^{\dagger} L_{f} + \left(R_{f}
ight)^{\dagger} R_{f} \right]$$

Correction

$$\Delta K = \sum_{f,g} \left[L_f^{\dagger} P_{fg} L_g + R_f^{\dagger} Q_{fg} R_g \right]$$

rightharpoonup Consider infinitesimal change parametrized by x

$$\Delta \mathcal{K}_L = -2xP$$

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Rotate to canonically normalized fields

$$L' \rightarrow L \simeq (1-xP) L'$$

extstyle ext

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Rotate to canonically normalized fields

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Corrections to leads to a change

$$\mathcal{W}_{v} = \frac{1}{2} (L' \cdot H_{u})^{T} \kappa_{v} L' \cdot H_{u}$$

$$\simeq \frac{1}{2} [(\mathbb{1} + xP) L \cdot H_{u}]^{T} \kappa_{v} [(\mathbb{1} + xP) L \cdot H_{u}]$$

$$\simeq \frac{1}{2} (L \cdot H_{u})^{T} \kappa_{v} L \cdot H_{u} + x (L \cdot H_{u})^{T} (P^{T} \kappa_{v} + \kappa_{v} P) L \cdot H_{u}$$

extstyle ext

$$\Delta \mathcal{K}_L = -2xP$$

Rotate to canonically normalized fields

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Corrections to leads to a change

$$\mathcal{W}_{\nu} \simeq \frac{1}{2} (L \cdot H_u)^T \kappa_{\nu} L \cdot H_u + x (L \cdot H_u)^T (P^T \kappa_{\nu} + \kappa_{\nu} P) L \cdot H_u$$

Differential equation

$$\frac{d}{dx}\kappa_{\nu}(x) = P^{T}\kappa_{\nu}(x) + \kappa_{\nu}(x)P$$

Kähler vs. RG corrections

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Same structure as RG evolution of neutrino mass operator

cf. Antusch et al. (2003)

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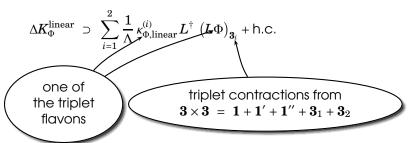
Same structure as RG evolution of neutrino mass operator

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 However, size of K\u00e4hler corrections can be substantially larger (no loop suppression)

Back to the A_4 example

Kähler potential may contain



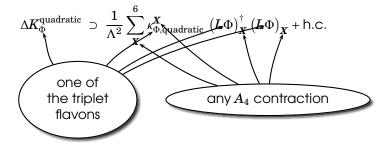
Back to the A_4 example

Kähler potential may contain

$$\Delta K_{\Phi}^{
m linear} \supset \sum_{i=1}^2 rac{1}{\Lambda} \kappa_{\Phi,
m linear}^{(i)} L^{\dagger} (L\Phi)_{{f 3}_i} + {
m h.c.}$$

However, such terms may be forbidden by additional symmetries

"Quadratic" K\u00e4hler corrections



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$$\Delta K_{\Phi}^{\mathrm{quadratic}} \;\supset\; \frac{1}{\Lambda^2} \sum_{\boldsymbol{X}}^{6} \kappa_{\Phi,\mathrm{quadratic}}^{\boldsymbol{X}} \; \left(L\Phi\right)_{\boldsymbol{X}}^{\dagger} \left(L\Phi\right)_{\boldsymbol{X}} + \mathrm{h.c.}$$

 Such terms cannot be forbidden by any (conventional) symmetry

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- Such terms cannot be forbidden by any (conventional) symmetry
- → Kähler corrections when flavon fields attain their VEVs
- riangleq Additional parameters κ_{Φ}^{X} reduce the predictivity of the scheme

Linear independent flavon corrections

From $\langle \Phi_e \rangle$

$$P_{\rm II} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, P_{\rm II} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$P_{\rm III} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

rianglerightarrow From $\langle \Phi_{\nu} \rangle$

$$P_{\text{IV}} = \left(egin{array}{ccc} 0 & 1 & 1 \ 1 & 0 & 1 \ 1 & 1 & 0 \end{array}
ight) \, , \quad P_{\text{V}} = \left(egin{array}{ccc} 0 & \mathsf{i} & -\mathsf{i} \ -\mathsf{i} & 0 & \mathsf{i} \ \mathsf{i} & -\mathsf{i} & 0 \end{array}
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- ilder Consider change induced by $P_{
 m V}$ correction
- rightharpoonup K Kähler metric of the form $\mathcal{K}_L = 1 2xP$ with

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The analytic formula evaluated at tri–bi–maximal mixing reads ($m_e \ll m_\mu \ll m_ au$)

$$\Delta\theta_{13} = \kappa_{\rm V} \cdot \frac{v^2}{\Lambda^2} \cdot 3\sqrt{\frac{3}{2}} \left(\frac{2m_1}{m_1 + m_3} + \frac{m_e^2}{m_\mu^2 - m_e^2} + \frac{m_e^2}{m_\tau^2 - m_e^2} \right)$$

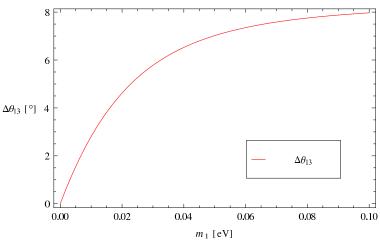
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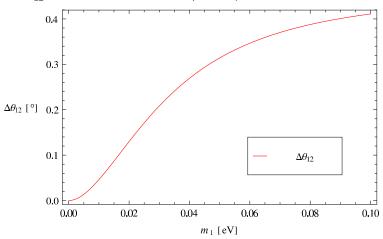
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 $rightharpoonup \Delta heta_{13}$ for Kähler coefficient $\kappa_{
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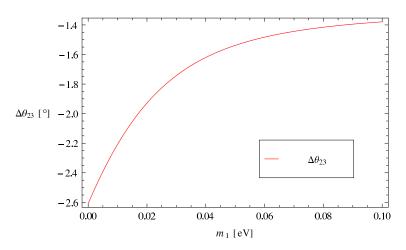
For comparison: change of θ_{12}

 $\Rightarrow \Delta\theta_{12}$ for Kähler coefficient $\kappa_{\rm V}=1$, $v/\Lambda=0.2$



For comparison: change of θ_{23}

 $rightharpoonup \Delta heta_{23}$ for Kähler coefficient $\kappa_{
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e.g. Antusch et al. (2003)

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Araki (2007) ; Araki et al. (2008)

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Araki (2007); Araki et al. (2008)

However, such corrections are exponentially suppressed

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 - ... expected to dominate

Kähler corrections induced by flavon VEVs

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- ilder Example: realistic $heta_{13}$ from tri-bi-maximal mixing scheme

Large uncertainties in a large class of popular constructions

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- In certain schemes (such as string compactifications) one may compute the coefficients of the higher order terms in the K\u00e4hler potential
- More effort both on theoretical and experimental side required to attack the flavor problem

Thank you very much!

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